

Coherent Bremsstrahlung Beam at MAX-lab Facility

The MAX-Tagg Collaboration

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Some History

Firstly the Coherent Bremsstrahlung (CB) beams were produced in Frascati by Diambri and his coworkers about **50 years ago (1960)**.

Since that time the Coherent Bremsstrahlung beams were created practically on all electron accelerators with energies $E_0 \sim > 1 \text{ GeV}$ (**SLAC, Erevan, Charkov, Tomsk, Tokyo**) and widely used for polarized photonuclear investigations.

At present the CB beams are successfully used for photonuclear experiments at **Mainz, Bonn, Jlab**.

principal possibility to produce the CB beam for
nuclear researches at MAX-lab was appeared
due to modernization of the MAX-lab facility
increasing maximal energy of the electron beam
 $E_0 \sim 200-250$ MeV.

Calculations of the CB beam expected parameters
for MAX-lab conditions have been produced in
2000-2001 years (*J-O.Adler, V.Ganenko,
Lindgren, Report 01/01 LUNF
/FFR-3086)1-31/2001, Lund 2001*)

The proposition was supported by Program
Advisory Committee of the MAX-lab in the 2001
OI-15 for NP-experiments at MAX-lab
*B.Ganenko et al. A coherent bremsstrahlung
beam at the MAX-lab facility. Lund, May 30
2002*).

When it was required *5 years* in order that to produce necessary equipment and develop experimental methods for the project realization.

007 of September - the first test experiment on CB beam production. It was performed:

- the goniometer placing on the beam line preliminary adjustment;
- the crystals orientation;
- preliminary CB spectra measurements.

008 of April - the second test experiment:

- checking of the crystals orientation;
- the CB spectra measurements in more wide energy range and orientations;
- measurements of the channeling radiation;
- test measurements of the deuteron disintegration by polarized photon at $E_\gamma \sim 50-60\text{MeV}$.

Production of the CB beam

The MAX-lab Facility

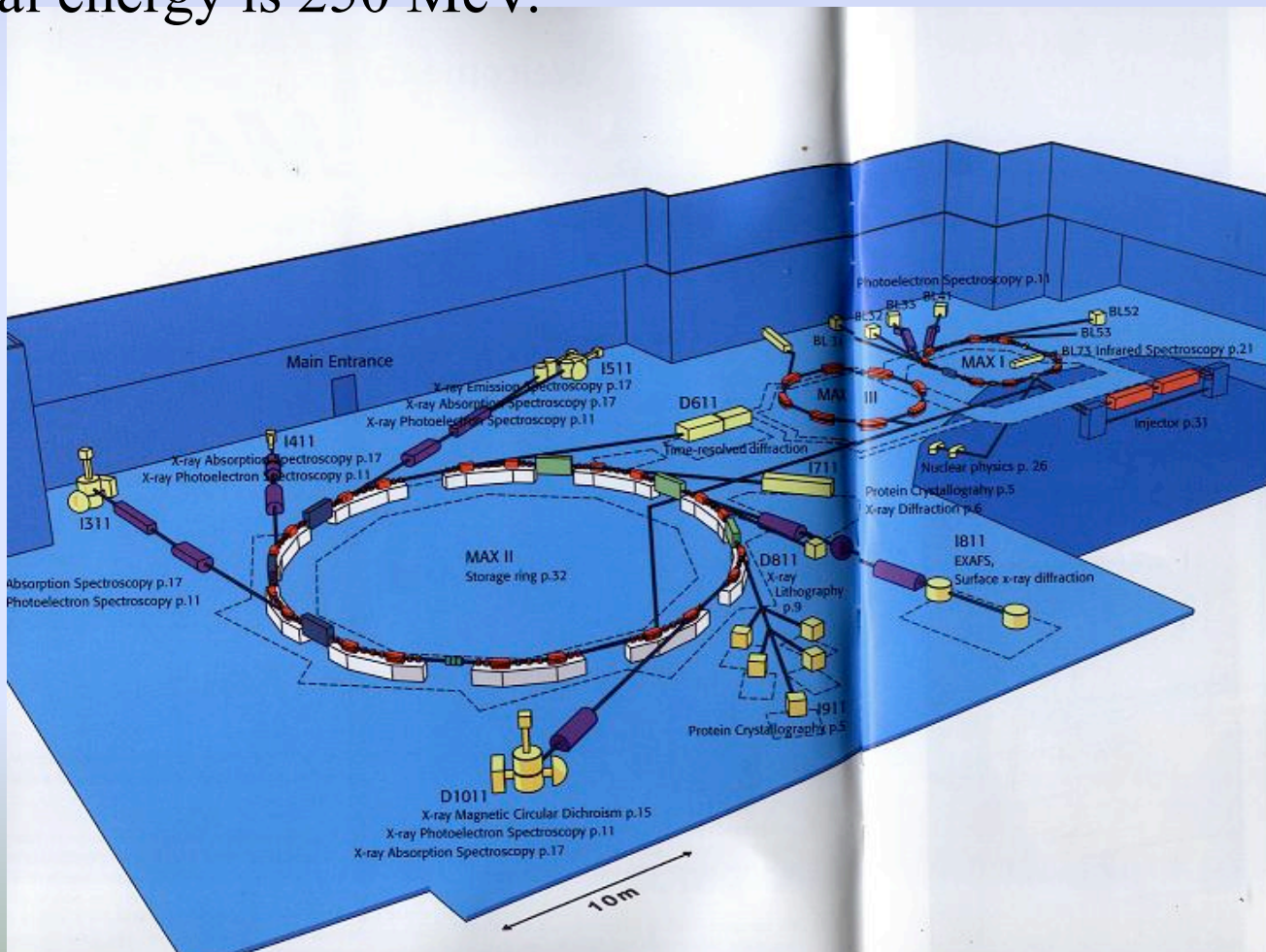
MAX-lab is national laboratory of Sweden. It is center of synchrotron radiation investigations. There are three storage rings at present time

MAX-I - energy 550 MeV.

MAX-II - energy 1.5 GeV.

MAX-III - energy 700 MeV.

There is beam line for photonuclear researches. The beam comes from storage ring MAX-I, which in this case works as stretcher. Final energy is 250 MeV.



Expected CB beam parameters at MAX-lab

Calculations of the expected CB parameters have been made for expected MAX-lab electron beam parameters
-in 2001 (old code);

-repeated in the 2004-2007 years using the code
developed by Peter Grabmayr.

J.-O.Adler, V.Ganenko, L.-J.Lindgren. *Report 01/01 LUN*
(NFFR-3086)1-31/2001, Lund 2001

V.Ganenko, G.Vashchenko, D.Burdeinii. *The expected polarized
photon beam parameters at MAX_lab facility // MA*
Activity Report 2005-2006, p.454

V.Ganenko, G.Vashchenko, D.Burdeinii. *The linearly polarized
photon beams for photonuclear investigations at low energy
facilities // Journal of Kharkov University, No.781, 2007.Paper
series "Nuclei, Particles, Fields", issue 3/35/, p.88-92.*

Calculations have shown that it is possible to get a beam of
the polarized photons in the range $E_\gamma \sim 20-100$ MeV
use:

- diamond crystal;
- electron beam with energy of $E_0 \sim 200-250$ MeV;
- collimation of the radiation in the angle $\theta_c \sim (1 - \frac{1}{\gamma})$
($\theta_\gamma = mc^2/E_0$)

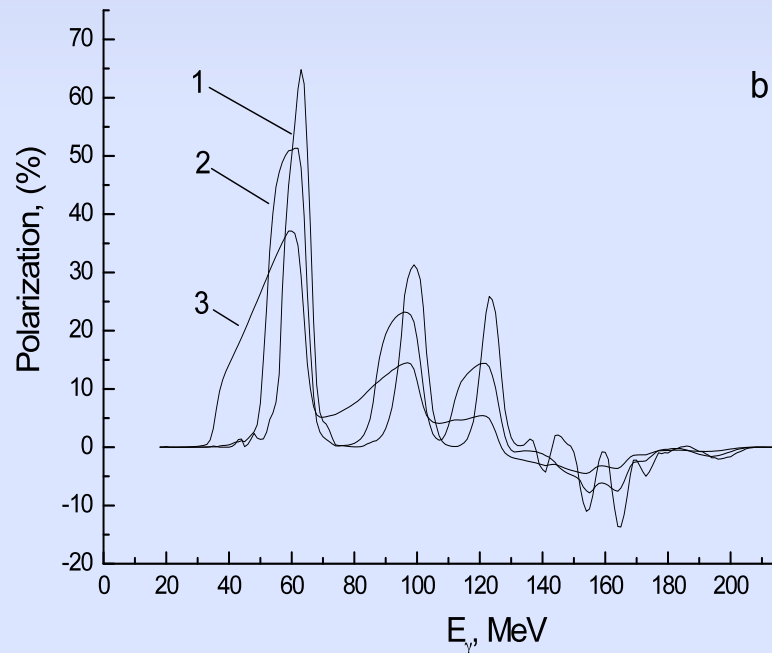
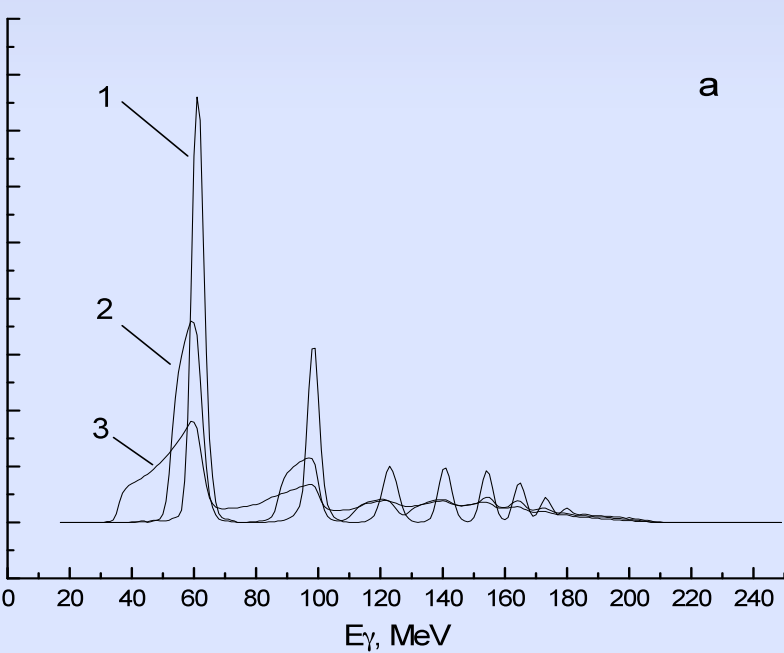
ere it is shown calculated spectra of the CB intensity and polarization for electron energy $E_0=250$ MeV and several collimation angles θ_c (in θ_γ units).

The intensity is presented as ratio

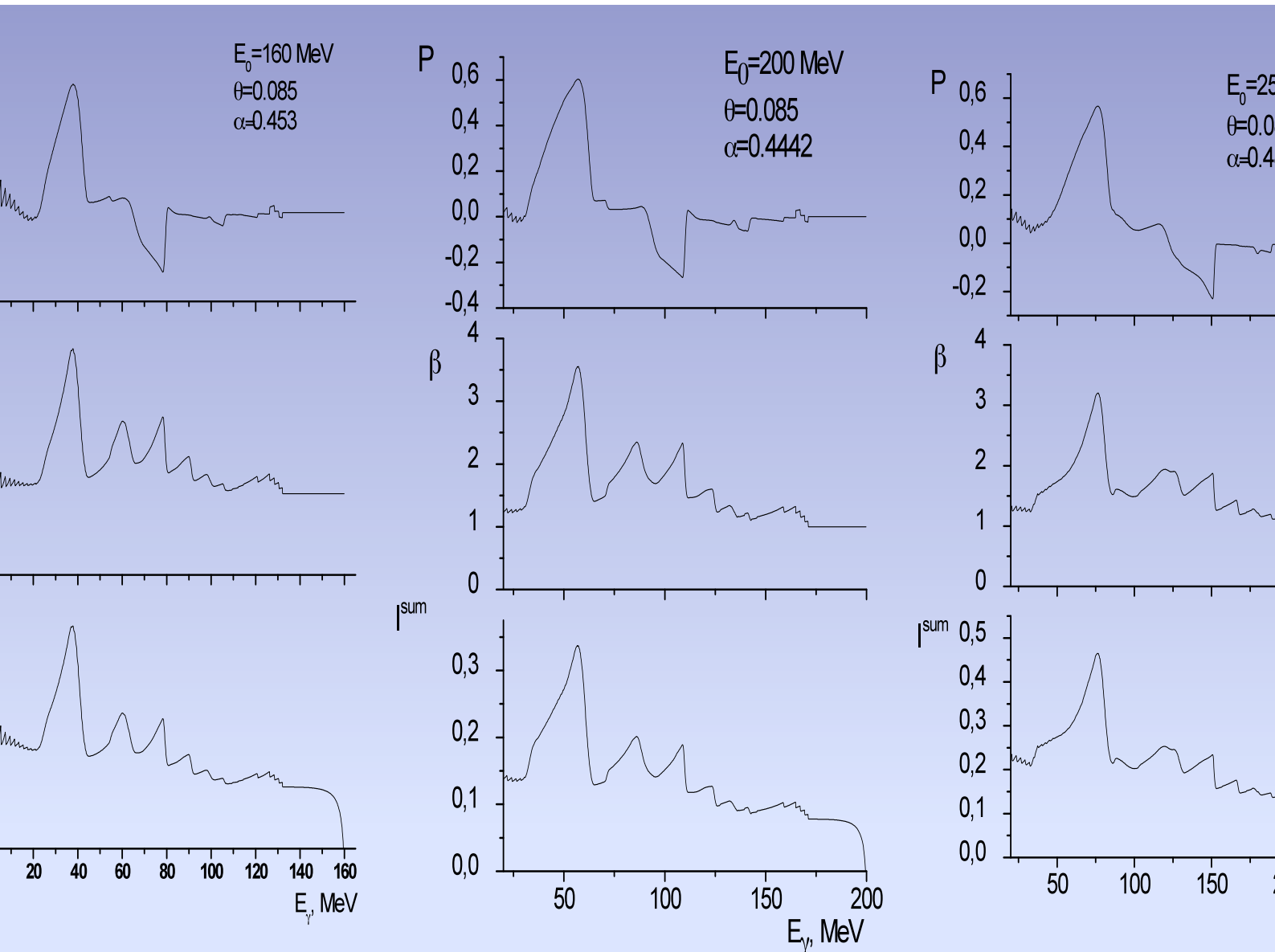
$$\beta = (I_{\text{coh}} + I_{\text{in}}) / I_{\text{in}}$$

where I_{coh} (I_{in}) is intensity of the coherent (incoherent) part of the CB.

As we can see, that intensity and polarization in the maximum are high enough ($\beta_{\text{max}} \sim 2 \dots 4.5$ and $P_{\gamma, \text{max}} \sim 30 \dots 60\%$) for the collimation angles θ_c in the interval $(0.25 - 1)\theta_\gamma$.



Intensity (a) and polarization (b) of the CB of electron and positron from a diamond crystal $L=0.1\text{mm}$ thick. $E_0=250$ MeV, $E_{\gamma,p}=60\text{MeV}$, collimation angles $\theta_c/\theta_\gamma=0.25(1), 0.5(2), 1(3)$.

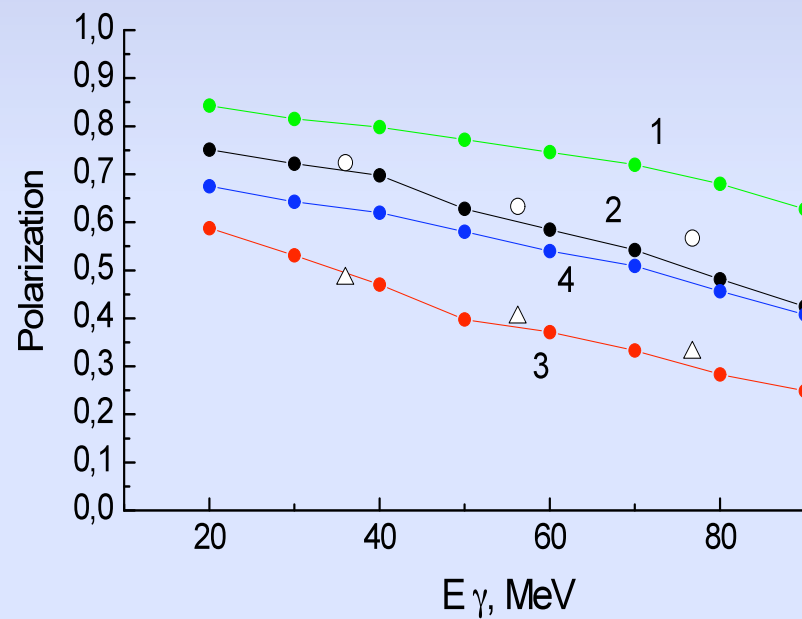
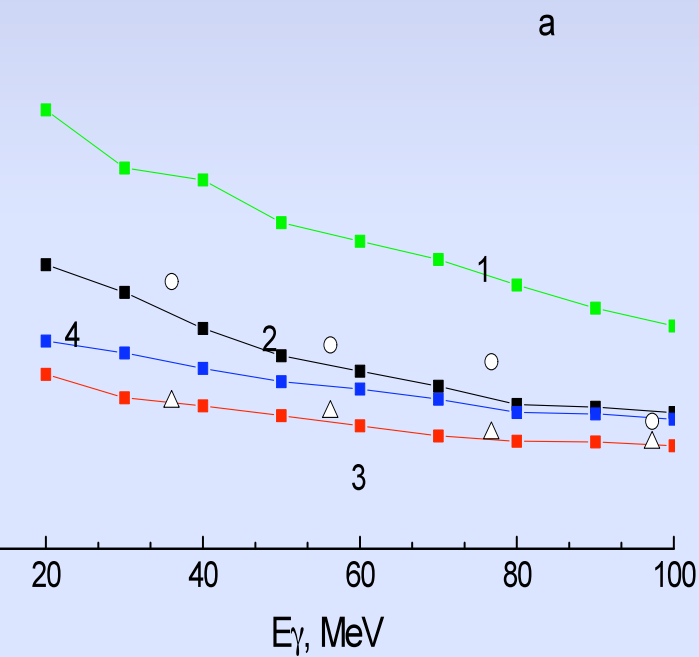


2. Polarization, coherent effect and intensity of the CB
electron energies $E_0=160 \text{ MeV}$ (collimation $\theta_c \sim 0.31\theta_\gamma$), 200 MeV ($0.36\theta_\gamma$) u 250 MeV ($\theta_c \sim 0.45\theta_\gamma$). Diamond is 0.1 mm thick.

These calculations have been made with the code developed by P. Grabmayr with colleagues. In this code experimental factors were taken into account more accurately. The calculations in a whole confirmed the results of the previous calculations.

the values of the CB beam parameters at which the beam is suitable for nuclear physics experiments are: $\beta_{\max} \geq 1.4$ and $P_{\gamma, \max} \geq 20\%$. Calculations have shown the energy ranges available for experiments:

E_0	$\theta_c \sim \theta_\gamma$	$\theta_c \sim 0.5\theta_\gamma$
150 MeV	up to 40 MeV	up to 60 MeV
200 MeV	up to 60 MeV	up to 80 MeV
250 MeV	up to 80 MeV	up to 100 MeV



Dependence of the coherent effect (a) and polarization (b) on the CB peak energy vs. the CB peak energy for $E_0=250\text{MeV}$, diamond thickness $L=0.1\text{mm}$.

*1- $\theta_c=0.25\theta_\gamma$; 2- $\theta_c=0.5\theta_\gamma$; 3- $\theta_c=\theta_\gamma$; 4- $\theta_c=0.5\theta_\gamma$ $L=0.3\text{mm}$.
 open circles - $\theta_c=0.46\theta_\gamma$, triangles - $\theta_c=0.91\theta_\gamma$.*

Experimental set up

The electron beam for photonuclear research is extracted from the storage ring MAX-I, which in this case works in a stretcher mode.

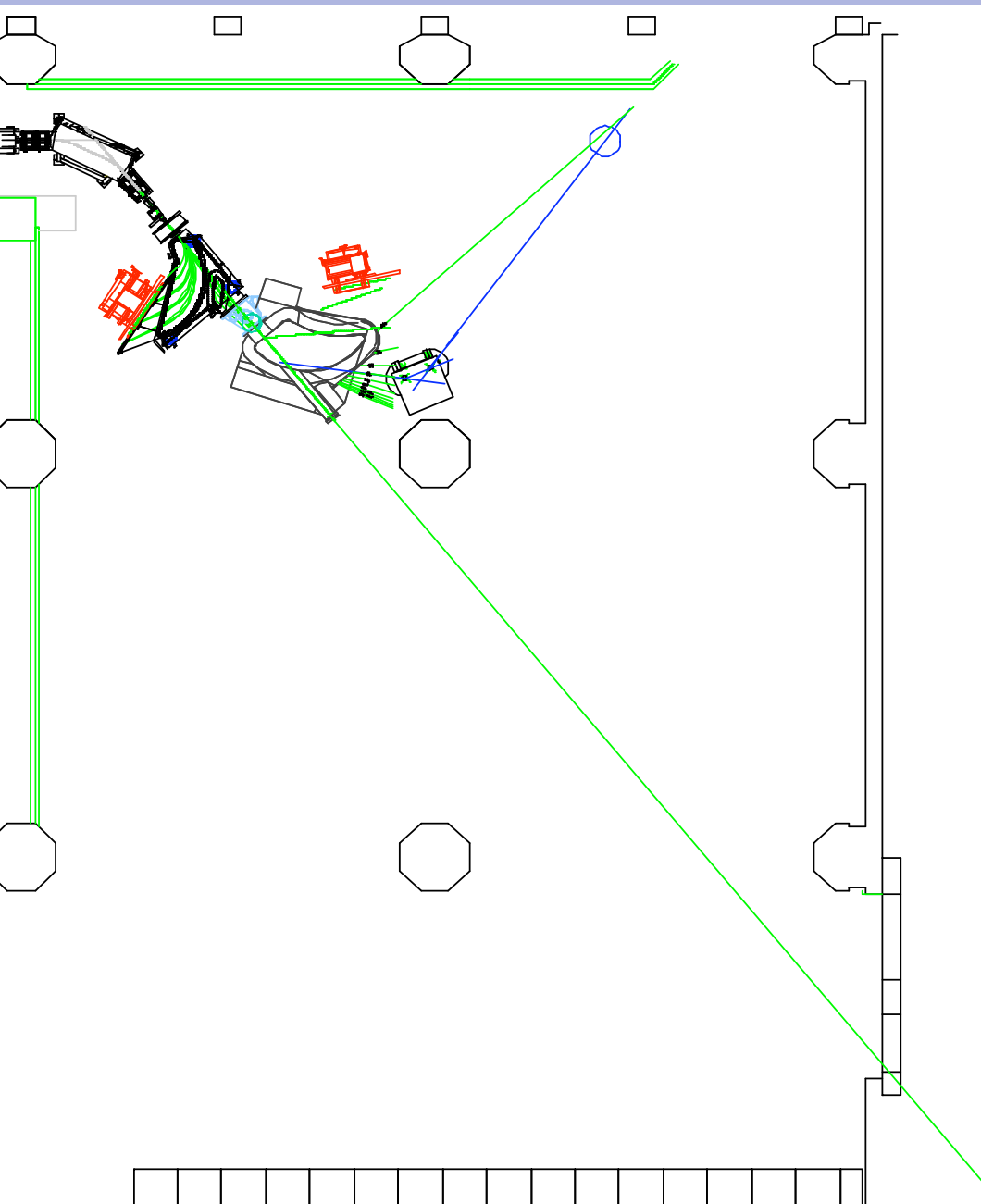


Fig.5. Scheme of MAX-lab beam line for photonuclear researches.

The electron beam comes into the experimental hall and is directed to the goniometer, which is fixed in the goniometer chamber. The goniometer is placed in a vacuum chamber between the magnets of the electron transport (ET) and main tagging systems.

After the passage of the photon radiators, the non-interacting part of the electron beam is directed to the beam dump. The beam dump consists of a paraday cup which measures the beam current.

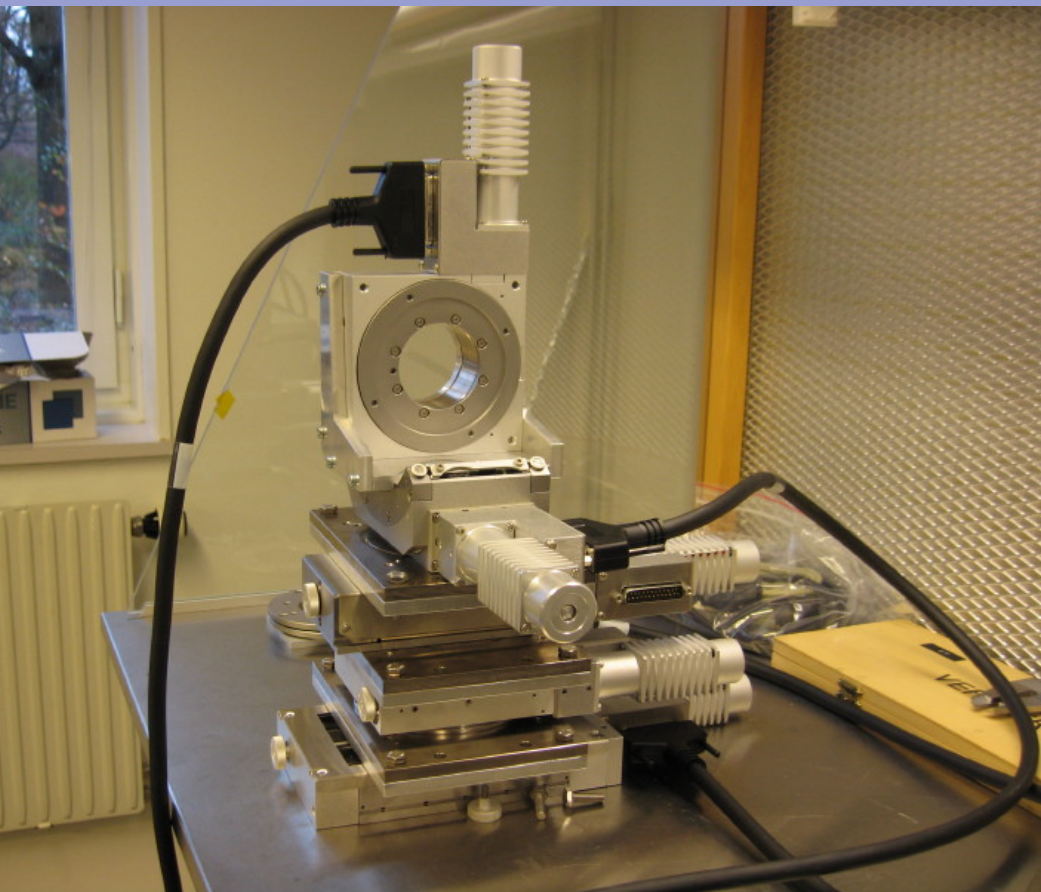
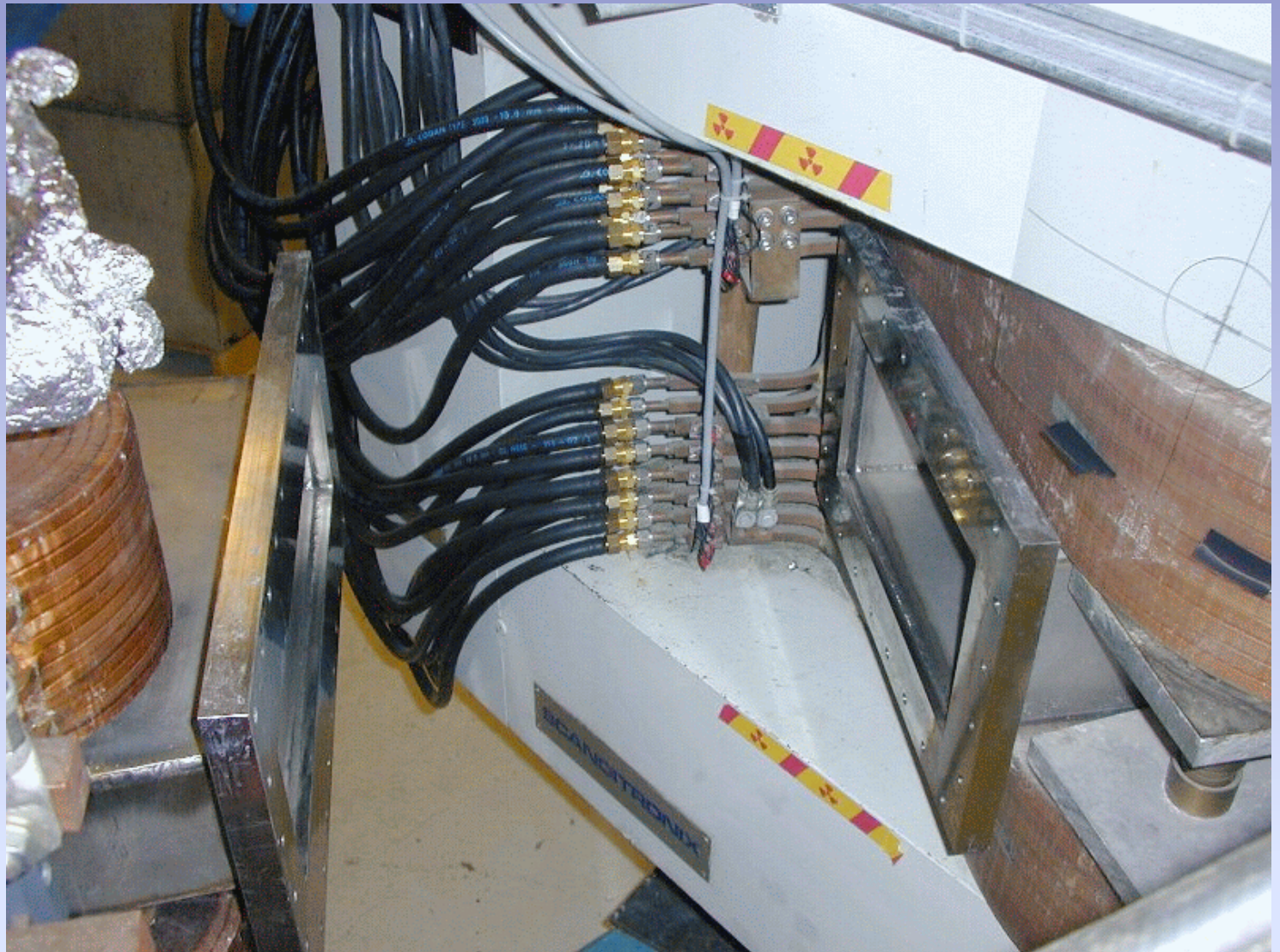


Fig.6. View of the multi-axis goniometer. In the prototype the M-4006 motion controller was used.

The goniometer consists of three rotation stages providing orientation of the crystal and two translation stages providing vertical adjustment of the target and the movement perpendicular to the neutron beam direction. Control of the goniometer is provided by the M-4006 motion controller connected to RS-232 port of the PC.

Table 2: Characteristics of the moving stages

Model, M-	Travel Range	Unidirectional Repeatability	Hysteresis	Accuracy	Resolution
RM100APEV6	360°	0.003°	0.006°	0.023°(abs)	0.001°
GM80PEV6	±45°	0.004°	0.02°	0.05°	0.001°
ZM80PE.1V6	4 mm	0.2 μm	3 μm	4 μm	0.1 μm
TM50PE1V6	50 mm	1.5 μm	3 μm	5 μm	1 μm

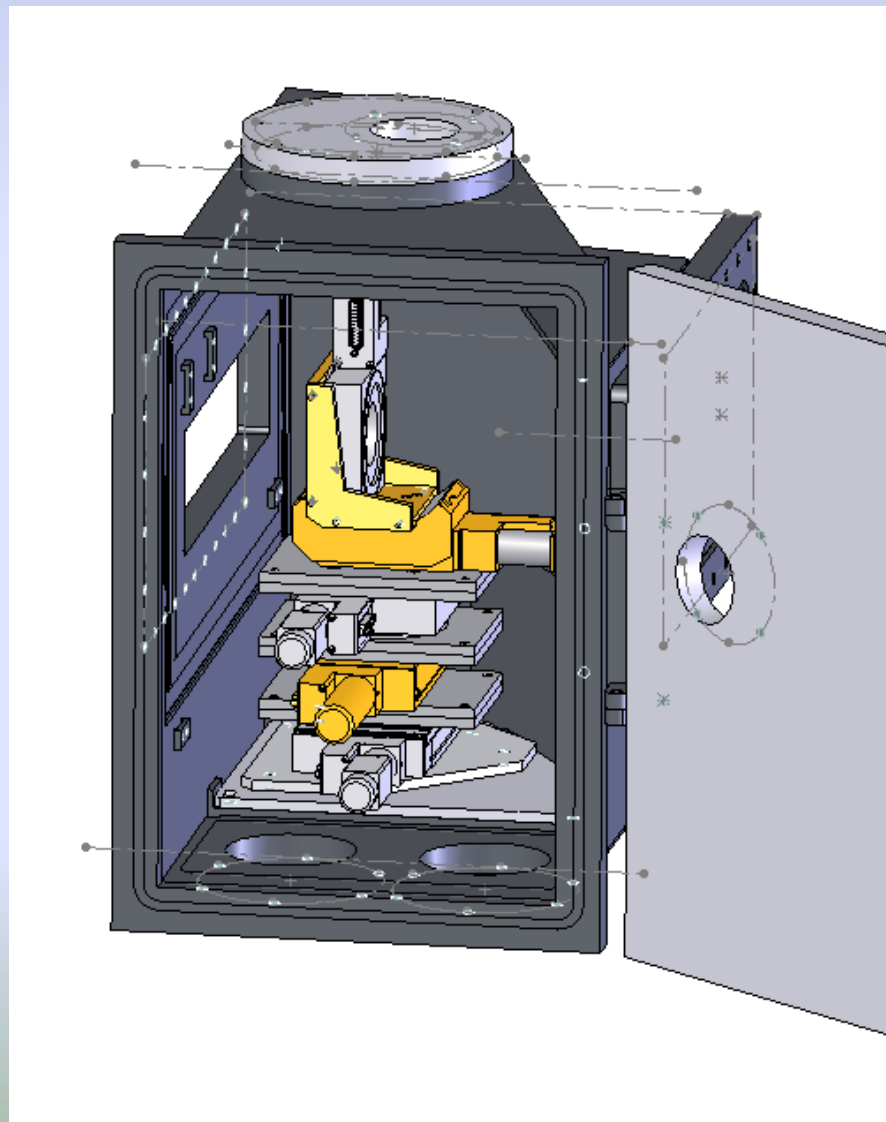


The goniometer was placed into vacuum chamber between MT magnets. The place is not very suitable for this purpose because the distance between flanges of these magnets is very small, about 100 mm. Secondly, the yoke of the MT magnet hangs over the front of the magnetic chamber. These reasons result to complicate for the design of the chamber, limit its sizes.

raft and working drawings were developed by Dmitriy Machev (from MAX-lab) with participation the Kharkov group. The chamber corps is made from (stainless) the not-magnetic material. It is of itself an empty parallelepiped with two truncated corners that repeats lines of the MT magnet and has size 526×677 mm.

The chamber was made in MAX-lab and placed on the beam line.

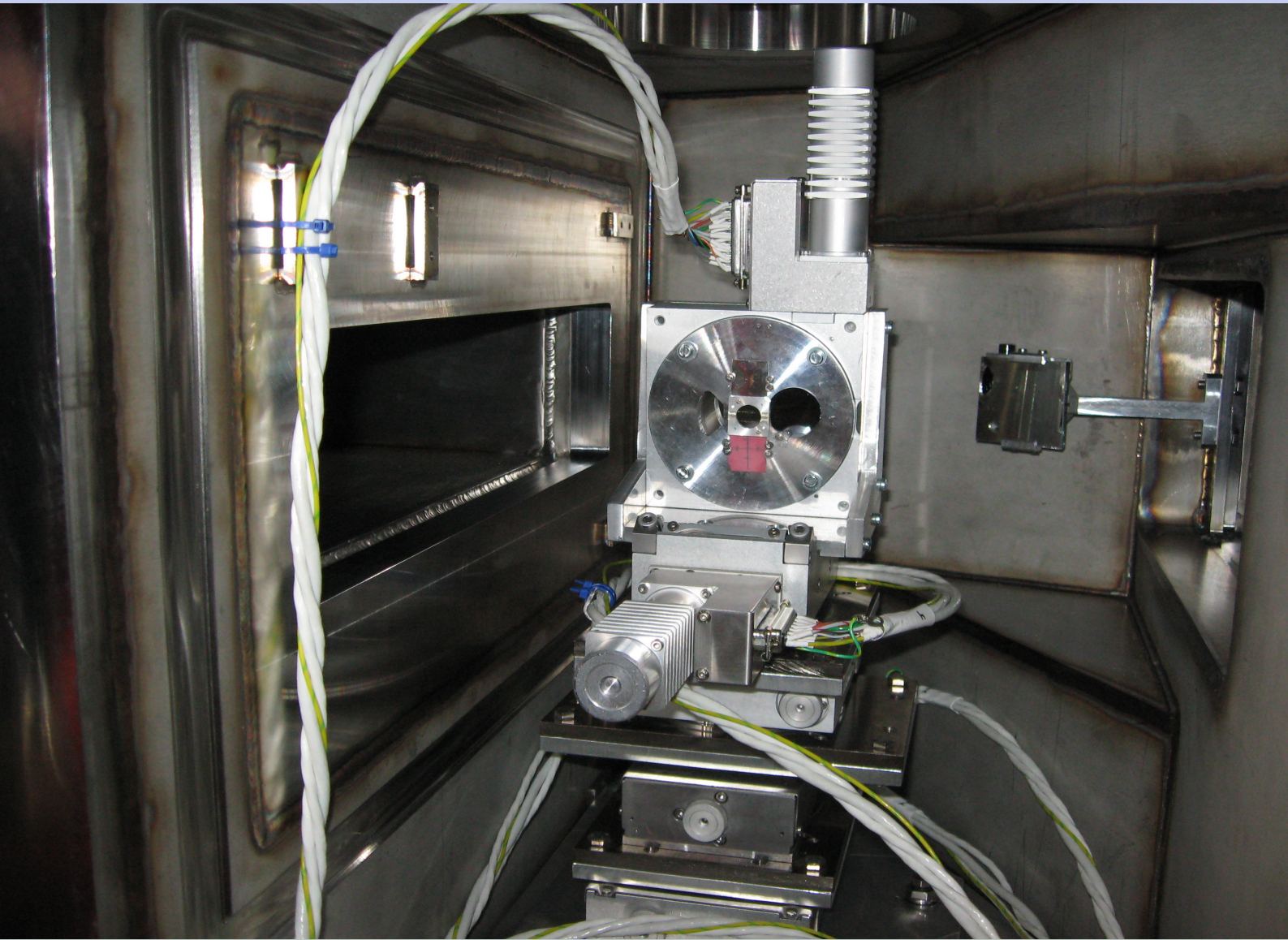
The model of the vacuum chamber. It is shown the position of the detector in the chamber.



Assembling of the goniometer on the beam line



rotating goniometer target holder has five positions for radiation
diamond crystal (100 μm) was fixed at center of the holder
100 μm crystal, Al radiator (50 μm), screen, empty slot
positioned on the circumference of the target holder.



Crystal orientation

The procedure of the crystal initial orientation includes finding angles of goniometer rotations, Φ_{v0} , Φ_{h0} , Φ_{a0} , along vertical, horizontal and azimuthal axes at which one of the crystal axes is directed along the axis of electron beam (along the impinging electrons \mathbf{P}_0), and the other two axes \mathbf{b}_2 and \mathbf{b}_3 are directed along the axes of goniometer rotations \mathbf{f}_v and \mathbf{f}_h . The diamond was oriented in such a way that axis $\langle 001 \rangle$ was perpendicular to the crystal plane.

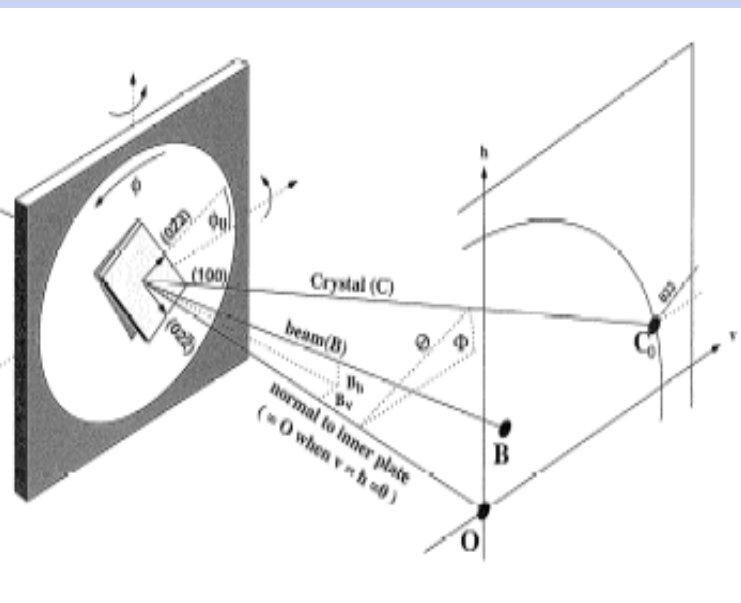
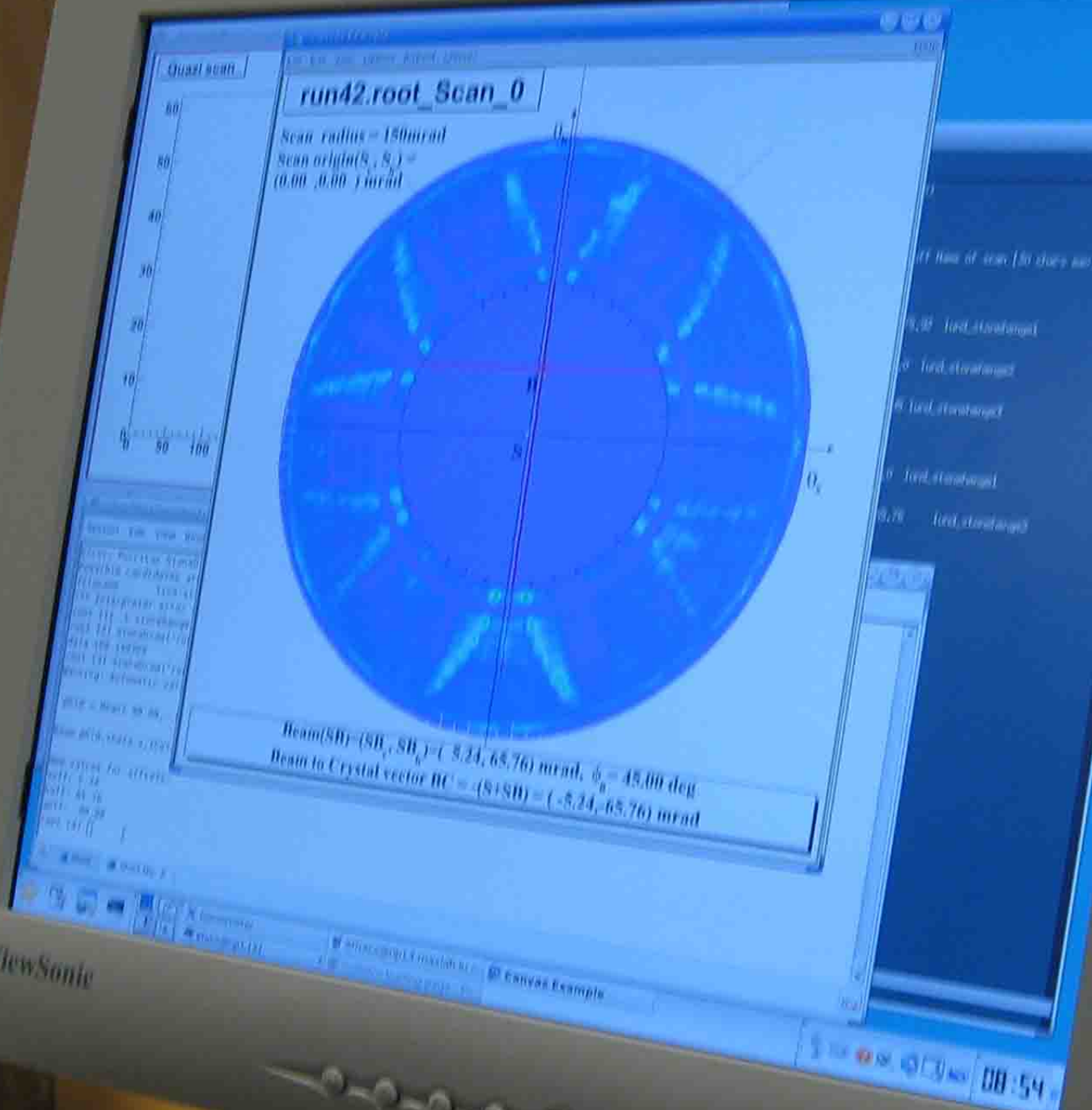


Fig. 9. Schematic of the goniometer and the definition of the angles. Figure courtesy Ken Livingston (Summary of the period 04.14-04.30.2008).

Two methods were used at MAX-lab, based on measurements of the orientation dependence of gamma radiation intensity.

“Stonehenge” technique (developed by Ken Livingston) was adapted for facilities with tagging systems. In this method it is the orientation dependence of the gamma radiation intensity as a function of the azimuthal rotation which was registered with the tagger hodoscope in a narrow energy interval ($\sim 20 \dots 80$ MeV).



ViewSonic

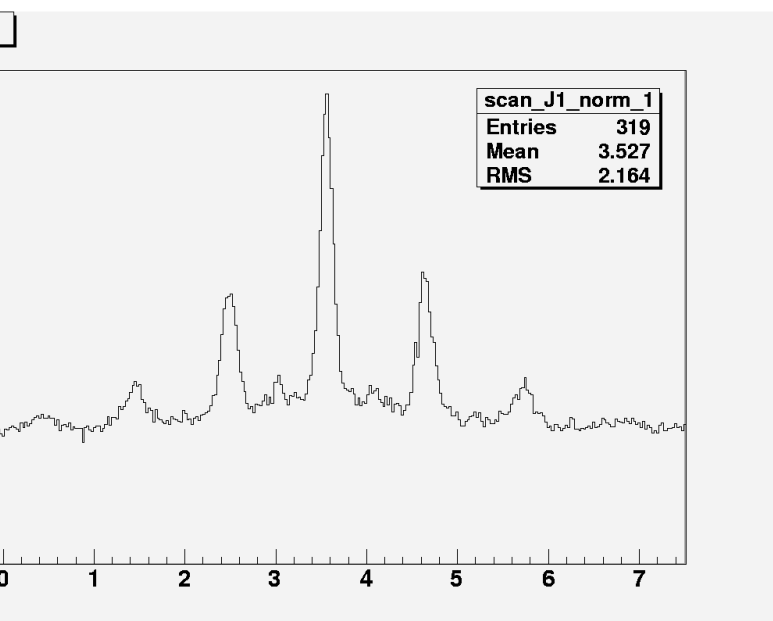
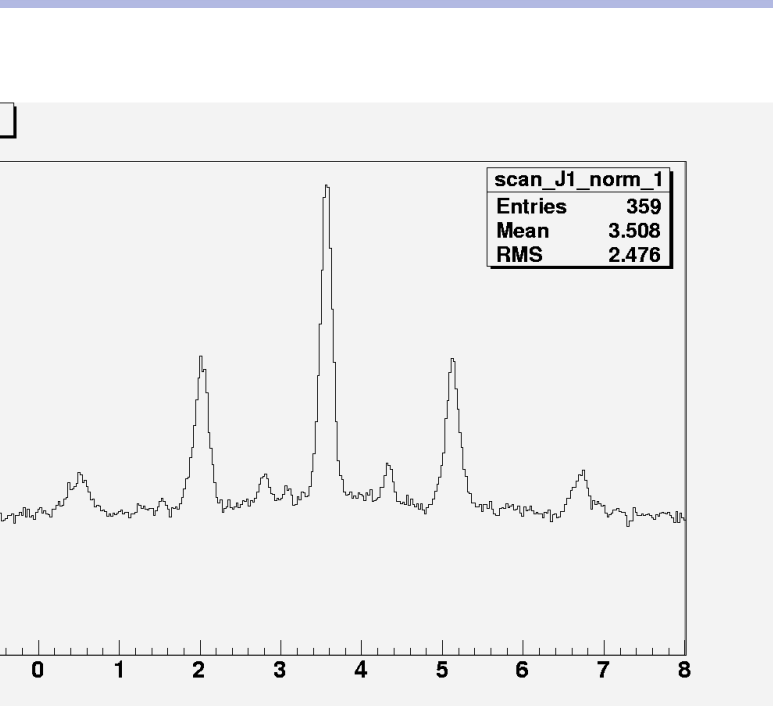
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The second method was traditional. It involved the measurement of orientation dependence of total photon flux with any photon beam monitor (e.g. ionization chamber or calorimeter).

In the MAX-lab experiments the scintillation gamma monitor was applied in front of which a 2-6 mm thick metal converter was placed for increasing the secondary charged particle yield.

The monitor worked in counting rate mode. The particle yield was normalized to electron beam current measured with a Faraday cup.

the measurements of gamma radiation intensity are carried out by rotating the crystal round one of the axis (for example, around the vertical axis) and for some fixed angles of horizontal rotation (Φ_V).



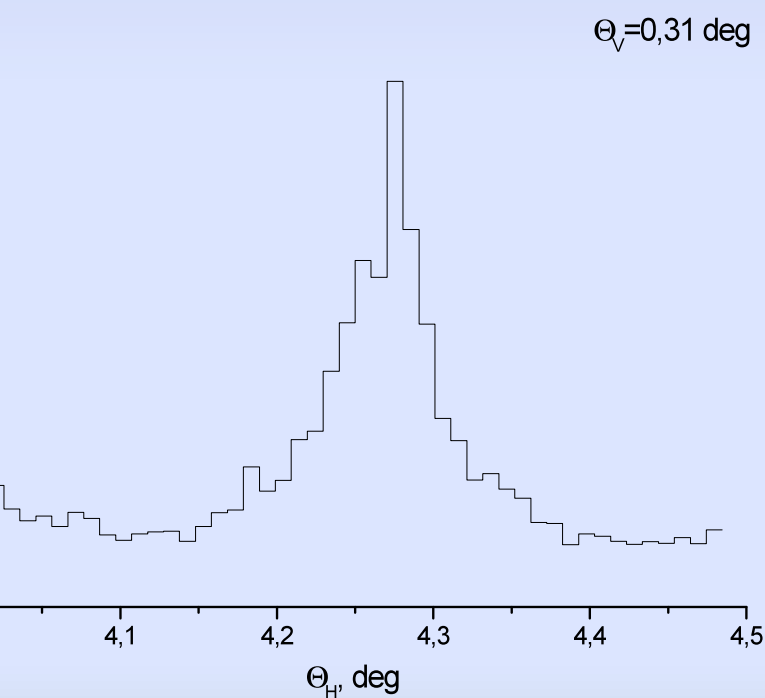
One can see the maximum intensity is observed in the central, corresponding to the strong plane (022), and other peaks corresponding to the planes (026), (062), (040).

*Fig.10. Orientation dependences
 $E_0=193.72$ MeV
 $\theta_V=-1.2^\circ$ and $\theta_V=-0.7^\circ$.*

When the crystal axis $\langle 001 \rangle$ coincides with direction of electron beam the large maximum is observed.

Intensity of the maximum is $\sim 3-5$ times larger than intensity in disoriented crystal. Large intensity is result of strong radiation of the electrons moving along crystal axis.

Channeling or above the barrier types of electron motion in this case are possible, as a result, the intensive photoemission in low energies is produced.



The width of the maximum is $\Delta\theta \sim 2(\psi_c + \theta_{msc})$ for diamond and silicon crystals

For diamond:

$\psi_c \sim 0.043^\circ (0.75 \text{ mrad})$ $\theta_{msc} \sim 0.0$

Fig. 11. Orientation dependence of the intensity of the radiation from a diamond crystal at the electron energy $E_0 = 193.72 \text{ MeV}$.

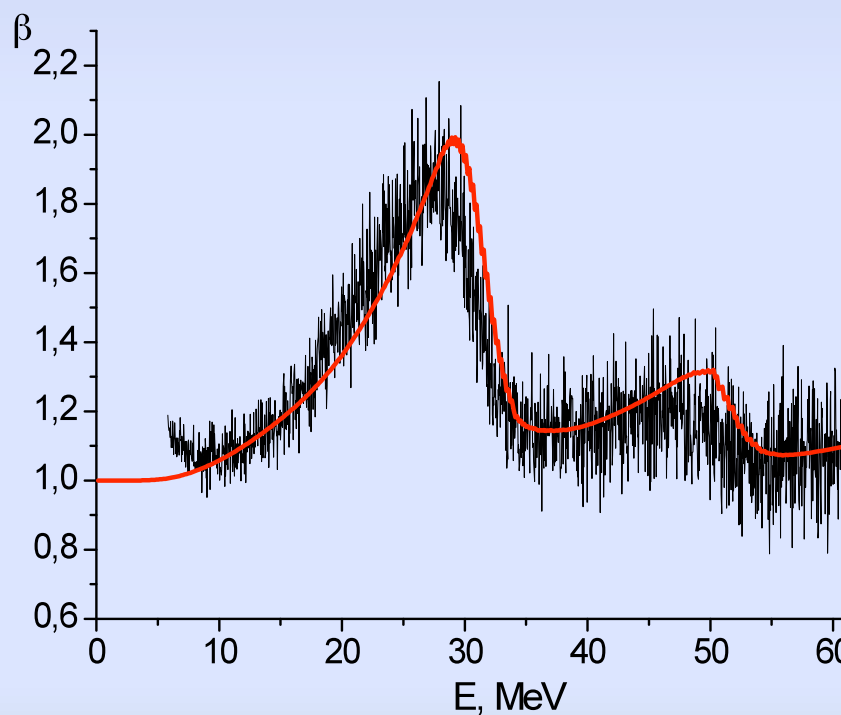
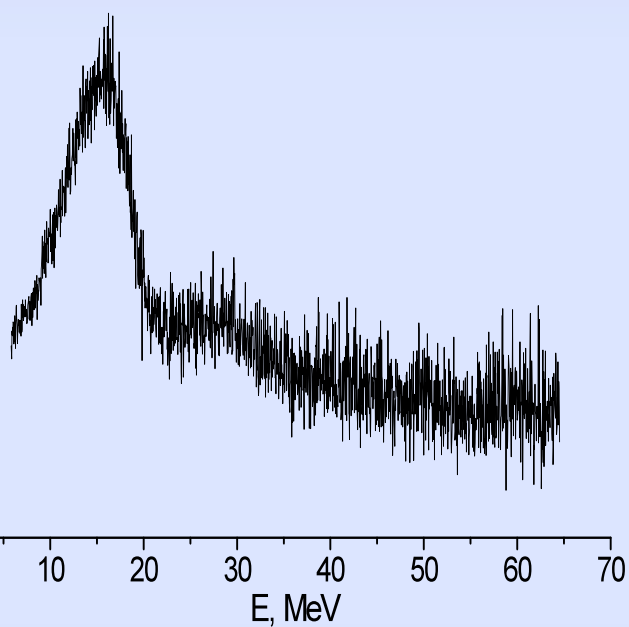
The CB spectra

The gamma radiation spectra were measured with NaI detector (25x25x25 cm³) at electron energies:

$E_0=143.87$ MeV (first run)

$E_0=192.66$ MeV (second run)

of the coherent effect is large enough $\beta_{\max} \sim 2$ at the CB energies $x_d < 0.2$ even for $E_0 \sim 150$ MeV. These values are in agreement with results of previous calculations and corresponds to a polarization of the CB maximum of $\sim 40\%$. The calculation of the CB spectrum is in good agreement with experiment.

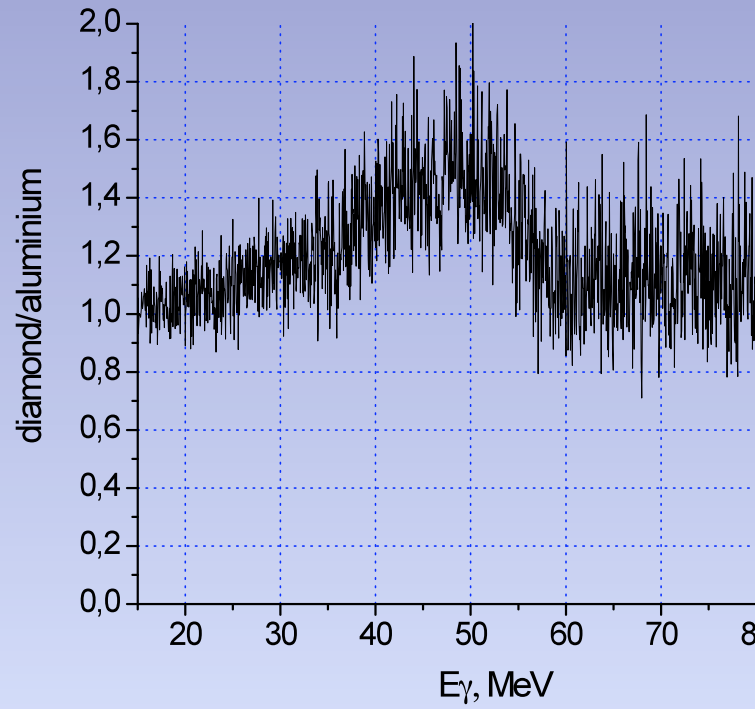
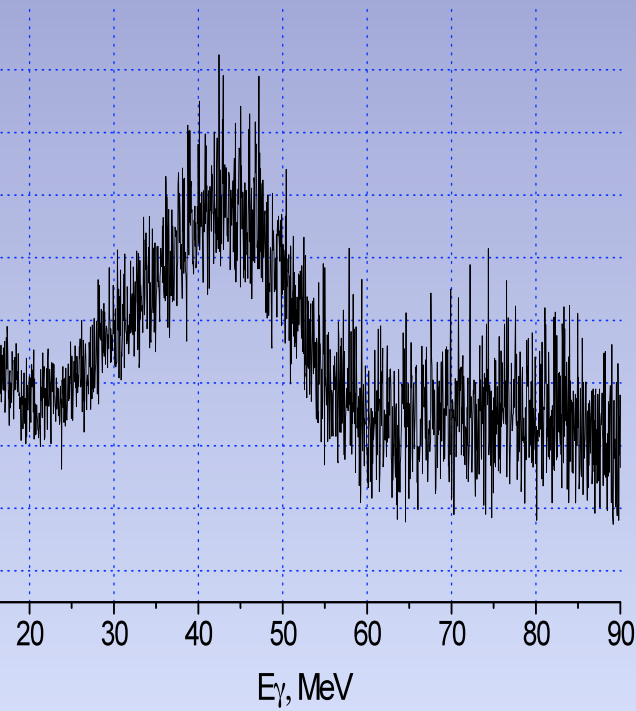


0. Spectra of the CB from diamond crystal 0.1 mm thick normalized to the bremsstrahlung spectrum from Al target. Electron energy is 143.9 MeV, $\theta_c \sim 0.73\theta_\gamma$. The curve is CB calculation. First

$\beta_{\max} \sim 1.7$

preliminary

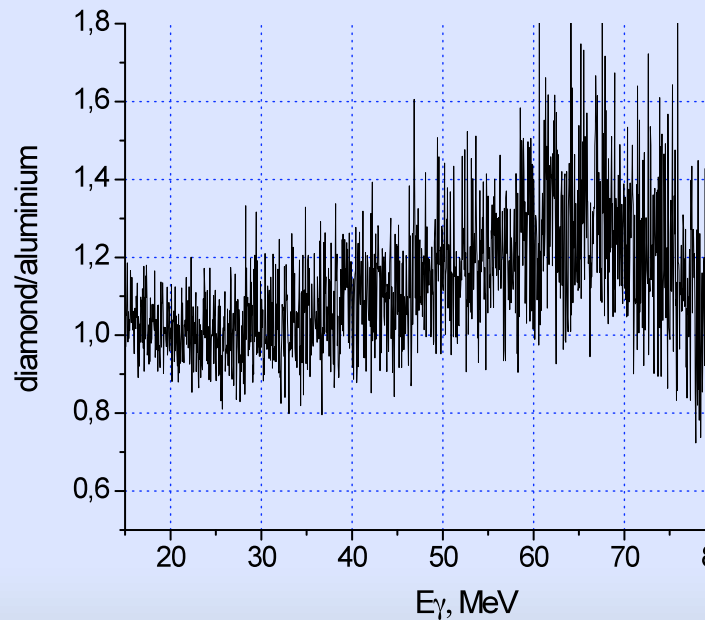
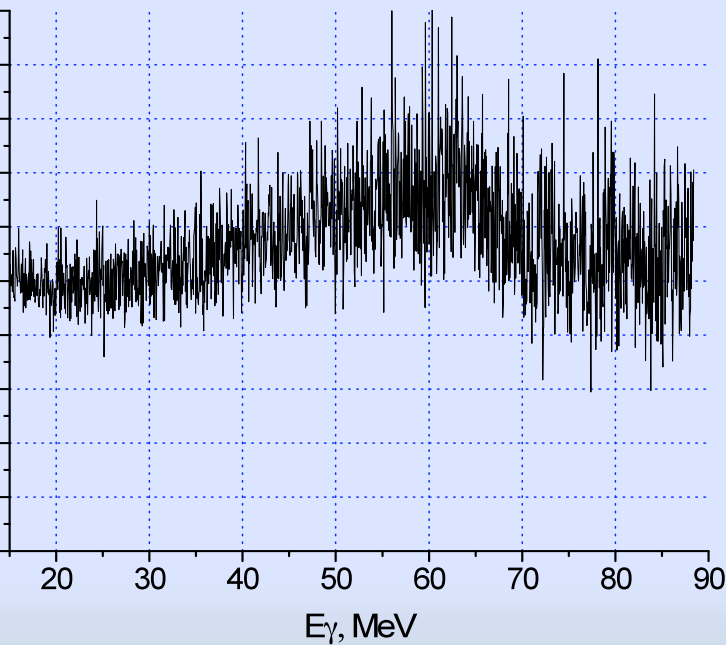
$\beta_{\max} \sim 1.5$



$\beta_{\max} \sim 1.4$

— 224-134

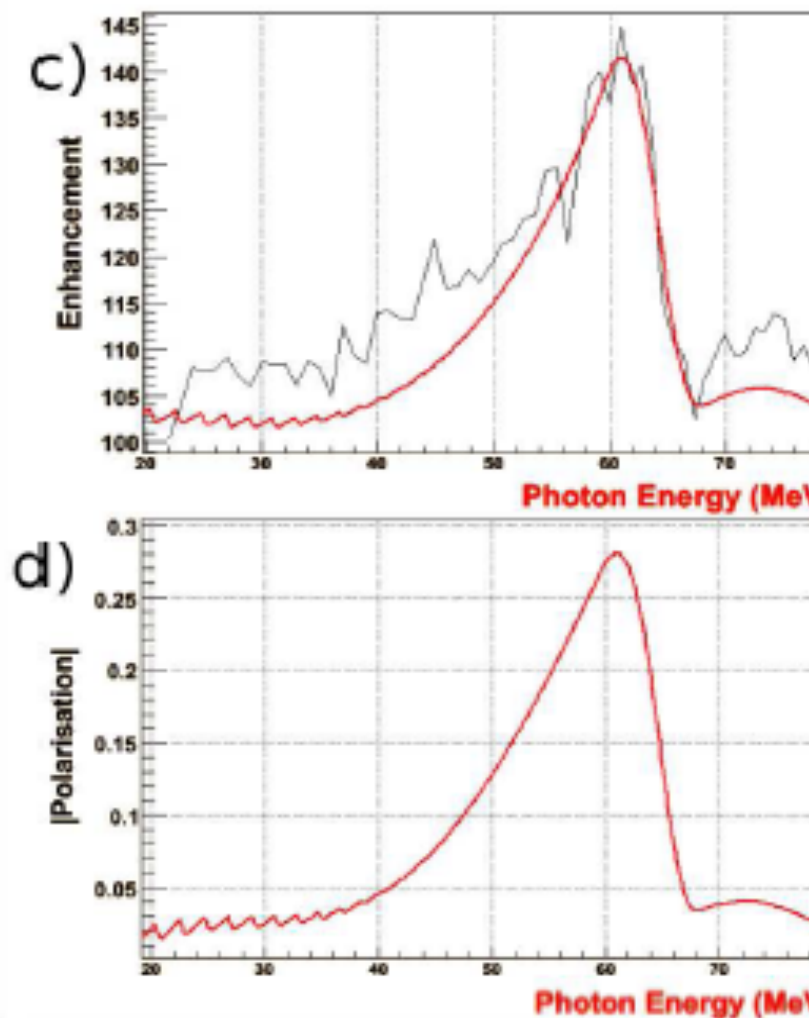
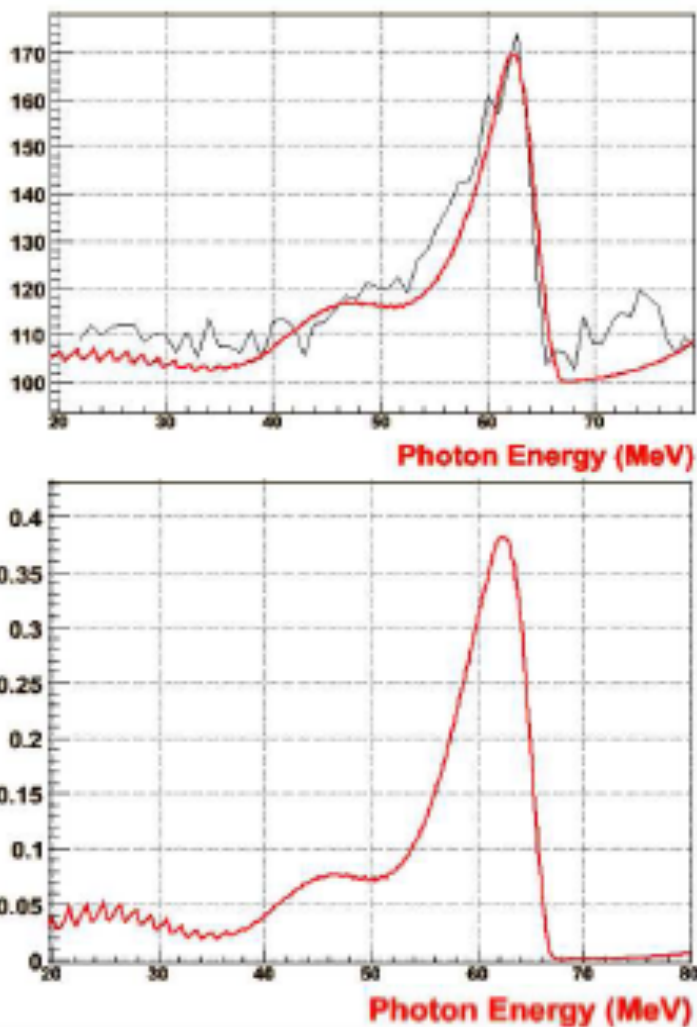
$\beta_{\max} \sim 1.3$



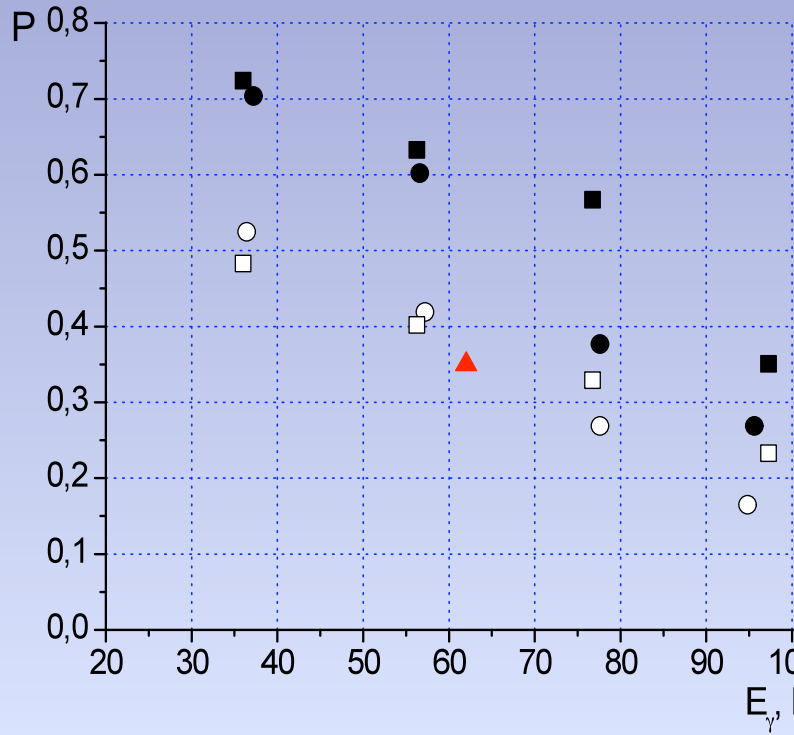
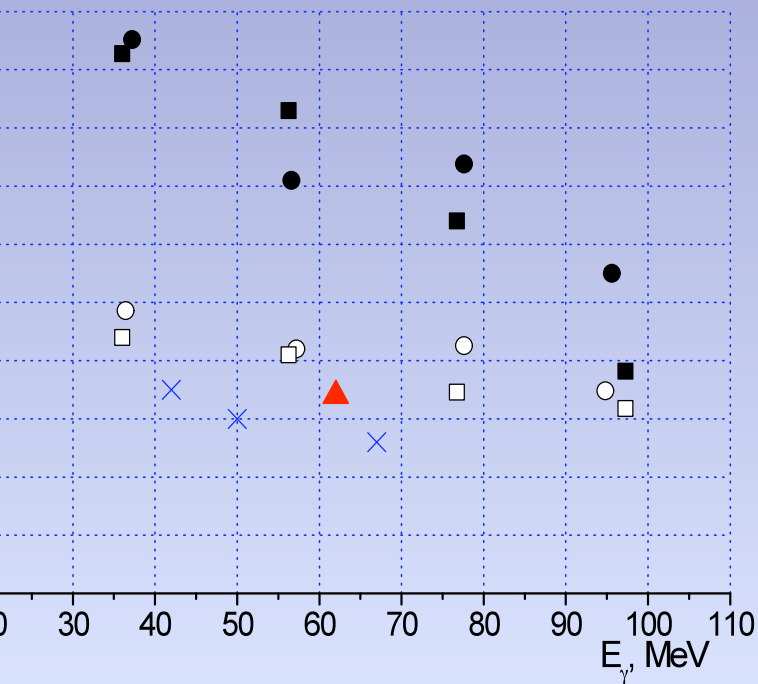
1. Spectra of the CB from a diamond crystal of thickness 0.1 cm, detection angle $\theta_c \sim 1.7\theta_\gamma$, electron energy $E_0 = 192.66 \text{ MeV}$.

Here it is presented the measured CB spectra for collimators 12mm and 4 mm (corresponds to collimation angles $\sim 1.1\theta_\gamma$) and $\theta_c \sim 0.37\theta_\gamma$) and results of the calculations. The photon energy is $E_\gamma \sim 62$ MeV. Polarization $\sim 37\%$.

Strong collimation does not improve the CB parameters (due to multiple scattering and comparable dimensions of the electron beam spot on the crystal (~ 3 mm) and the collimator hole).



12. The spectra (a) and polarization (b) of the CB for 12 mm collimator. The same plots for 4 mm collimator can be seen in (c), (d).



13. Dependencies of the coherent effect and polarization in maximum as a function of the peak energy.

50 MeV (squares: filled - $\theta_c=0.46\theta_\gamma$, empty - $\theta_c=0.91\theta_\gamma$).

100 MeV (circles: filled - $\theta_c=0.36\theta_\gamma$, empty - $\theta_c=0.54\theta_\gamma$).

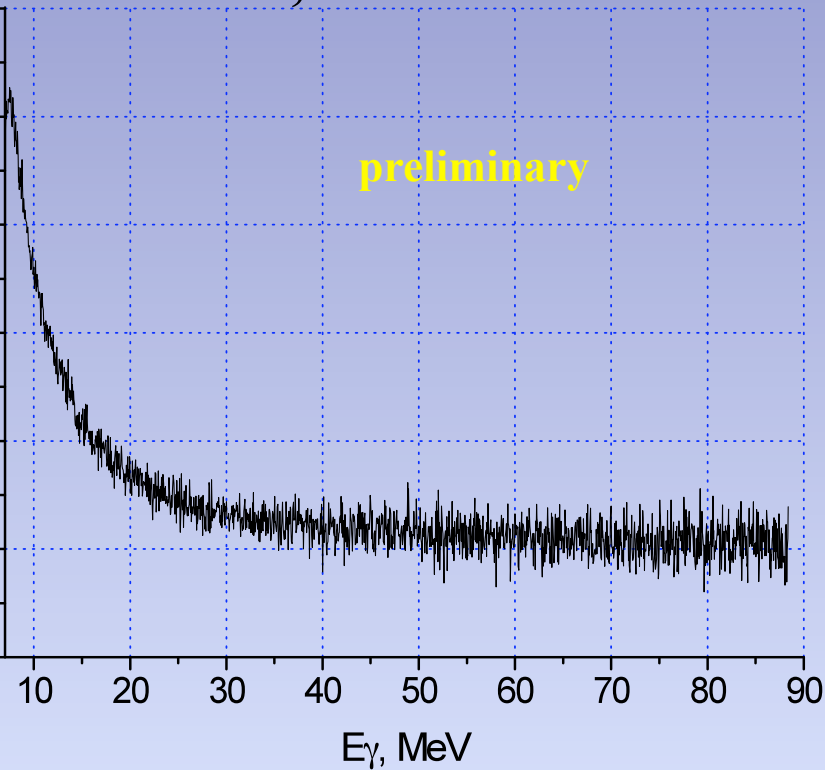
Measurements at $E_0=192.66$ MeV: crosses - $\theta_c\sim 1.7\theta_\gamma$, red triangle - $\theta_c\sim 1.7\theta_\gamma$.

$1\theta_\gamma$

Channeling radiation

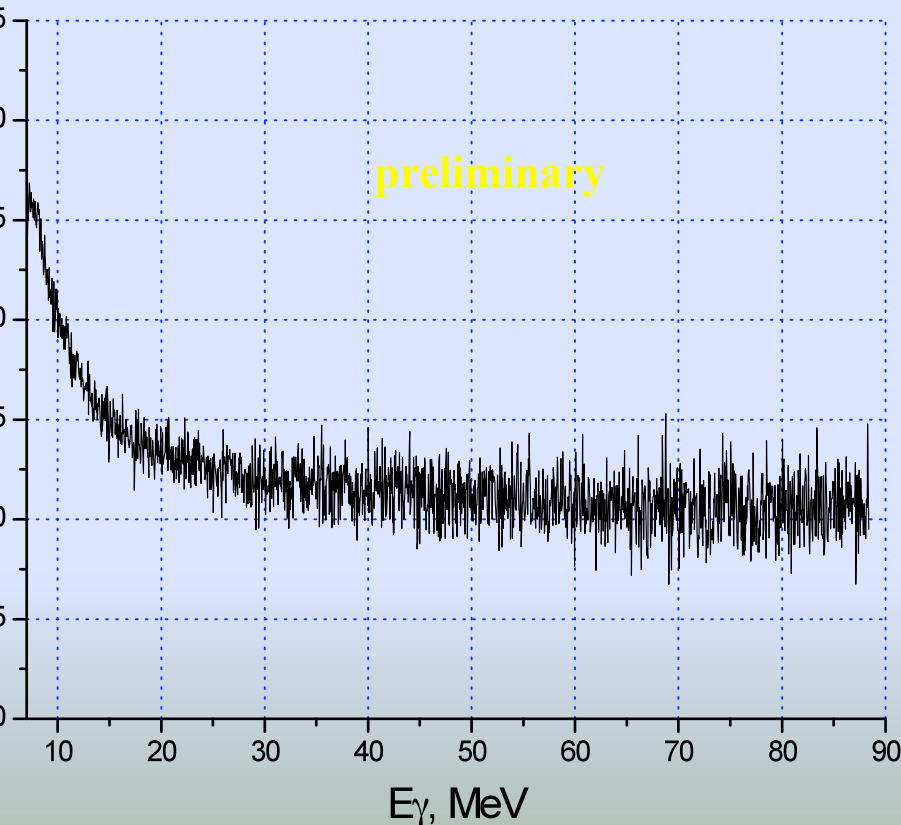
diamond, axis $\langle 001 \rangle$

— E.dat176



silicon, axis $\langle 111 \rangle$

— E.dat196



Strong increasing of radiation intensity at energies, at $\sim 6-7$ MeV almost 5 times (for diamond exceeds intensity of electron radiation in the amorphous matter of the same thickness).

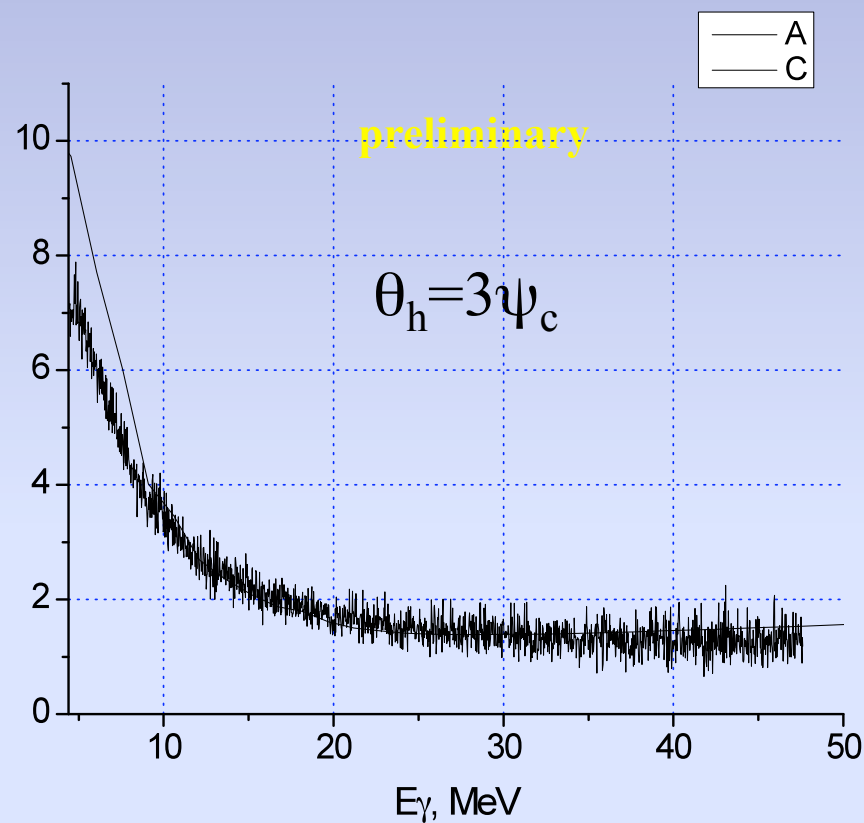
Critical angle of the axis channeling:

diamond $\psi_c \sim 0.75$ mrad

silicon $\psi_c \sim 0.91$ mrad

Fig.14. Spectra radiation of electrons with energy $E_0 = 192.66$ MeV conditions of axis channeling ($\theta_h = 0, \theta_v \sim \theta_c \sim 1.7 \theta_\gamma$).

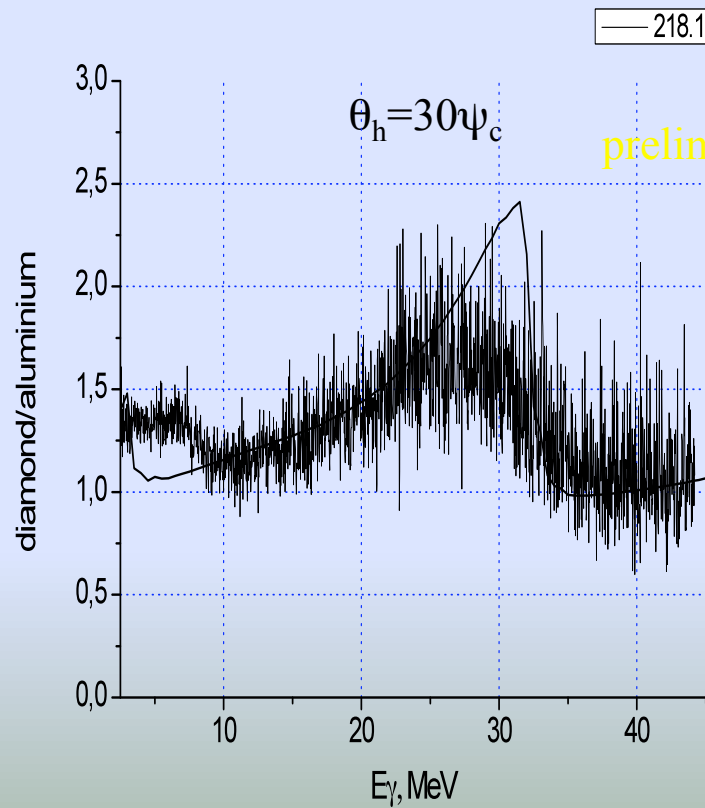
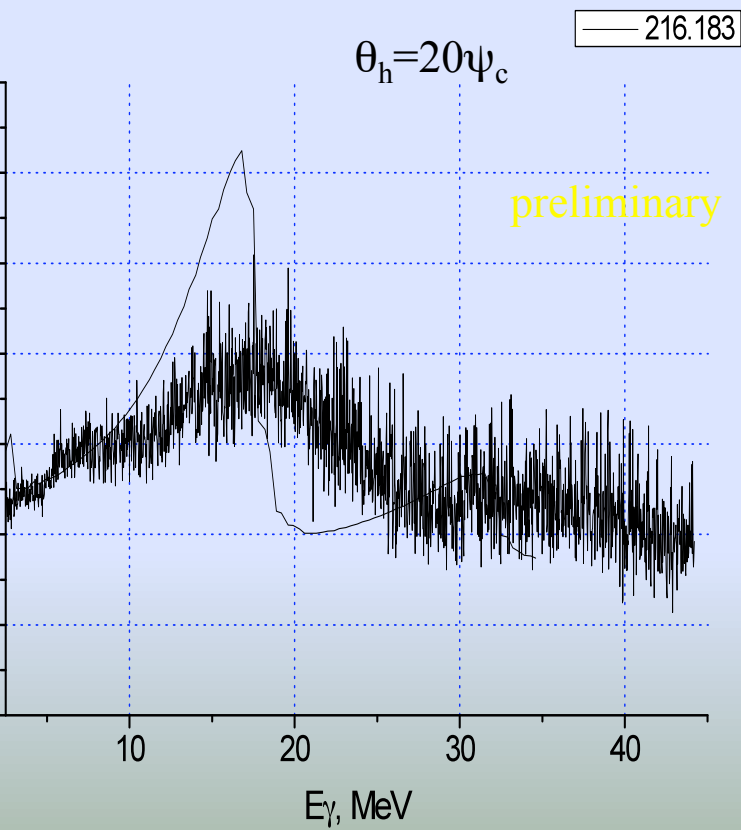
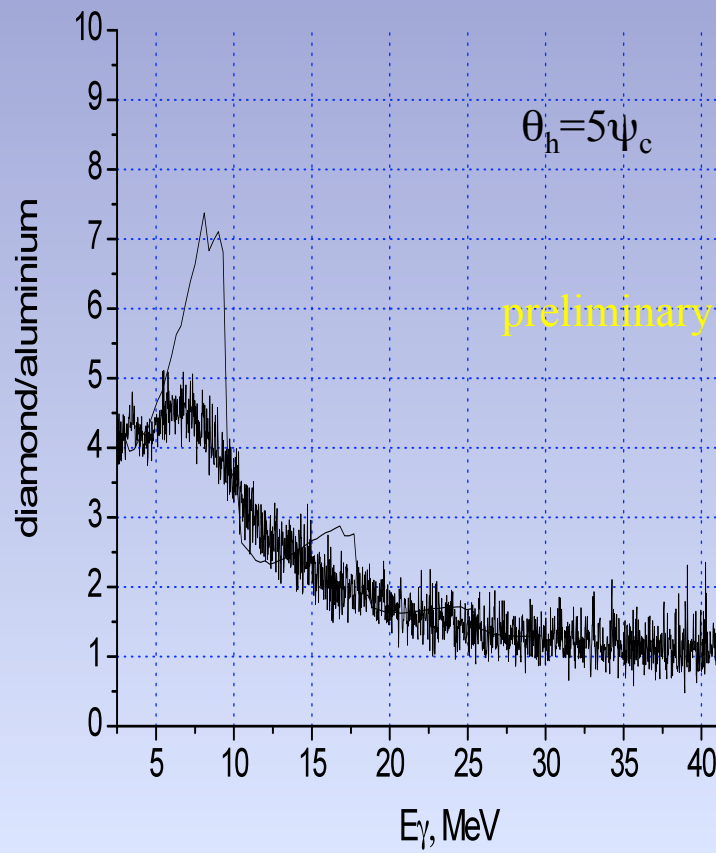
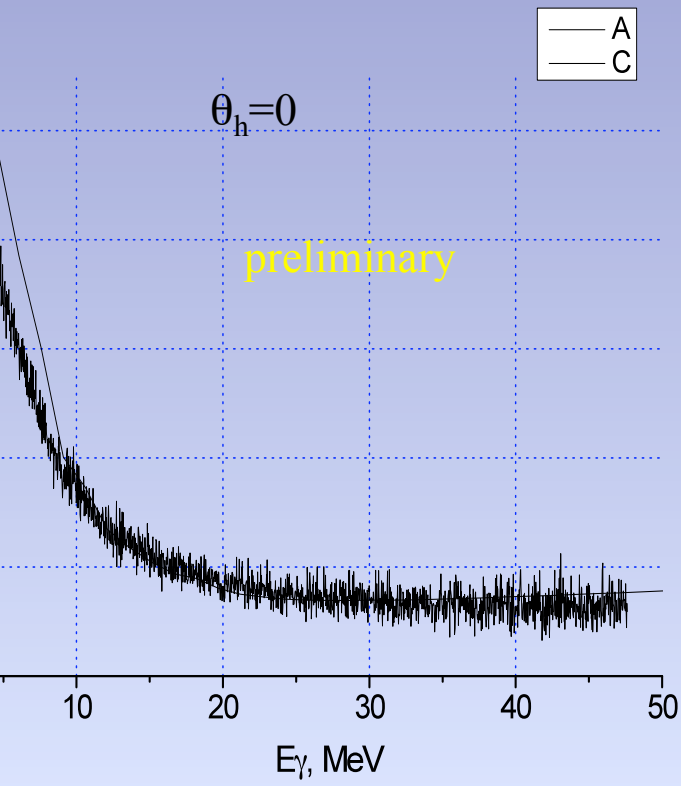
is shown the expected spectrum of electron radiation
ions move near axis $\langle 110 \rangle$ of diamond crystal 0.3 mm
ion energy is $E_0=200\text{MeV}$. At $\sim 3\text{MeV}$ it almost 10 times ex
ity of electron radiation in the amorphous matter of the
ness.

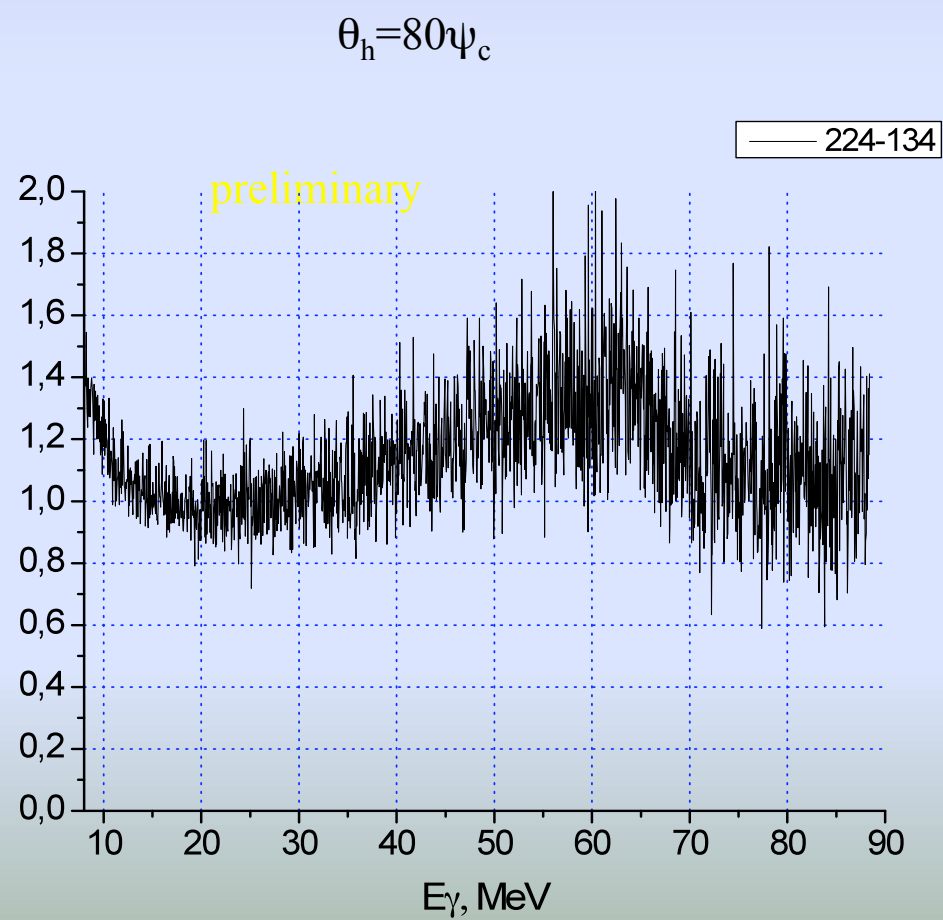
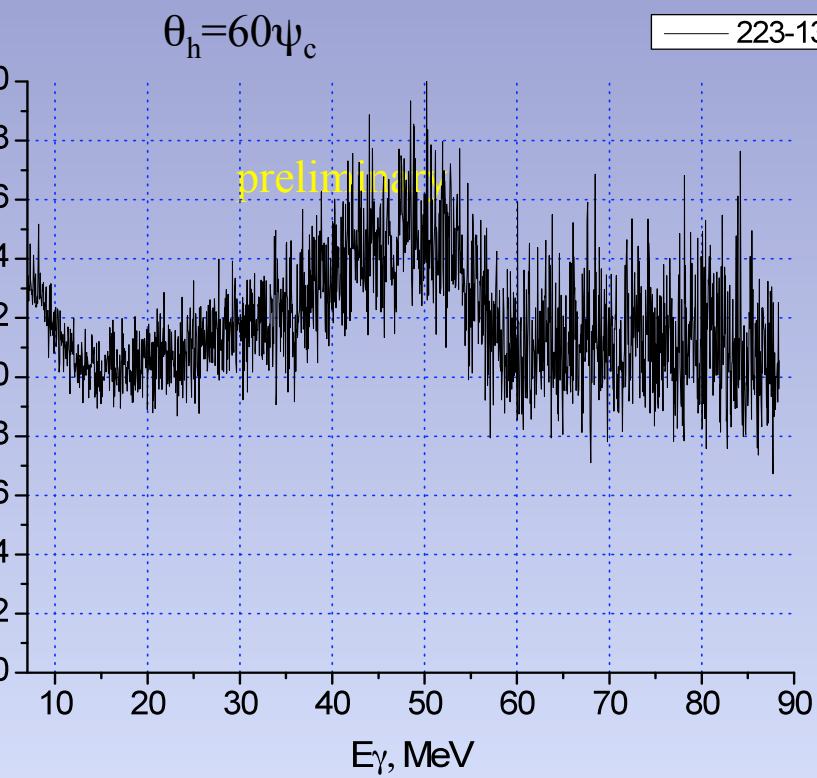


*16. Radiation spectrum of electron with energy $E_0=193.72$
ing under zero and three critical angle of axial channeling
 $\langle 001 \rangle$ along plane (220) of diamond crystal 0.1 mm
imation angle $\theta_c \sim 1.7\theta_\gamma$.*

*ve is the calculation in the model based on the semi-cla
roximation developed by Kharkov group (N.F. Shul'ga, V.I Tr
Greenenko. Nucl.Instr.Meth. 1998, v. B145)*

under angle θ_h to $\langle 100 \rangle$ in the plane (220)





Preliminary results

The test measurements have shown that the beam at MAX-lab was produced and the previous computations were confirmed.

The electron beam with energy $E_0=200$ MeV provides the photo-nuclear investigation in the range up to $E_\gamma \sim 60$ MeV (with beam collimation $\theta_c \sim \theta_\gamma$). The CB peak energies $x_d \sim 0.1-0.2$, ($E_{\gamma,p} \sim 20-30$ MeV). The polarization can achieve values $P_{\gamma,\max} \approx 50\%$.

To increase the polarization and coherent effect, the electron beam energy increasing up to 250 MeV is most effective. It will allow us to get the CB parameters suitable for nuclear physics experiments in the energy range up to 100 MeV.

Another possibility to improve the CB parameters is to use thinner crystal, ~ 50 μm .