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A. Bramon and M. Greco : THE REACTION  $e^+e^- \rightarrow \pi^+\pi^-\pi^+\pi^-$   
AND THE  $\rho^+$ -MESON.

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## The Reaction $e^+e^- \rightarrow \pi^+\pi^-\pi^+\pi^-$ and the $\rho'$ -Meson.

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In the preceding letter BARBARINO *et al.* (1) have reported on the observation of the process

$$(1) \quad e^+e^- \rightarrow \pi^+\pi^-\pi^+\pi^-$$

at total c.m. energies of the colliding beams  $2E = \sqrt{s}$  ranging from 1.2 to 2.4 GeV. The quantum numbers of the  $4\pi$  final state are well defined:  $J^{PC} = 1^{--}$ ,  $G = +$  and consequently  $I = 1$ , the same ones of the  $\rho$ -meson. The actual results however can be hardly understood under the assumption that reaction (1) proceeds only through a  $\rho$  intermediate state, because of the smallness of the cross-section at the total energies of 0.99 GeV (the ACO results (2)), 1.2 and  $\sim 2.0$  GeV. The threshold in fact for  $4\pi$  production is 0.6 GeV roughly, while the  $A_1\pi$  (3,4) mode has a smoother energy behaviour than the experiment actually shows.

On the other hand, an attractive and at the same time natural explanation of the data can be given assuming the existence of a new vector meson, the  $\rho'$ . This meson has been suggested by several theoretical models (5) and introduced by the authors (6) some time ago as a possible explanation of the rather high cross-sections observed in the multi-hadron production at ADONE.

In the present note we reexamine the previous attempt in light of new experimental information and analyse the implications concerning process (1).

The starting point of our model is to get an estimate of the coupling constants of the  $\rho'$  with the photon,  $g_{\rho'\gamma} = em_{\rho'}^2/f_{\rho'}$ , and with the mesons involved in its different decay modes. This was achieved (see ref. (6) for more details) by fitting the available experimental results concerning the decay rates of some pseudoscalar and vector

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(1) G. BARBARINO, F. CERADINI, M. CONVERSI, M. GRILLI, E. IAROCCI, M. NIGRO, L. PAOLUZI, R. SANTONICO, P. SPILLANTINI, L. TRASATTI, V. VALENTE, R. VISENTIN and G. T. ZORN: *Lett. Nuovo Cimento*, **3**, 689 (1972).

(2) J. LEFRANCOIS: talk given at the 1971 Intern. Symp. on Electron and Photon Interactions at High Energies, Cornell University (Ithaca, 1971).

(3) J. LAYSSAC and F. M. RENARD: *Lett. Nuovo Cimento*, **1**, 197 (1971).

(4) M. T. VAUGHN and P. J. POLITO: *Lett. Nuovo Cimento*, **1**, 74 (1971).

(5) G. VENEZIANO: *Nuovo Cimento*, **57A**, 190 (1968); J. A. SHAPIRO: *Phys. Rev.*, **179**, 1345 (1969).

(6) A. BRAMON and M. GRECO: *Lett. Nuovo Cimento*, **1**, 739 (1971).

mesons in the spirit of a vector-dominance model generalized with the inclusion of these new states. In particular, for our actual purposes, it suffices to consider the chain of processes (7)

$$(2) \quad \omega \rightarrow \pi^+ \pi^- \pi^0, \quad \omega \rightarrow \pi^0 \gamma, \quad \pi^0 \rightarrow \gamma \gamma,$$

which are simply related by VMD and whose experimental values are firmly established (8). By using the old ORSAY results (9) for the photon–vector-meson couplings, AUGUSTIN *et al.* (10) showed that a disagreement between the VMD predictions and the experimental values, of roughly a factor of two, was obtained at each step of the chain (2). We showed (6) that the existence of a new vector nonet could restore the agreement if the coupling constants of the new mesons satisfy the relation

$$(3) \quad \lambda = \frac{g_{\rho' \omega \pi}}{g_{\rho \omega \pi}} \frac{f_{\rho'}}{f_{\rho}} \simeq -0.2.$$

If however the new ORSAY data (11), now in agreement with those of NOVOSIBIRSK (12), where used, the disagreement in ref. (10) would be reduced to a factor 1.4 roughly at each step. Correspondingly, in our extended vector-dominance model the parameter  $\lambda$  introduced in (3) should be given by

$$(4) \quad \lambda \simeq -0.12.$$

From this result and by assuming  $m_{\rho'} \simeq 1.6$  GeV, we predict the partial width of the  $\rho' \rightarrow \omega \pi \rightarrow \pi^+ \pi^- \pi^0 \pi^0$  decay mode to be

$$(5) \quad \Gamma(\rho' \rightarrow \omega \pi) = 2.4 \left( \frac{f_{\rho'}}{f_{\rho}} \right)^2 \Gamma(\omega \rightarrow \pi^+ \pi^- \pi^0),$$

and therefore a factor of 3 smaller than our previous result (6). By means of usual  $SU_3$  and phase-space evaluation techniques and introducing the experimental value  $\Gamma(\omega \rightarrow \pi^+ \pi^- \pi^0) \simeq 10$  MeV we obtain for the main  $\rho' \rightarrow VP$  decay modes

$$(6) \quad \Gamma(\rho' \rightarrow \omega \pi) = 24 \left( \frac{f_{\rho'}}{f_{\rho}} \right)^2, \quad \Gamma(\rho' \rightarrow \rho \eta) = 7 \left( \frac{f_{\rho'}}{f_{\rho}} \right)^2, \quad \Gamma(\rho' \rightarrow K^* K) = 4 \left( \frac{f_{\rho'}}{f_{\rho}} \right)^2,$$

where all the widths are expressed in MeV. We have also estimated for the pseudoscalar-pseudoscalar decays modes  $\Gamma(\rho' \rightarrow PP) \simeq 8(f_{\rho'}/f_{\rho})^2$  MeV, which is compatible with the experimental information available (13). By adding all these rates for the decays  $\rho' \rightarrow VP, PP$ , we obtain  $\Gamma_{VP,PP} \simeq 43(f_{\rho'}/f_{\rho})^2$  MeV  $\simeq 215$  MeV where the value  $(f_{\rho'}/f_{\rho})^2 \approx 5$ ,

(7) A similar analysis of the three related processes has been done by R. A. BRANDT and G. PREPARATA: *Phys. Rev. Lett.*, **25**, 1530 (1970), by using operator product expansions techniques.

(8) PARTICLE DATA GROUP: *Rev. Mod. Phys.*, **43** (April 1971).

(9) See for instance ref. (8).

(10) J. E. AUGUSTIN, D. BENAKSAS, J. C. BIZOT, J. BUON, B. DELCOURT, V. GRACCO, J. HAISSINSKI, J. JEAN JEAN, D. LALANNE, F. LAPLANCHE, J. LEFRANÇOIS, P. LEHMANN, P. MARIN, H. NGUYEN NGOC, J. PEREZ-Y-JORBA, F. RICHARD, F. RUMPF, E. SILVA, S. TAVERNIER and D. TREILLE: *Phys. Lett.*, **28 B**, 503 (1969).

(11) See ref. (8).

(12) See, for example, R. MARSHALL: *Proceedings of the Daresbury Study Weekend* (June 1970).

(13) A. BRAMON and M. GRECO: *Journ. Phys.*, **C 3**, 93 (1971).

estimated in ref. (6), has been tentatively used. Further decay modes are therefore required if the total  $\rho'$  width is about 0.35 GeV. This value is roughly confirmed by the new results (14) on the reaction  $\gamma p \rightarrow p\pi^+\pi^-\pi^+\pi^-$ , besides those of the preceding letter, both experiments suggesting a sizeable  $\rho' \rightarrow \pi^+\pi^-\pi^+\pi^-$  decay mode.

More in detail, DAVIER *et al.* (14) have studied the just mentioned reaction at photon energies up to 18 GeV, showing evidence for a  $4\pi^\pm$  enhancement centered at  $(1.6 \pm 0.1)$  GeV with a width of  $(0.5 \pm 0.1)$  GeV and the following properties: *a*) the corresponding cross-section of  $(0.9 \pm 0.2) \mu b$  is energy independent above 6 GeV; the slope parameter of  $d\sigma/dt$  is  $(5.8 \pm 0.4) (\text{GeV})^{-2}$ , very near to the value  $(6.6 \pm 0.4) (\text{GeV})^{-2}$  for the analogous slope in the reaction  $\gamma p \rightarrow \rho^0 p$  at the same energies.

This effect can be understood in terms of diffractive production of a  $\rho'$  which subsequently decays into four pions. By assuming the equality of the elastic  $\rho$ -nucleon and  $\rho'$ -nucleon cross-sections, as can be expected from the quark model, and neglecting nondiagonal terms  $\rho p \rightarrow \rho' p$  with respect to the diagonal ones  $\rho p \rightarrow \rho p$  and  $\rho' p \rightarrow \rho' p$  we easily obtain

$$(7) \quad \frac{\Gamma(\rho' \rightarrow 4\pi^\pm)}{\Gamma_{\rho'}} \simeq (0.08 \pm 0.02) \left( \frac{f_{\rho'}}{f_\rho} \right)^2,$$

where we have used for the  $\rho$ -photoproduction cross-section the value  $\sigma_\rho = 12 \mu b$  (15).

If is easy to evaluate the implications of this result in the  $4\pi^\pm$  production in  $e^+e^-$  annihilation, by using the expression for the  $e^+e^- \rightarrow \rho' \rightarrow f$  cross-section

$$(8) \quad \sigma_f(s) = 16\pi^2 \alpha^2 \frac{m_{\rho'}}{f_{\rho'}^2} \frac{1}{s^{\frac{3}{2}}} \frac{\Gamma_{(\rho' \rightarrow f)}}{(s - m_{\rho'}^2)^2 + m_{\rho'}^2 \Gamma_{\rho'}^2}.$$

We find

$$(9) \quad \sigma_{e^+e^- \rightarrow 4\pi^\pm} = (1.6 \pm 0.5) \cdot 10^{-32} \text{ cm}^2.$$

The energy dependence is shown in Fig. 1 (\*) and it is compared with the experiments. Looking at this Figure it should be realized also the presence of a background essentially coming from the  $A_1\pi$  contribution of the  $\rho$  tail previously discussed, which can be roughly estimated to be about 5 nb in the average.

In order to have more precise information on the four pions  $\rho'$ -decay modes, we have studied the reaction  $\rho' \rightarrow \rho\pi\pi$  with techniques of finite-energy sum rules, similar to the approach used in  $\rho\pi$  scattering. We omit the details, which will be discussed elsewhere. The preliminary results of this analysis are  $\Gamma(\rho' \rightarrow \rho\pi^+\pi^-) \simeq 30(f_{\rho'}/f_\rho)^2$  MeV and  $\Gamma(\rho' \rightarrow \rho\pi\pi^0) \simeq 12(f_{\rho'}/f_\rho)^2$  MeV in rough agreement with eq. (7).

We are now able to estimate the strength of the  $\rho'$ -photon coupling by adding all the different decay rates to a total width of  $\simeq 0.35$  GeV. We obtain  $(f_{\rho'}/f_\rho)^2 \simeq 4$ . This value is consistent with the typical limits given in the literature, using however realistic assumption on the  $\rho'$  width and on the branching ratio to two pions (13).

(14) M. DAVIER, I. DERADO, D. C. FRIES, F. F. LIU, R. F. MOZLEY, A. C. ODIAN, J. PARK, W. P. SWANSON, F. VILLA and D. YOUNT: SLAC Preprint (June 1971).

(15) J. PARK, M. DAVIER, I. DERADO, D. C. FRIES, F. F. LIU, R. F. MOZLEY, A. C. ODIAN, W. P. SWANSON, F. VILLA and D. YOUNT: SLAC-PUB-972 (1971).

(\*) The experimental points are taken from ref. (1).

We discuss now briefly the implications of our last result on proton Compton scattering and vector-meson photoproduction, where discrepancies of the order of factors of two are observed for the  $\rho$ -photon coupling constants when compared to the values found in  $e^+e^-$  annihilation.

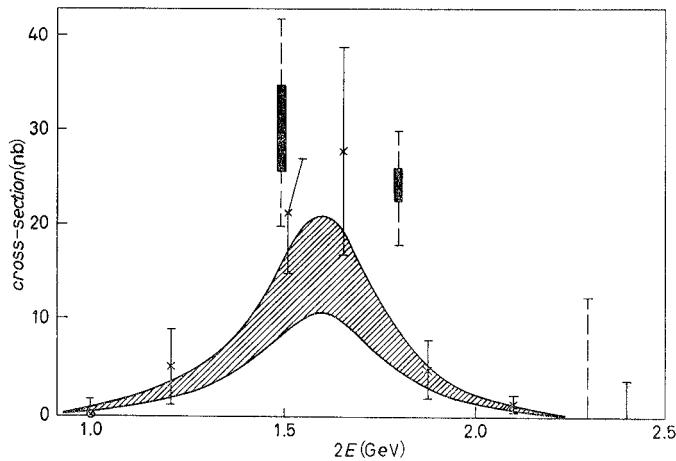


Fig. 1. —  $e^+e^- \rightarrow \pi^+\pi^-\pi^+\pi^-$ ;  $\odot$  ACO, ■ Boson group,  $\times$   $\mu\pi$  » experiment.

From the VDM model follows the relation

$$(10) \quad \sigma_{\text{tot}}(\gamma p) = \sqrt{4\pi\alpha} \sum_{v=\rho,\omega,\varphi} \sqrt{\frac{4}{f_v^2/4\pi} \frac{1}{1 + \eta_v^2} \frac{d\sigma^0}{dt} (\gamma p \rightarrow vp)},$$

where  $d\sigma^0/dt$  denotes the forward cross-section and  $\eta_v$  the ratio of the real to the imaginary part. It is well known that the value of  $f_\rho^2/4\pi$  obtained through eq. (10), from the total  $\gamma p$  cross-section and from photoproduction data is smaller by a factor of about two than the results from  $e^+e^-$  colliding beams. More explicitly one gets  $f_\rho^2/4\pi = 1.4 \pm 0.2$  to be compared with the value  $2.56 \pm 0.22$  <sup>(1)</sup>. Since eq. (10) is based on VMD alone the disagreement has to be attributed to a failure of the model. The situation substantially changes in the presence of higher vector mesons coupled to the photon. Let us in fact extend the sum over the vector mesons in eq. (10), by including as next step the  $\rho'$  and its  $SU_3$  partners. By using the above value  $(f_{\rho'}/f_\rho)^2 \approx 4$  and assuming  $d\sigma^0/dt(\gamma p \rightarrow \rho' p) = (f_{\rho'}/f_\rho)^2 d\sigma^0/dt(\gamma p \rightarrow \rho p)$  the r.h.s. of eq. (10) rises by a factor 1.25, to be compared with 1.35 approximately.

This improvement of the VMD results suggests that moderate but significant corrections can arise from the continuum of higher-mass states, and underlines the fact that the hadronic properties of the photon cannot only be described by the nearby vector mesons, namely  $\rho$ ,  $\omega$  and  $\varphi$ . Similar conclusions have been obtained by ALTARELLI and PREPARATA <sup>(16)</sup> who have estimated the contribution from the continuum by using the operator expansions near the light-cone.

It is very tempting to speculate on what can be the overall contribution of the vector mesons coupled to the photon. Should the naive relation  $f_\rho(n)/f_\rho = 2^n$  be true

<sup>(16)</sup> G. ALTARELLI and G. PREPARATA: University of Rome preprint n. 300 (1971).

for every state  $n$ , the eq. (10) has to be replaced by

$$(11) \quad \sigma_{\text{tot}}(\gamma p) = \sqrt{4\pi\alpha} \sum_{v=\rho,\omega,\phi} \sqrt{\frac{4}{f_v^2/4\pi} \frac{1}{1 + \eta_v^2} \frac{d\sigma^0}{dt} (\gamma p \rightarrow V p)} \left\{ 1 + \frac{1}{4} + \frac{1}{4^2} + \dots \right\} = \\ = \frac{4}{3} \sqrt{4\pi\alpha} \sum_{v=\rho,\omega,\phi} \sqrt{\frac{4}{f_v^2/4\pi} \frac{1}{1 + \eta_v^2} \frac{d\sigma^0}{dt} (\gamma p \rightarrow V p)},$$

in very good agreement with experiment. This leads to an effective  $\rho$ -photon coupling measured in photon-nucleon reactions

$$(12) \quad \left( \frac{\gamma_\rho^2}{4\pi} \right)_{\text{eff}} = \frac{1}{4} \left( \frac{f_\rho^2}{4\pi} \right)_{\text{eff}} = 0.36,$$

where the value  $f_\rho^2/4\pi = 2.56$  has been used. The value (12) is in surprising agreement with different high-energy results<sup>(17)</sup> on photoabsorption and Compton scattering experiments.

Finally, we observe that experimental evidence in favour of the  $1^-$  ( $\sim 1680$ ) but against the  $1^-$  ( $\sim 1250$ ) is predicted in some Veneziano-type form factors<sup>(18)</sup>, which allow the photon to couple only to the even daughters of the  $\rho$ -meson.

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<sup>(17)</sup> D. O. CALDWELL, V. B. ELINGS, W. P. HESSE, R. J. MORRISON, F. V. MURPHY, B. W. WORSTER and D. E. YOUNT: *Phys. Rev. Lett.*, **25**, 609 (1970); G. BUSCHHORN, L. CRIEGEE, L. DUBAL, G. FRANKE, G. GEWENIGER, P. HEIDE, R. KOTTHAUS, G. POELZ, U. TIMM, K. WEGENER, H. WERNER, M. WOUG and W. ZIMMERMANN: *Phys. Lett.*, **33 B**, 241 (1970).

<sup>(18)</sup> See, for example, H. SUURA: *Phys. Rev. Lett.*, **23**, 551 (1969).