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M. Basile et al. : A MEASUREMENT OF TWO RESONANT  
CONTRIBUTIONS IN THE  $\Lambda_c^+$  BRANCHING RATIOS

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## A Measurement of Two Resonant Contributions in the $\Lambda_c^+$ Branching Ratios.

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**Summary.** — The contribution of two resonant states  $\Delta^{++}$  and  $\bar{K}^{*0}$  to the  $\Lambda_c^+$  decay has been measured. The world average is given.

### 1. - Introduction.

The purpose of this paper is to report on a study of the contributions to the decay  $\Lambda_c^+ \rightarrow pK^-\pi^+$  of the two well-known resonant states

$$(1) \quad \bar{K}^{*0} \rightarrow K^-\pi^+ \quad \text{and} \quad \Delta^{++} \rightarrow p\pi^+.$$

The experiment has been performed by using, as the  $\Lambda_c^+$  production reaction, high-energy pp collisions at the CERN Intersecting Storage Rings (ISR). The reaction investigated was

$$(2) \quad pp \rightarrow \Lambda_c^+ + e^- + \text{anything},$$

at a total pp c.m. energy of  $\sqrt{s} = 62$  GeV.

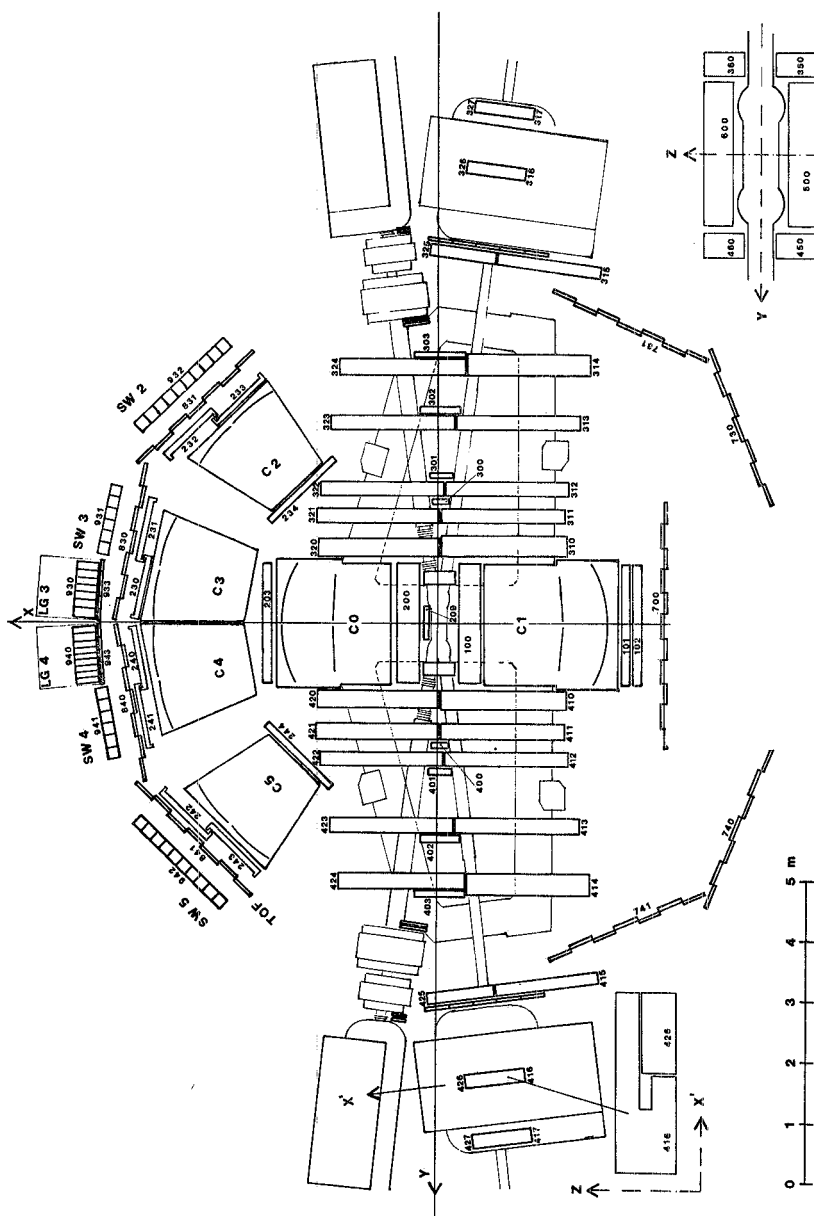


Fig. 1. - Top view of the SFM detector, showing the MWPCs and the external apparatus for particle identification: i)  $C_i$  ( $i = 0, \dots, 5$ ) are gas threshold Čerenkov counters, ii)  $SW_i$  ( $i = 2, \dots, 5$ ) are lead/scintillator sandwiches, iii)  $LG_i$  ( $i = 3, 4$ ) are lead-glass counters. Notice that only coincidences  $C_0 C_3$  or  $C_0 C_4$  associated to a minimum energy release ( $\geq 500$  MeV) in the electromagnetic shower detectors ( $LG_3, SW_3$  or  $LG_4, SW_4$ ) were used in the «electron» trigger.

## 2. - Experimental set-up.

The experimental set-up consists of a system of multiwire proportional chambers (MWPCs) in a large-volume, 10 kG magnetic field, produced by the Split-Field Magnet (SFM). There are also gas Čerenkov counters and electromagnetic shower detectors for «electron» trigger, plus a large assembly of time-of-flight (TOF) counters for particle identification (p, K,  $\pi$ ). Figure 1 shows a schematic lay-out of the set-up, which has been described in detail elsewhere <sup>(1)</sup>.

## 3. - Data analysis.

From a sample of proton-proton collisions with an electron produced at  $90^\circ$  we obtained <sup>(2)</sup> a clear signal of the charmed baryon  $\Lambda_c^+$  decaying via the channel  $pK^-\pi^+$ .

Firstly, the events were analysed in order to reduce the background contribution to genuine electron events. The contamination from charged hadrons simulating electrons was reduced to the 2% level. Background electrons from Dalitz  $\pi^0$  and  $\eta$  decays and from external  $\gamma$ -conversions were rejected to the 50% level via pulse-height analysis with a  $dE/dx$  chamber placed near the intersection region.

From the sample of events left, we selected those with at least one positive particle with  $x = 2p_x/\sqrt{s} \geq 0.3$  and an error on the measured momentum  $\Delta p/p \leq 30\%$ . In these events the positive particle with the highest  $x$  has been assumed to be a proton. In the invariant mass,  $pK^-\pi^+$ ,  $K^-(\pi^+)$  is any negative (positive) track not identified as an  $\bar{p}$  or  $\pi^-$  ( $p$  or  $K^+$ ) by the time of flight. All possible combinations are entered into the mass plot, provided they satisfy the following conditions:

<sup>(1)</sup> R. BOUCLIER, R. C. A. BROWN, E. CHESI, L. DUMPS, H. G. FISCHER, P. G. INNOCENTI, G. MAURIN, A. MINTEN, L. NAUMANN, F. PIUZ and O. ULLALAND: *Nucl. Instrum. Methods*, **125**, 19 (1975); M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: *Nucl. Instrum. Methods*, **163**, 93 (1979); M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. DI CESARE, B. ESPOSITO, L. FAVALE, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: *Nucl. Instrum. Methods*, **179**, 477 (1981); H. FREHSE, F. LAPIQUE, M. PANTER and F. PIUZ: *Nucl. Instrum. Methods*, **156**, 87 (1978); H. FREHSE, M. HEIDEN, M. PANTER and F. PIUZ: *Nucl. Instrum. Methods*, **156**, 97 (1978).

<sup>(2)</sup> M. BASILE, G. CARA ROMEO, L. CIFARELLI, A. CONTIN, G. D'ALI, P. DI CESARE, B. ESPOSITO, P. GIUSTI, T. MASSAM, F. PALMONARI, G. SARTORELLI, G. VALENTI and A. ZICHICHI: submitted to *Nuovo Cimento*.

a) all particles are fitted to the vertex and have a momentum uncertainty  $\Delta p/p \leq 30\%$ ,

b) both  $K^-$  and  $\pi^+$  are in the same rapidity hemisphere as that of the proton,

c) the rapidity of both the  $K^-$  and the  $\pi^+$  is greater than 1.0.

In order to enhance the  $\Lambda_c^+$  signal, we have imposed one of the following conditions on the hemisphere opposite to that of the  $\Lambda_c^+$ :

i) either the presence of a leading system of charged particles with  $x_{\text{total}} = \sum_{i=1}^n x_i \geq 0.5$ , where  $x = 2p_L/\sqrt{s}$ ;

ii) or a signature for a leading system escaping detection, *i.e.*  $x_{\text{total}} \leq 0.1$ .

The above sum  $\sum_{i=1}^n$  is extended to all particles fitted to the vertex and with  $\Delta p/p \leq 30\%$ . The results are shown in fig. 2.

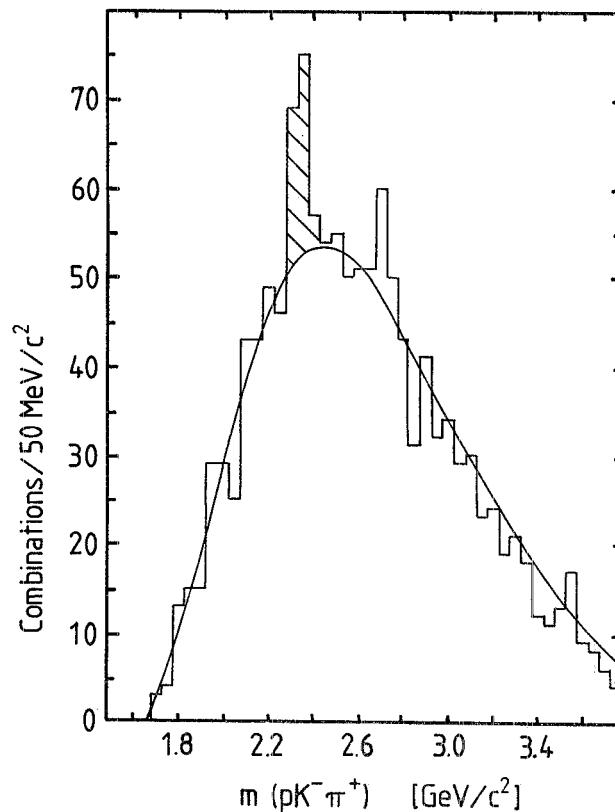


Fig. 2. - Mass distribution of  $pK^-\pi^+$  combinations selected as described in the text. The solid line is a fit to the mass spectrum with  $e^+$  trigger. This shape represents the background.

In order to study the resonant contributions to the decay  $\Lambda_c^+ \rightarrow pK^-\pi^+$  from

$$\Lambda_c^+ \rightarrow \bar{K}^{*0}p \quad \text{and} \quad \Lambda_c^+ \rightarrow \Delta^{++}K^-,$$

the mass spectra for  $K^-\pi^+$  and  $p\pi^+$  have been studied in three different mass ranges for the  $pK^-\pi^+$  mass spectrum. These are indicated in table I.

Figure 3 shows the results. Above (fig. 3a) and below (fig. 3e) the  $\Lambda_c^+$  mass range, the  $K^-\pi^+$  and  $p\pi^+$  mass spectra do not show any enhancement.

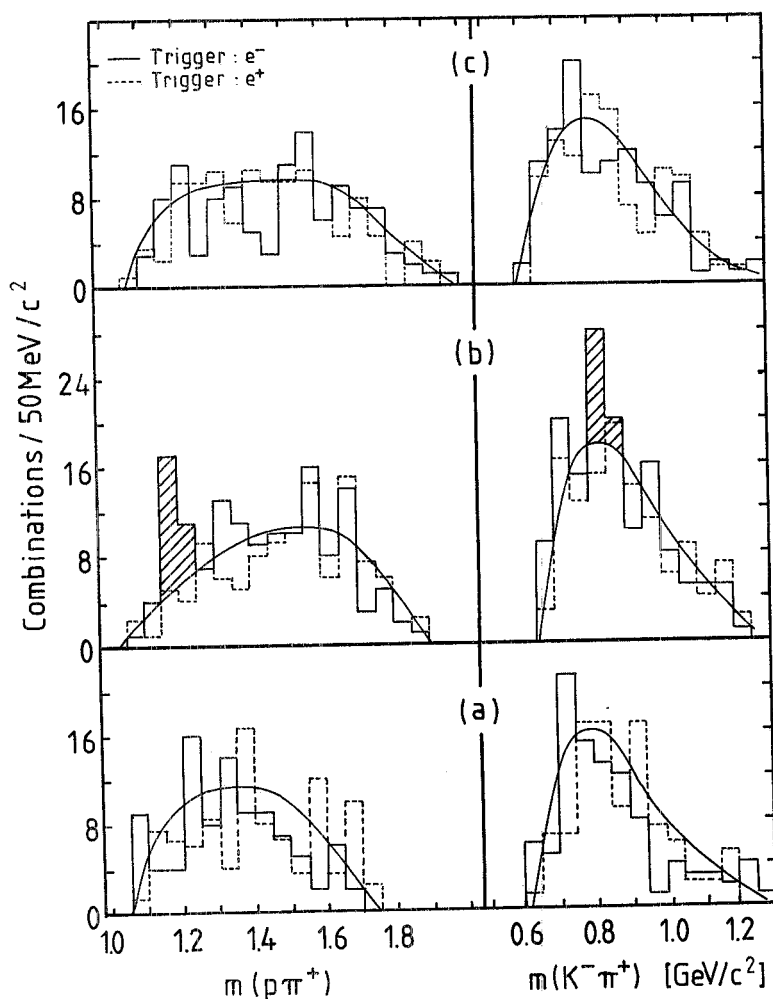


Fig. 3. - Invariant-mass spectra for  $K^-\pi^+$  and  $p\pi^+$  combinations as a function of three different cuts in the  $pK^-\pi^+$  mass: a)  $2.18 \leq m(pK^-\pi^+) < 2.28 \text{ GeV}/c^2$  (below  $\Lambda_c^+$ ), b)  $2.28 \leq m(pK^-\pi^+) < 2.38 \text{ GeV}/c^2$  (in the  $\Lambda_c^+$  peak), c)  $2.38 < m(pK^-\pi^+) \leq 2.48 \text{ GeV}/c^2$  (above  $\Lambda_c^+$ ). The full-line histograms refer to  $e^-$  triggers; the dashed-line histograms refer to  $e^+$  triggers. The solid lines are fits to these mass spectra with  $e^+$  triggers.

This is not the case in fig. 3b), which corresponds to the  $\Lambda_c^+$  mass range. Here we observe a peak in the mass spectrum of both the  $K^-\pi^+$  and the  $p\pi^+$ , at mass values corresponding to  $\bar{K}^{*0}$  and  $\Delta^{++}$ , respectively.

TABLE I. - *Mass ranges of  $pK^-\pi^+$  ( $\text{GeV}/c^2$ ).*

below $\Lambda_c^+$	2.18 ÷ 2.28
in the $\Lambda_c^+$ peak	2.28 ÷ 2.38
above $\Lambda_c^+$	2.38 ÷ 2.48

Notice that the dashed-line histograms correspond to mass spectra with  $e^+$ -triggered events. Here there is no evidence for any resonant state. This is another proof that, in order to have a clean  $\Lambda_c^+$  signal, the  $e^-$  trigger is needed. In fact, the associated production of  $\Lambda_c^+$  with an anticharmed particle requires the presence of an  $e^-$  and not of an  $e^+$ . The  $e^+$ -triggered events were used to draw the shape of the background in the  $e^-$ -triggered spectra.

#### 4. - Results.

From the number of events above background in the  $\bar{K}^{*0}$  and  $\Delta^{++}$  peaks and from the total number of observed  $\Lambda_c^+$ , the following branching ratios are obtained:

$$\frac{\Lambda_c^+ \rightarrow \bar{K}^{*0}p}{\Lambda_c^+ \rightarrow pK^-\pi^+} = 0.28 \pm 0.16,$$

$$\frac{\Lambda_c^+ \rightarrow \Delta^{++}K^-}{\Lambda_c^+ \rightarrow pK^-\pi^+} = 0.40 \pm 0.17.$$

TABLE II. - *Branching ratios of  $\bar{K}^{*0}$ ,  $\Delta^{++}$  resonant contributions to the  $\Lambda_c^+ \rightarrow pK^-\pi^+$  decay channel.*

	$\bar{K}^{*0}p/pK^-\pi^+$	$\Delta^{++}K^-/pK^-\pi^+$
SPEAR (Mark II) <sup>(3)</sup>	$0.12 \pm 0.07$	$0.17 \pm 0.07$
ISR (SAJOT) <sup>(4)</sup>	$0.41 \pm 0.18$	$0.53 \pm 0.22$
ISR (this experiment)	$0.28 \pm 0.16$	$0.40 \pm 0.17$
world average	unweighted	$0.27 \pm 0.08$
	weighted	$0.18 \pm 0.06$
$\frac{\Lambda_c^+ \rightarrow (\Delta^{++}K^- + \bar{K}^{*0}p)}{\Lambda_c^+ \rightarrow pK^-\pi^+} = \begin{cases} 0.64 \pm 0.13 & \text{(unweighted)} \\ 0.41 \pm 0.08 & \text{(weighted)} \end{cases}$		

<sup>(3)</sup> G. GOLDHABER: *Topical Workshop on the Production of New Particles in Superhigh Energy Collisions* (Madison, Wis., 1979); LBL-10428 (January 1980).

<sup>(4)</sup> G. SAJOT: Ph. D. Thesis, University of Paris VI (April 1980).

The interest in these measurements is due to the fact that only two results <sup>(3,4)</sup> are available so far. Table II summarizes the present knowledge of these branching ratios. The world average is also given.

● RIASSUNTO

Analizzando un campione di interazioni protone-protone ad un'energia totale nel centro di massa di 62 GeV, è stata studiata la produzione associata del barione incantato  $\Lambda_c^+$  con una particella antincantata che decade semileptonicamente. È stato misurato il contributo dei due stati risonanti  $\Delta^{++}$  e  $\bar{K}^{*0}$  al decadimento del  $\Lambda_c^+$  nel canale  $\Lambda_c^+ \rightarrow pK^-\pi^+$ .

**Измерение двух резонансных вкладов в отношения ветвей  $\Lambda_c^+$ .**

**Резюме (\*).** — Измеряется вклад двух резонансных состояний  $\Delta^{++}$  и  $\bar{K}^{*0}$  в распад  $\Lambda_c^+$ . Приводится « мировое » среднее.

(\*) *Переведено редакцией.*