

# Radiative $K_{e3}$ decays revisited



Bastian Kubis



HISKP (Theory), Universität Bonn



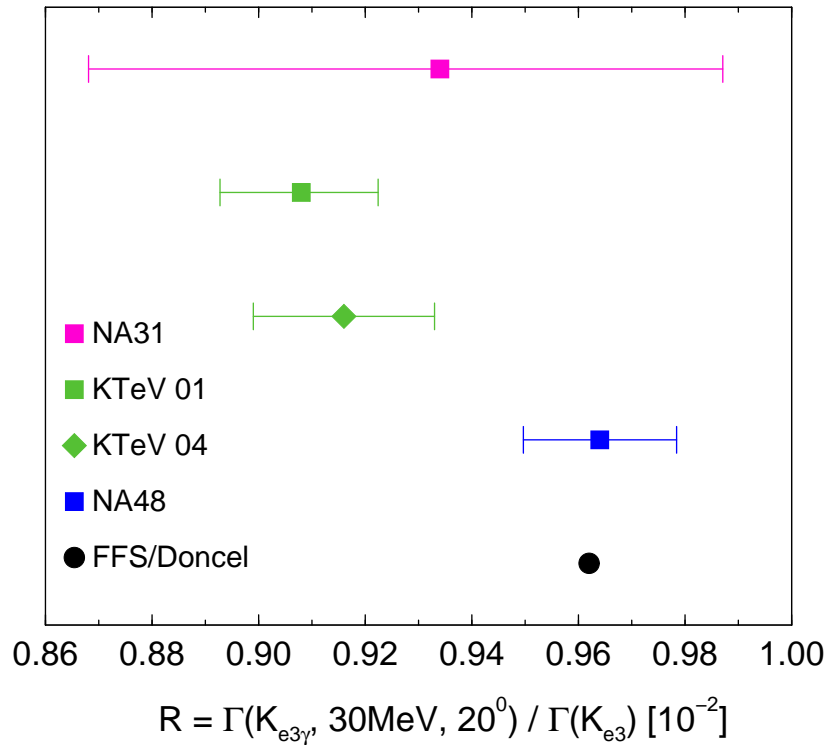
in collaboration with Jürg Gasser, Nello Paver, Michela Verbeni

EURIDICE Collaboration Meeting  
INFN Frascati, 9/2/2005

$$K_L \rightarrow \pi^\mp e^\pm \nu_e \gamma \quad [K_{e3\gamma}]$$

- Experimental status
- Introduction: ChPT and Low's theorem
- Formalism:  $K_{e3}$ ,  $K_{e3\gamma}$  amplitudes
- Observables: decay widths, the ratio  $R$
- Electromagnetic corrections
- Structure dependent terms from differential rates
- Summary

## Experimental status



KTeV 2001:

⇒ "... significantly lower than all published theoretical predictions."

KTeV, Phys. Rev. D64 (2001) 112004

latest results confirm this...

KTeV, Phys. Rev. D71 (2005) 012001

... or don't:

NA48, Phys. Lett. B605 (2005) 247

# Introduction

## Chiral Perturbation Theory

- expansion in small masses ( $M_\pi, M_K$ ) and momenta
- equivalent to loop expansion;  $K \rightarrow \pi \ell \nu_e \gamma$  [ $K_{\ell 3 \gamma}$ ] up to one loop:

J. Bijnens, G. Ecker, J. Gasser, Nucl. Phys. B396 (1993) 81

## Low's Theorem

- radiative processes  $X \rightarrow Y \gamma$  involving charged particles are **infrared divergent** for photon momentum  $q \rightarrow 0$ :

$$T(X \rightarrow Y \gamma) = \underbrace{\frac{T_{-1}}{q} + T_0}_{\text{inner bremsstrahlung (IB)}} + \underbrace{T_1 q + \dots}_{\text{structure dependent (SD)}}$$

- **Low's theorem**:  $T_{-1}, T_0$  given in terms of the **non-radiative** amplitude (and derivatives thereof):

$$T_{-1}, T_0 \propto T(X \rightarrow Y), \partial T(X \rightarrow Y)$$

F.E. Low, Phys. Rev. 110 (1958) 974

## Strategy: ChPT + Low's Theorem

- chiral representation fulfils Low's theorem
- expect bremsstrahlung to dominate due to infrared singularity  
 $\Rightarrow$  improve on chiral representation by using **phenomenological amplitudes** for  $T_{-1}, T_0$
- ChPT for structure dependent contributions

## Formalism

- amplitude for  $K_{e3}$  described in terms of single form factor  $f_+(t)$ ,  
parametrisation  $f_+(t) = f_+(0) \left\{ 1 + \lambda_+ \frac{t}{M_\pi^2} \left( + \lambda_+'' \frac{t^2}{M_\pi^4} \right) \right\}$
- amplitude for  $K_{e3\gamma}^0$  ( $K^0(p) \rightarrow \pi^-(p') e^+(p_e) \nu_e(p_\nu) \gamma(q)$ ):

$$\begin{aligned}
 T(K_{e3\gamma}^0) = & e G_F V_{us}^* \epsilon^\mu(q)^* \left[ (V_{\mu\nu} - A_{\mu\nu}) \underbrace{\bar{u}(p_\nu) \gamma^\nu (1 - \gamma_5) v(p_e)}_{\text{leptonic weak current}} \right. \\
 & \left. + \underbrace{f_+(t) \frac{(p + p')_\nu}{2p_e q} \bar{u}(p_\nu) \gamma^\nu (1 - \gamma_5) (m_e - \not{p}_e - \not{q}) \gamma_\mu v(p_e)}_{\text{photon radiation from positron}} \right]
 \end{aligned}$$

## Characterisation of $V_{\mu\nu}$ , $A_{\mu\nu}$ :

- decomposition  $V_{\mu\nu} = V_{\mu\nu}^{\text{IB}} + V_{\mu\nu}^{\text{SD}}$ ,

$V_{\mu\nu}^{\text{IB}} \simeq$  photon radiation from pion, depends *only* on  $f_+$

H.W. Fearing, E. Fischbach, J. Smith, Phys. Rev. D2 (1970) 542

- Ward identities:

$$q^\mu V_{\mu\nu}^{\text{IB}} = f_+(t) (p + p')_\nu, \quad q^\mu V_{\mu\nu}^{\text{SD}} = q^\mu A_{\mu\nu} = 0$$

- $V_{\mu\nu}^{\text{SD}}$ ,  $A_{\mu\nu}$  decomposed in terms of 8 functions  $V_{1-4}$ ,  $A_{1-4}$  :
  - $\Rightarrow V_3, A_3$  suppressed by  $m_e^2/M_K^2 \approx 10^{-6}$ , not observable in  $K_{e3\gamma}$
  - $\Rightarrow V_4, A_4$  suppressed by two orders in the chiral expansion
  - $\Rightarrow$  (essentially) 4 functions  $V_{1/2}, A_{1/2}$  for SD contribution

## Chiral one-loop prediction for structure dependent part:

$$V_1 = -\frac{8}{F_\pi F_K} L_9 + (\pi, K, \eta)\text{-loops}$$

$$V_2 = \frac{4}{F_\pi F_K} (L_9 + L_{10}) + (\pi, K, \eta)\text{-loops}$$

$$A_1 = 0$$

$$A_2 = -\frac{1}{8\pi^2 F_\pi F_K}$$

- low-energy constants:  $L_9 \leftrightarrow \langle r^2 \rangle_\pi^V$  or  $\lambda_+$ ,  $L_{10} \leftrightarrow (\pi \rightarrow e\nu\gamma)$
- axial form factors  $A_i$  given in terms of **chiral anomaly**

J. Wess, B. Zumino, Phys. Lett. B37 (1971) 95

E. Witten, Nucl. Phys. B223 (1983) 422



## Observables: decay widths

- for the non-radiative and radiative widths:

$$\Gamma(K_{e3}) = \frac{M_K^5 G_F^2 |V_{us}|^2}{128\pi^3} f_+(0)^2 \times I$$

$$\Gamma(K_{e3\gamma}) = \frac{\alpha M_K^5 G_F^2 |V_{us}|^2}{16\pi^7} f_+(0)^2 \times I^\gamma(E_\gamma^{\text{cut}}, \theta_{e\gamma}^{\text{cut}})$$

in the following: “standard cuts”  $E_\gamma^{\text{cut}} = 30 \text{ MeV}$ ,  $\theta_{e\gamma}^{\text{cut}} = 20^\circ$

- accessible in experiments:

$$R = \frac{\Gamma(K_{e3\gamma})}{\Gamma(K_{e3})} = \frac{8\alpha}{\pi^4} \frac{I^\gamma}{I}$$

$\Rightarrow$  all sorts of constants cancel in the ratio

- $I, I^\gamma$  depend on  $\lambda_+$ :  $I^{(\gamma)} = a_0^{(\gamma)} + a_1^{(\gamma)} \lambda_+ + a_2^{(\gamma)} \lambda_+^2$  ,

$$\frac{I^{(\gamma)}(\lambda_+ = \lambda_+^{\text{exp}})}{I^{(\gamma)}(\lambda_+ = 0)} - 1 \approx 0.1 \quad \text{but} \quad \frac{R(\lambda_+ = \lambda_+^{\text{exp}})}{R(\lambda_+ = 0)} - 1 = \mathcal{O}(10^{-4})$$

form factor dependence  **cancels completely**  in  $R$ !

- in **ChPT**: only visible one-loop effect from **structure dependent** contributions:

$$\frac{R(\text{IB} + \text{SD})}{R(\text{IB})} - 1 \approx -0.01$$

⇒ **inner bremsstrahlung** completely **dominant**!

**tree prediction** that is accurate at the **1%**-level!

## Electromagnetic corrections

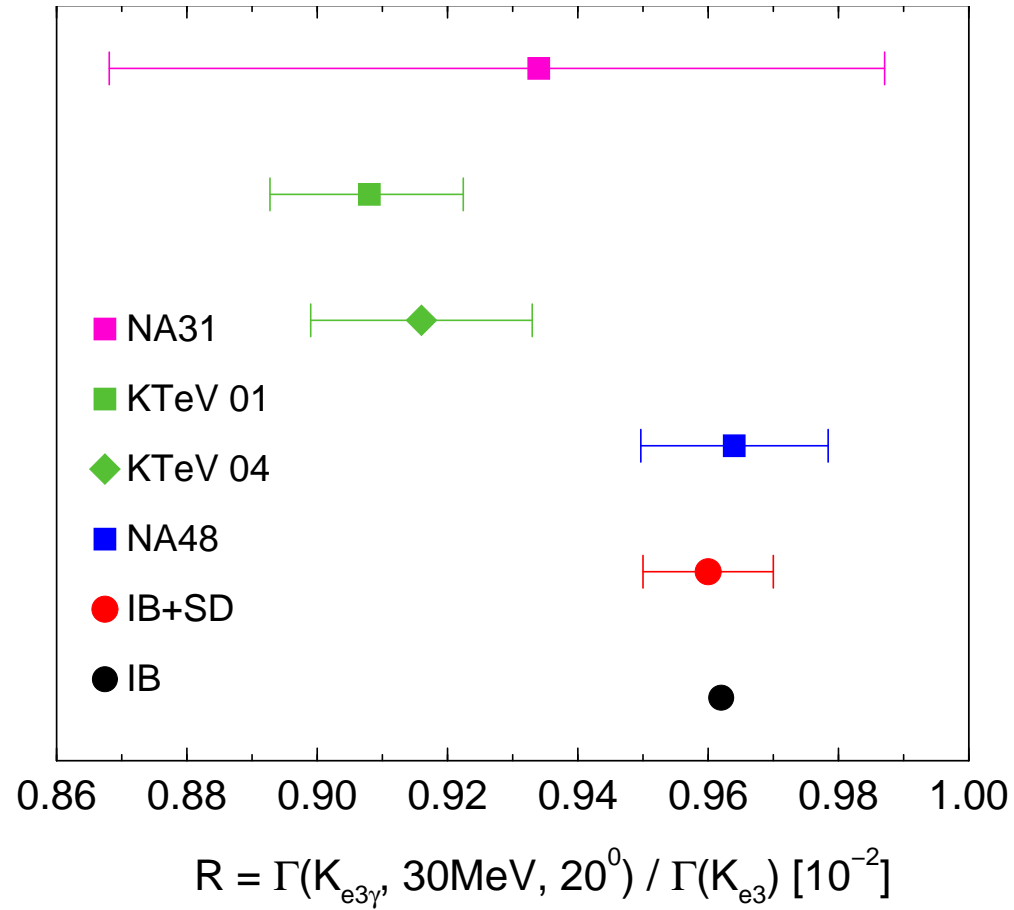
Hadronic prediction of  $\sim 1\%$  accuracy – radiative corrections?

$$R = \frac{\Gamma_{\text{incl}}(K_{e3\gamma}, E_\gamma^* > E_\gamma^{\text{cut}}, \theta_{e\gamma}^* > \theta_{e\gamma}^{\text{cut}})}{\Gamma_{\text{incl}}(K_{e3})}$$

- denominator: see [H. Neufeld's talk](#), net effect:  $I$  increased by  $0.2\%$
  - numerator: potential large logarithms  $\log \frac{m_e}{M_K}$  known (KLN theorem) largely absorbed in  $f_+(0)$ , remaining effect also  $0.2\%$  only
- $\Rightarrow$  assume non-enhanced corrections to be small,  $\pm 1.0\%$  uncertainty
- $\Rightarrow$  our estimate [Low + ChPT](#):

$$R = (0.96 \pm 0.01) \times 10^{-2}$$

# Comparison experiment/theory:



## A second look at structure dependent terms

- Idea: access information independent on branching ratios by studying (unnormalised) differential rates  $\Rightarrow V_{1/2}, A_{1/2}$
  - $d\Gamma/dE_\gamma^*$  most promising due to special role of  $E_\gamma^*$  in IB vs. SD
  - KTeV: study  $d\Gamma/dE_\gamma^*$  with
    1. real and constant structure functions
    2. two SD terms, neglect other two in “soft kaon approximation”
- $\Rightarrow$  ChPT: first assumption reasonable?
- $\Rightarrow$  check validity of the second assumption in distributions
- $\Rightarrow$  study also other distributions

# 1. Real & constant? – Analytic structure of $V_i, A_i$

- $V_i, A_i$  functions of **three** kinematical variables, e.g.

$$s = (p' + q)^2, \quad t = (p - p')^2, \quad u = (p - q)^2$$

$\Rightarrow$  three-body decay with variable “mass”  $W^2 = (p_e + p_\nu)^2$

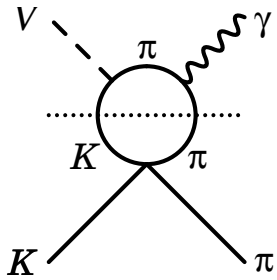
- Reminder:  $V_i$  at  $\mathcal{O}(p^4)$

$$V_1 = \underbrace{-\frac{8L_9}{F_\pi F_K}}_{\text{constant}} + \underbrace{\text{loops}}_{\text{non-trivial}}, \quad V_2 = \underbrace{\frac{4(L_9 + L_{10})}{F_\pi F_K}}_{\text{constant}} + \underbrace{\text{loops}}_{\text{non-trivial}}$$

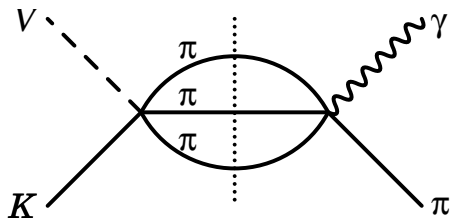
$V_i$  at  $\mathcal{O}(p^6)$ : linear terms in  $s, t, u, M_K^2, M_\pi^2$  + two-loop diagrams

- WZW–anomaly  $\mathcal{O}(p^4) \Rightarrow$  constant  $A_i$   
at  $\mathcal{O}(p^6)$  linear terms in  $s, t, u, M_K^2, M_\pi^2$  + one-loop diagrams

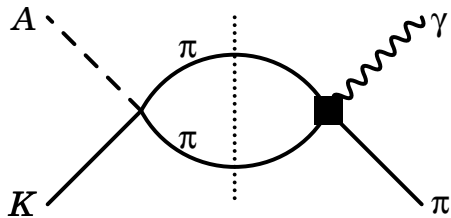
Diagrams corresponding to the most important cuts:



$\mathcal{O}(p^4)$ : only  $t$ -channel cut in  $V_i$   
 far outside physical region  $\Rightarrow$  smooth & flat



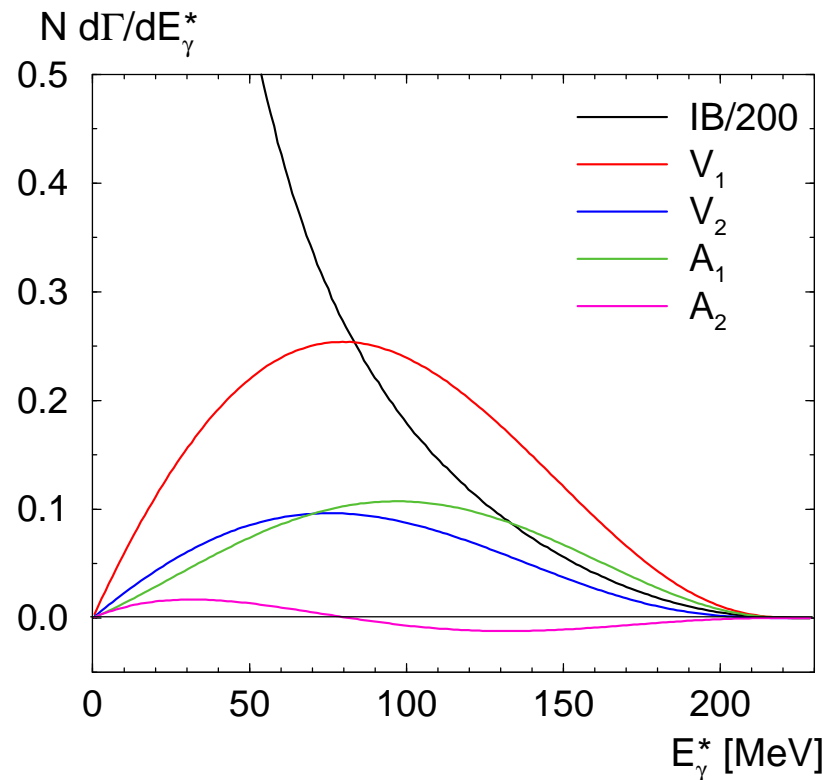
$\mathcal{O}(p^6)$ : “physical”  $s$ -channel cut in  $V_i$   
 $\hat{=}$  intermediate  $3\pi$  state  $\Rightarrow$  two-loop kinematically suppressed (phase space!)



$\mathcal{O}(p^6)$ : “physical”  $s$ -channel cut in  $A_i$   
 $\hat{=}$  intermediate  $2\pi$  state  $\Rightarrow$  one-loop + anomalous rescattering  $\rightarrow$  calculated  $\Rightarrow$  small!

dominant  $\mathcal{O}(p^6)$  are  $\mathcal{O}(M_K^2)$  corrections  $\Rightarrow$  still smooth & flat!

## 2. “Soft kaon” ? – structure dependent terms in $d\Gamma/dE_\gamma^*$



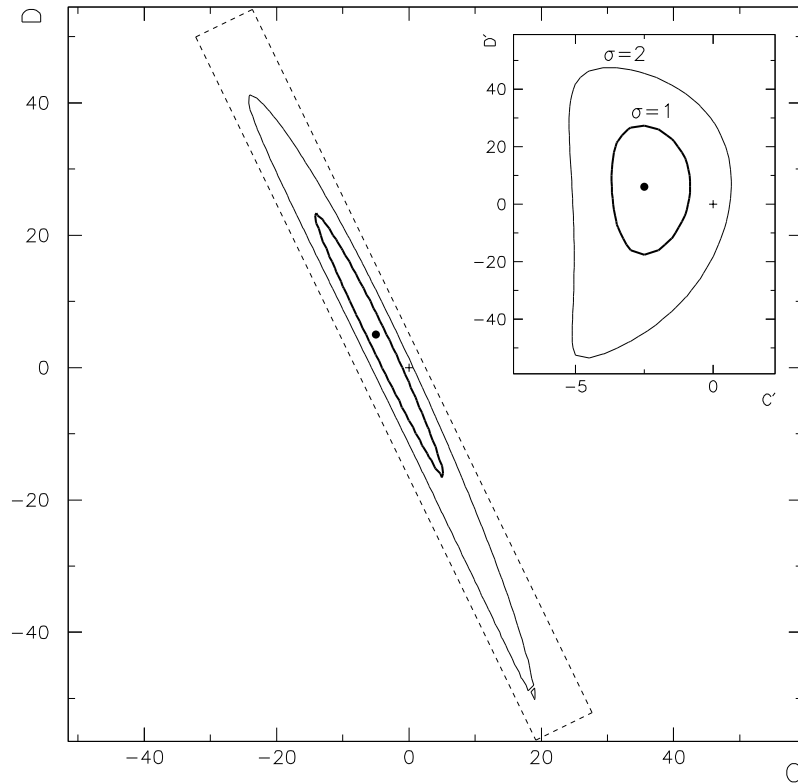
$d\Gamma/dE_\gamma^*$  is essentially sensitive to *one* linear combination:

$$\begin{aligned}
 C' &= 0.9 V_1 + 0.4 V_2 + 0.4 A_1 \\
 &= 0.9(V_1 - V_2) + \underbrace{1.3 V_2}_{\text{neglected?!}} \\
 &+ 0.4(A_1 + A_2) - \underbrace{0.4 A_2}_{\text{neglected?!}}
 \end{aligned}$$

⇒ “soft kaon” not a good approximation!



## Compare KTeV result to ChPT prediction:

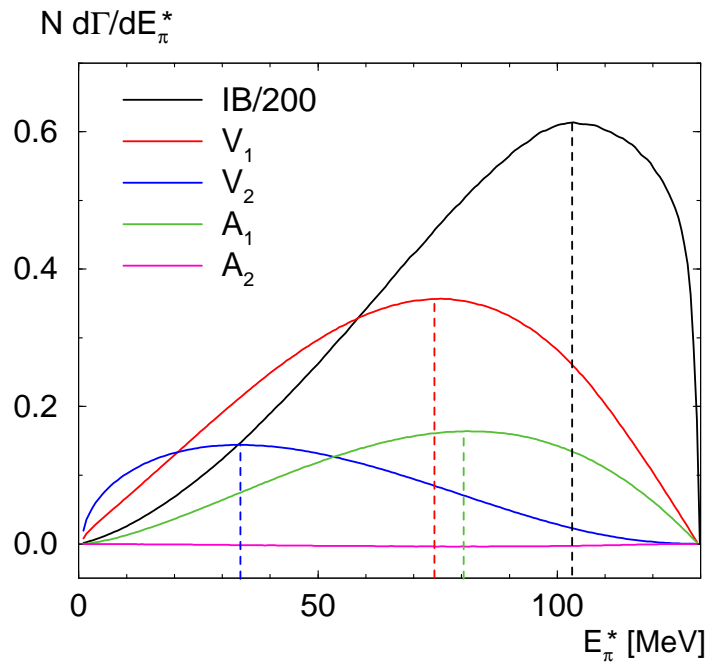


KTeV, Phys. Rev. D64 (2001) 112004

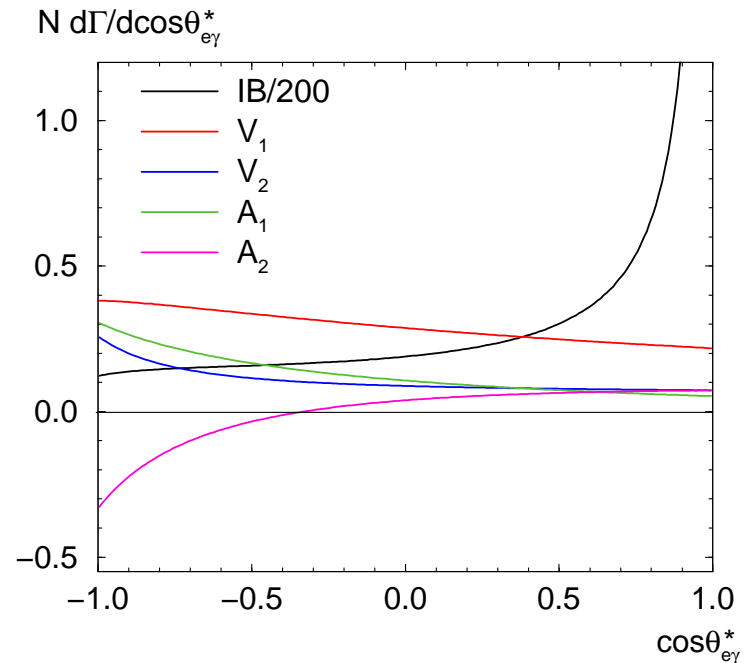
- KTeV:  $C' = -2.5_{-1.0}^{+1.5} \pm 1.5$
  - ChPT:  $C' = -1.6 \pm 0.4$
- $\Rightarrow$  good numerical agreement within  $1\text{-}\sigma$  error
- $\Rightarrow$  interpretation of  $C'$  in terms of structure functions different
- $\Rightarrow$  serious constraints on SD terms feasible!

Other distributions:  $d\Gamma/dE_\pi^*$ ,

$d\Gamma/d\cos\theta_{e\gamma}^*$



⇒ potentially sensitive to  $V_2$  :  
peaked at low pion energies



⇒ potentially sensitive to  $A_2$  :  
strong backward variation

## Summary (1)

Low's theorem + ChPT allow for a precise prediction of  $\Gamma(K_{e3\gamma})/\Gamma(K_{e3})$

- completely insensitive to details of  $K_{e3}$  form factor
- structure dependent terms very small
- constants like  $G_F$  or  $|V_{us}|$  cancel

Precision limited by radiative corrections:

$$R = \Gamma(K_{e3\gamma})/\Gamma(K_{e3}) = (0.96 \pm 0.01) \times 10^{-2}$$

vs.  $R = (0.908 \pm 0.008_{\text{stat}} \begin{smallmatrix} + 0.013 \\ - 0.012 \end{smallmatrix}_{\text{syst}}) \times 10^{-2}$  (KTeV 2001)

$R = (0.916 \pm 0.017_{\text{stat+syst}}) \times 10^{-2}$  (KTeV 2004)

$R = (0.964 \pm 0.008_{\text{stat}} \begin{smallmatrix} + 0.011 \\ - 0.009 \end{smallmatrix}_{\text{syst}}) \times 10^{-2}$  (NA48 2004)

## Summary (2)

Experimental extraction of **structure dependent terms** from  $d\Gamma/dE_\gamma^*$  shown to be **feasible** by KTeV collaboration:

$$C' = -2.5_{-1.0}^{+1.5} \text{stat} \pm 1.5_{\text{syst}} \text{ (KTeV)} \quad \text{vs.} \quad C' = -1.6 \pm 0.4 \text{ (ChPT)}$$

**ChPT**: these structure functions are fairly smooth  
imaginary parts and momentum dependence suppressed

Tentative ideas for experimentalists and theorists:

- extract  $V_2$  from  $d\Gamma/dE_\pi^*$  in order to disentangle  $V_1$  and  $V_2$  in  $C'$
- chiral anomaly in  $A_2$  at most accessible in  $d\Gamma/d\cos\theta_{e\gamma}^*$
- perform same study for **other  $K_{\ell 3\gamma}$  channels!**