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E. Etim, A.F. Grillo, G. Pancheri-Srivastava and Y. Srivastava:
COMMENTS ON THE LEPTONIC COUPLINGS OF ψ (3.1 GeV)
RESONANCE. -

E. Etim^(x), A. F. Grillo, G. Pançheri-Srivastava and Y. Srivastava^{(o)(+)}:
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ABSTRACT. -

We study in detail some bounds imposed on the leptonic coupling of the recently discovered ψ (3.1 GeV) resonance. These are based on the magnetic moment anomaly for the muon and the purely leptonic neutral current experiments. If the dip-like structure that seems to be present both in the SPEAR and in the Adone $e^+e^- \rightarrow e^+e^-$ data is taken seriously, then a possible explanation in terms of Lee-Wick type negative metric boson is also presented.

The recently discovered ψ (3.1 GeV) resonance⁽¹⁾ seems to exhibit some peculiarities :

- i) Under the assumption that the resonance is much narrower than the experimental energy resolution ($\Delta E \approx 1.9$ MeV), it has been estimated that $\Gamma(\psi \rightarrow e^+e^-) \simeq \Gamma(\psi \rightarrow \mu^+\mu^-) \simeq 2$ to 5 keV and $\Gamma(\psi \rightarrow \text{hadrons}) \simeq 25$ to 60 keV^(1, 2).
- ii) SPEAR as well as Adone groups^(2, 3) reports seem to show interfe-

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rence effect in the e^+e^- channel - which is destructive before the resonance. The peak in this channel seems shifted (from that seen in the hadronic channel) and the dip is ~ 1 MeV to the left of the peak. In the $\mu^+\mu^-$ decay channel a dip before the peak (again approximately ~ 1 MeV) seems also to exist.

iii) MEA group⁽⁴⁾ at ADONE has also presented some evidence for a charge asymmetry in the $\mu^+\mu^-$ channel which appears to change sign in approximately 2 MeV around the resonance. If confirmed this would imply a strong parity violating component in the resonance decay.

If the interference effects in ii) and iii) (in the $\mu^+\mu^-$ channel) are taken at face value we are forced to say that the spin J of the resonance is ≥ 1 . Theoretically it is appealing then to entertain the possibility that $J=1$ and that it is indeed the intermediate vector boson for neutral weak current⁽⁵⁾. But one immediately runs into severe obstacles with this nice interpretation:

- a) If one accepts that Γ_e and $\Gamma_T \ll \Delta E$, it is practically impossible to see an interference effect (dips and so on) at an energy interval comparable to ΔE .
- b) One obtains an enhancement instead of a dip before the resonance in the e^+e^- decay channel in this scheme.
- c) The sign of the charge asymmetry as given by the MEA group is opposite to that obtained with a standard V-A coupling.

Regarding point a), one possible way out is to suppose that $\Gamma_T \gg \Gamma_{\text{hadron}}$ - the rest coming mainly from decays into neutrinos. In addition, $\Gamma_{e^+e^-}$ and $\Gamma_{\mu^+\mu^-}$ also have to be much larger than the earlier estimates of $\lesssim 5$ keV. Below, we try to bound $\Gamma_{\mu^+\mu^-}$ from the $(g-2)$ for the muon and $\Gamma_{\nu\bar{\nu}}$ from the available $e\nu$ data.

I. $(g-2)$ for the muon.

Under the above hypothesis that the resonance is a pure V-A, one finds for the contribution to a_μ

$$a_{\mu}^{(\psi)} \approx -\frac{4}{\pi} \frac{\Gamma(\psi \rightarrow \mu \bar{\mu})}{m_{\psi}} \left(\frac{m_{\mu}}{m_{\psi}}\right)^2 \approx -(1.5 \times 10^{-9}) (\Gamma_{\mu \bar{\mu}} / \text{keV}).$$

If one compares the experimental data with the QED plus the hadronic contribution one finds the possible anomaly⁽⁶⁾

$$\Delta a_{\mu} \approx -(50 \times 10^{-9}).$$

Thus, in absence of any other cancelling contribution, $\Gamma_{\mu \bar{\mu}} \lesssim 100 \text{ keV}$. If instead of V-A, one had pure V coupling (like in the "charm" or "color" classifications) the bound is much worse, i. e. $\Gamma_{\mu \bar{\mu}}^{(V)} \lesssim 175 \text{ keV}$. Note that in this case $a_{\mu}^{(V)}$ is opposite in sign.

II. Elastic $e\nu$ and $e\bar{\nu}$ scattering.

Regarding some speculations which admit of the possibility that the neutrino decay modes may be anomalously large (for example, as large as $\sim 700 \text{ keV}$!), we present some bounds which rule out this way of escape. Our observations stem from considering $\nu_{\mu} e^{-}$ and $\bar{\nu}_{\mu} e^{-}$ elastic scattering, for which experimental bounds exist. Adler and Tuan⁽⁷⁾ have given expressions for these cross-sections taking into account the experimental cuts. In terms of the $e^{+}e^{-}$ and $\nu_{\mu} \bar{\nu}_{\mu}$ decay widths, these cross-sections become (in the pure V-A form):

$$\sigma(\bar{\nu}_{\mu} e^{-} \rightarrow e^{-} \bar{\nu}_{\mu}) \approx (0.053) \left(\frac{\Gamma_e \Gamma_{\nu_{\mu}}}{\text{keV}^2}\right) (0.035 \times 10^{-41} \text{ cm}^2 \frac{E_{\bar{\nu}}}{\text{GeV}}),$$

$$\sigma(\nu_{\mu} e^{-} \rightarrow e^{-} \nu_{\mu}) \approx (0.16) \left(\frac{\Gamma_e \Gamma_{\nu_{\mu}}}{\text{keV}^2}\right) (0.054 \times 10^{-41} \text{ cm}^2 \frac{E_{\bar{\nu}}}{\text{GeV}}),$$

to be compared with CERN Gargamelle results

$$(0.03) \times 10^{-41} \text{ cm}^2 \frac{E_{\bar{\nu}}}{\text{GeV}} < \sigma(\bar{\nu}_{\mu} e^{-} \rightarrow e^{-} \bar{\nu}_{\mu}) < (0.3) \times 10^{-41} \text{ cm}^2 \frac{E_{\bar{\nu}}}{\text{GeV}}$$

4.

and

$$\sigma(\nu_{\mu} e^{-} \rightarrow e^{-} \nu_{\mu}) < 0.26 \times 10^{-41} \text{ cm}^2 E_{\bar{\nu}}/\text{GeV}.$$

Thus, one obtains

$$18 \lesssim \Gamma_e \Gamma_{\nu_{\mu}} \lesssim 30 \text{ keV}^2.$$

Experimentally, from $\gamma\gamma$ group at Adone, $\Gamma_e > 2.4 \text{ keV}$, which limits

$$\Gamma_{\nu_{\mu}} < 12.5 \text{ keV}.$$

Similar bounds can also be obtained for Γ_{ν_e} . Saturating the experimental upper bound

$$\sigma(\bar{\nu}_e e^{-} \rightarrow e^{-} \bar{\nu}_e) < (1.3 \times 10^{-41} \text{ cm}^2) (E_{\bar{\nu}}/\text{GeV}),$$

we obtain

$$\Gamma_{\nu_e} \lesssim 20 \text{ keV}.$$

These bounds are not much different for Weinberg-Salam type models. Thus, the total decay width cannot be substantially increased by appealing to the neutrino channels.

Conclusion

As discussed previously, taking the dip which is seen in the e^+e^- data about 1 MeV below the peak at its face value, one is led to demand a large total width which however cannot now be attributed to known leptonic channels. The problem of the destructive interference before the resonance can be partially understood provided one employs a negative metric boson of the Lee-Wick type. Also, the charge asymmetry in $\mu^+\mu^-$ channel as observed requires one to relax the pure V-A structure and thus give up chirality. Lastly the possibility of spin other than 1 is still open. We hope to return to these problems elsewhere.

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