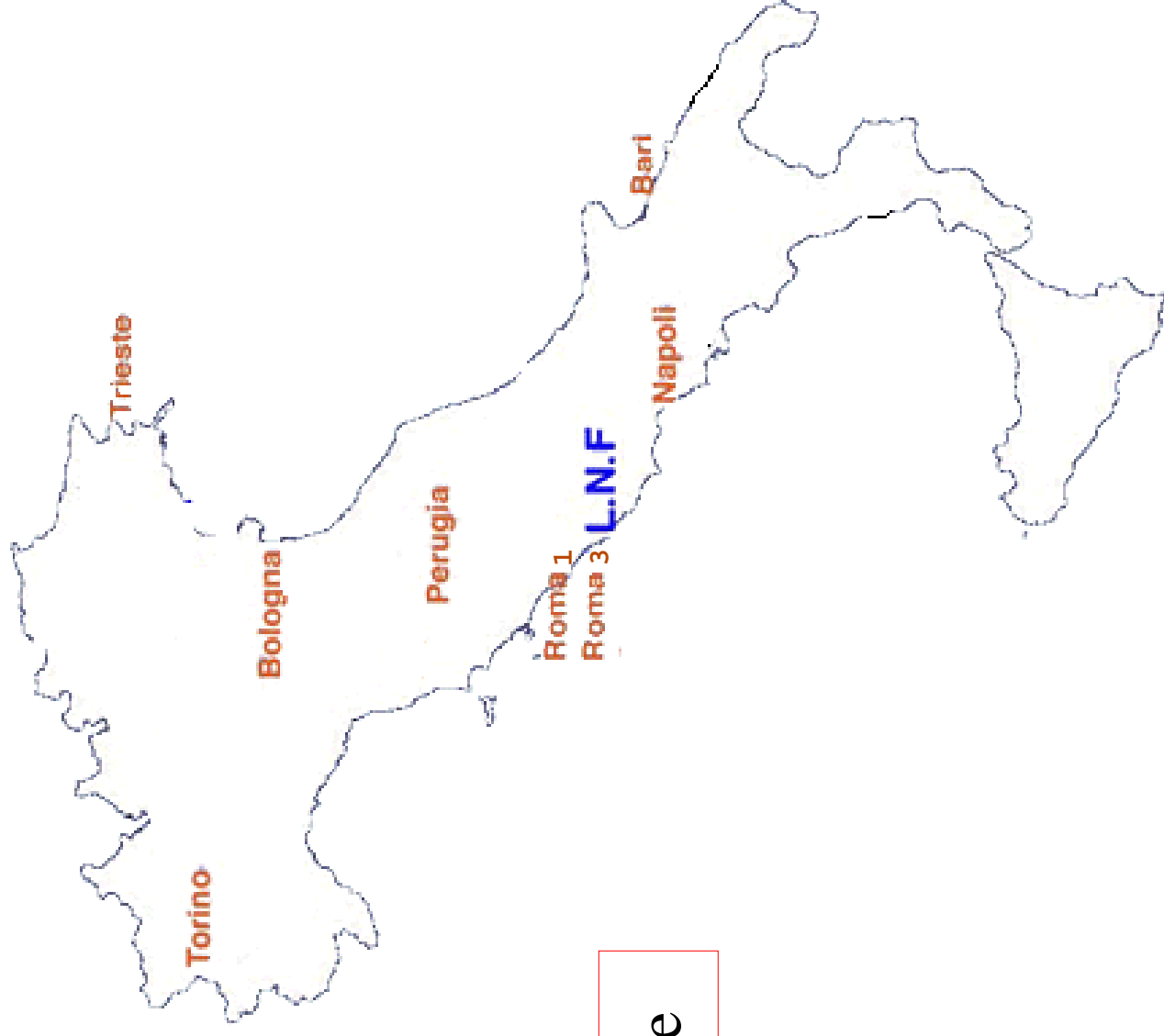


*Gino Isidori (LNF)*

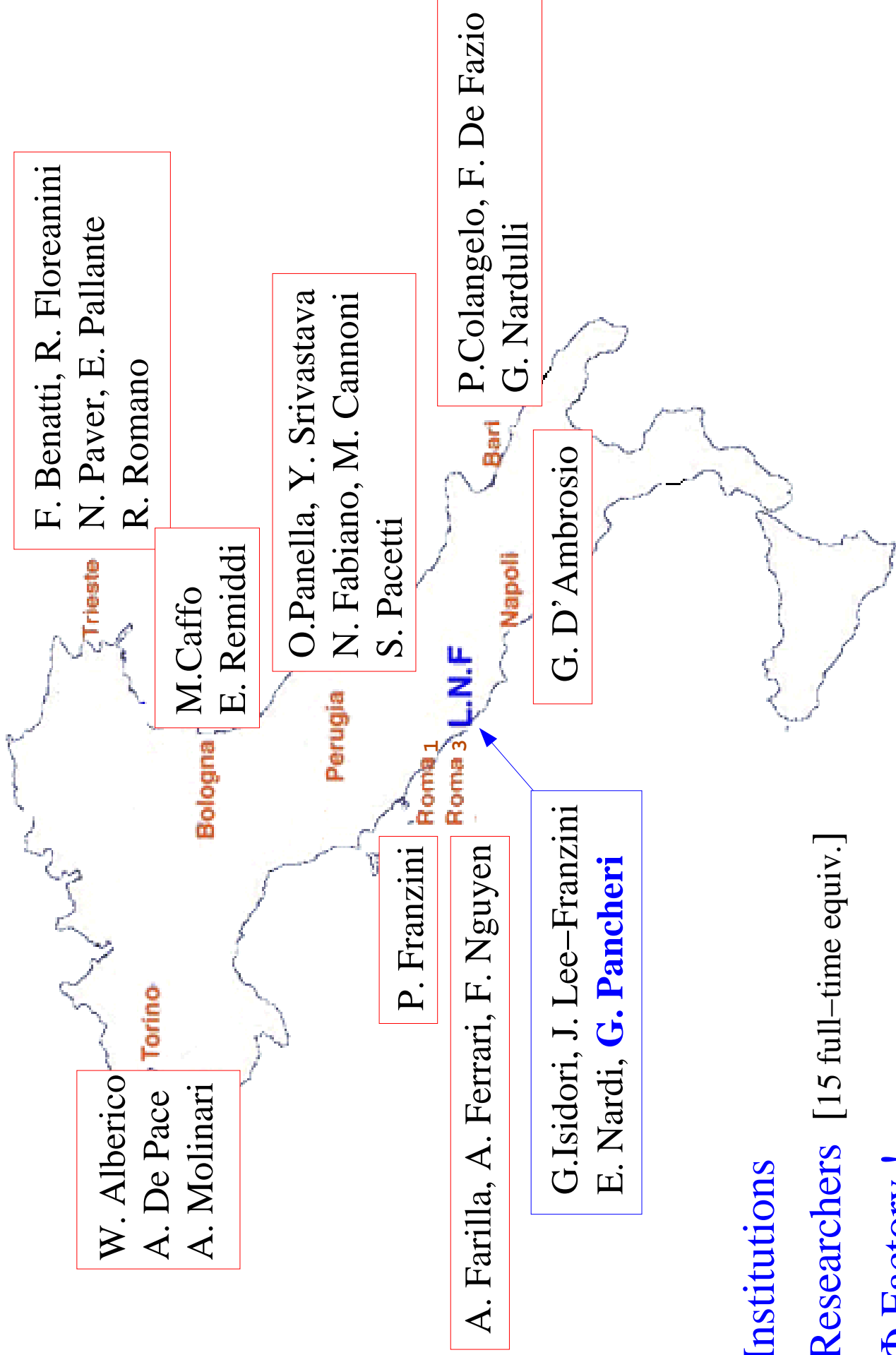
1<sup>st</sup> Euridice Collaboration Meeting

[Frascati, 18–20 October 2002]



The **EURIDICE** – **INFN** node



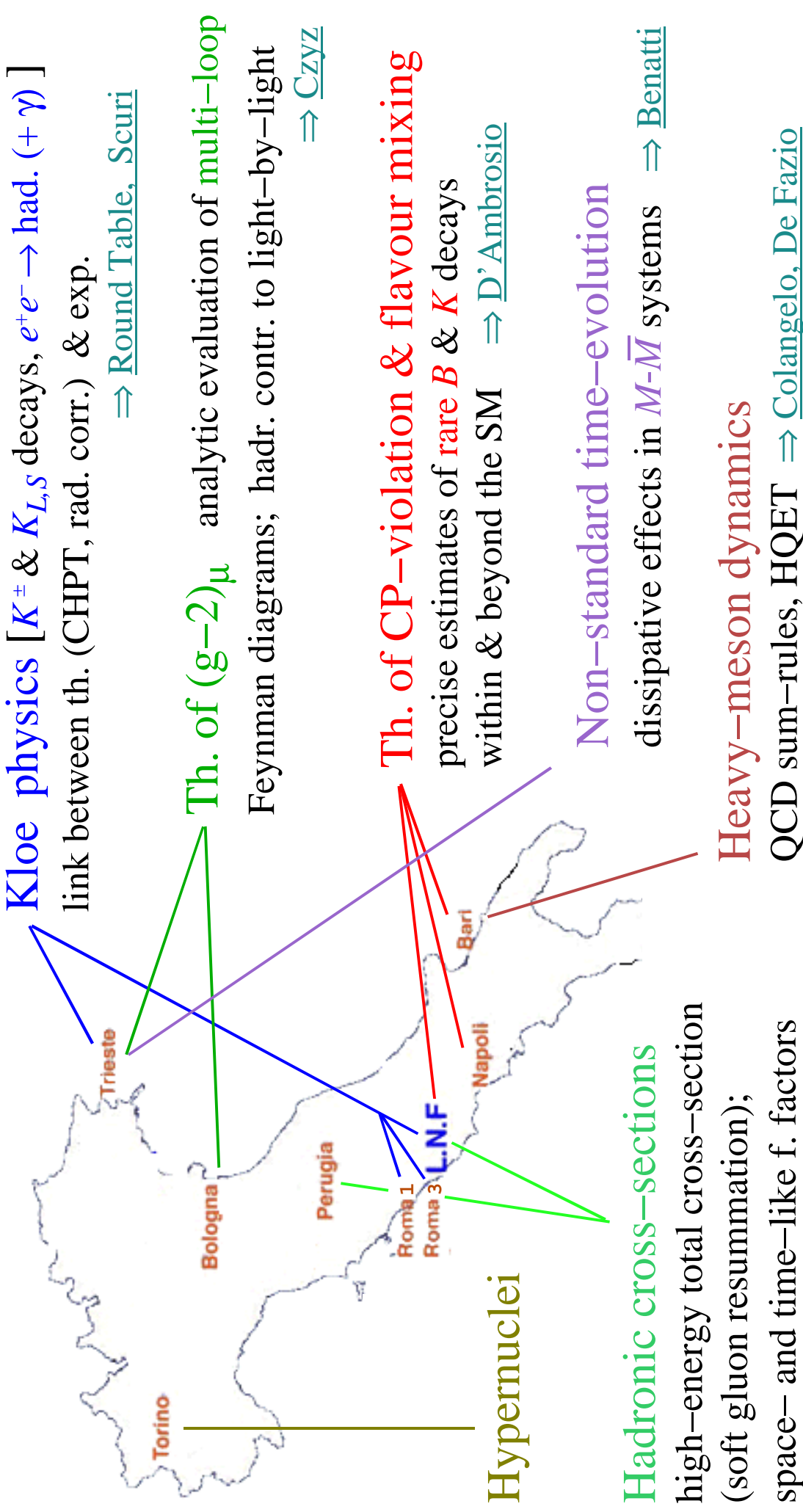


9 Institutions

27 Researchers [15 full-time equiv.]

1  $\Phi$  Factory!

# • The EURIDICE – INFN physics program

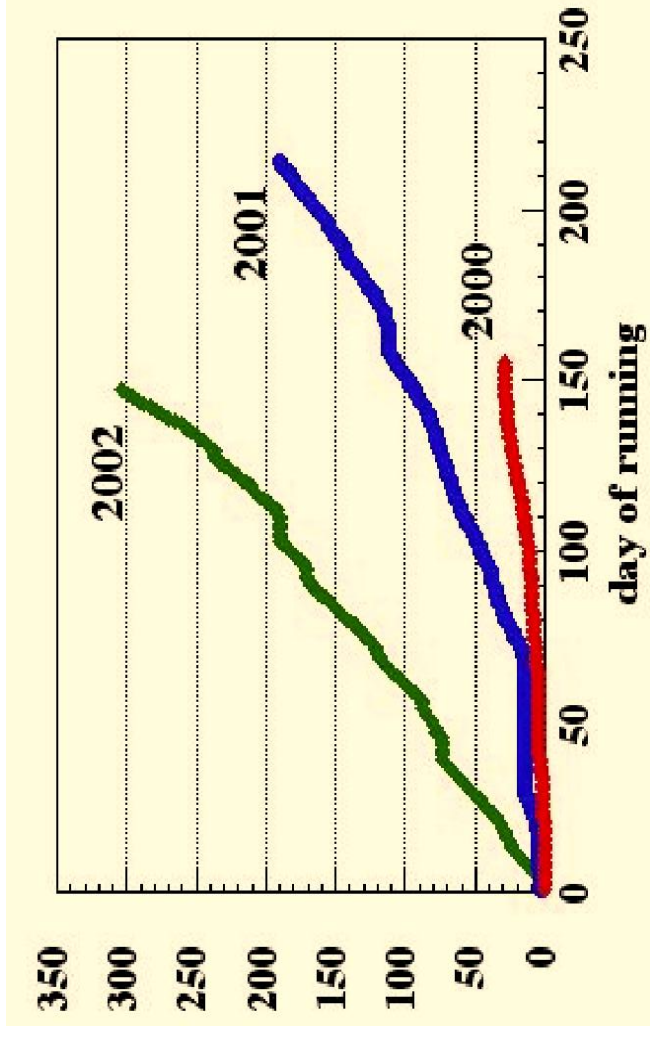


## • KLOE physics

Very efficient **DAΦNE**  $\oplus$  **KLOE**  
run in 2002:  $\int \mathcal{L} ('01+'02) = 500 \text{ pb}^{-1}$



$\sim 5 \times 10^8$   $K_L K_S$  (untagged) pairs



Wide list of on-going activities [ **$\Phi$** ,  **$\eta$**  and  **$K$**  decays]

Three outstanding research lines with  $500 \text{ pb}^{-1}$  :  
(personal point of view)

several non-INFN links:

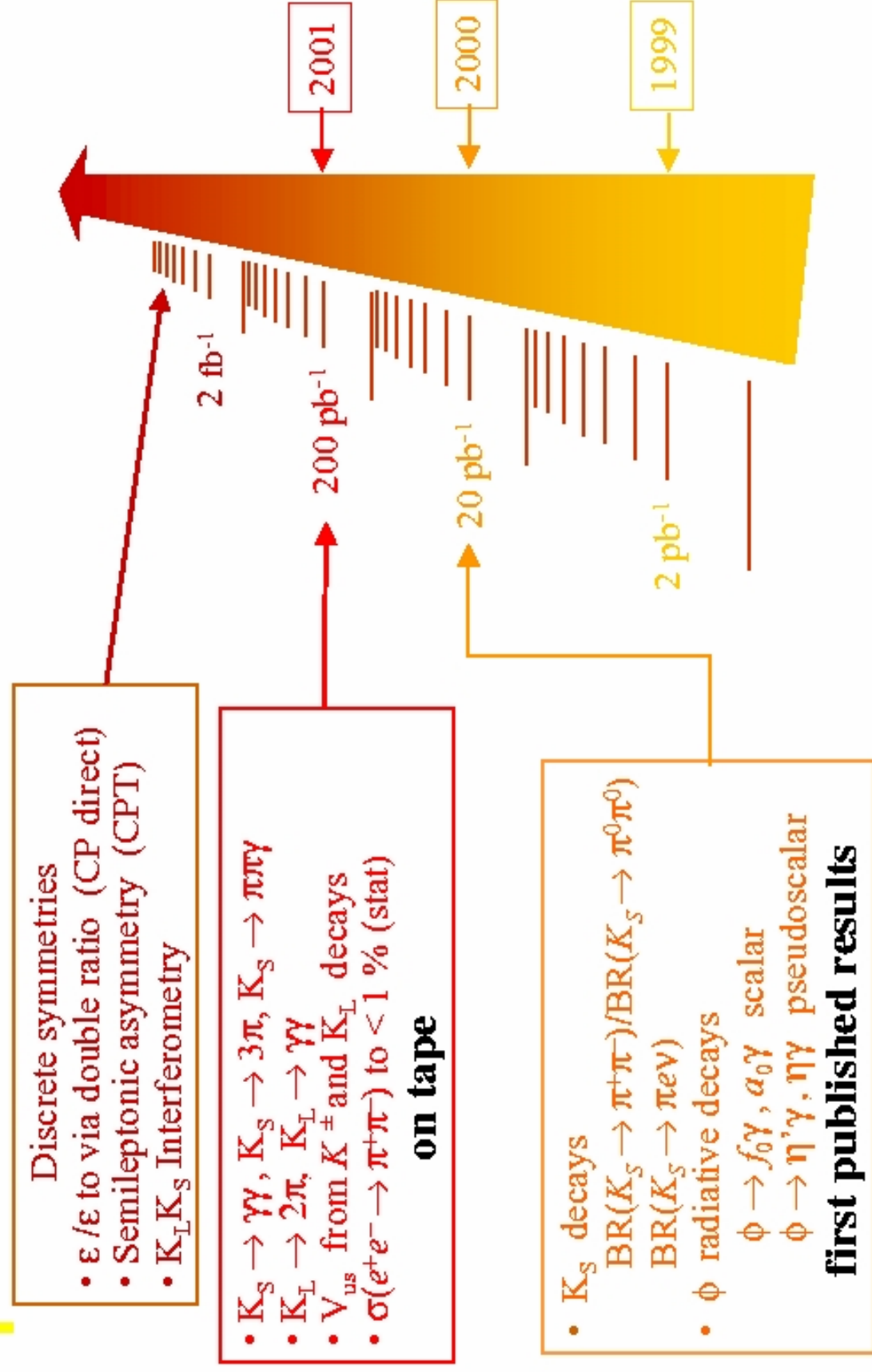
•  $\sigma[e^+e^- \rightarrow \text{had.} (+\gamma)]$  @  $10^{-2}$  level (or better)  $\Leftrightarrow \Pi_\gamma(q^2)$  [DESY, Karlsruhe, Warsaw]

•  $\Gamma(K_{l3}^\pm)$  &  $\Gamma(K_{l3}^0)$  @  $10^{-3}$  level (+  $\gamma$  spectrum)  $\Leftrightarrow V_{us}$ , CHPT [Bern, Marseille, Valencia, Wien]

• Rare  **$K_S$**  decays [ **$K_S \rightarrow \pi l \nu$** ,  **$3\pi$** ,  **$\pi l l$**  down to  $10^{-8}$  SES]  $\Leftrightarrow$  **CPT, CP, CHPT** [Valencia, Wien]



# KLOE physics program

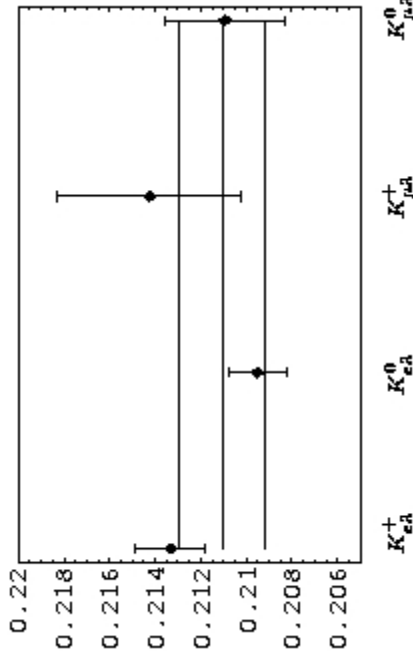


Present status of  $V_{us}$   $f(0) \times V_{us}$

large spread after SU(2) corr.

[ $\gamma$ -inclusive measurements?]

Measuring  $V_{us}$  at KLOE



data after inclusion  
of SU(2) breaking  
corrections

Cirigliano *et al.* '02

$\Rightarrow$  talk by Neufeld

KLOE can improve the accuracy in the knowledge of  $\Gamma(e3)$  measuring:

- $\triangleright$  both charged and neutral kaons absolute BR's with the **same detector**
- $\triangleright$  directly the kaon partial decay width

Traditional method:

- $\diamond$  measurement of  $\Gamma = 1/\tau$
- $\diamond$  measurement of  $BR(e3)$



$$\Gamma(e3) = BR(e3) \times \Gamma$$

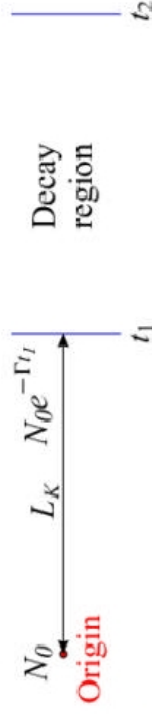
$$\frac{\delta \Gamma(e3)}{\Gamma(e3)} = \sqrt{\left(\frac{\delta \tau}{\tau}\right)^2 + \left(\frac{\delta(BR(e3))}{BR(e3)}\right)^2}$$

KLOE:

- $\diamond$  by the tag count the number of K produced,  $N_{kl}$
- $\diamond$  count the number  $N_{e3}$  of semileptonic decays in the decay region

$$\Gamma(e3) = (\Delta N_{e3} / \Delta t) / N_{kl}$$

$\Gamma$  is a correction &  $\delta\tau/\tau$   
dependence reduced by a  
factor  $\approx 5$



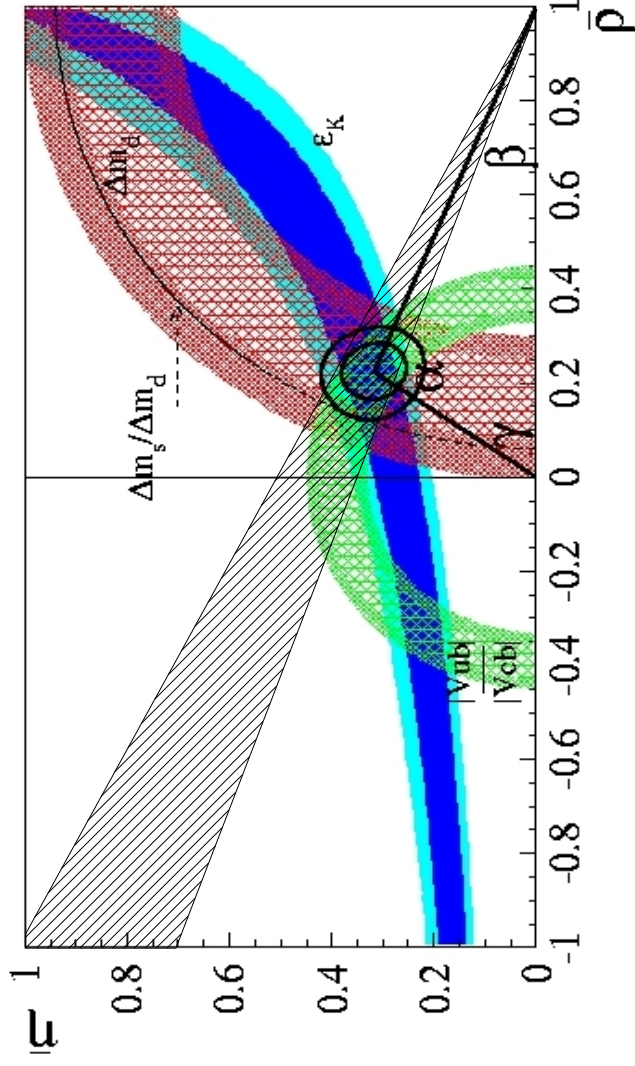
E. De Lucia,  
Sept. 2002

# ● Theory of CP–violation and flavour mixing

My main activity  
in '01–'02

## The Flavour Problem

Available data on  $\Delta F=2$  FCNC amplitudes (meson–antimeson mixing) provides very serious constraints on the scale of New Physics...



E.g.:  $\bar{K}^0 - K^0$  mixing



$\Lambda \gtrsim 100 \text{ TeV}$  for  $O^{(6)} \sim (\bar{s}d)^2$

much more stringent that the bounds on the scale of flavour–conserving op. from ew precision data (  $\sim$  few TeV)

...while a natural stabilization of the Higgs potential  $\Rightarrow \Lambda \sim 1 \text{ TeV}$

After the recent precise data from  $B$  factories, it is very difficult to believe this is an accident...

Two possible solutions:

- *pessimistic* [very unnatural]:  $\Lambda > 100 \text{ TeV}$   
 $\Rightarrow$  almost nothing to learn from other FCNC processes
- *natural*:  $\Lambda \sim 1 \text{ TeV}$  + flavour-mixing protected by additional symmetries  
 $\Rightarrow$  still a lot to learn from FCNCs [especially from  $\Delta F=1$  processes]



### Minimal Flavour Violation (MFV) hypothesis:

The breaking of the flavour symmetry occurs at very high scales and is mediated at low energies only by terms proportional to SM Yukawa couplings.

- natural implementation in many consistent scenarios

[SUSY, technicolour, extra dimensions,...]

[wide literature]

- possible to build a predictive low-energy EFT including

only SM fields  $\Rightarrow$  model-independent approach

[D'Ambrosio, Giudice,  
G.I., Strumia, '02]

# Minimal Flavour Violation

The maximal group of unitary field transf. allowed by  $\mathcal{L}_{\text{gauge}}^{\text{SM}}$  is:

$$G_F = \text{U}(3)^5 = \text{SU}(3)^2_{\textcolor{red}{l}} \times \text{SU}(3)^3_{\textcolor{red}{q}} \times \text{U}(1)_{\textcolor{red}{PQ}} \times \text{U}(1)_{\textcolor{red}{E_R}} \times \text{U}(1)_{\textcolor{red}{B}} \times \text{U}(1)_{\textcolor{red}{L}} \times \text{U}(1)_{\textcolor{red}{Y}}$$

subgroup broken by  $\mathcal{L}_{\text{Yukawa}}^{\text{SM}}$

$$\text{SU}(3)^2_{\textcolor{red}{l}} = \text{SU}(3)_{\textcolor{red}{L_L}} \times \text{SU}(3)_{\textcolor{red}{E_R}}$$

$$\text{SU}(3)^3_{\textcolor{red}{q}} = \text{SU}(3)_{\textcolor{red}{Q_L}} \times \text{SU}(3)_{\textcolor{red}{U_R}} \times \text{SU}(3)_{\textcolor{red}{D_R}}$$

subgroup responsible for quark mixing  
[Yukawa structure, CKM matrix]

$\text{U}(1)_{\textcolor{red}{PQ}}$ : glob. phase of  $D_R$  &  $E_R$

$\text{U}(1)_{\textcolor{red}{E_R}}$ : glob. phase of  $E_R$

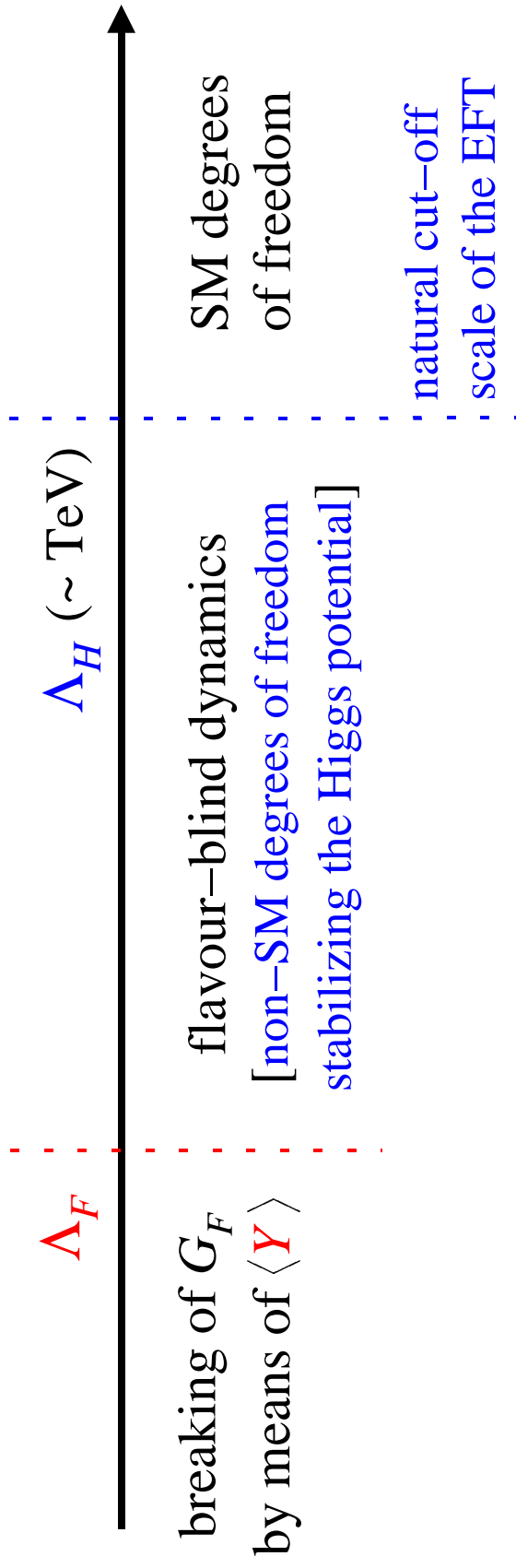
groups relevant in multi-Higgs models [overall Yukawa norm.]

$$\mathcal{L}_{\text{Yukawa}} = \bar{Q}_L Y_D D_R H + \bar{Q}_L Y_U U_R (H)_c + \bar{L}_L Y_L E_R H + \text{h.c.}$$

Since  $G_F$  is already broken within the SM, it is not consistent to impose it as an exact symmetry beyond the SM.

However, we can (formally) promote  $G_F$  to be an exact symmetry, assuming the Yukawa matrices are the vacuum expectation values of appropriate auxiliary fields:

$$\text{e.g.: } Y_D \sim (3, \bar{3}, 1) \quad \& \quad Y_U \sim (3, 1, \bar{3}) \quad \text{under} \quad \text{SU}(3)_{Q_L} \times \text{SU}(3)_{U_R} \times \text{SU}(3)_{D_R}$$



A low-energy EFT (including only SM fields) satisfies the criterion of MFV if all higher-dimensional operators, constructed from SM and  $Y$  fields, are (formally) invariant under  $G_F$

Chivukula & Georgi, '89  
DGIS, '02

$$Y_D = \text{diag}(y_d, y_s, y_b) \quad Y_U = V^+ \times \text{diag}(y_u, y_c, y_t)$$

→ the CKM matrix is the only source of quark mix.

$$(\lambda_{\text{FC}})_{ij} = (Y_U Y_U^\dagger)_{ij} \approx y_t^2 V_{3i}^* V_{3j} \quad \lambda_{\text{FC}} \text{ is the effective coupl. ruling all FCNCs with external d quarks}$$

↓ (as within the SM for s.d. dominated amplitudes)

- all FCNC amplitudes have the same CKM structure as in the SM

$$[\text{e.g.: } A(b \rightarrow s \gamma) \propto V_{bt} V_{ts}, \quad A(s \rightarrow d \gamma) \propto V_{st} V_{td}, \quad \dots]$$

- only the flavour-independent magnitude of FCNC amplitudes can be modified by (non-standard) dimension-six operators → Rare decays

- "phase measurement" [e.g.:  $a(B \rightarrow \psi K_S)$ ,  $\Delta M_{B_d}/\Delta M_{B_s}$ ] are completely unaffected by (non-standard) dimension-six operators

# $B \rightarrow X_s l^+ l^-$ @ NNLO

Different LL counting than in  $B \rightarrow X_s \gamma$

[  $Q_9 \leftrightarrow Q_{1-6}$  starts @ 1 loop  $\Rightarrow$  NNLO simpler ]

- partial NNLO available for the dilepton spectrum

[Bobeth *et al.* '01, Asatrian *et al.* '01]

- full NNLO available for the lepton FB asymmetry

[Ghinchulov *et al.* '02, Asatrian *et al.* '02]

$$\int \frac{d^2 \Gamma(B \rightarrow X_s \mu^+ \mu^-)}{ds d \cos \vartheta} \operatorname{sgn}(\cos \vartheta) \propto \Re \left[ C_{10}^* \left( s C_9^{\text{eff}}(s) + r(s) C_7 \right) \right]$$

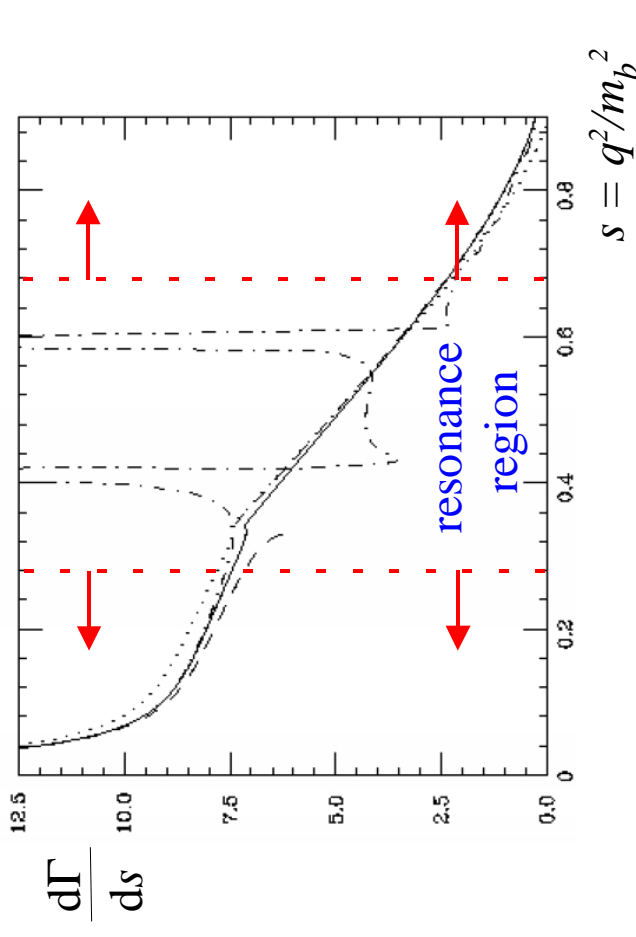
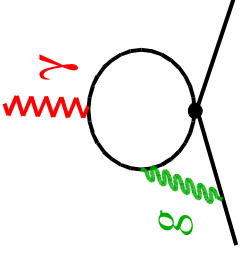
$\theta$  = angle between

$\mu^+$  &  $B$  momenta

- direct access to the relative phases of the  $C_i$

- axial current ( $C_{10}$ )  $\Rightarrow$  no  $4q$ -mixing

$\Rightarrow$  short-distance not diluted  $\Rightarrow$  excellent sensitivity to  $\Lambda_{\text{NP}}$



th. error  
 $\lesssim 5\%$



Possible extensions/generalizations of the MFV approach:

I) Inclusion of extra Higgs doublets  $\Leftrightarrow$  breaking of the  $U(1)$ 's

The breaking of the  $SU(3)$  groups is not necessarily related to the breaking of the  $U(1)$ 's

The CKM matrix remains the only source of quark mixing, but we are free to modify the overall normalization of the  $Y \Leftrightarrow$  large  $\tan\beta$



Model-indep. approach to large  $\tan\beta$  effects (SUSY and non-SUSY) in rare decays (all orders resummation of non-decoupl. terms)

G.I. & A. Retico, '01, '02;  
DGIS '02

II) Inclusion of SUSY degrees of freedom  $\Rightarrow$  EFT valid to higher scales

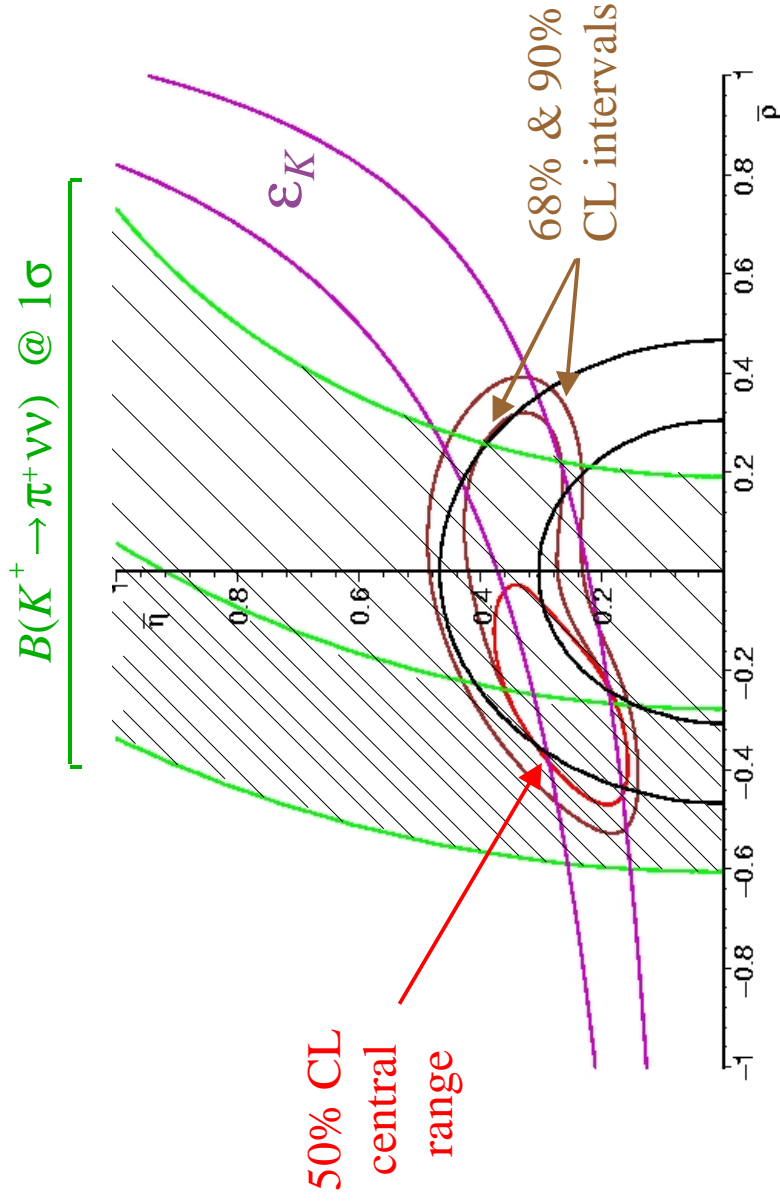
$\vdots$

Interesting developments still to be investigated

## Is there still room for non-MFV models ?

The key prediction of MFV is the link between  $b \rightarrow s$ ,  $b \rightarrow d$ , &  $s \rightarrow d$  and this is certainly not well tested yet

E.g.: UT fit ignoring  $\Delta B=2$  data [i.e. allowing a non-MFV contr. to  $(\bar{b}d)^2$ ]:



There is still room for non-MFV operators, even in the  $\Delta F=2$  sector, but these can at most induce  $O(1)$  effects with respect to the SM amplitude

important to perform new precision tests of the  $b \rightarrow s$ ,  $b \rightarrow d$ ,  $s \rightarrow d$  link

## ● Conclusions

The INFN node has

- several well-established links inherited from **EURODAPHNE**
- several on-going projects on the wider physics program of **EURIDICE** ...

[many activities have not been presented in this short overview  
⇒ see individual talks today & tomorrow]

...which could naturally lead to new productive collaborations