# Status and perspectives of PHOTOS and TAUOLA comments on KKMC and BHLUMI

# Z. Was

### Institute of Nuclear Physics, PAN, Kraków, Poland

- What is new or important? Time limits → not a coherent presentation.
- especially for interpretators of new experimental data (low and high energy):
- TAUOLA: extra weights for multidimensional fits with MC. T. Przedzinski, V. Cherepanov, P. Golonka.
- MC-TESTER for automated comparisons and its database, P. Golonka, N. Davidson
- universal interface of TAUOLA: N. Davidson, E. Richter-Was, T. Przedzinski
- PHOTOS for radiative correction in decays: NLO in W-decays. G. Nanava.
- KKMC S. Banerjee, B. Pietrzyk, J.M. Roney

This talk in slightly shorter version was presented at TAU08 Novosibirsk conference.

My web page is at http://home.cern.ch/wasm

# TAUOLA: basic structure

- Phase space × Matrix elements, it is a must
- Electroweak vertex can stay in theorist hands only.
- Semileptonic decays are difficult: Hadronic current need to remain experiments' property, in cases expoeriment wish so.
- The last point enforces constraint for program organization and requests good communication between experimentalists, model builders and TAUOLA authors.
- This topic is being developed now with contributions from Vladimir Cherepanov (Novosibirsk), Tomasz Przedzinski (Krakow) and Piotr Golonka (CERN).
- Larger activity in the near future?
- TMVA Toolkit for Multivariate Data Analysis with ROOT http://tmva.sourceforge.net/ can be useful?
- People (and/or data) from Berlin, Beijing, Charkov, Frascati, Katowice, Karlsruhe,
   KEK, Novosibirsk, SLAC. Big topic for private discussions (too many directions).

### Semileptonic decays: Phase-space × weak-current × hadronic-current

The differential partial width for the channel under consideration reads

$$d\Gamma_X = G^2 \frac{v^2 + a^2}{4M} d\text{Lips}(P; q_i, N) (\omega + \hat{\omega} + (H_\mu + \hat{H}_\mu) s^\mu)$$

• The phase space distribution is given by the following expression where a compact notation with  $q_5=N$  and  $q_i^2=m_i^2$  is used

$$\begin{split} d \mathrm{Lips}(P;q_1,q_2,q_3,q_4,q_5) &= \frac{1}{2^{23}\pi^{11}} \int_{Q_{min}^2}^{Q_{max}^2} dQ^2 \int_{Q_{3,min}^2}^{Q_{3,max}^2} dQ_3^2 \\ & \int_{Q_{2,min}^2}^{Q_{2,max}^2} dQ_2^2 \quad \times \quad \int d\Omega_5 \frac{\sqrt{\lambda(M^2,Q^2,m_5^2)}}{M^2} \int d\Omega_4 \frac{\sqrt{\lambda(Q^2,Q_3^2,m_4^2)}}{Q^2} \\ & \quad \times \quad \int d\Omega_3 \frac{\sqrt{\lambda(Q_3^2,Q_2^2,m_3^2)}}{Q_3^2} \int d\Omega_2 \frac{\sqrt{\lambda(Q_2^2,m_2^2,m_1^2)}}{Q_2^2} \\ & \quad Q^2 &= (q_1+q_2+q_3+q_4)^2, \quad Q_3^2 = (q_1+q_2+q_3)^2, \quad Q_2^2 = (q_1+q_2)^2 \\ Q_{min} &= m_1 + m_2 + m_3 + m_4, \quad Q_{max} = M - m_5 \, Q_{3,min} = m_1 + m_2 + m_3, \qquad Q_{3,max} = Q - m_4 \\ Q_{2,min} &= m_1 + m_2, \qquad Q_{2,max} = Q_3 - m_3 \end{split}$$

 These formulas if used directly, are inefficient for a Monte Carlo algorithm if sharp peaks due to resonances in the intermediate states are present. The changes affect the program efficiency, but the actual density of the phase space remains intact. No approximations are introduced.

## **General formalism for semileptonic decays**

Matrix element used in TAUOLA for semileptonic decay

$$\tau(P,s) \to \nu_{\tau}(N)X$$

$$\mathcal{M} = \frac{G}{\sqrt{2}}\bar{u}(N)\gamma^{\mu}(v+a\gamma_5)u(P)J_{\mu}$$

- ullet  $J_{\mu}$  the current depends on the momenta of all hadrons.
- I can provide only prototypes for  $J_{\mu}$ . Here TAUOLA must be open to experiment interior. difficult in mixed C++ Fortran environment, and of different experiments software.

$$|\mathcal{M}|^{2} = G^{2} \frac{v^{2} + a^{2}}{2} (\omega + H_{\mu} s^{\mu})$$

$$\omega = P^{\mu} (\Pi_{\mu} - \gamma_{va} \Pi_{\mu}^{5}), \qquad H_{\mu} = \frac{1}{M} (M^{2} \delta_{\mu}^{\nu} - P_{\mu} P^{\nu}) (\Pi_{\nu}^{5} - \gamma_{va} \Pi_{\nu})$$

$$\Pi_{\mu} = 2 [(J^{*} \cdot N) J_{\mu} + (J \cdot N) J_{\mu}^{*} - (J^{*} \cdot J) N_{\mu}]$$

$$\Pi^{5\mu} = 2 \operatorname{Im} \epsilon^{\mu\nu\rho\sigma} J_{\nu}^{*} J_{\rho} N_{\sigma}, \qquad \gamma_{va} = -\frac{2va}{v^{2} + a^{2}}$$

• If  $\tau$  coupling  $v+a\gamma_5$  and  $m_{\nu_{\tau}}\neq 0$  is allowed, one has to add to  $\omega$  and  $H_{\mu}$ :

$$\hat{\omega} = 2 \frac{v^2 - a^2}{v^2 + a^2} m_{\nu} M(J^* \cdot J)$$

$$\hat{H}^{\mu} = -2 \frac{v^2 - a^2}{v^2 + a^2} m_{\nu} \operatorname{Im} \epsilon^{\mu\nu\rho\sigma} J_{\nu}^* J_{\rho} P_{\sigma}$$

### **Main references:**

- R. Decker, S.Jadach, M.Jeżabek, J.H.Kuhn, Z. Was, Comput. Phys. Commun. 76 (1993) 361, ibid. 70 (1992) 69, ibid. 64 (1990) 275
- 2. P. Golonka, B. Kersevan ,T. Pierzchala, E. Richter-Was, Z. Was, M. Worek, Comput.Phys.Commun.174:818-835,2006
- 3. J.H.Kuhn, Z. Was, hep-ph/0602162, Acta Phys. Polon. 39, (2008) 47 (5-pions)



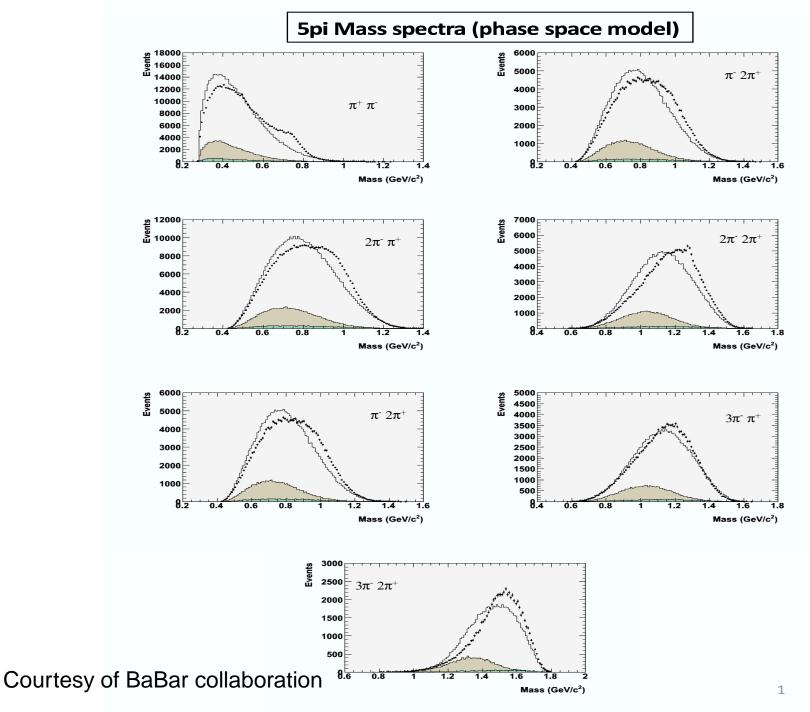
- 1. Alain Weinstein www home page: http://www.cithep.caltech.edu/~ajw/korb\_doc.html#fi les
- 2. B. Bloch, private communications.
- 3. R. Decker, M. Finkemeier, P. Heiliger and H.H. Jonsson, Z. Phys. C **70** (1996) 247, now standard  $4\pi$  formfactors.
- 4. A. E. Bondar, S. I. Eidelman, A. I. Milstein, T. Pierzchala, N. I. Root, Z. Was and M. Worek, Comput. Phys. Commun. **146**, 139 (2002)
- 5. P. Abreu et al., Phys. Lett. B426 (1998) 411 (alternative  $3\pi$  formf.)
- 6. Sherry Towers alternative formf. in K $\pi\pi$  modes, hep-ex/9908013, Eur. Phys. J. **C13** (2000) 197.

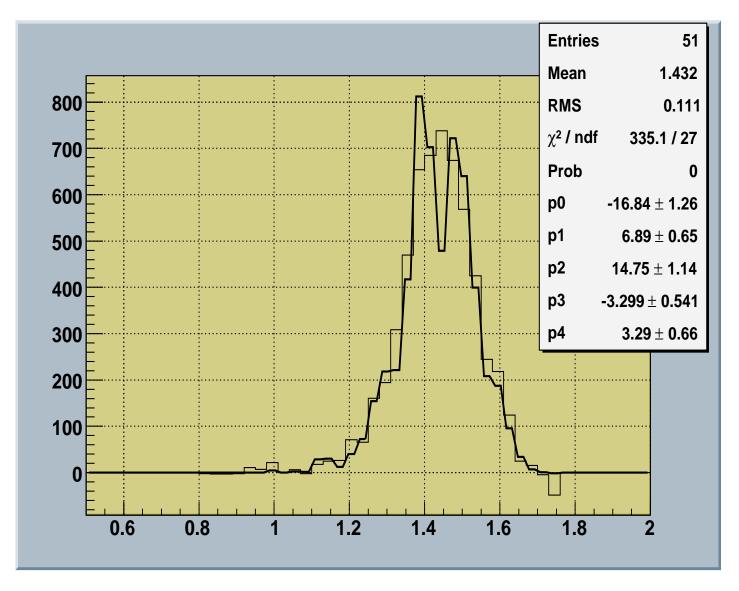
## Formfactors secret life

The studies within collaborations often rely on private form-factors, wealth of versions were/are regularly created for more general, or specific purposes. I have seen some.

### Arrangements for multidimensional fits with MC.

- for two hadron final states fits are easy, for three, one can separate contributions of different resonances chains using angle dependent projection operators ( J. H. Kuhn, E. Mirkes Z.Phys.C 1992), but loss of sensitivity and systematic errors. For more hadrons nobody even tried. Too many interferences ...
- For each tau decay calculate wector of weights for alternative decay models
- Choose the best one comparing simulated sample with data. If dependence is linear in current it is qadratic in weight. Fit function can be analyticall all over parameter space.
   Otherwise linearization/iteration necessary.
- Arrangement works for TAUOLA standalone and also for KKMC as installed in Belle software. Thanks to A. Bozek for help.
- First result point to control of relation between model assumptions and unitarity.
- Exact phase space brings consequences. Unitarity discipline is a must:



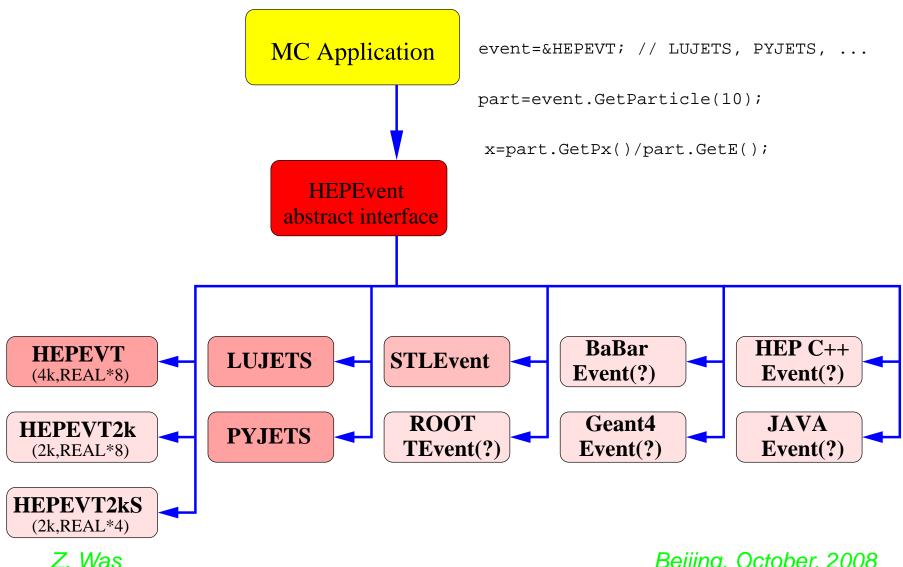


Technical test from the work with V. Cherepanov. Good example of collecting experience.  $\rho$  system propagator has to have proper phase space dependent width:  $\tau \to \eta \pi \pi \nu$ .

# MC-TESTER was developped to automate comparisons

- It was essential in our work on TAUOLA, and we expect it to be even more essential in the future, also for debugging.
- The same is true for our projects on PHOTOS developments.
- necessary tool for migration from Fortran to C++.
- Now installed in ATLAS collaboration ATHENA system and used by LCG MC group at CERN.
- Thanks to P. Golonka, T. Przedzinski and N. Davidson for efforts!
- Now we can use it with C++ HepMC event record and all C++ generators as well.

This tool can be used for any MC storing events in standard common blocks: HEPEVT, PYJETS, ... It may also be extended to adopt new event-record data-structures (i.e. in C++). Recent attempt to have standard: T. Sjostrand et al. A standard format for Les Houches Event Files, hep-ph/0609017. Is it going to be useful and accepted (this time)?



MC-TESTER RESULTS

http://mc-tester.web.cern.ch/MC-TESTER/mc-tester\_results/results.html

#### MC-TESTER RESULTS

Here you will find some of the MC-TESTER output comparing various Monte-Carlo Generators. To expand this list, we encourage you to share MC-TESTER validation results or suggestions.

Click on the pdf files in the comparison matrix to get the results booklet.

Click on the generator name to get the ROOT output file from the generation step of MC-TESTER.

#### **Tau Decay Results**

Note: The ROOT files were produced with ROOT Version 5.18, histogram range 0-2 GeV in 60 bins. The comparisons were done with User Analysis MCTest01

Generator	Pythia 6.4.14	Pythia 8.1	Tauola - Cleo
Pythia 6.4.14	-		
Pythia 8.1	tester 6.4vs8.1.pdf	-	
Tauola - Cleo	tauola cleo vs pythia 6.4.pdf	tauola cleo vs pythia 8.1.pdf	-

#### **B+ Decay Results**

Note: The ROOT files were produced with ROOT Version 5.18, histogram range 0-6 GeV in 60 bins. The comparisons were done with User Analysis MCTest01

Generator	Pythia 8.1
EvtGenLHC 5.15	evtgenlhc vs pythia 8.pdf

web-page http://mc-tester.web.cern.ch/MC-TESTER/mc-tester\_results/results.html Most of the test results are hidden in Atlas repository. In particular one can find there tests of tau decays in HERWIG Sherpa etc.

# MC-TESTER results for decays of particle $\tau^-$ (PDG code 15).

Piotr Golonka

Tomasz Pierzchala

Zbigniew Was

May 22, 2004

#### **Results from generator 1.**

tauola-cleo starting point no modifications in any case May 19 2004.

- From directory: /home/wasm/y2004/TAUOLA-all/nowa-tauola/TAUOLA/tauola-old/demo-standalone/prod
- Total number of analyzed decays: 5000000
- Number of decay channels found: 32

#### **Results from generator 2.**

tauola-cleo new version new channels installed, brs=\*0.001 May 22 2004.

- From directory:
- /home/wasm/y2004/TAUOLA-all/nowa-tauola/TAUOLA/tauola-new/demo-standalone/prod
- Total number of analyzed decays: 5000000
- Number of decay channels found: 32 + 8

1

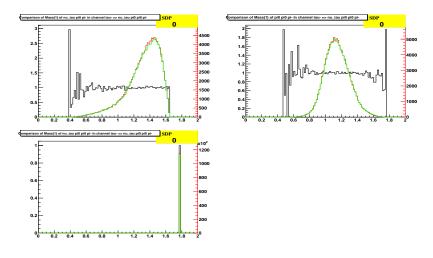
MC-TESTER booklet: Page 1. Run like for installation tests.

#### Found decay modes:

Decay channel	Branching Ratio ± Rough Errors  Generator #1 Generator #2		Max. shape dif. param.
$ au^-  ightarrow  u_{ au} K^-$	$4.5460 \pm 0.0095\%$	$4.5500 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  au_ au \pi^0 \pi^0 \pi^+ \pi^- \pi^-$	$4.5460 \pm 0.0095\%$	$4.5425 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^+ \pi^+ \pi^- \pi^- \pi^-$	$4.5457 \pm 0.0095\%$	$4.5303 \pm 0.0095\%$	0.00000
$\tau^- \to \nu_{\tau} \pi^0 \pi^0 \pi^0 \pi^+ \pi^- \pi^-$	$4.5449 \pm 0.0095\%$	$4.5271 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  au_ au \pi^0 \pi^0 \pi^0 \pi^-$	$4.5416 \pm 0.0095\%$	$4.5366 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 \pi^+ \pi^- \pi^-$	$4.5392 \pm 0.0095\%$	$4.5371 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \gamma \pi^0 \pi^-$	$4.5368 \pm 0.0095\%$	$4.5160 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 \pi^0 K^-$	$4.5268 \pm 0.0095\%$	$4.5468 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 \pi^- \eta$	$4.5236 \pm 0.0095\%$	$4.5154 \pm 0.0095\%$	0.00000
$ au^-  ightarrow \mu^- \widetilde{ m V}_{\mu} { m V}_{ au}$	$4.3942 \pm 0.0094\%$	$4.3919 \pm 0.0094\%$	0.00000
$ au^-  ightarrow e^- \widetilde{ m V}_e { m V}_{ m T}$	$3.8276 \pm 0.0087\%$	$3.8245 \pm 0.0087\%$	0.00000
$ au^-  ightarrow  au_ au \pi^0 \pi^0 \pi^-$	$2.2907 \pm 0.0068\%$	$2.2669 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au K_S^0 K^-$	$2.2832 \pm 0.0068\%$	$2.2582 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  au_ au \pi^0 K_L^0 K^-$	$2.2825 \pm 0.0068\%$	$2.2698 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_{ au} K_L^0 K^-$	$2.2795 \pm 0.0068\%$	$2.2725 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 K_L^0 \pi^-$	$2.2756 \pm 0.0067\%$	$2.2680 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au K_L^0 \pi^- K_S^0$	$2.2756 \pm 0.0067\%$	$2.2667 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 K_S^0 K^-$	$2.2717 \pm 0.0067\%$	$2.2606 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 \pi^- K_S^0$	$2.2582 \pm 0.0067\%$	$2.2663 \pm 0.0067\%$	0.00000
$ au^-  ightarrow  u_ au \pi^+ \pi^- \pi^-$	$2.2449 \pm 0.0067\%$	$2.2822 \pm 0.0068\%$	0.00000
$ au^-  ightarrow  u_{ au} \pi^0 K^-$	$1.5545 \pm 0.0056\%$	$1.5441 \pm 0.0056\%$	0.00000
$ au^-  ightarrow  u_{ au} \pi^- K_S^0$	$1.5047 \pm 0.0055\%$	$1.4819 \pm 0.0054\%$	0.00000
$ au^-  ightarrow  u_{ au} K_L^0 \pi^-$	$1.5019 \pm 0.0055\%$	$1.4915 \pm 0.0055\%$	0.00000
$ au^-  ightarrow  u_ au \pi^- K^+ K^-$	$4.5561 \pm 0.0095\%$	$4.5349 \pm 0.0095\%$	0.00000
$\tau^- \to \nu_\tau \pi^-$	$4.5501 \pm 0.0095\%$	$4.5291 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_{ au} \pi^+ \pi^- K^-$	$4.5465 \pm 0.0095\%$	$4.5461 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^0 \pi^-$	$4.5528 \pm 0.0095\%$	$4.5405 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  au_ au K_L^0 K_L^0 \pi^-$	$1.1407 \pm 0.0048\%$	$1.1324 \pm 0.0048\%$	0.00000
$\tau^- \rightarrow \nu_\tau \pi^0 \pi^+ \pi^+ \pi^- \pi^- \pi^-$	$4.5557 \pm 0.0095\%$	$4.5381 \pm 0.0095\%$	0.00000
$ au^-  ightarrow  u_ au \pi^- K^0_S K^0_S$	$1.1340 \pm 0.0048\%$	$1.1404 \pm 0.0048\%$	0.00000
$ au^-  ightarrow e^- \widetilde{ u_e}  u_{ au} \gamma$	$0.7181 \pm 0.0038\%$	$0.7164 \pm 0.0038\%$	0.00000
$ au^-  ightarrow \mu^- \widetilde{ m V}_\mu { m V}_{ m T} \gamma$	$0.1507 \pm 0.0017\%$	$0.1489 \pm 0.0017\%$	0.00000

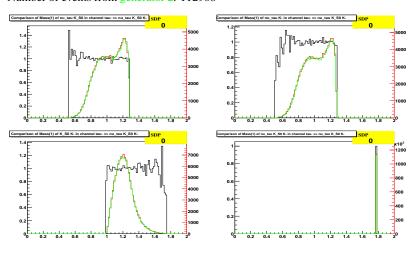
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MC-TESTER booklet: Page 2. Run like for installation tests.



#### 13 Decay Channel: $\tau^- \rightarrow \nu_{\tau} K_S^0 K^-$

Number of events from generator 1: 114161 Number of events from generator 2: 112908



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MC-TESTER booklet: Page 49. Run like for installation tests.

#### MC-TESTER results for decays of particle $B^+$ (PDG code 521).

Piotr Golonka

Tomasz Pierzchala April 18, 2008

Zbigniew Was

#### **Results from generator 1.**

EvtGenLHC demo.

- Total number of analyzed decays: 500000
- Number of decay channels found: 12876 + 20016

#### **Results from generator 2.**

Pythia 8.1 demo; p-p at 14 TeV gg->bbbar. *B*+ decay analysed

- Total number of analyzed decays: 484029
- Number of decay channels found: 12876 + 13012

taken from http://mc-tester.web.cern.ch/MC-TESTER/mc-tester\_results/results.html web page. Nice results for B+ decays. Compared MC are EvtGenLHC 5.15 and Pythia 8.1. In total 50 k-channels found, 1000+ pages booklet created.

# PHOTOS for bremsstrahlung in decays and precision

E. Barberio, B. van Eijk, Z. Was, Comput. Phys. Commun.(1991) ibid. (1994)

See also: P. Golonka et al. hep-ph/0312240, Comput.Phys.Commun.174 (2006) 818.

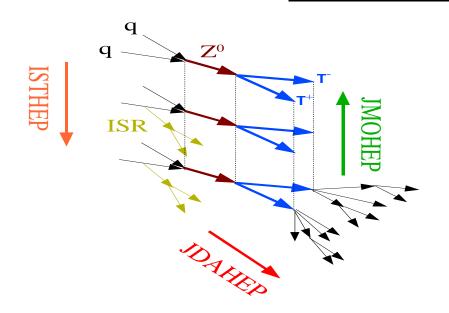
P. Golonka and Z. Was hep-ph/0604232 EPJC (2007), G. Nanava and Z. Was hep-ph/0607019 EPJC 2007

- ullet It was developed as single photon emission. starting from MUSTRAAL (F. Berends, R. Kleiss, S. Jadach, Comput. Phys. Commun. (1982)) option for fi nal state bremsstrahlung in Z decay only.
- Factorization of phase space for photonic variables and two-body decay phase space was studied. Similarly for matrix element: process independent kernel was found. Phase space is exact.
- ullet Interference between emission from  $\mu^+$  and  $\mu^-$  is dropped and re-introduced later.
- The algorithm of the antenna type was created: full phase space NLO ready
- Works for single emission (orthodox ME Monte Carlo)
- Fixed order (up to double emission), useful for tests with second order ME.
- Fixed order (up to quatric emission) or multiphoton.
- Nice environment to study options of factorization schemes, relations between exponentiation structure function evolution, etc.



- Generally kernels in PHOTOS, are not better than LL. To improve, process dependent weights are needed. Complications for users, but otherwise straightforward!
- Special weights with complete matrix elements are available now for:  $Z \to \mu^+\mu^-$  (2005), and for  $B^+ \to K^+\pi^0$ ,  $B^0 \to K^+\pi^-$ , etc. (2006), for  $B^0$  and  $B^\pm$  decays scalar QED)
- ullet W o l 
  u (2008)
- We will see that effects are small, it is sufficient to keep them for tests only.
- Lots of numerical tests.
- For other decay modes such exercises easy to repeat, if matrix elements are available.
- As consequence: PHOTOS is ready for improvements with measured data as well!
- PHOTOS uses mother-daughter relations in HEPEVT.
- C++ version is prepared but not distributed, need to migrate to HepMC event record.
- Program analyze whole event record and may add bremsstrahlung at any branching.
- Appropriately modifi es particles momenta of the whole cascade!
- Algorithm is vulnerable on the way how HEPEVT is filled in. Any new inconistency and ...

### Problems With Event Record



- 1. Hard process
- 2. with shower
- 3. after hadronization
- Event record overloaded with physics beyond design → gramar problems.
- 5. Here we have basically LL phenomenology only.

# This Is Physics Not F77!

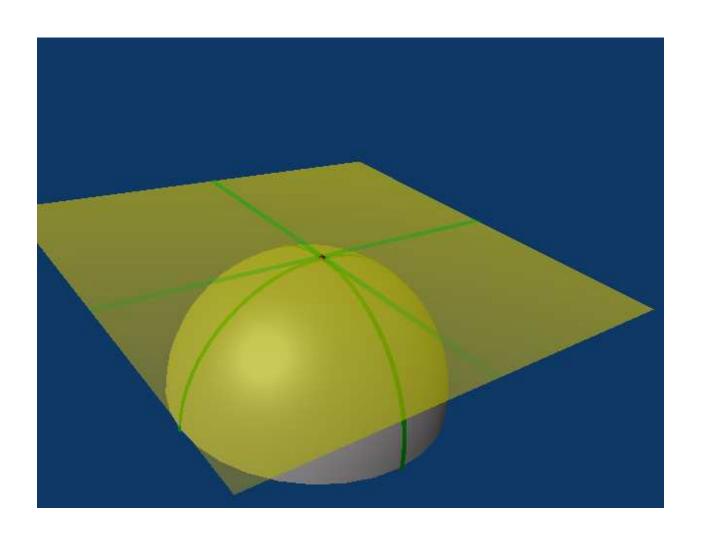
Similar problems are in any use of full scale Monte Carlos, lots of complaints at MC4LHC workshop, HEPEVTrepair utility (C. Biscarat and ZW) being probed in D0.

Design of event structure WITH some grammar requirements AND WITHOUT neglecting possible physics is needed NOW to avoid large problems later.

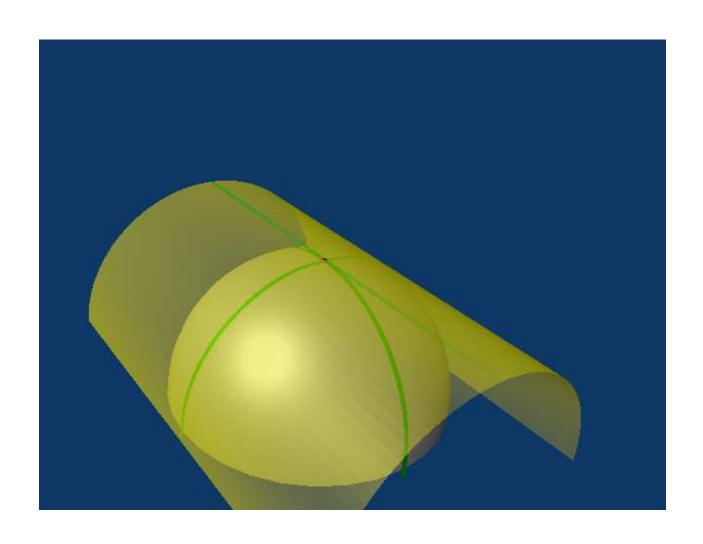
### G. Nanava and Z. Was hep-ph/0607019 EPJC 2007

- It is not just about computer technology.
- I should provide discussion of phase-space in PHOTOS here
- It is always refreshing and rewarding subject.
- Question of matching two NLO 'emission showers' both of full (thus overlapping)
  phase-space. Clean environment where soft singularity can be conrolled by YFS
  exponentiation to all orders
- sigularity structure must be consistent with triangulation. CW complexes and homotopy grup.
- Instead I will just show one plot which brings more doubts than explains.
- I will skip NLO matching of consecutive emissions (1994 done on basis of  $Z \to l^+ l^-$ ) too.

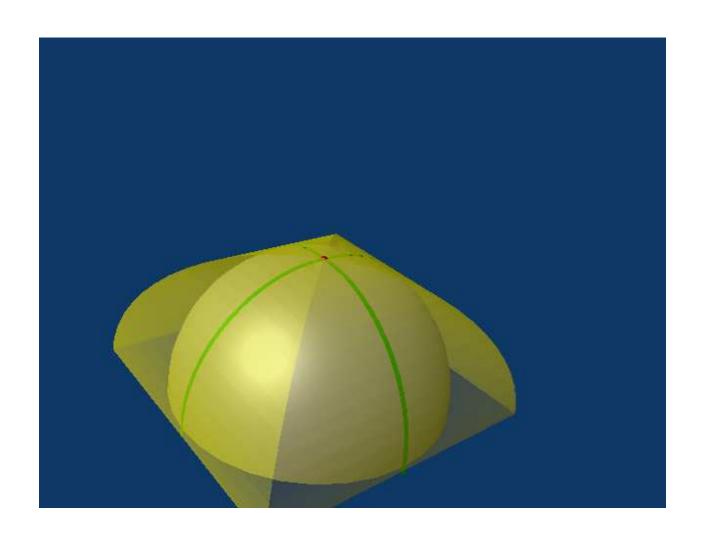
Heuristic plots



Heuristic plots



Heuristic plots



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### Results of P. Golonka and Z. Was hep-ph/0604232

- ullet From the point of view of matrix element and choice of internal angular variables in case of Z decays, PHOTOS is very close to MUSTRAAL MC by F. Berends S. Jadach and R. Kleiss, CPC 29 (1983) 185.
- The first step to re-introduce NLO terms for Z decay into PHOTOS was to check relations between 4-vectors and angles, see some last transparencies.
- For PHOTOS those angles are defined in complicated way. Incoming beams (defining Z polarisations) have to be along z axis, to install old (1983) NLO formula.
- To match this requirement PHOTOS has to read in events from appropriately orienting Born level events generator.
- The fully differential distribution from MUSTRAAL (used also in KORALZ for single photon mode) reads:

$$X_f = \frac{Q'^2 \alpha (1 - \Delta)}{4\pi^2 s} s^2 \left\{ \frac{1}{(k'_+ k'_-)} \left[ \frac{\mathrm{d}\sigma_B}{\mathrm{d}\Omega}(s, t, u') + \frac{\mathrm{d}\sigma_B}{\mathrm{d}\Omega}(s, t', u) \right] \right\}$$

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Here:

$$s = 2p_{+} \cdot p_{-}, \quad s' = 2q_{+} \cdot q_{-},$$
  
 $t = 2p_{+} \cdot q_{+}, \quad t' = 2p_{+} \cdot q_{-},$   
 $u = 2p_{+} \cdot q_{-}, \quad u' = 2_{-} \cdot q_{+},$   
 $k'_{\pm} = q_{\pm} \cdot k, \quad x_{k} = 2E_{\gamma}/\sqrt{s}$ 

- The  $\Delta$  term is responsable for final state mass dependent terms,  $p_+$ ,  $p_-$ ,  $q_+$ ,  $q_-$ , k denote four-momenta of incoming positron, electron beams, outcoming muons and bremsstrahlung photon.
- after trivial manipulation it can be written as:

$$X_{f} = \frac{Q'^{2}\alpha(1-\Delta)}{4\pi^{2}s}s^{2} \left\{ \frac{1}{(k'_{+}+k'_{-})} \frac{1}{k'_{-}} \left[ \frac{d\sigma_{B}}{d\Omega}(s,t,u') + \frac{d\sigma_{B}}{d\Omega}(s,t',u) \right] + \frac{1}{(k'_{+}+k'_{-})} \frac{1}{k'_{+}} \left[ \frac{d\sigma_{B}}{d\Omega}(s,t,u') + \frac{d\sigma_{B}}{d\Omega}(s,t',u) \right] \right\}$$

• In PHOTOS the following expression is used:

$$\begin{split} X_f^{PHOTOS} &= \frac{Q'^2\alpha(1-\Delta)}{4\pi^2s}s^2 \Bigg\{ \\ &\frac{1}{k'_+ + k'_-}\frac{1}{k'_-} \quad \left[ (1+(1-x_k)^2)\frac{\mathrm{d}\sigma_B}{d\Omega} \left(s,\frac{s(1-\cos\Theta_+)}{2},\frac{s(1+\cos\Theta_+)}{2}\right) \right] \frac{(1+\beta\cos\Theta_\gamma)}{2} \\ &+ \frac{1}{k'_+ + k'_-}\frac{1}{k'_+} \quad \left[ (1+(1-x_k)^2)\frac{\mathrm{d}\sigma_B}{d\Omega} \left(s,\frac{s(1-\cos\Theta_-)}{2},\frac{s(1+\cos\Theta_-)}{2}\right) \right] \frac{(1-\beta\cos\Theta_\gamma)}{2} \Bigg\} \\ &\qquad \qquad \qquad \qquad \qquad \qquad \qquad \\ \text{where} : \Theta_+ &= \angle(p_+,q_+), \; \Theta_- = \angle(p_-,q_-) \\ &\qquad \qquad \qquad \qquad \qquad \qquad \\ \Theta_\gamma &= \angle(\gamma,\mu^-) \; \text{are defined in } (\mu^+,\mu^-) \text{-pair rest frame} \end{split}$$

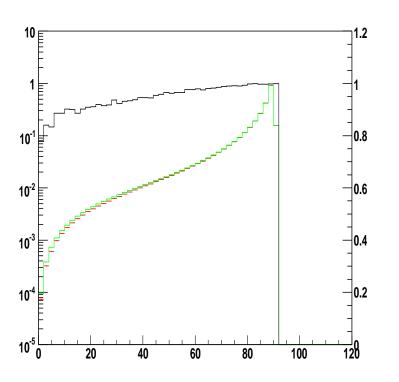
- The two expressions define weight to make with PHOTOS complete first order.
- Here we give up the results of 1991 work!

### The differences are important

- The two expressions define weight to make out of PHOTOS complete first order.
- The PHOTOS expression separates (i) Final state bremsstrahlung (ii) electroweak parameters of the Born Cross section (iii) Initial state bremsstrahlung that is orientation of the spin quantization axix for Z.
- The constraints can be overcomed but the price is discipline of event record.
- Of course all this has to be understood in context of Leading Pole approximation. For example initial-final state interference breaks the simplification. Limitations need to be controlled: Phys.Lett.B219:103,1989

Now results of tests:

# $Z \to \mu^+ \mu^-$ standard PHOTOS vs. Matrix Element.



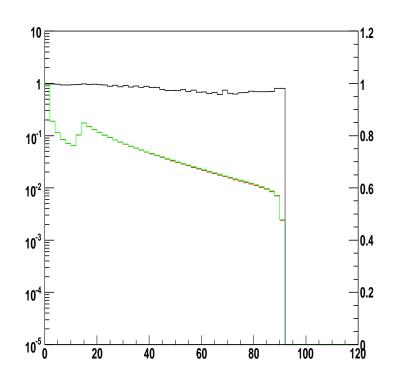
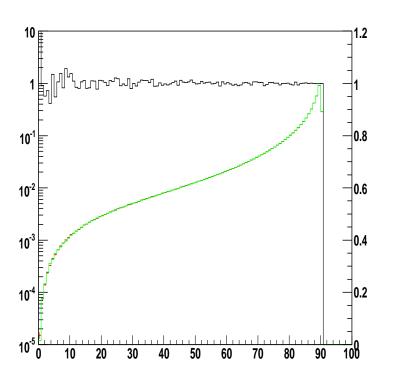


Figure 1: Comparison of standard PHOTOS and KORALZ: single photon emission level. On the left hand side the invariant mass of the  $\mu^+\mu^-$  pair; SDP=0.00534. On right hand side the invariant mass of  $\mu^-\gamma$ ; SDP=0.00296. The distributions for  $\mu^+\gamma$  are identical to  $\mu^-\gamma$ . The histograms produced by the two programs (logarithmic scale) and their ratio (linear scale, black line) are plotted on both figures. Test1, as defined in Section 3, is used, overall SDP=0.00534, fraction of events with hard photon was 17.4863  $\pm$  0.0042% for KORALZ and 17.6378  $\pm$  0.0042% for PHOTOS.

# $Z \to \mu^+ \mu^-$ PHOTOS vs. Matrix Element.



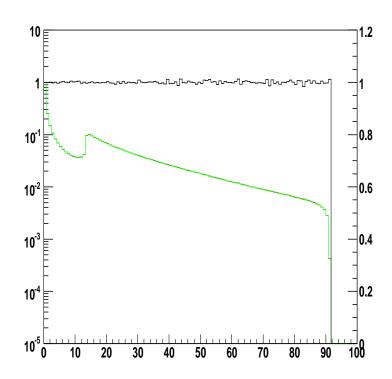
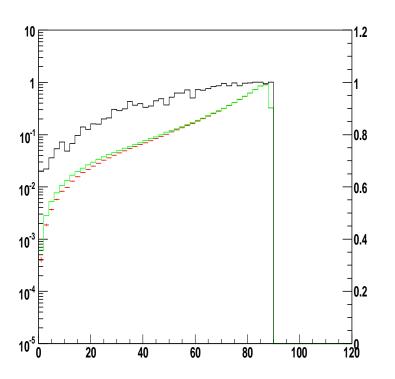


Figure2 Comparisons of improved PHOTOS and KORALZ: single photon emission level. On the left hand side the invariant mass of the  $\mu^+\mu^-$  pair. On right hand side the invariant mass of  $\mu^-\gamma$  is shown. In both cases differences between PHOTOS and KORALZ are below statistical error. As in Fig 1 distributions for  $\mu^+\gamma$  are skipped. Test1, as defined in Section 3, is used, overall SDP=0.0, fraction of events with hard photon was 17.4890  $\pm$  0.0042% for KORALZ and 17.4926  $\pm$  0.0042% for PHOTOS.

# $Z \to \mu^+ \mu^-$ PHOTOS vs. Matrix Element. $\gamma, \gamma$ conf.



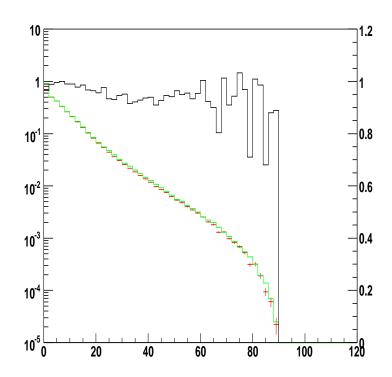
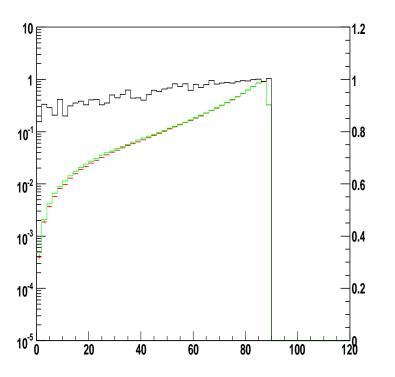


Figure 5: Comparisons of standard PHOTOS with multiple photon emission and KKMC with second order matrix element and exponentiation. Two comparison figures of worst agreement were selected from 2 hard photon configurations. On the left hand side the invariant mass of the  $\mu^+\mu^-$  pair is shown; SDP= 0.00918. On the right hand side the invariant mass of the  $\gamma\gamma$  pair; SDP=0.00268. Test2, as defined in Section 3, is used, overall SDP=0.00918, fraction of events with two hard photons was 1.2659  $\pm$  0.0011% for KORALZ and 1.2952  $\pm$  0.0011% for PHOTOS.

# $Z \to \mu^+ \mu^-$ NLO PHOTOS vs. Matrix Element $\gamma, \gamma$ conf.



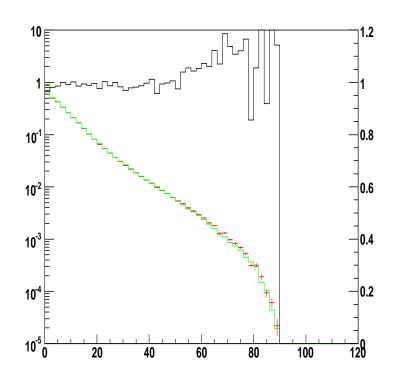


Figure 6: Comparisons of improved PHOTOS with multiple photon emission and KKMC with second order matrix element and exponentiation. Two comparison figures of worst agreement were selected from two-hard-photon configurations. On the left hand side the invariant mass of the  $\mu^+\mu^-$  pair is shown; SDP= 0.00142. On the right hand side the invariant mass of the  $\gamma\gamma$ ; SDP=0.00293. Test2, as defined in Section 3, was used, overall SDP= 0.00293, fraction of events with two hard photons was 1.2659  $\pm$  0.0011% for KORALZ and 1.2868  $\pm$  0.0011% for PHOTOS.

Figure 1: Results from PHOTOS, standard version, and SANC for  $B^- \to \pi^0 K^-(\gamma)$  decay are superimposed on the consecutive plots. Standard distributions, as defined in the text, are used. Logarithmic scales are used. The distributions from the two plots overlap almost completely. Samples of  $10^9$  events were used.

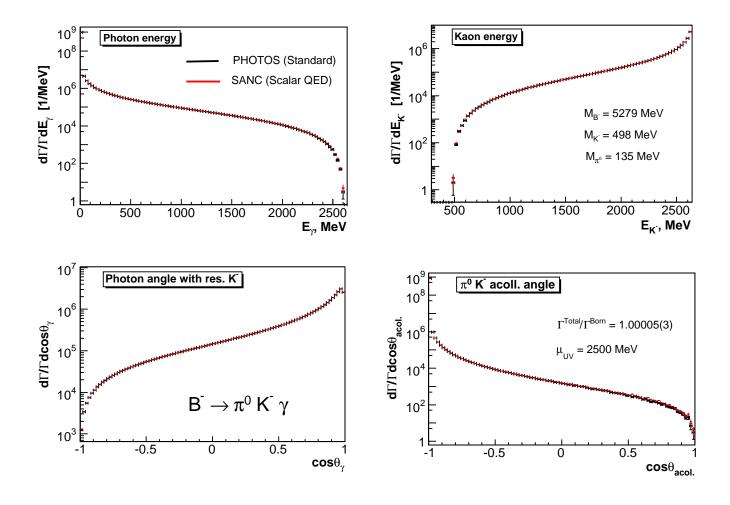


Figure 2: Results from PHOTOS, standard version, and SANC for ratios of the  $B^- \to \pi^0 K^-(\gamma)$  distribution in fi g.1 are presented. Differences between PHOTOS and SANC are small, but are clearly visible now.

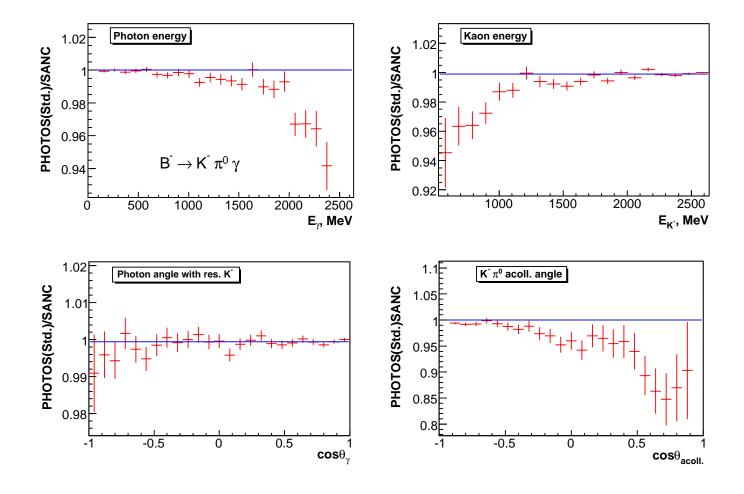
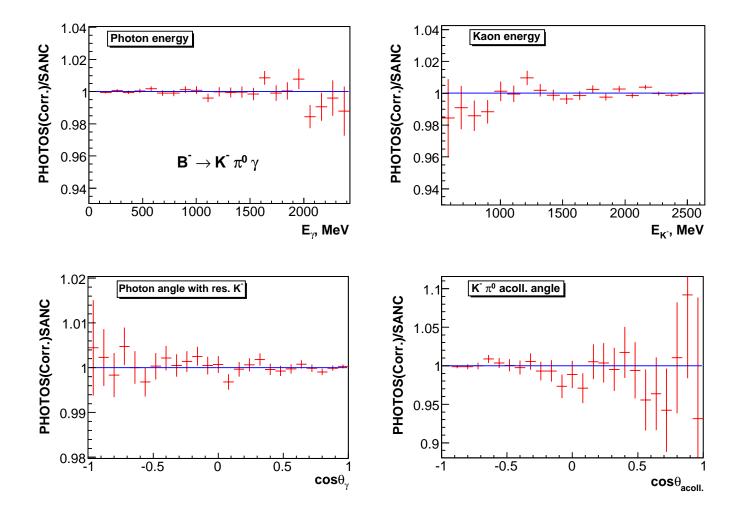
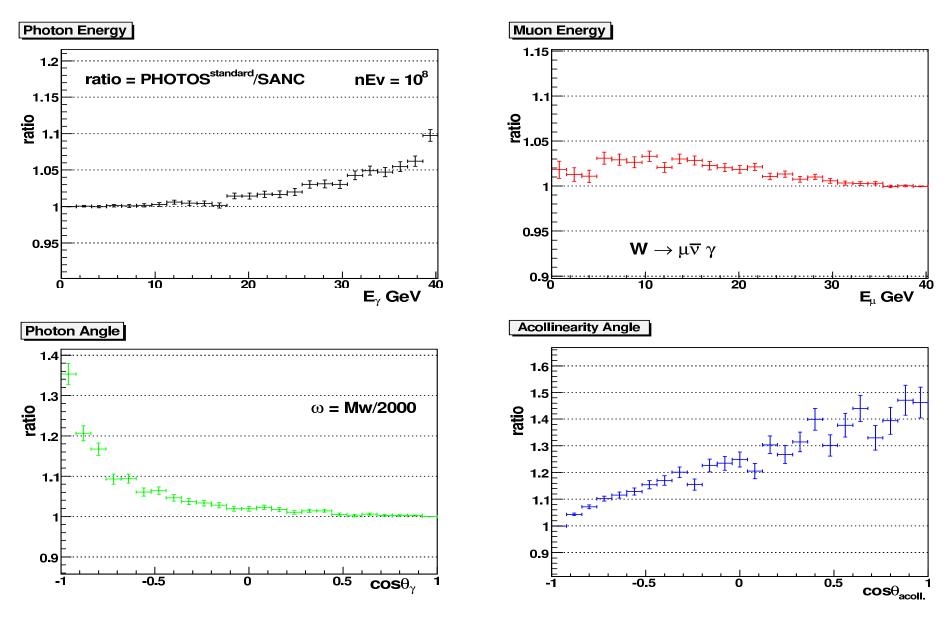
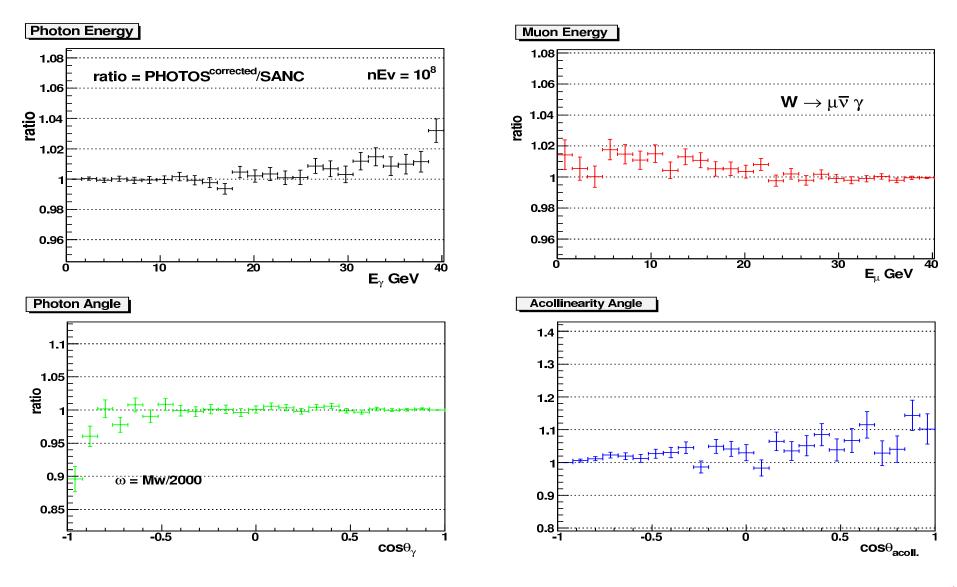


Figure 3: Results from PHOTOS with the exact matrix element, and SANC for ratios of the  $B^- \to \pi^0 K^-(\gamma)$  distributions. Differences between PHOTOS and SANC are below statistical error for samples of  $10^9$  events.

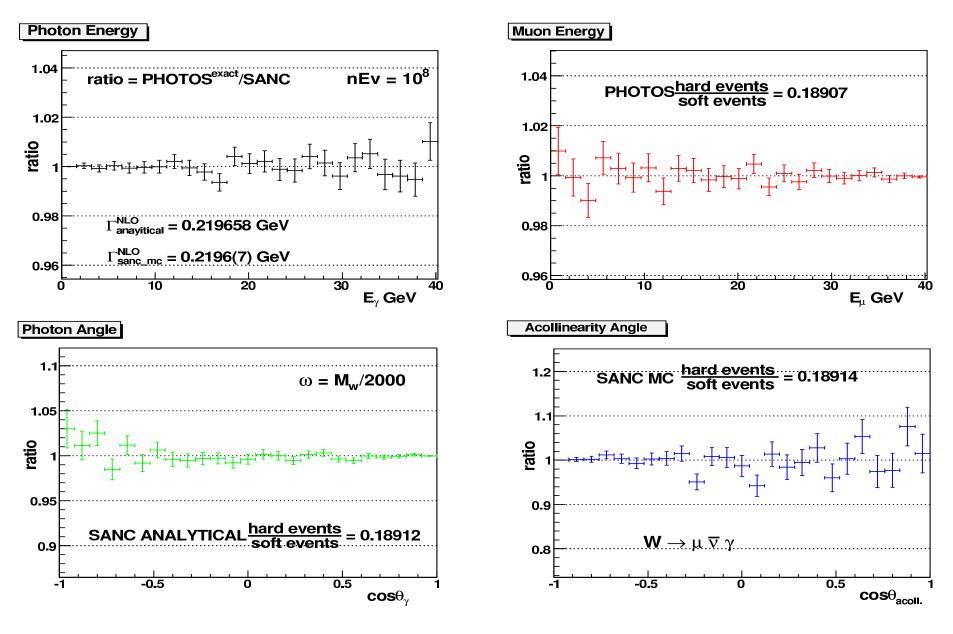




Results from PHOTOS without any correcting weight, and SANC for ratios of the  $W \to l\nu(\gamma)$  distributions. Distribution shapes are similar to those of B- deays (and we skip them).



Results from PHOTOS with the corecting weight of 2003, and SANC for ratios of the  $W \to l\nu(\gamma)$  distributions.



Results from PHOTOS with the exact matrix element, and SANC for ratios of the  $W \to l\nu(\gamma)$  distributions. Differences are below statistical error of  $10^8$  events..

# **Conclusion on PHOTOS**

- ullet PHOTOS work exellently for statistical samples of up to  $10^9$  events
- Matrix element used in generation is explicitly localized
- Even though in case of scalars as decay products, results means not much direct progress in itself, possibility to play with measurements of form-factors opens.
- Standard PHOTOS is most probably sufficient for many years to go.
- In case of doubts or 'big' measuremens technique of validation is prepared.
- Extensions to QCD is an open possibility now.

# **TAUOLA** universal interface

- To run, generator for tau decays must be combined with the part for tau production.
- In cases of our packages such as KORALB, KORALZ, KKMC host programs provide environment for TAUOLA use.
- I will concentrate on physics points in case when only information from event records is used.
- I will skip technicalities related to the way how HEPEVT common block is filled in 3 versions of PYTHIA conventions and HERWIG.
- also I will skip new developments in domain of event records.
- ullet TAUOLA universal interface reads information from HEPEVT common block, there au leptons to be decayed are found,
- ullet and their spin states are calculated from kinematical configurations of hard processes leading to au's.

# Formalism for $\tau^+\tau^-$

• Because narrow  $\tau$  width approximation can be obviously used for phase space, cross section for the process  $f\bar{f}\to \tau^+\tau^-Y; \tau^+\to X^+\bar{\nu}; \tau^-\to \nu\nu$  reads:

$$d\sigma = \sum_{spin} |\mathcal{M}|^2 d\Omega = \sum_{spin} |\mathcal{M}|^2 d\Omega_{prod} \ d\Omega_{\tau^+} \ d\Omega_{\tau^-}$$

- This formalism is fine, but because of over 20  $\tau$  decay channels we have over 400 distinct processes. Also picture of production and decay are mixed.
- but (only  $\tau$  spin indices are explicitly written):

$$\mathcal{M} = \sum_{\lambda_1 \lambda_2 = 1}^2 \mathcal{M}^{prod}_{\lambda_1 \lambda_2} \, \mathcal{M}^{ au^+}_{\lambda_1} \mathcal{M}^{ au^-}_{\lambda_2}$$

Formula for the cross section can be re-written

$$d\sigma = \left(\sum_{spin} |\mathcal{M}^{prod}|^2\right) \left(\sum_{spin} |\mathcal{M}^{\tau^+}|^2\right) \left(\sum_{spin} |\mathcal{M}^{\tau^-}|^2\right) wt \ d\Omega_{prod} \ d\Omega_{\tau^+} \ d\Omega_{\tau^-}$$

where

$$wt = \left(\sum_{i,j=0,3} R_{ij}h^{i}h^{j}\right)$$
  
 $R_{00} = 1, < wt >= 1, 0 \le wt \le 4.$ 

 $R_{ij}$  can be calculated from  $\mathcal{M}_{\lambda_1\lambda_2}$  and  $h^i,h^j$  respectively from  $\mathcal{M}^{\tau^+}$  and  $\mathcal{M}^{\tau^-}$ .

ullet Bell inequalities tell us that it is impossible to re-write wt in the following form

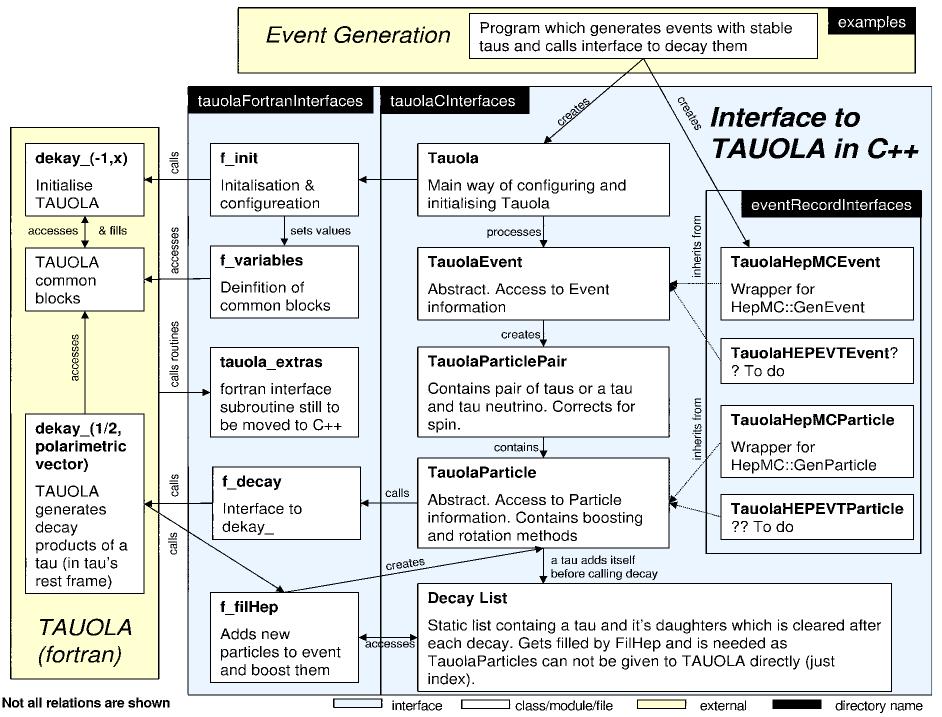
$$wt \neq \left(\sum_{i,j=0,3} R_i^A h^i\right) \left(\sum_{i,j=0,3} R_j^B h^j\right)$$

that means it is impossible to generate first  $\tau^+$  and  $\tau^-$  first in some given 'quantum state' and later perform separately decays of  $\tau^+$  and  $\tau^-$ 

- It can be done only if approximations are used !!!
- May be often reasonable, but nonetheless approximations.

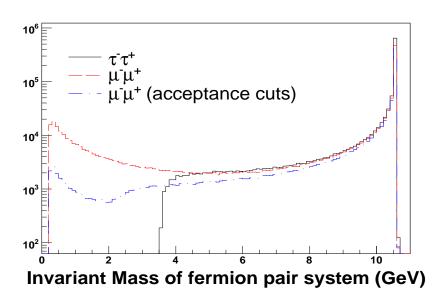
# Also in C++

- 1. Many young people prefer that language ...
- 2. Many new LHC generators use C++ HepMC structure to store events ...
- 3. Universal interface is mostly about helping experimental physicist. It must remain in their hands, but it is not their main worry or work direction.



web-page http://www.ph.unimelb.edu.au/ ndavidson/tauola/doxygen/index.html  $\beta$  version. No spin effects yet.

# For 0.1% precision at low energies ...



- 1. vacuum polarization
- 2. mass terms also in interference
- 3. pair corrections
- 4. S. Banerjee et al. Phys.Rev.D77:054012,2008. Precision tag for 10 GeV CMS ebe is just 0.35 %  $(\tau^+\tau^-)$  and 0.44 %  $(\mu^+\mu^-)$ . But it was quickfi x.
- 5. Coulomb interaction
- 6. hadronic currents instead of  $\mu^+\mu^-$  final states?

.. one need technical work

This is mostly question of manpower and organization of workshare.

Nothing of what in principle is OK at high energies will work without attention

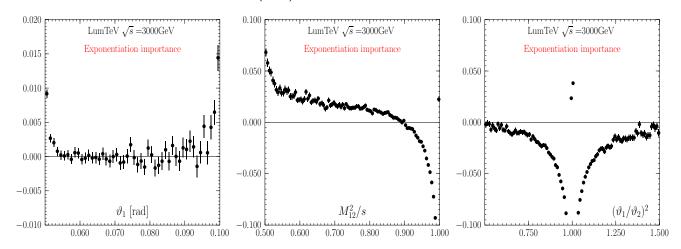
We are basically as in: http://home.cern.ch/jadach/public/frascati\_nov\_2001.pdf

Luminosity at LC

14

#### Is exponentiation important for photonic QED r.c.'s?

The difference  $\mathcal{O}(\alpha^3)_{exp} - \mathcal{O}(\alpha^2)$  in LL approximation (ISR only) gives us hint how bad the calculation in  $\mathcal{O}(\alpha^2)$  without exponentiation actually would be:



Conclusion: Exponentiation of photonic QED is absolute necessity!

Note that  $M_{12}/s=z_1z_2$  and  $\vartheta_1/\vartheta_2\simeq z_1/z_2$  are basic variables for determination of the luminosity distribution. Effects close to  $\vartheta$ -edges are due to soft ISR photons.

S. Jadach

SLAC October 24-th, 2002

**from** http://home.cern.ch/jadach/public/LumLCslac.pdf Warning on  $lpha^3$  and cut-off dependence

Z. Was



- We have reviewed news on the following tools for simulation of  $\tau$  physics:
  - TAUOLA as generator for  $\tau$  decays: additinal weights to be used for fi ts.
  - PHOTOS as generator for radiative corrections in decays: new precision results for W decays.
  - benchmark distribution strategies and MC-TESTER now also for C++ applications
  - Universal interface of TAUOLA, migration to C++ is on its way
  - Uncertainty of total cross section at lower energies as predicted by KKMC is now reduced to 0.3-0.5 %, if appropriate imrovements of photon vacuum polarzation are installed (S. Banerjee, Phys.Rev. D77:054012,2008).
  - stability of C++ event records, need verification still
  - encapsulation and inconvenience for end users, problems seem to be under control.



- TAUOLA and associated programs seem to be a living project
- As in the past different parametrizations will be developed within collaborations.
   Also, as in the past, will function as private code. Huge machinery.
- Some "cross talk" may be useful. Non-tau experiments like LHC may profit.
- We have prepared, updated TAUOLA version, with open slots for many new currents to be studied simultaneously.
- High precison of PHOTOS verified for W decays (NLO kernel). We need to enter into phenomenological project where analysis of data is essential.
- Fortran to C++ shift is mainly community issue, technicaly not a problem.
- There is little activity on KKMC and BHLUMI. Some work to control perecision at lower energies is a must. Especially for vacuum polarization and effects like Coulomb interaction.
- Let us have private discussions now.