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Isospin Breaking in $K \rightarrow 3\pi$ Decays

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Overview

- *K → 3π Decays in Chiral Perturbation Theory*, J. Bijnens, P. Dhonte and F. Persson, hep-ph/0205341, Nucl. Phys. B648 (2003) 317-344.
- *Isospin Breaking in K → 3π Decays I: Strong Isospin Breaking*, J. Bijnens and F. Borg, hep-ph/0405025, Nuclear Physics B697 (2004) 319-342.
- *Isospin Breaking in K → 3π Decays II: Radiative Corrections*, J. Bijnens and F. Borg, hep-ph/0410333, accepted in Eur. Phys. J. C.
- *Isospin Breaking in K → 3π Decays III: Bremsstrahlung and Fit to Experiment*, J. Bijnens and F. Borg, hep-ph/0501163.

Note: Fredrik Borg = Fredrik Persson, obtained PhD 28/1/2005

Overview

- ChPT in the nonleptonic mesonic sector
- Lagrangians
- $K \rightarrow 3\pi$ kinematics and isospin
- Overview of calculations and results
- Data and Fits
- Conclusions

ChPT in the nonleptonic sector

- some earlier work: especially on decays with photons
- Kambor, Missimer, Wyler (KMW) : Constructed \mathcal{L} and ∞ 1990
- G. Esposito-Farese: Checked \mathcal{L} and ∞ 1991
- KMW : Calculated $K \rightarrow 2\pi$ and $K \rightarrow 3\pi$ 1991
- Donoghue + Holstein + KMW : clarified the relations between observables 1992
- BUT: explicit formulas lost (US Mail)
- Kambor, Ecker, Wyler : simplified octet \mathcal{L} 1993
- $K \rightarrow 2\pi$ redone: Bijnens, Pallante, Prades 1998

ChPT in the nonleptonic sector

- First paper: redo $K \rightarrow 3\pi$
- Ecker Isidori Muller Neufeld Pich: Electromagnetic octet \mathcal{L} Lagrangian plus ∞ 2000
- applications to $K \rightarrow 2\pi$: Several papers
- Isospin breaking in $K \rightarrow 3\pi$: remaining papers
 $\mathcal{O}(p^4, p^2(m_u - m_d), e^2 p^2)$.

Lagrangians: p^2 and e^2

$$\mathcal{L}_2 = \mathcal{L}_{S2} + \mathcal{L}_{W2} + \mathcal{L}_{E2}$$

$$\mathcal{L}_{S2} = \frac{F_0^2}{4} \langle u_\mu u^\mu + \chi_+ \rangle$$

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$$u_\mu = i u^\dagger D_\mu U u^\dagger = u_\mu^\dagger, \quad u^2 = U, \quad \chi_\pm = u^\dagger \chi u^\dagger \pm u \chi^\dagger u,$$

U contains the Goldstone boson fields

$$U = \exp \left(\frac{i\sqrt{2}}{F_0} M \right), \quad M = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi_3 + \frac{1}{\sqrt{6}}\eta_8 & \pi^+ & K^+ \\ \pi^- & \frac{-1}{\sqrt{2}}\pi_3 + \frac{1}{\sqrt{6}}\eta_8 & K^0 \\ K^- & \overline{K^0} & \frac{-2}{\sqrt{6}}\eta_8 \end{pmatrix}.$$

$$\text{Here } \chi = 2B_0 \begin{pmatrix} m_u & & \\ & m_d & \\ & & m_s \end{pmatrix} \text{ and } D_\mu U = \partial_\mu U - ie [Q, U].$$

Lagrangians: p^2 and e^2

$$\mathcal{L}_{W2} = CF_0^4 \left[G_8 \langle \Delta_{32} u_\mu u^\mu \rangle + G'_8 \langle \Delta_{32} \chi_+ \rangle + G_{27} t^{ij,kl} \langle \Delta_{ij} u_\mu \rangle \langle \Delta_{kl} u^\mu \rangle \right]$$

$$\Delta_{ij} \equiv u \lambda_{ij} u^\dagger, \quad (\lambda_{ij})_{ab} \equiv \delta_{ia} \delta_{jb}, \quad t^{21,13} = t^{13,21} = 1/3, \dots$$

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$$C = -\frac{3}{5} \frac{G_F}{\sqrt{2}} V_{ud} V_{us}^*: \text{chiral and large } N_c \text{ limits } G_8 = G_{27} = 1$$

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$$\text{Electromagnetic: } \mathcal{L}_{E2} = e^2 F_0^4 Z \langle \mathcal{Q}_L \mathcal{Q}_R \rangle + (e^2 C F_0^6 G_E \langle \Delta_{32} \mathcal{Q}_R \rangle)$$

$$\mathcal{Q}_L = u Q u^\dagger, \quad \mathcal{Q}_R = u^\dagger Q u \quad \text{with} \quad Q = \text{diag} (2/3, -1/3, -1/3)$$

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Note: F_0 not well known, fit $CF_0^4 G_8, \dots$

use *numerically* $F_0 = F_\pi$ to quote G_8, \dots

Lagrangians: p^4 and $e^2 p^2$

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\mathcal{L}_{S4} values of L_i^r use Amoros, Bijnens, Talavera p^4 fit

\mathcal{L}_{W4} thirteen N_i^r (octet) and twelve D_i^r might contribute
(two N_i^r and two D_i^r extra for photon-reducible)

\mathcal{L}_{S2E2} eleven K_i^r can contribute

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Done here :

Isospin conserved: 11 combinations of N_i^r, D_i^r relevant: \tilde{K}_i
7 coefficients order m_K^4 , 4 order $M_K^2 m_\pi^2$
2 (virtually) indistinguishable from G_8 and G_{27} .

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Isospin Broken: 30 combinations of N_i^r, D_i^r, Z_i^r relevant
Results from isospin breaking with $K_i^r = Z_i^r = 0$

$K \rightarrow 3\pi$ Kinematics and Isospin

$$K_L(k) \rightarrow \pi^0(p_1)\pi^0(p_2)\pi^0(p_3), \quad [A_{000}^L],$$

$$K_L(k) \rightarrow \pi^+(p_1)\pi^-(p_2)\pi^0(p_3), \quad [A_{+-0}^L],$$

$$K_S(k) \rightarrow \pi^+(p_1)\pi^-(p_2)\pi^0(p_3), \quad [A_{+-0}^S],$$

$$K^+(k) \rightarrow \pi^0(p_1)\pi^0(p_2)\pi^+(p_3), \quad [A_{00+}^L],$$

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plus charge conjugate ones

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even under $p_1 \leftrightarrow p_2$ except A_{+-0}^S odd (CP and Bose symmetry)

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Kinematics: $s_1 = (k - p_1)^2, s_2 = (k - p_2)^2, s_3 = (k - p_3)^2.$

Dalitz plot variables:

$$y = (s_3 - s_0)/m_{\pi^+}^2, x = (s_2 - s_1)/m_{\pi^+}^2, s_0 = (s_1 + s_2 + s_3)/3.$$

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Bremsstrahlung problem: s_i from E_i or $m_{\pi\pi}^2$ different

Kinematics and Isospin

Valid also at p^6

$$\begin{aligned} A_{000}^L &= M_0(s_1) + M_0(s_2) + M_0(s_3), \\ A_{+-0}^L &= M_1(s_3) + M_2(s_1) + M_2(s_2) + M_3(s_1)(s_2 - s_3) + M_3(s_2)(s_1 - s_3), \\ A_{+-0}^S &= M_4(s_1) - M_4(s_2) + M_5(s_1)(s_2 - s_3) - M_5(s_2)(s_1 - s_3) + M_6(s_3)(s_1 - s_2), \\ A_{00+} &= M_7(s_3) + M_8(s_1) + M_8(s_2) + M_9(s_1)(s_2 - s_3) + M_9(s_2)(s_1 - s_3), \\ A_{++-} &= M_{10}(s_3) + M_{11}(s_1) + M_{11}(s_2) + M_{12}(s_1)(s_2 - s_3) + M_{12}(s_2)(s_1 - s_3). \end{aligned}$$

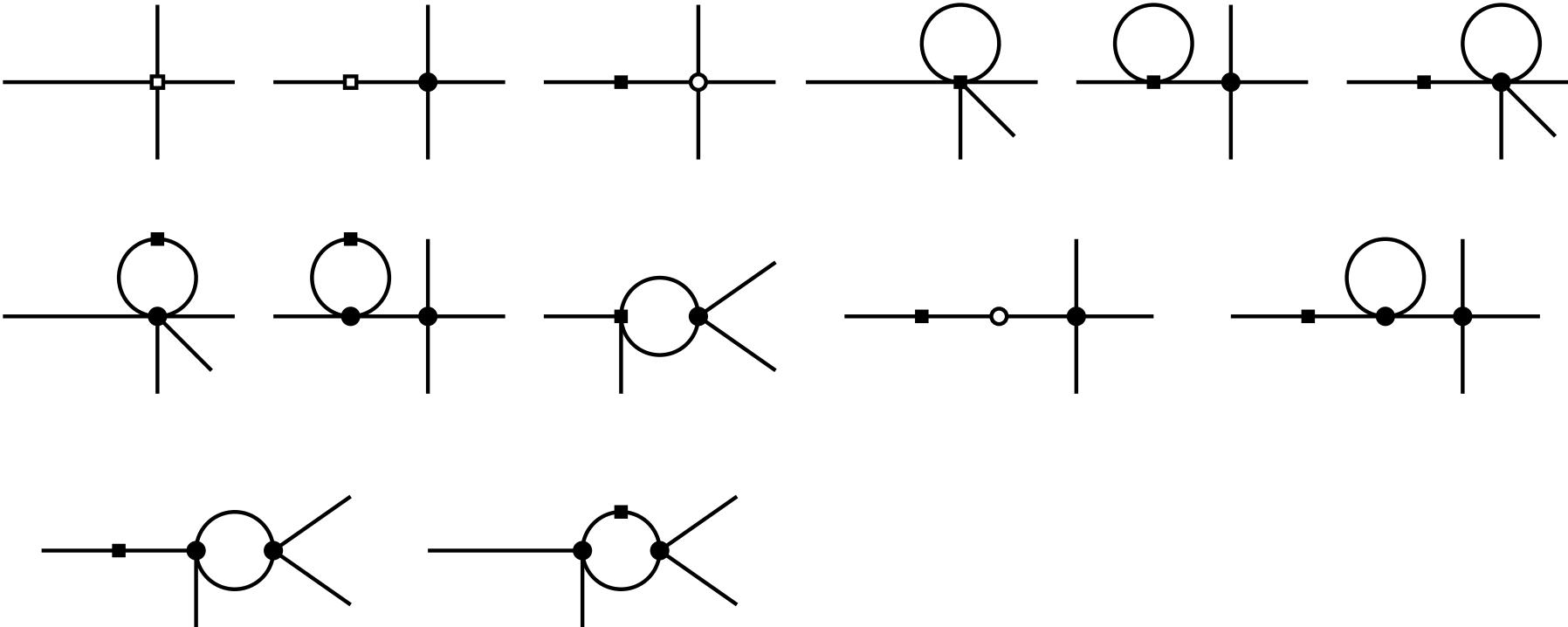
polynomial ambiguity from $s_1 + s_2 + s_3 = \sum_i m_i^2$

Isospin:

$$\begin{aligned} M_0(s) &= M_1(s) + 2M_2(s), \\ 2M_7(s) + 4M_8(s) &= M_{10}(s) + 2M_{11}(s), \\ M_4(s) &= \frac{1}{3} (M_7(s) - M_8(s) + M_{10}(s) - M_{11}(s)) \\ M_5(s) - M_6(s) &= M_9(s) + M_{12}(s). \end{aligned}$$

Calculations: isospin limit

The diagrams of order p^4



Expressions for the $M_i(s)$: see first paper

Confirmed by Gamiz, Prades, Scimemi, hep-ph/0305164
Lashin, hep-ph/0308200 (how not clear)

Calculations: I: Strong Isospin Breaking

Includes:

- $m_u - m_d$ and the effects of \mathcal{L}_{E2} , \mathcal{L}_{E2S2} and $\mathcal{L}_{W2E2}(G_8)$
- π^0 - η mixing (η' via LECs)
- Same diagrams as before
- Formulas immediately a lot longer:
see <http://www.thep.lu.se/~bijnens/chpt.html>

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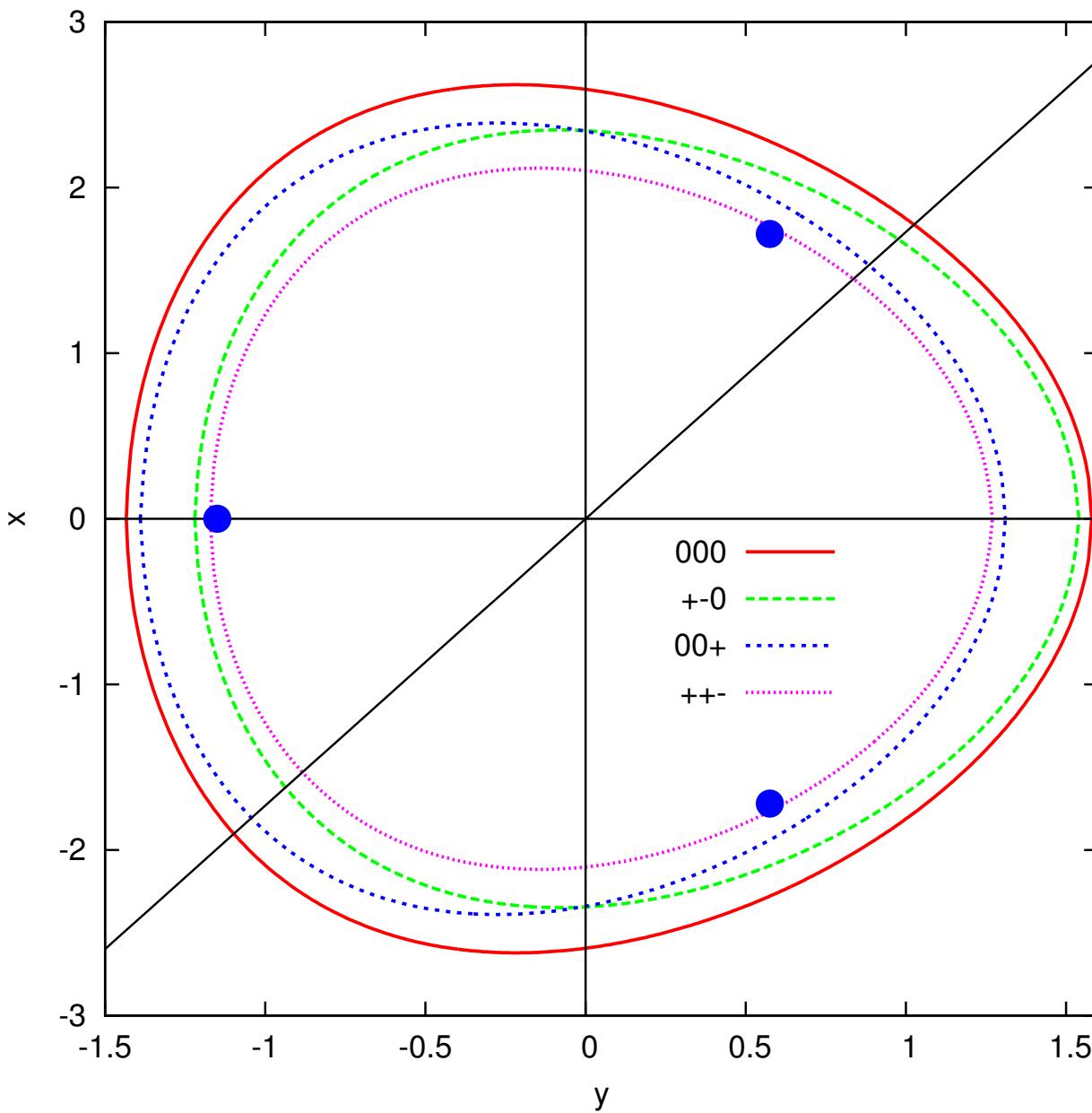
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Effects:

- F_{π^+} , F_{K^+} overall factor
- Mass difference $\pi^+ - \pi^0$
- Others typically a few %

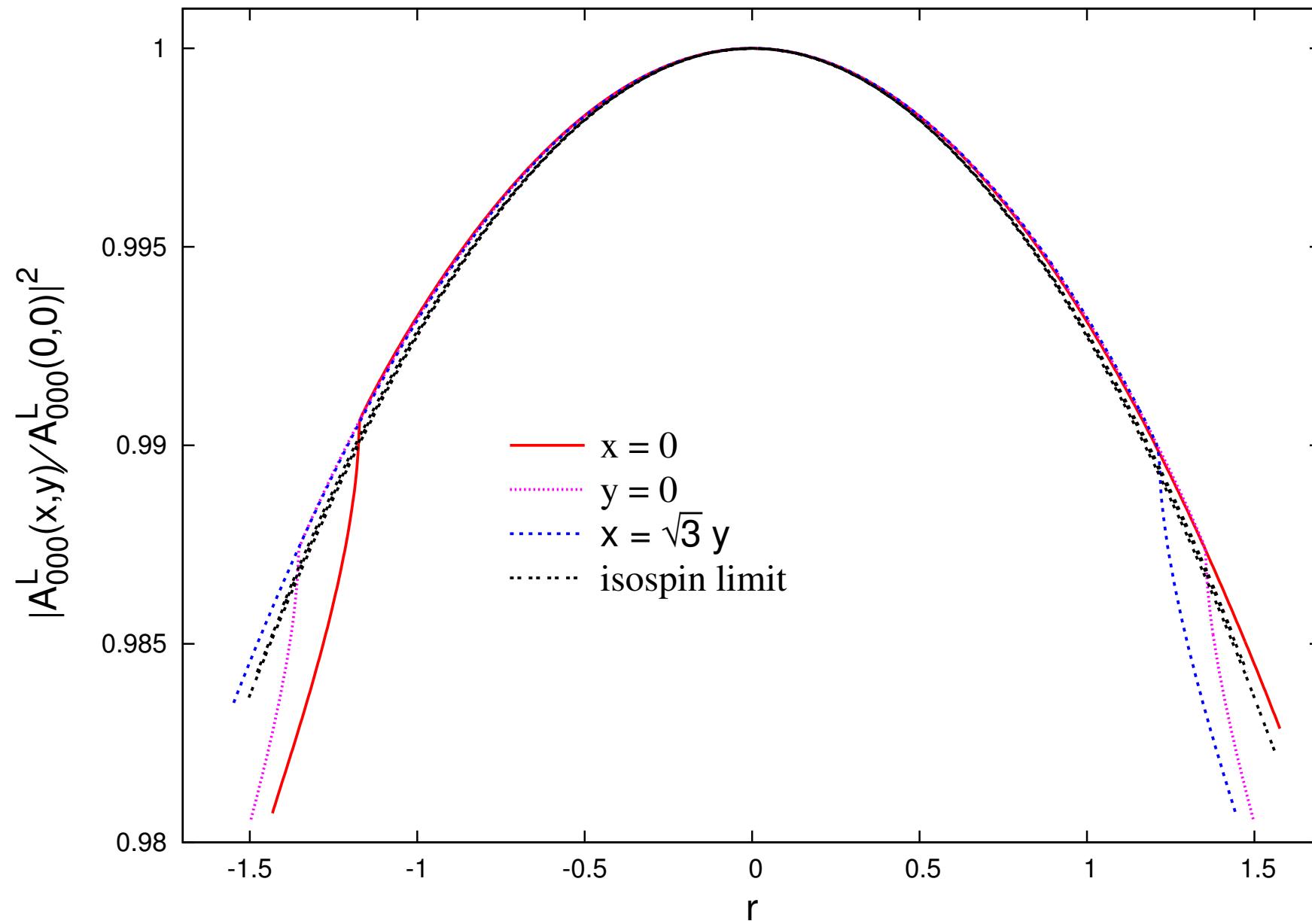
Phasespace



Curves:
Phasespace
boundaries

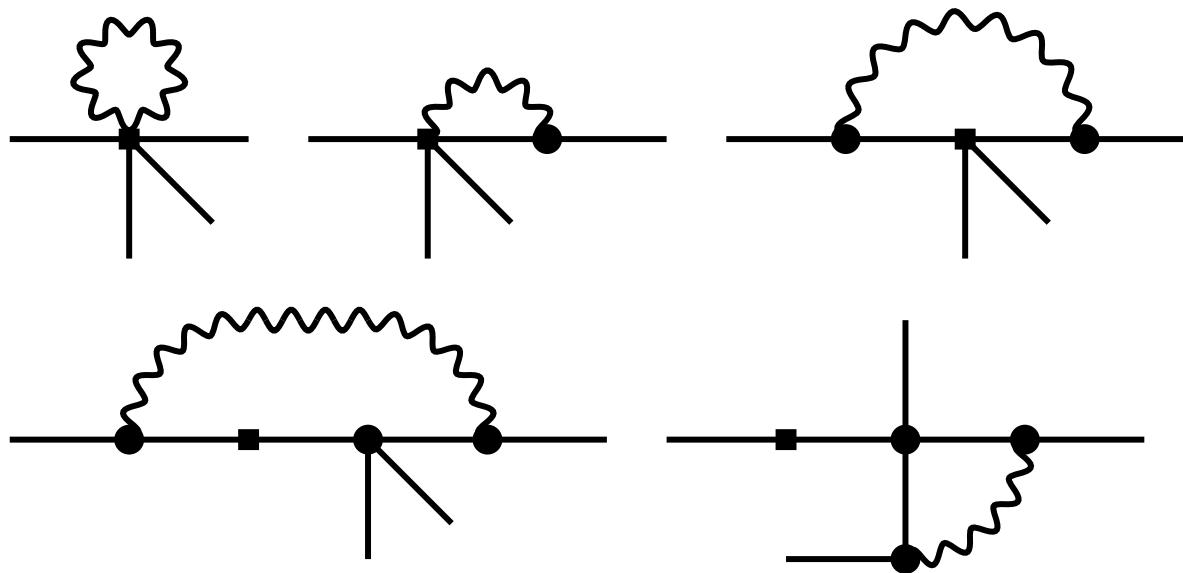
Dots:
 $\pi\pi$ pair
at rest

Lines:
Show $|A|^2$
along these

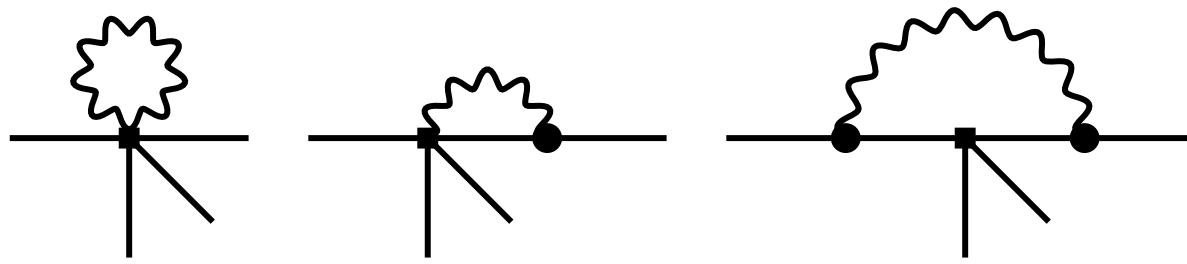


Calculations: II: Radiative Corrections

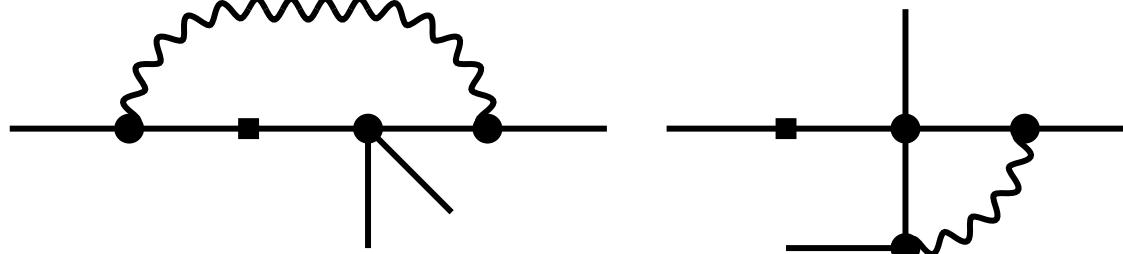
- Photon Loops



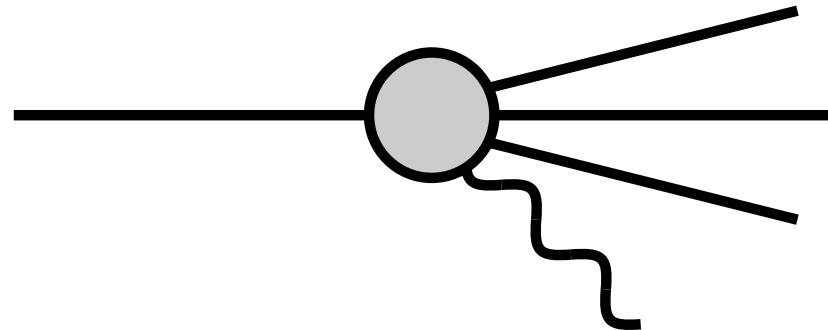
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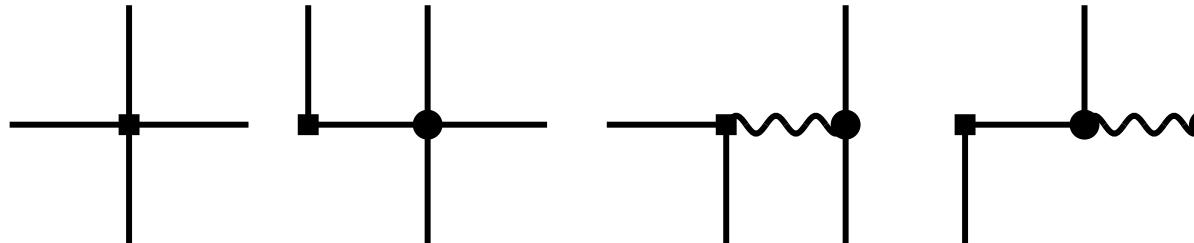
- Soft Bremsstrahlung



Calculations: II: Radiative Corrections

- Photon reducible diagrams

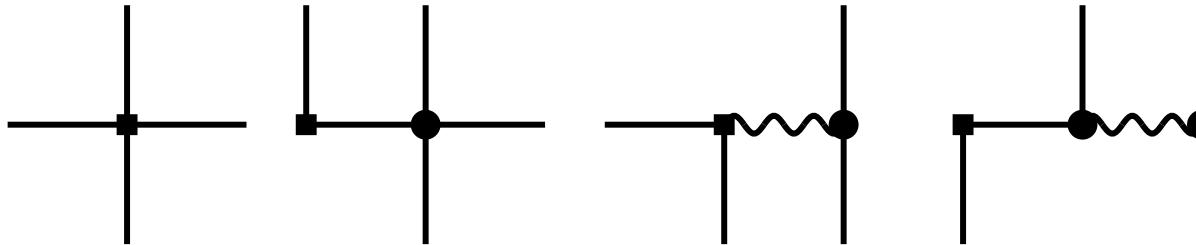
- Tree Level:



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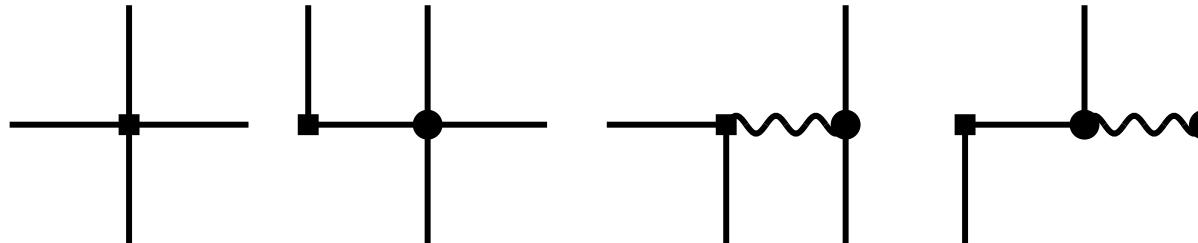


Extra vanish exactly

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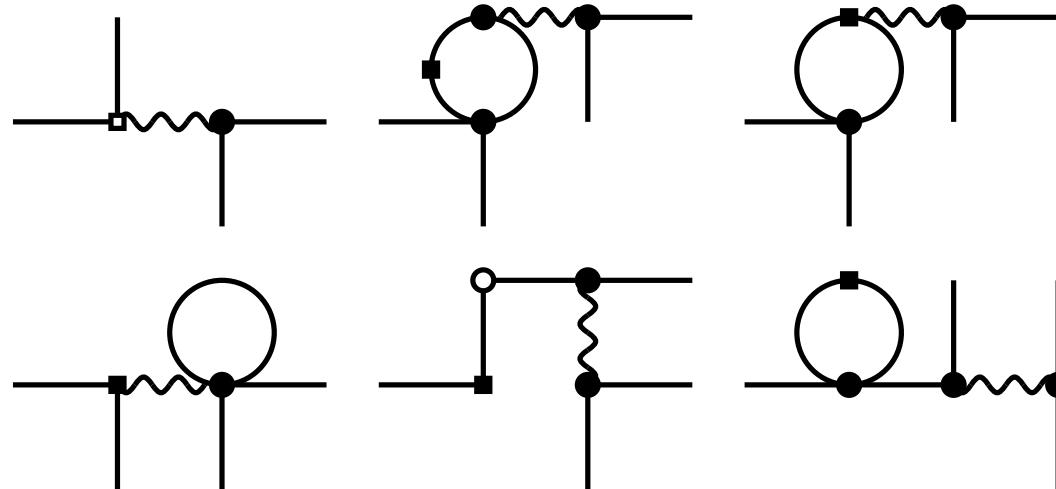
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- One-loop Level:

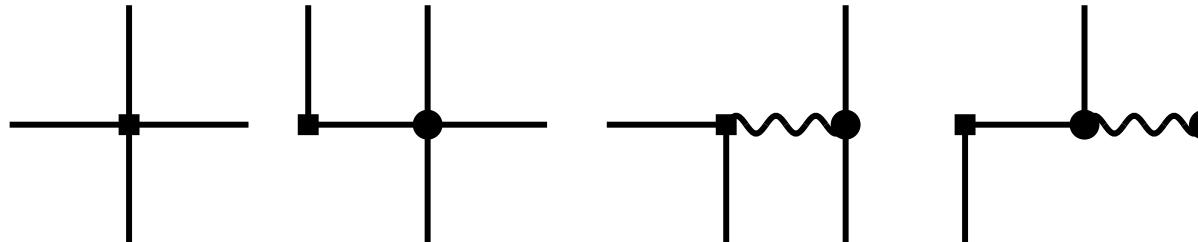


Extra N_i^r, D_i^r from $K \rightarrow \pi \ell^+ \ell^-$

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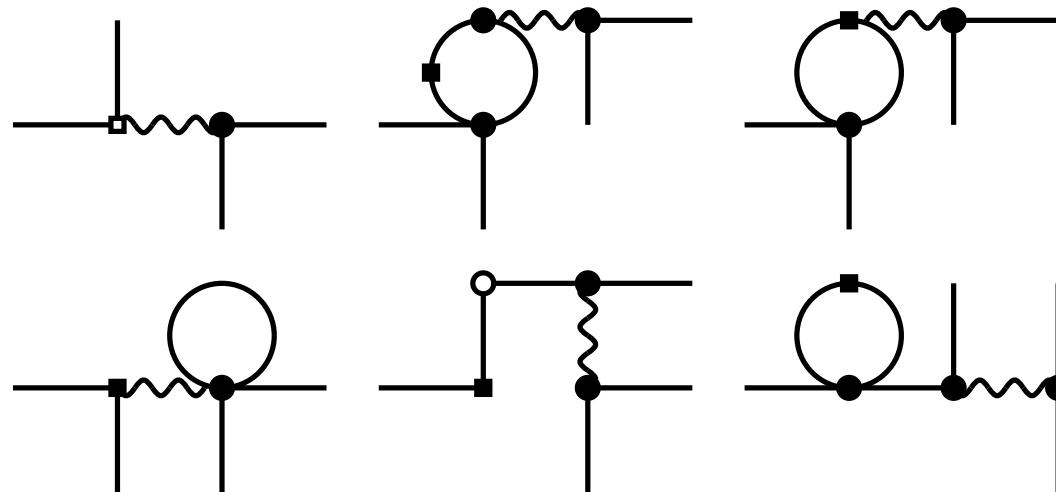
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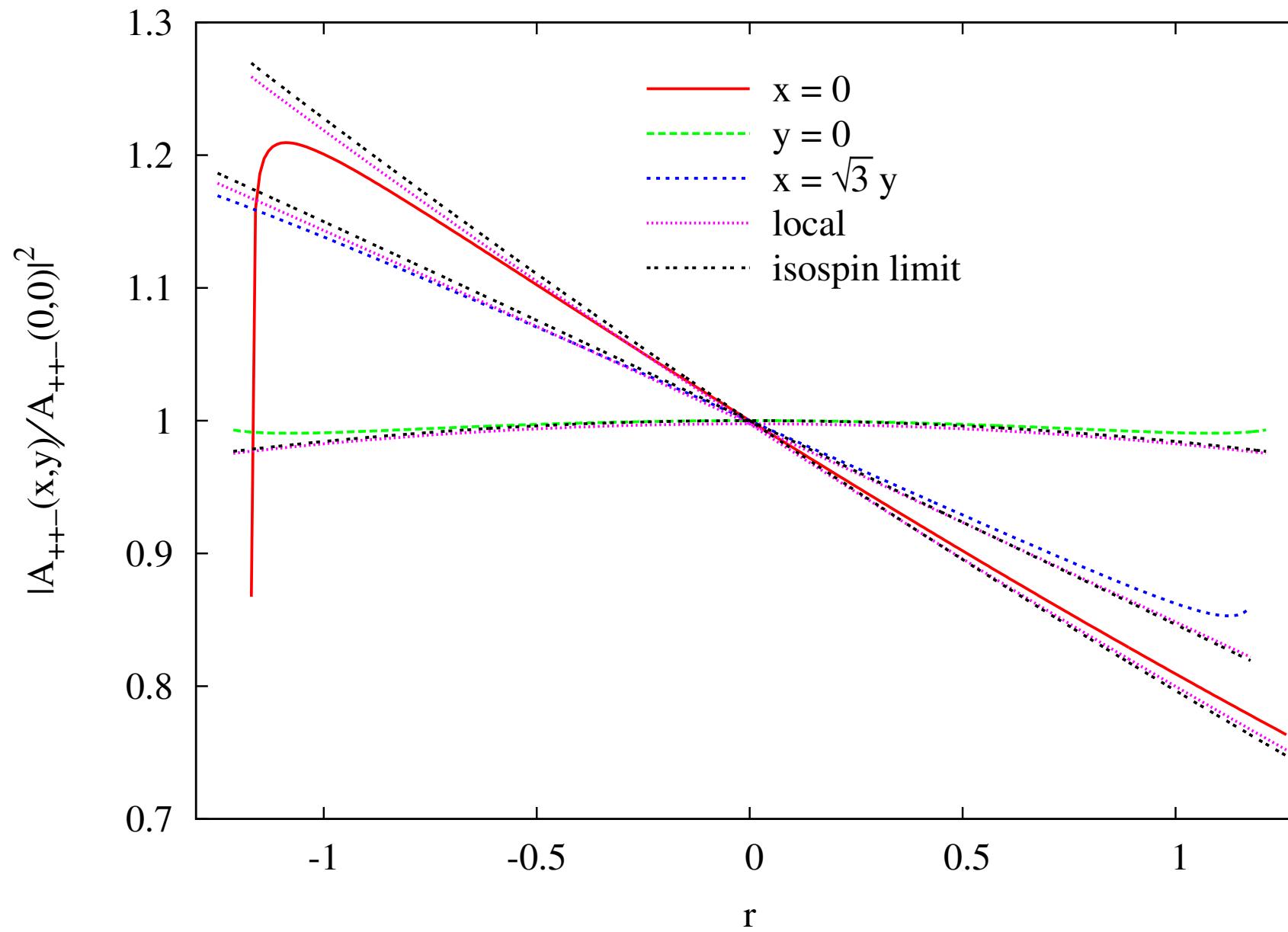
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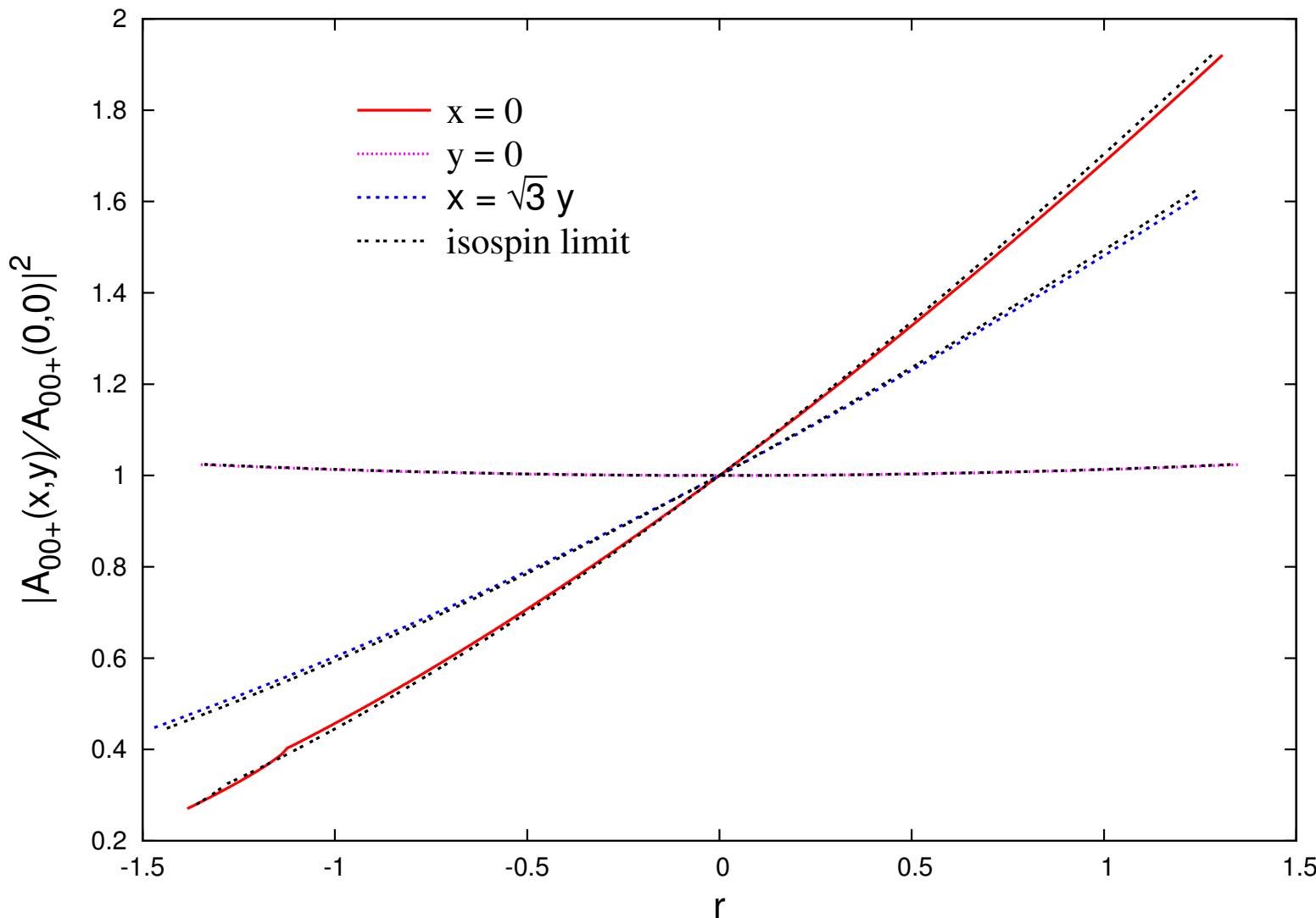
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Numerically Negligible





Note: Disagree numerically with Nehme, hep-ph/0406209

Calculations: III: Hard Bremsstrahlung

- Include hard photon Bremsstrahlung to lowest order
Full amplitude from Low's theorem
- Do everything also for $K \rightarrow 2\pi$ to have same treatment
Note: some new pieces, e.g., $e^2 p^2 G_{27}$

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- Checks:
- Everything done twice independently
 - UV infinities cancel analytically: no $1/(d - 4)$
 - IR infinities cancel analytically: no m_γ
 - soft-hard Bremsstrahlung match: no E_γ
(except for F_π, F_K parts)

Experimental Inputs

- Decay Rates
- Dalitz Plot Parameters

$$\left| \frac{A(s_1, s_2, s_3)}{A(s_0, s_0, s_0)} \right|^2 = 1 + gy + hy^2 + kx^2$$

g, h, k for $K_L \rightarrow \pi^+ \pi^- \pi^0$,

$K^+ \rightarrow \pi^0 \pi^0 \pi^+$,

$K^+ \rightarrow \pi^+ \pi^+ \pi^-$

h for $K_L \rightarrow \pi^0 \pi^0 \pi^0$

$(g = 0 \text{ and } k = h/3)$

$$A_{+-0}^S = \gamma_S x - \xi_S xy \text{ for } K_S \rightarrow \pi^+ \pi^- \pi^0$$

- All input taken from PDG2004

Treatment of Bremsstrahlung

Old experiments: what to do?

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- We adopted:
- Subtract hard Bremsstrahlung from Decay Rate
 - From remainder get $|A(s_0, s_0, s_0)|^2$ using **experimental** g, h, k
 - Fit $|A(s_0, s_0, s_0)|^2, g, h, k$ with soft Bremsstrahlung and loops, ...

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We really need to know what Phasespace/Photons covered

Results

$K \rightarrow \pi\pi$ only: essentially same as Ecker et al.

Tree level: $G_8 = 10.36$ $G_{27} = 0.550$

Full: $G_8 = 5.39$ $G_{27} = 0.359$

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Full fit:

μ	0.77 GeV	1.0 GeV	0.6 GeV	0.77 GeV
G_8	5.39(1)	4.60(1)	6.43(1)	5.39(1)
G_{27}	0.359(2)	0.301(1)	0.438(2)	0.359(2)
$\delta_2 - \delta_0$	-57.9(1.5) ^o	-57.3(1.4) ^o	-58.9(1.4) ^o	-57.9(1.4) ^o
$10^3 \tilde{K}_1/G_8$	$\equiv 0$	$\equiv 0$	$\equiv 0$	$\equiv 0$
$10^3 \tilde{K}_2/G_8$	48.5(2.4)	56.5(2.4)	41.2(1.9)	46.6(1.6)
$10^3 \tilde{K}_3/G_8$	2.6(1.2)	-1.7(1.1)	6.7(1.0)	3.5(0.8)
$10^3 \tilde{K}_4/G_{27}$	$\equiv 0$	$\equiv 0$	$\equiv 0$	$\equiv 0$
$10^3 \tilde{K}_5/G_{27}$	-41.2(16.9)	-52.0(17.7)	-31.1(12.0)	-27.0(8.3)
$10^3 \tilde{K}_6/G_{27}$	-102(105)	-114(105)	-93(76)	$\equiv 0$
$10^3 \tilde{K}_7/G_{27}$	78.6(33)	78.0(33.5)	79.6(22.7)	50.0(13.0)
χ^2/DOF	29.3/10	27.2/10	33.0/10	30.5/11

vary
 \tilde{K}_1, \tilde{K}_4
 see first
 paper

Octet and μ -variation

μ	octet 0.77 GeV	μ variation 1.0 GeV	μ variation 0.6 GeV
G_8	4.84(1)	–	–
G_{27}	0.430(1)	–	–
$\delta_2 - \delta_0$	$-57.9(0.2)^o$	–	–
$10^3 \tilde{K}_1/G_8$	2.0(1)	–5.88	5.61
$10^3 \tilde{K}_2/G_8$	63.0(1.5)	–2.69	2.57
$10^3 \tilde{K}_3/G_8$	–6.0(7)	0.159	–0.152
$10^3 \tilde{K}_4/G_{27}$	$\equiv 0$	–9.93	9.48
$10^3 \tilde{K}_5/G_{27}$	$\equiv 0$	0	0
$10^3 \tilde{K}_6/G_{27}$	$\equiv 0$	27.0	–25.8
$10^3 \tilde{K}_7/G_{27}$	$\equiv 0$	–21.5	20.5
$10^3 \tilde{K}_8/G_8$	20.4(1)	–0.546	0.521
$10^3 \tilde{K}_9/G_8$	9.1(1)	–2.92	2.79
$10^3 \tilde{K}_{10}/G_8$	$\equiv 0$	11.6	–11.1
$10^3 \tilde{K}_{11}/G_8$	$\equiv 0$	–1.66	1.58
χ^2/DOF	33.3/10	–	–

Octet: $G_{27} = 0$
 keep also \tilde{K}_i
 with $m_K^2 m_\pi^2$ factors

μ variation
 $\tilde{K}_i(\mu) - \tilde{K}_i(0.77 \text{ GeV})$

Fit about same as
 before, some exper-
 iments fit better, oth-
 ers smaller error

Models

See Ecker, Kambor, Wyler 1993 for explanations

- Vector octet dominance $\implies \tilde{K}_3 = -\frac{1}{2}\tilde{K}_2$: Not at all

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$$N_1^r = 2k_f(32/3 L_1^r + 4 L_3^r + 2/3 L_9^r),$$

$$N_2^r = 2k_f(16/3 L_1^r + 4 L_3^r + 10/3 L_9^r),$$

$$N_3^r = 2k_f(8L_2^r - 2L_9^r),$$

$$N_4^r = 2k_f(-16/3 L_1^r - 8/3 L_3^r - 4/3 L_9^r),$$

$$N_5^r = 2k_f(-L_5^r),$$

$$N_6^r = 2k_f(2/3 L_5^r),$$

$$N_7^r = 2k_f(L_5^r),$$

$$N_8^r = 2k_f(4 L_4^r + 2 L_5^r),$$

$$N_9^r = N_{10}^r = N_{11}^r = N_{12}^r = N_{13}^r = 0$$

Models

See Ecker, Kambor, Wyler 1993 for explanations

- Vector octet dominance $\implies \tilde{K}_3 = -\frac{1}{2}\tilde{K}_2$: Not at all
- Weak Deformation and Factorization

$$\begin{aligned}
 N_1^r &= 2k_f(32/3 L_1^r + 4 L_3^r + 2/3 L_9^r), \\
 N_2^r &= 2k_f(16/3 L_1^r + 4 L_3^r + 10/3 L_9^r), \\
 N_3^r &= 2k_f(8L_2^r - 2L_9^r), \\
 N_4^r &= 2k_f(-16/3 L_1^r - 8/3 L_3^r - 4/3 L_9^r), \\
 N_5^r &= 2k_f(-L_5^r), \\
 N_6^r &= 2k_f(2/3 L_5^r), \\
 N_7^r &= 2k_f(L_5^r), \\
 N_8^r &= 2k_f(4 L_4^r + 2 L_5^r), \\
 N_9^r &= N_{10}^r = N_{11}^r = N_{12}^r = N_{13}^r = 0
 \end{aligned}$$

μ	0.77 GeV	0.9 GeV	0.842 GeV
G_8	4.18(1)	4.42	4.22(1)
G_{27}	0.360(2)	0.326(10)	0.339(10)
k_F	2.61(1)	4.94(2)	3.60(5)
χ^2/DOF	109/14	182/14	60.4/13

WDM or $k_f = 1/2$ No
 Factorization: surprisingly OK
 But k_f not naive one

Data and Fits

Decay	Width [GeV]	ChPT [GeV]	Fact. [GeV]
$K^+ \rightarrow \pi^+ \pi^0$	$(1.1231 \pm 0.0078) \cdot 10^{-17}$	$1.123 \cdot 10^{-17}$	$1.127 \cdot 10^{-17}$
$K_S \rightarrow \pi^0 \pi^0$	$(2.2828 \pm 0.0104) \cdot 10^{-15}$	$2.282 \cdot 10^{-15}$	$2.283 \cdot 10^{-15}$
$K_S \rightarrow \pi^+ \pi^-$	$(5.0691 \pm 0.0108) \cdot 10^{-15}$	$5.069 \cdot 10^{-15}$	$5.069 \cdot 10^{-15}$
$K_L \rightarrow \pi^0 \pi^0 \pi^0$	$(2.6748 \pm 0.0358) \cdot 10^{-18}$	$2.618 \cdot 10^{-18}$	$2.698 \cdot 10^{-18}$
$K_L \rightarrow \pi^+ \pi^- \pi^0$	$(1.5998 \pm 0.0271) \cdot 10^{-18}$	$1.658 \cdot 10^{-18}$	$1.711 \cdot 10^{-18}$
$K^+ \rightarrow \pi^0 \pi^0 \pi^+$	$(9.195 \pm 0.0213) \cdot 10^{-19}$	$8.934 \cdot 10^{-19}$	$8.816 \cdot 10^{-19}$
$K^+ \rightarrow \pi^+ \pi^+ \pi^-$	$(2.9737 \pm 0.0174) \cdot 10^{-18}$	$2.971 \cdot 10^{-18}$	$2.933 \cdot 10^{-18}$

Decay	Quantity	Experiment	ChPT	Fact.
$K_L \rightarrow \pi^0 \pi^0 \pi^0$	h	-0.0050 ± 0.0014	-0.0062	-0.0025
$K_L \rightarrow \pi^+ \pi^- \pi^0$	g	0.678 ± 0.008	0.678	0.654
	h	0.076 ± 0.006	0.088	0.083
	k	0.0099 ± 0.0015	0.0057	0.0068
$K_S \rightarrow \pi^+ \pi^- \pi^0$	γ_S	$(3.3 \pm 0.5) \cdot 10^{-8}$	$3.0 \cdot 10^{-8}$	$2.9 \cdot 10^{-8}$
$K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm$	g	0.638 ± 0.020	0.636	0.648
	h	0.051 ± 0.013	0.077	0.080
	k	0.004 ± 0.007	0.0047	0.0069
$K^+ \rightarrow \pi^+ \pi^+ \pi^-$	g	-0.2154 ± 0.0035	-0.215	-0.226
	h	0.012 ± 0.008	0.012	0.019
	k	-0.0101 ± 0.0034	-0.0034	-0.0033

Conclusions

- Calculated $K \rightarrow 3\pi$ to first nontrivial order in isospin breaking fully
- New fit done and tested a few models
- Fit with/without isospin breaking seems similar
- Need Photon/Phasespace information from experiment
- CP violation: not started yet
- Corrections for previous talk: higher order than those included here