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MOMENTUM PRODUCTION DISTRIBUTION IN pp INTERAC-
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Measurement of the Λ_c^+ Transverse-Momentum Production Distribution in pp Interactions at $\sqrt{s} = 62$ GeV.

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Introduction. - Detailed studies of charmed-particle production in high-energy proton-proton interactions are of particular interest. The experiments performed so far have measured the total cross-section for inclusive ($pp \rightarrow D^+X$, $pp \rightarrow \Lambda_c^+X$) (¹⁻⁴) or associated ($pp \rightarrow \Lambda_c^+e^-X$, $pp \rightarrow e^-D^0X$, $pp \rightarrow \bar{\Lambda}_c^0e^+X$) production (^{5,6}). In these processes no direct measurement is available on the production distributions of the charmed particle.

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(²) K. L. GIBONI, D. DIBITONTO, M. BARONE, M. M. BLOCK, A. BÖHM, R. CAMPANINI, F. CERADINI, J. EICKMEYER, D. HANNA, J. IRON, A. KERNAN, H. LUDWIG, F. MULLER, B. NAROSKA, F. NAVACH, M. NISSBAUM, J. O'CONNOR, C. RUBBIA, D. SCHINZEL, H. SEEBRUNNER, A. STAUDE, R. TIRLER, G. VAN DALEN, R. VOSS and R. WOJSLAW (ACHMNR COLLABORATION): *Phys. Lett. B*, **85**, 437 (1979).

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In the field of $D\bar{D}$ production there was only one indirect measurement (coming from μ -production studies) of the transverse-momentum distribution of the charmed mesons in the reaction $pp \rightarrow DDX$ at $\sqrt{s} = 27$ GeV⁽⁷⁾.

The purpose of this paper is to report the first measurement of the transverse-momentum production distribution of the Λ_c^+ charmed baryon in the reaction

$$(1) \quad pp \rightarrow \Lambda_c^+ + e^- + \text{anything}$$

at a centre-of-mass energy of $\sqrt{s} = 62$ GeV, the Λ_c^+ being identified via its decay mode $\Lambda_c^+ \rightarrow pK^-\pi^+$.

The experimental set-up. – The experiment was performed at the CERN Intersecting Storage Rings (ISR), by using the split-field magnet (SFM) detector⁽⁸⁾, with the addition of Čerenkov counters and electromagnetic shower detectors (EMSD)⁽⁹⁾ to trigger on electrons produced at 90°, and a large area time-of-flight system (TOF)⁽¹⁰⁾ for particle identification. A schematic plan of the experimental set-up is shown in fig. 1.

A multiwire proportional chamber (MWPC) with analog read-out (« dE/dx » chamber⁽¹¹⁾), placed near the intersection region, was used at the software level to reduce the background, from electrons produced in the π^0 and η Dalitz decay and in γ conversions, down to about 50% level. The contamination from charged hadrons simulating electrons was reduced to the 2% level, after a software refinement of the trigger conditions.

Data analysis. – From the sample of selected events, a clear signal of Λ_c^+ decaying into $pK^-\pi^+$ ⁽⁵⁾ was obtained. Let us briefly recall how the analysis was performed. (For full details, see ref. (5).)

The events were first required to have at least one positive particle with $x_L \simeq 2p_L/\sqrt{s} > 0.3$ and an error on the measured momentum $\Delta p/p < 30\%$. The positive particle, with the highest x_L , was assumed to be a proton (the π^+ contamination on the proton sample ranges from 20% to 2%, exponentially decreasing with increasing x_L). In the invariant mass $pK^-\pi^+$, the $K^-(\pi^+)$ is any negative (positive) track not identified as \bar{p} or π^- (p or K^+) by the time of flight. All possible combinations (the ratio combinations-to-events is only 1.15) were entered into the mass plot, provided they satisfied the following conditions:

- a) all particles were fitted to the vertex (within ± 5 cm) and had a momentum uncertainty $\Delta p/p < 30\%$;
- b) both K^- and π^+ were in the same rapidity hemisphere as that of the proton;
- c) the rapidity of both the K^- and the π^+ was greater than 1.0.

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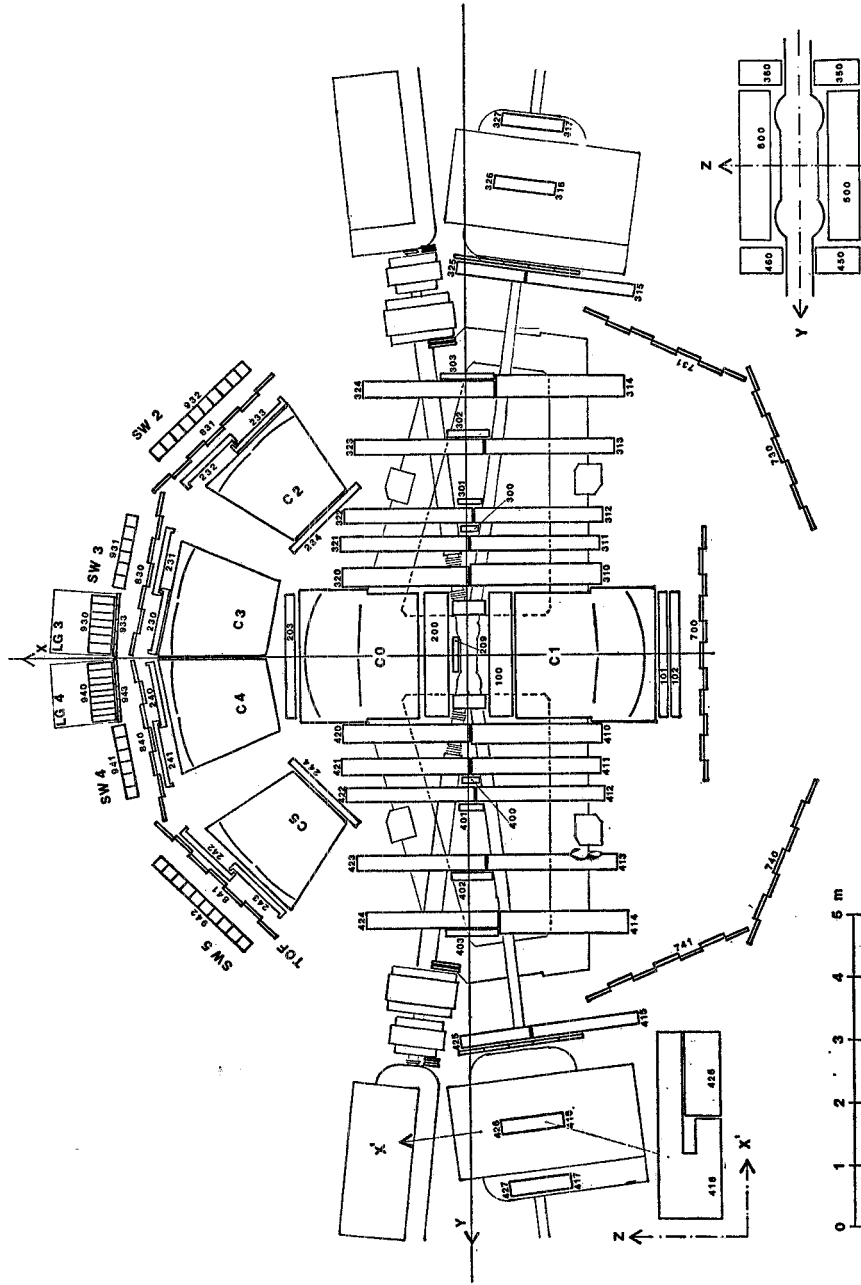


Fig. 1. - Top view of the SFM detector, showing the MWPCs and the external apparatus for particle identification (TOF). C_i ($i=0, \dots, 5$) are the gas threshold Čerenkov counters. The electromagnetic shower detectors (EMSD) are of two types, lead/scintillator sandwiches and lead glass, indicated as SW $_i$ ($i=2, \dots, 5$) and LG $_i$ ($i=3, 4$), respectively. Only the coincidences C_0C_3 or C_0C_4 , associated with a minimum energy release greater than 500 MeV in the EMSDs, were used in the «electron» trigger. The «dE/dx» chamber is indicated by the number 209.

In order to enhance the Λ_c^+ signal, the following conditions were imposed on the hemisphere opposite to that of the Λ_c^+ : either the presence of a leading system of charged particles with $x_{\text{tot}} = \sum_{i=1}^n (x_{iL}) > 0.5$, or a signature for a leading system escaping detection, *i.e.* $x_{\text{tot}} < 0.1$. The above sum was extended to all particles fitted to the vertex and with $\Delta p/p < 30\%$. The results are shown in fig. 2. Here the mass spectrum

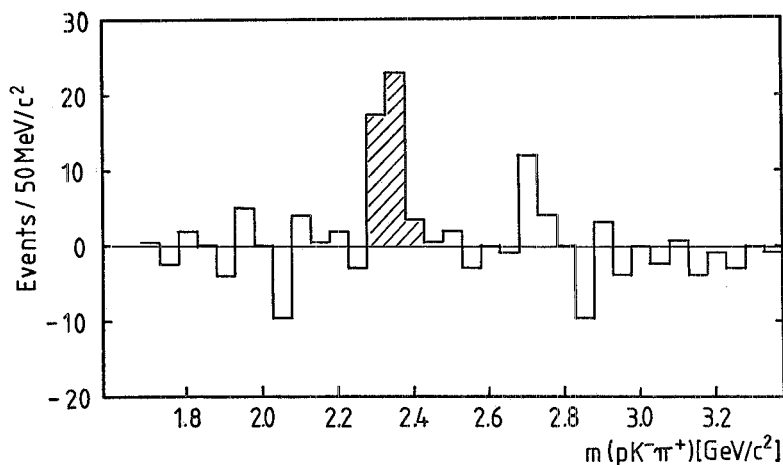


Fig. 2. - Difference between the $pK^- \pi^+$ invariant mass spectrum associated with e^- -triggered events and the same mass combination relative to e^+ -triggered events.

for e^- -triggered events was subtracted from the mass spectrum for e^+ -triggered events. The presence of a « signal » in only the e^- -triggered events is the proof that the enhancement in the $pK^- \pi^+$ mass plot corresponds to a charmed baryon produced in association with an anticharmed particle decaying semi-leptonically. This is the source of the e^- triggers. The « signal » is therefore identified as the charmed baryon Λ_c^+ , first discovered by SAMIOS and collaborators ⁽¹²⁾.

Results. - To study the transverse-momentum (p_T) dependence of the produced Λ_c^+ , the $pK^- \pi^+$ mass spectrum was divided into two different mass ranges: the « IN » region ($2.28 \leq m(pK^- \pi^+) < 2.38 \text{ GeV}/c^2$) and the « OUT » region ($2.18 \leq m(pK^- \pi^+) < 2.28 \text{ GeV}/c^2$, $2.38 \leq m(pK^- \pi^+) < 2.48 \text{ GeV}/c^2$). The events in the « OUT » region have been used to give the shape of the background to the events in the « IN » region. The p_T distribution of the Λ_c^+ was thus determined by subtracting, the « OUT » p_T distribution from the « IN » p_T distribution. The « OUT » p_T distribution was normalized to the total number of background events in the « IN » region. The corrections for the apparatus acceptance were applied to the real events, using single-track acceptance tables computed by means of a Monte Carlo simulation program.

Figure 3 shows the apparatus acceptance, as a function of p_T , for $pK^- \pi^+$ events in the invariant mass range between $2.18 \text{ GeV}/c^2$ and $2.48 \text{ GeV}/c^2$.

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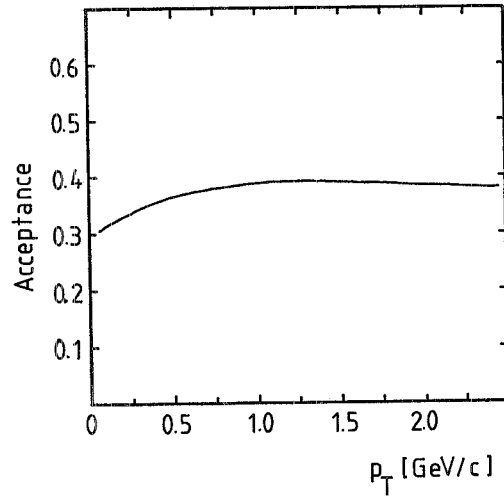


Fig. 3. - Apparatus acceptance as a function of p_T for $pK^-\pi^+$ combinations in the mass range between 2.18 GeV/c² and 2.48 GeV/c².

The experimental data are shown in fig. 4. They have been fitted with an exponential function of the form: $(1/p_T)(\Delta N/\Delta p_T) \propto \exp[-bp_T]$. The best fit gives for the exponent the following result:

$$b = 2.5 \pm 0.4.$$

The fit is indicated by the dashed line in fig. 4.

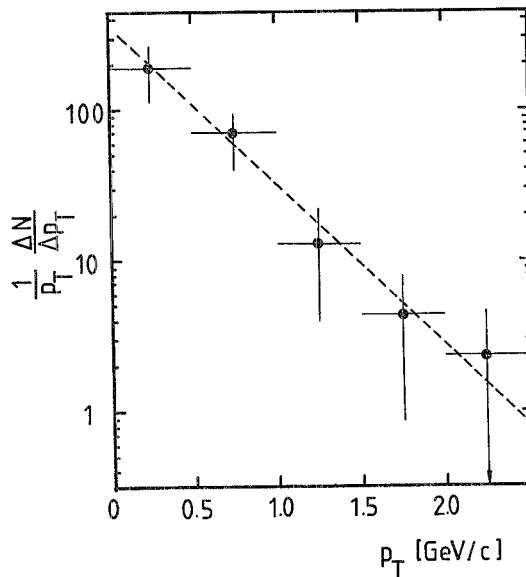


Fig. 4. - Experimental p_T distribution of the Λ_c^+ events. The dashed line is the best fit.

Recent QCD calculations by MARGOLIS and collaborators⁽¹⁸⁾, using two-gluon annihilation as the main process for open charm production in pp collisions, lead to a value

$$b = 2.28,$$

which is in excellent agreement with our experimental result.

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