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G. Parisi and F. Zirilli : HARD BREMSSTRAHLUNG IN  $e^+e^-$   
COLLISIONS

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**Hard Bremsstrahlung in  $e^+e^-$  Collisions.**

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It is well known that there are important radiative corrections to  $e^+e^-$  scattering due to bremsstrahlung. It is usually assumed that these corrections are important only for the emission of soft photons; in the case they are hard, the angle between them and the direction of at least one of the incoming or outgoing electrons should be of order  $m/E_\gamma$ .

The aim of this letter is to show that there is an important contribution to the total production cross-section due to the process  $e^+e^- \rightarrow e^+e^-\gamma$  where the final photon is hard and emitted with a large angle. An approximate evaluation brings to the following typical results. We denote by  $\bar{\theta}$  the angle between the incoming and the outgoing  $e^-$  in the centre-of-mass frame, by  $\theta_{in}$  and  $\theta_{out}$  the angles between the photon and the incoming and outgoing  $e^-$ , by  $\eta$  the ratio between the photon energy  $E_\gamma$  and the beam energy  $E$ . With this notation the cross-section for the production of  $e^+e^-\gamma$  with  $135^\circ > \bar{\theta} > 45^\circ$ ,  $\theta_{in}, \theta_{out} > 20^\circ$  is  $\sim 0.6\sigma_\mu$ , where  $\sigma_\mu = \pi\alpha^2/3E^2$  is the total cross-section for the production of a  $\mu^+\mu^-$  pair.

This result is brought out in the following way: using an almost real approximation, one gets for the differential cross-section, in the case of zero-mass electrons,

$$(1) \quad d\sigma = \sigma(\bar{\theta}) d\Omega_{e^-} \frac{\alpha}{4\pi^2} \left[ \frac{1}{\sin^2 \theta_{in}} + \frac{1}{\sin^2 \theta_{out}} \right] \frac{1 + \eta^2}{\eta} d\Omega_\gamma d\eta,$$

where

$$(2) \quad \sigma(\bar{\theta}) = \frac{\alpha^2}{8E^2} \left[ \frac{1 + \cos^4(\theta/2)}{\sin^4(\theta/2)} - \frac{2\cos^4(\theta/2)}{\sin^2(\theta/2)} + \frac{1 + \cos^2\theta}{2} \right]$$

is the differential cross-section for the Bhabha scattering.

This formula is obtained if one treats  $\eta, \theta_1, \theta_2$  as small quantities. If we try to use this formula at all angles and energies, we find that the ratio between  $\sigma(\bar{\theta})$  and the cross-section for the production of  $e^+e^-\gamma$  with the same  $\bar{\theta}$  and  $\eta > 0.1, \theta_1, \theta_2 > 20^\circ$  is  $\sim 40$ .

By a simple integration over  $\bar{\theta}$  we obtain the required cross-section. In order to test the reliability of this result, we have to check the validity of formula (1) in the large-angle region.

To do this we have performed the exact calculation of the differential cross-section,

not using the trace theorems, but calculating the helicity amplitudes and summing their squared absolute value<sup>(1)</sup>. The results are in perfect agreement with the trace calculations of ref. (2).

The ratio between the exact value and formula (2) is plotted in Figs. 1 and 2.

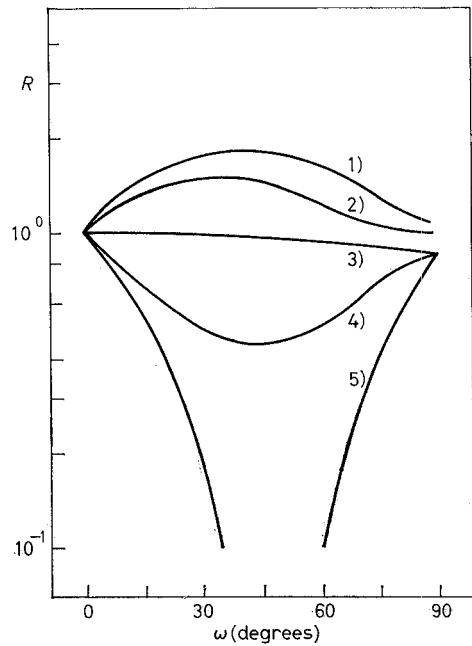


Fig. 1.

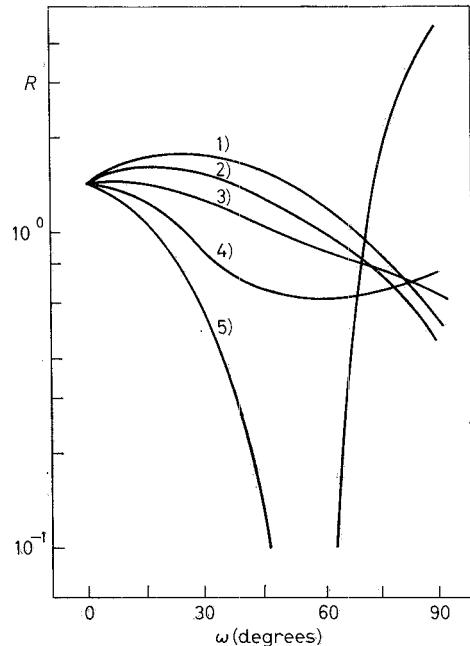


Fig. 2.

Fig. 1. —  $R$  is the ratio between formula (1) and the exact cross-sections. The angle  $\varphi$  has the following values: curve 1)  $\varphi = 0$ , curve 2)  $\varphi = 45^\circ$ , curve 3)  $\varphi = 90^\circ$ , curve 4)  $\varphi = 135^\circ$ , curve 5)  $\varphi = 180^\circ$ .  $\eta = 0.1$ ,  $\bar{\theta} = 90^\circ$ .

Fig. 2. —  $R$  is the ratio between formula (1) and the exact cross-sections. The angle  $\varphi$  has the following values: curve 1)  $\varphi = 0$ , curve 2)  $\varphi = 45^\circ$ , curve 3)  $\varphi = 90^\circ$ , curve 4)  $\varphi = 135^\circ$ , curve 5)  $\varphi = 180^\circ$ .  $\eta = 0.4$ ,  $\bar{\theta} = 90^\circ$ .

The cross-section is calculated at  $\bar{\theta} = 90^\circ$ ,  $\omega$  and  $\varphi$  are the latitude and longitude of a polar co-ordinate system where the  $\omega = 0$ ,  $\varphi = 0$  axis is in the beam direction and the  $\omega = 90^\circ$ ,  $\varphi = 0$  axis is in the direction of the outgoing electron.

One can easily see that our approximation is not too bad: eq. (2) gives, in the average, the correct result. The most prominent departure is given by a destructive interference which causes a zero at  $\varphi \sim 180^\circ$ ,  $\omega \sim 45^\circ$ .

We can conclude that our result is essentially correct and, therefore, the process  $e^+e^- \rightarrow e^+e^-\gamma$  gives a relevant contribution to the multiple production of noncollinear particles at large angle in  $e^+e^-$  colliding beams.

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(1) The details of this method can be found in G. PARISI and F. ZIRILLI: to be published.  
(2) S. M. SWANSON: *Phys. Rev.*, **154**, 1601 (1967).