

ISTITUTO NAZIONALE DI FISICA NUCLEARE

Laboratori Nazionali di Frascati

LNF -81/45

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AT  $\sqrt{s} = 62$  GeV IN PROTON-PROTON COLLISIONS

Estratto da :

Lett. Nuovo Cimento 30, 487 (1981)

Servizio Documentazione  
dei Laboratori Nazionali di Frascati  
Cas. Postale 13 - Frascati (Roma)

**The Leading Effect in  $\Lambda_c^+$  Production at  $\sqrt{s} = 62$  GeV  
in Proton-Proton Collisions.**

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(ricevuto il 10 Marzo 1981)

*Introduction* – In this paper the first measurement of the longitudinal momentum distribution of the charmed baryon  $\Lambda_c^+$  produced in the reaction

$$(1) \quad pp \rightarrow \Lambda_c^+ + e^- + \text{anything}$$

is reported. The experiment has been carried out at a centre-of-mass energy of  $\sqrt{s}=62$  GeV.

The purpose of this study was to establish whether in the production of a heavy flavour in pp interactions, the «leading hadron effect» was present.

*Experimental set-up* – The data were obtained using the split-field-magnet (SFM) spectrometer of the CERN Intersecting Storage Rings (ISR). Čerenkov counters were used together with electromagnetic shower detectors (EMSD) to trigger on electrons produced in a solid angle of about 0.2 sr around 90°. A large-area time-of-flight system was used for p, K and  $\pi$  identification up to a momentum of about 1.5 GeV/c.

Figure 1 shows the top view of the apparatus. Details of the apparatus and of the single-electron trigger can be found elsewhere (1,2).

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(<sup>1</sup>) R. BOUCLEIR, R. C. A. BROWN, E. CHESI, L. DUMPS, H. G. FISCHER, P. G. INNOCENTI, G. MAURIN, A. MINTEN, L. NAUMANN, F. PIUZ and O. ULLALAND: *Nucl. Instrum. Methods*, **125**, 19 (1975); M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: *Nucl. Instrum. Methods*, **163**, 93 (1979); M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. DI CESARE, B. ESPOSITO, L. FA-VALE, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: *Nucl. Instrum. Methods*, **179**, 477 (1981); H. FREHSE, F. LAPIQUE, M. PANTER and F. PIUZ: *Nucl. Instrum. Methods*, **156**, 87 (1978); H. FREHSE, M. HEIDEN, M. PANTER and F. PIUZ: *Nucl. Instrum. Methods*, **156**, (1978).

(<sup>2</sup>) M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. DI CESARE, B. ESPOSITO, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: *Measurement of associated charm production in pp interactions at  $\sqrt{s} = 62$  GeV*, to be published in *Nuovo Cimento*.

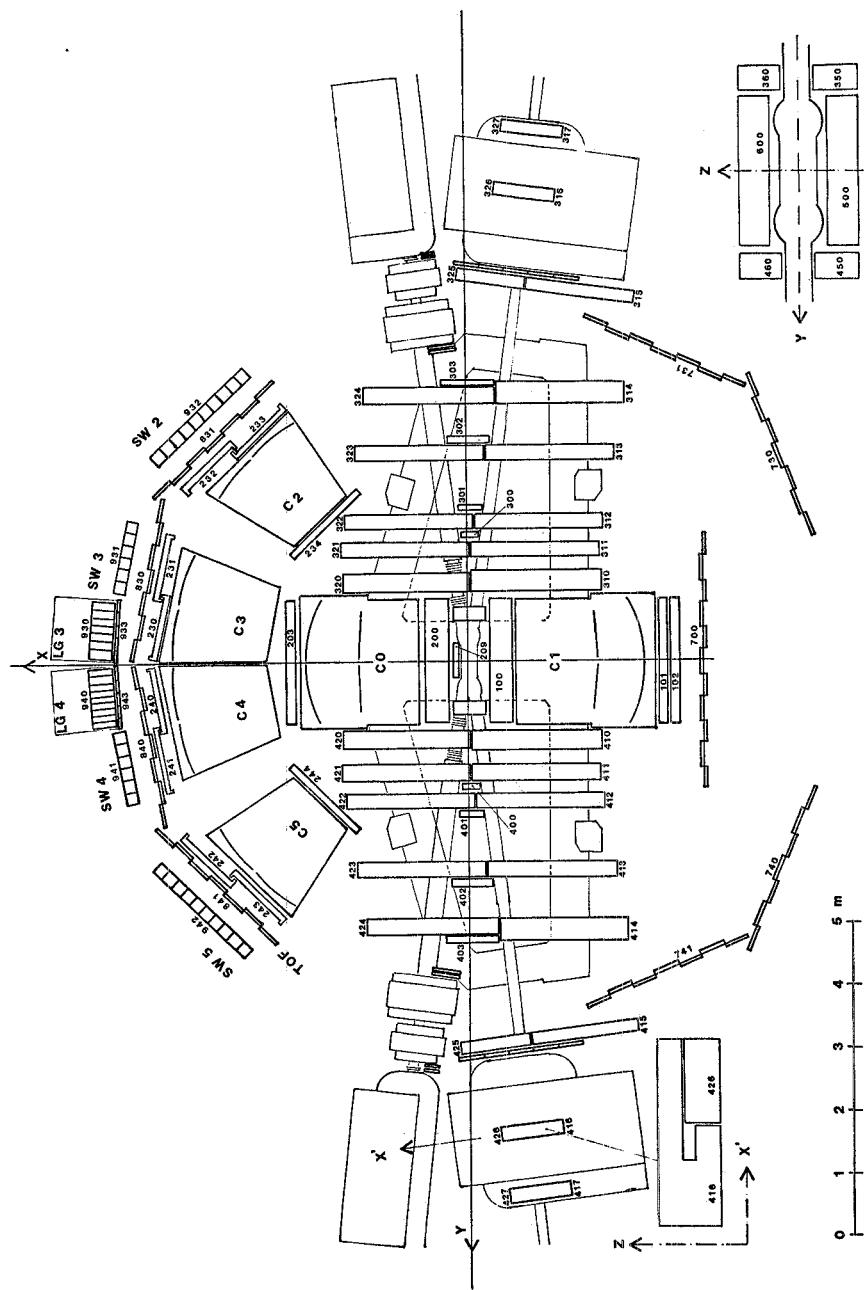


Fig. 1. — Top view of the SMT detector, showing the MWPCs and: a) the time-of-flight system (TOF); b) the electromagnetic shower detectors (EMSD) which are either lead-glass (LG) or lead-glass sandwiches (SW) or lead-glass sandwiches (LG); c) the gas threshold Čerenkov counters (C); d) the  $dE/dx$  chamber (209).

*Data Analysis* — In order to obtain a clean sample of directly produced electrons, the software analysis proceeded in two steps.

First, a software refinement of the trigger conditions, plus a request of a minimum energy release in the EMSDs ( $E > 500$  MeV), reduced the charged-hadron contamination to the 2% level. Secondly the analysis of the pulse-height of a  $dE/dx$  chamber, placed very near to the intersection region, rejected electrons from  $\pi^0$  and  $\eta$  Dalitz decay and  $\gamma$  conversions. This allowed background reduction to the 50% level.

In the sample of single-electron events, a clean signal of  $\Lambda_c^+$  decaying into  $pK^-\pi^+$  was observed (2).

In order to reduce the combinatorial background, the proton was identified as the fastest positive particle with  $x_L \simeq 2p_L/\sqrt{s} > 0.3$  in the event. According to previous measurements (3) the  $\pi^+$  contamination ranges from 20% to 2%, decreasing exponentially with increasing  $x_L$ .

The  $K^-$  and the  $\pi^+$  were any negative or, respectively, positive tracks not identified as  $\bar{p}$  or  $\pi^-$  (respectively  $p$  or  $K^+$ ) by the time-of-flight system. Furthermore, they had to

- a) be fitted to the event vertex (within  $\pm 5$  cm) with a momentum uncertainty  $\Delta p/p < 30\%$ ;
- b) be in the same rapidity hemisphere as that of the proton;
- c) have a rapidity greater than 1.0.

In order to enhance the  $\Lambda_c^+$  signal, either the presence of a leading system of charged particles with  $x_{tot} = \sum_{i=1}^n (x_{iL}) > 0.5$ , or a signature for a leading system escaping detection, i.e.  $x_{tot} < 0.1$ , has been imposed on the hemisphere opposite to that of the  $\Lambda_c^+$  (the sum is extended to all particles fitted to the vertex and with  $\Delta p/p < 30\%$ ). The final combination-to-event ratio was about 1.15.

Figure 2 shows the  $pK^-\pi^+$  invariant mass spectrum for reaction (1). The back-

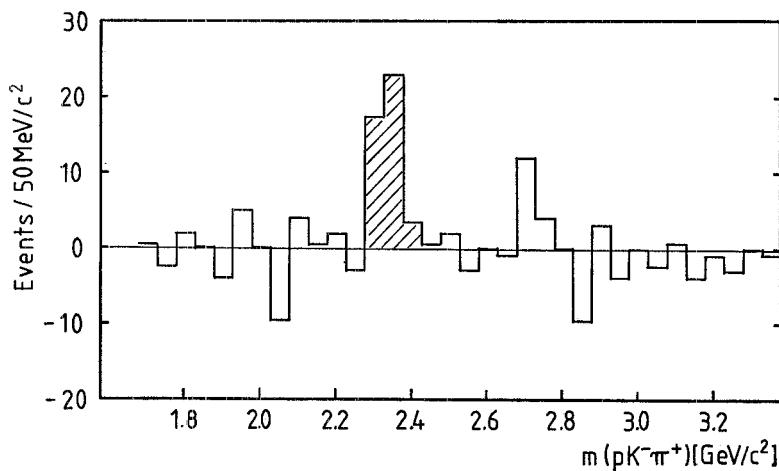


Fig. 2. — Difference between the  $pK^-\pi^+$  invariant mass spectrum associated with  $e^-$ -triggered events and the same mass combination relative to  $e^+$ -triggered events.

(3) P. CAPILUPPI, G. GIACOMELLI, A. M. ROSSI, G. VANNINI and A. BUSSIÈRE: *Nucl. Phys. B*, **70**, 1 (1974).

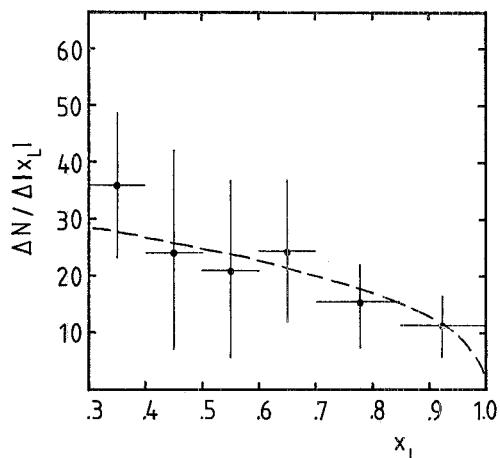


Fig. 3. — Experimental  $x_L$  distribution of the  $\Lambda_c^+$  events. The dashed line is the fit:  $\Delta N / \Delta |x_L| \propto \propto (1 - |x_L|)^{0.4}$ .

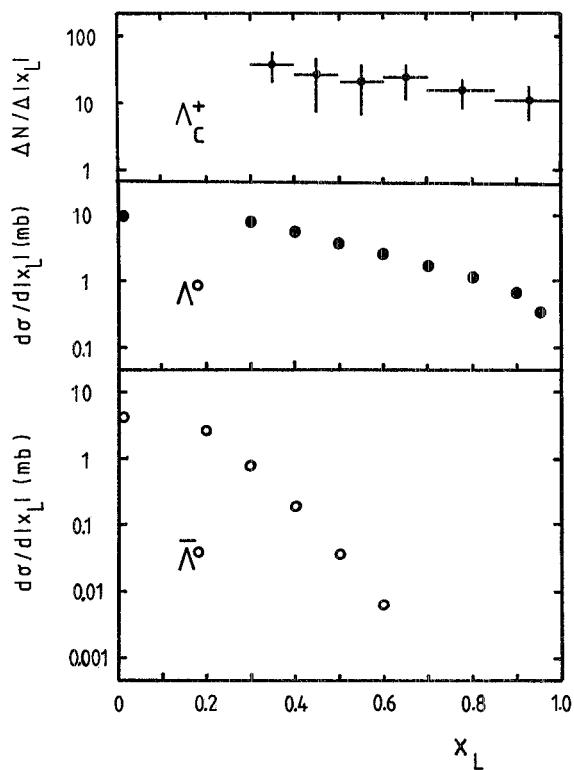


Fig. 4. — Experimental  $x_L$  distribution of  $\Lambda^0$  (black circles) and  $\bar{\Lambda}^0$  (open circles) produced in proton-proton collisions at  $\sqrt{s} = 53$  GeV. Also shown is the  $x_L$  distribution for  $\Lambda_c^+$  in this experiment.

ground was determined, in the same invariant mass range, using the mass spectrum associated to the  $e^+$  trigger, which, as expected for a charmed baryon, does not show any enhancements in the  $\Lambda_c^+$  region.

*Results* — The  $x_L$  distribution of the  $\Lambda_c^+$  events was computed using as reference the  $x_L$  distribution of the events above ( $2.38 < m(pK^-\pi^+) < 2.48 \text{ GeV}/c^2$ ) and below ( $2.18 < m(pK^-\pi^+) < 2.28 \text{ GeV}/c^2$ ) the  $\Lambda_c^+$  mass range. This was in order to determine the shape of the background events inside the  $\Lambda_c^+$  mass range. All events had been corrected for the apparatus acceptance.

The experimental data are shown in fig. 3. They have been fitted with a function:  $(\Delta N/\Delta|x_L|) \propto (1 - |x_L|)^\alpha$ . The dashed line is the best fit to the data. The exponent is found to be

$$\alpha = 0.40 \pm 0.25.$$

By comparing the longitudinal-momentum distribution of the  $\Lambda_c^+$  (quark composition: cud) with that of the  $\Lambda^0$  (quark composition: sud), both produced in high-energy proton-proton collisions, the following remarks are in order. These two particles are very similar, since they have two quarks (ud) which might be in common with the incoming proton and one (c and s, respectively) which is not present in the pp initial state.

Figure 4 shows the experimental  $x_L$  distributions of the  $\Lambda^0$  and  $\bar{\Lambda}^0$ <sup>(4)</sup> together with

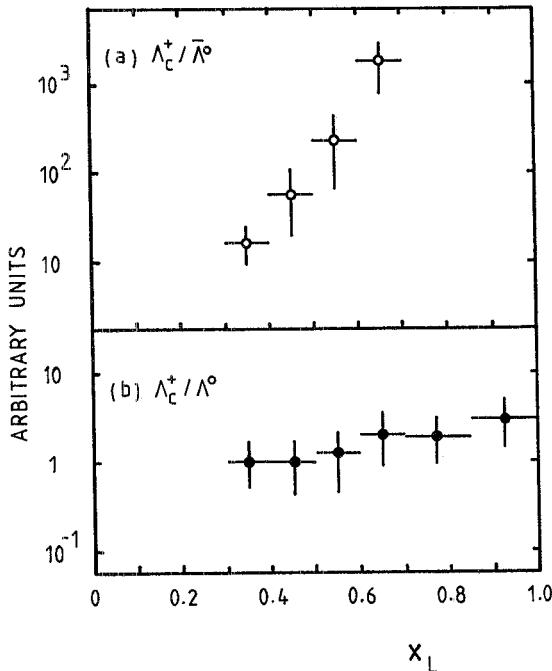


Fig. 5. — Quantitative comparison between  $(\Lambda_c^+/\bar{\Lambda}^0)$  (a)) and  $(\Lambda_c^+/\Lambda^0)$  (b))  $x_L$  distributions.

<sup>(4)</sup> S. ERHAN, W. LOCKMAN, T. MEYER, J. RANDER, P. SCHLEIN, R. WEBB and J. ZSEMBERY: *Phys. Lett. B*, **85**, 447 (1979).

the experimental  $x_L$  distribution of the  $\Lambda_c^+$  measured in the present experiment. The shapes of these distributions, indicate that the  $\Lambda_c^+$  is produced very similarly to the  $\Lambda^0$ , and much flatter in  $x$  than the  $\bar{\Lambda}^0$ .

The quantitative analysis is shown in fig. 5, where the ratios  $(\Lambda_c^+/\bar{\Lambda}^0)$  and  $(\Lambda_c^+/\Lambda^0)$  vs.  $x_L$  are presented. The clear enhancement with increasing  $x_L$  for the  $(\Lambda_c^+/\bar{\Lambda}^0)$  ratio (fig. 5a) is evidence for a leading effect in the  $\Lambda_c^+$  production. The flatness vs.  $x_L$  of the ratio  $(\Lambda_c^+/\Lambda^0)$  (fig. 5b), is evidence that the production mechanism of charmed and strange baryons, in proton-proton interactions, has analogous « leading hadron » effects.