

Laboratori Nazionali di Frascati

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**Total and Differential Cross-Sections
of Negative-Pion Photoproduction off Neutrons up to 1 GeV.
Possible Evidence for the P_{11} Resonance.**

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We report in this paper some preliminary results on the total and differential cross-sections of the reaction $\gamma+n \rightarrow p+\pi^-$ obtained by exposing a deuterium bubble chamber to the 1 GeV bremsstrahlung beam of the Frascati electron-synchrotron. In a previous paper ⁽¹⁾ we have reported some results concerning the double pion photoproduction off neutrons, using data obtained in the same experiment.

The main result of this investigation is a significant difference between the cross-sections of π^+ and π^- photoproduction which appears in the second resonance region, if the comparison is made for the same $\pi-N$ c.m. energy. Indeed, while the π^+ cross-section shows the well-known $D_{13}(1520)$ peak at 1480 MeV, for the π^- the bump is displaced backwards at about 1400 MeV.

The method of data analysis is very similar to the one described in ref. ⁽¹⁾. The actual measured reaction is $\gamma+d \rightarrow p_s+p+\pi^-$, where p_s is the spectator proton, which is assumed to be the one with the lower momentum.

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TABLE I.

	E^{*2} ((GeV) ²)	E_{γ} (GeV)	σ (μb)	10 \div	20 \div	30 \div	40 \div	50 \div	60 \div
				$\div 20$	$\div 30$	$\div 40$	$\div 50$	$\div 60$	$\div 70$
I	1.22 \div $\div 1.26$	0.179 \div $\div 0.201$	131.2 \pm ± 5.2	3.3	3.8	3.9	5.2 \pm ± 1.2	6.9 \pm ± 1.3	10.6 \pm ± 1.5
II	1.26 \div $\div 1.30$	0.201 \div $\div 0.222$	133.9 \pm ± 5.6	2.6	1.0	1.5 \pm ± 0.9	5.1 \pm ± 1.3	8.9 \pm ± 1.6	8.8 \pm ± 1.5
III	1.30 \div $\div 1.34$	0.222 \div $\div 0.243$	162.5 \pm ± 6.7	6.7	4.7	3.7 \pm ± 1.4	5.3 \pm ± 1.4	8.5 \pm ± 1.6	9.1 \pm ± 1.6
IV	1.34 \div $\div 1.38$	0.243 \div $\div 0.265$	187.3 \pm ± 7.5	4.0	3.4	3.9 \pm ± 1.5	8.9 \pm ± 1.9	7.7 \pm ± 1.6	15.2 \pm ± 2.2
V	1.38 \div $\div 1.42$	0.265 \div $\div 0.286$	193.4 \pm ± 8.1	9.4	2.2	3.6 \pm ± 1.5	5.0 \pm ± 1.5	6.5 \pm ± 1.6	10.7 \pm ± 1.9
VI	1.42 \div $\div 1.46$	0.286 \div $\div 0.307$	215.2 \pm ± 9.1	5.2	2.6	4.8 \pm ± 1.8	9.7 \pm ± 2.3	10.4 \pm ± 2.2	18.4 \pm ± 2.7
VII	1.46 \div $\div 1.50$	0.307 \div $\div 0.328$	214.0 \pm ± 9.5	0	3.6	5.3 \pm ± 2.0	12.4 \pm ± 2.7	10.2 \pm ± 2.3	23.0 \pm ± 3.3
VIII	1.50 \div $\div 1.54$	0.328 \div $\div 0.350$	177.6 \pm ± 9.0	0	1.3	7.7 \pm ± 2.6	10.6 \pm ± 2.6	15.1 \pm ± 2.9	17.7 \pm ± 3.0
IX	1.54 \div $\div 1.62$	0.350 \div $\div 0.392$	167.3 \pm ± 6.7	5.7	7.6 \pm ± 2.3	9.7 \pm ± 2.1	9.9 \pm ± 1.9	11.5 \pm ± 1.9	10.9 \pm ± 1.7
X	1.62 \div $\div 1.70$	0.392 \div $\div 0.435$	134.7 \pm ± 6.5	10.9	9.1 \pm ± 2.6	7.4 \pm ± 2.0	12.9 \pm ± 2.3	14.5 \pm ± 2.3	10.2 \pm ± 1.8
XI	1.70 \div $\div 1.78$	0.435 \div $\div 0.477$	102.0 \pm ± 6.1	15.0	5.0 \pm ± 2.0	7.2 \pm ± 2.1	9.6 \pm ± 2.1	6.2 \pm ± 1.6	9.0 \pm ± 1.8
XII	1.78 \div $\div 1.86$	0.477 \div $\div 0.520$	98.3 \pm ± 6.4	0	4.9 \pm ± 2.2	12.1 \pm ± 2.9	9.6 \pm ± 2.3	6.8 \pm ± 1.8	12.8 \pm ± 2.4
XIII	1.86 \div $\div 1.94$	0.520 \div $\div 0.563$	87.7 \pm ± 6.6	5.4	7.6 \pm ± 2.9	9.4 \pm ± 2.7	10.7 \pm ± 2.6	6.5 \pm ± 1.9	8.4 \pm ± 2.0
XIV	1.94 \div $\div 2.02$	0.563 \div $\div 0.605$	99.4 \pm ± 7.4	2.9	11.1 \pm ± 3.7	9.0 \pm ± 2.8	6.5 \pm ± 2.1	6.9 \pm ± 2.1	8.4 \pm ± 2.2
XV	2.02 \div $\div 2.10$	0.605 \div $\div 0.648$	92.3 \pm ± 7.6	7.1	7.5 \pm ± 3.4	5.6 \pm ± 2.5	11.6 \pm ± 3.2	7.7 \pm ± 2.4	6.3 \pm ± 2.1
XVI	2.10 \div $\div 2.18$	0.648 \div $\div 0.690$	80.1 \pm ± 7.5	3.0	1.6 \pm ± 1.6	8.3 \pm ± 3.1	7.6 \pm ± 2.7	10.7 \pm ± 3.0	3.7 \pm ± 1.7
XVII	2.18 \div $\div 2.26$	0.690 \div $\div 0.733$	67.7 \pm ± 7.2	3.1	10.3 \pm ± 4.2	5.1 \pm ± 2.5	5.2 \pm ± 2.3	9.8 \pm ± 3.0	3.2 \pm ± 1.6
XVIII	2.26 \div $\div 2.34$	0.733 \pm $\div 0.776$	56.3 \pm ± 6.7	3.1	5.2 \pm ± 3.0	10.3 \pm ± 3.6	6.2 \pm ± 2.5	2.7 \pm ± 1.6	4.9 \pm ± 2.0
XIX	2.34 \div $\div 2.54$	0.776 \div $\div 0.882$	43.0 \pm ± 3.9	2.7	3.1 \pm ± 1.6	3.5 \pm ± 1.4	5.6 \pm ± 1.6	2.8 \pm ± 1.1	5.1 \pm ± 1.4
XX	2.54 \div $\div 2.74$	0.882 \div $\div 0.988$	33.7 \pm ± 5.8	7.2	4.2 \pm ± 3.0	3.1 \pm ± 2.2	3.8 \pm ± 2.2	3.2 \pm ± 1.9	2.9 \pm ± 1.7

$\sigma/d\Omega^*$ ($\mu\text{b}/\text{sr}$).

θ^* (degrees)										
70 ÷ ÷80	80 ÷ ÷90	90 ÷ ÷100	100 ÷ ÷110	110 ÷ ÷120	120 ÷ ÷130	130 ÷ ÷140	140 ÷ ÷150	150 ÷ ÷160	160 ÷ ÷170	170 ÷ ÷180
9.8 ± 1.4	10.4 ± 1.4	11.9 ± 1.5	10.4 ± 1.4	10.0 ± 1.4	13.6 ± 1.8	13.6 ± 1.9	11.5 ± 1.9	11.6 ± 2.3	13.1 ± 3.1	15.2 ± 5.7
10.1 ± 1.5	14.6 ± 1.8	11.8 ± 1.6	11.7 ± 1.7	16.3 ± 2.0	15.2 ± 2.1	12.5 ± 2.0	17.4 ± 2.6	18.2 ± 3.1	11.4 ± 3.2	13.0 ± 5.8
13.6 ± 1.9	13.1 ± 1.8	13.7 ± 1.9	12.0 ± 1.8	18.7 ± 2.3	21.6 ± 2.6	17.8 ± 2.5	19.3 ± 2.9	17.6 ± 3.3	17.9 ± 4.2	11.8 ± 5.9
13.1 ± 2.0	19.6 ± 2.4	18.1 ± 2.3	18.6 ± 2.4	20.3 ± 2.5	16.5 ± 2.4	19.9 ± 2.8	12.5 ± 2.5	18.3 ± 3.5	22.2 ± 5.0	0
16.5 ± 2.3	13.8 ± 2.1	21.6 ± 2.6	20.4 ± 2.6	16.5 ± 2.4	19.4 ± 2.7	23.4 ± 3.2	21.8 ± 3.4	18.5 ± 3.7	13.3 ± 4.0	3.6 ± 3.6
14.6 ± 2.4	24.6 ± 3.0	25.7 ± 3.1	21.1 ± 2.9	19.7 ± 2.8	15.0 ± 2.6	21.5 ± 3.4	25.3 ± 4.0	20.2 ± 4.2	15.8 ± 4.8	8.5 ± 6.0
19.9 ± 2.9	20.9 ± 3.0	20.0 ± 2.9	18.1 ± 2.8	20.3 ± 3.1	21.9 ± 3.3	27.2 ± 4.0	15.3 ± 3.3	10.9 ± 3.3	22.6 ± 6.0	9.6 ± 6.8
14.2 ± 2.6	17.9 ± 2.9	15.6 ± 2.7	12.3 ± 2.4	17.7 ± 3.0	14.5 ± 2.8	20.0 ± 3.6	19.1 ± 3.9	17.3 ± 4.3	8.8 ± 4.0	5.2 ± 5.2
17.1 ± 2.1	13.5 ± 1.8	15.0 ± 1.9	15.2 ± 2.0	14.3 ± 2.0	15.8 ± 2.2	14.4 ± 2.3	15.0 ± 2.6	12.0 ± 2.7	11.8 ± 3.4	5.8 ± 4.1
12.0 ± 1.9	13.1 ± 2.0	8.7 ± 1.6	12.9 ± 2.0	10.9 ± 1.9	8.1 ± 1.7	7.4 ± 1.7	9.1 ± 2.1	9.6 ± 2.6	5.6 ± 2.5	6.6 ± 4.7
10.2 ± 1.9	8.5 ± 1.7	9.5 ± 1.8	8.4 ± 1.7	6.7 ± 1.6	8.3 ± 1.9	4.3 ± 1.4	8.3 ± 2.2	8.0 ± 2.5	7.9 ± 3.2	3.9 ± 3.9
9.5 ± 2.0	7.2 ± 1.7	8.0 ± 1.8	5.8 ± 1.5	8.8 ± 2.0	8.8 ± 2.1	7.9 ± 2.1	7.0 ± 2.2	4.7 ± 2.1	1.5 ± 1.5	4.6 ± 4.6
5.6 ± 1.6	10.3 ± 2.1	5.8 ± 1.6	7.9 ± 1.9	4.9 ± 1.6	5.5 ± 1.7	4.4 ± 1.7	3.9 ± 2.7	2.1 ± 1.5	3.5 ± 2.4	5.1 ± 5.1
8.5 ± 2.1	9.7 ± 2.2	10.2 ± 2.3	6.3 ± 1.8	5.6 ± 1.8	9.3 ± 2.4	7.2 ± 2.3	4.4 ± 2.0	4.8 ± 2.4	2.0 ± 2.0	5.9 ± 5.9
11.8 ± 2.8	9.5 ± 2.5	8.9 ± 2.4	7.2 ± 2.2	3.5 ± 1.6	8.5 ± 2.6	10.7 ± 3.1	3.3 ± 1.9	6.0 ± 3.0	4.9 ± 3.4	7.2 ± 7.2
1.2 ± 2.9	6.7 ± 2.1	6.7 ± 2.1	4.9 ± 1.8	5.2 ± 2.0	4.9 ± 2.0	10.4 ± 3.1	9.4 ± 3.3	1.6 ± 1.6	0	7.7 ± 7.7
7.5 ± 2.4	3.7 ± 1.6	5.9 ± 2.1	2.3 ± 1.3	3.2 ± 1.6	9.8 ± 3.0	4.1 ± 2.1	3.8 ± 2.2	6.9 ± 3.4	2.8 ± 2.8	8.4 ± 8.4
5.3 ± 2.0	2.2 ± 1.3	4.4 ± 1.8	6.1 ± 2.2	3.2 ± 1.6	1.8 ± 1.3	3.1 ± 1.8	3.8 ± 2.2	8.7 ± 3.9	2.8 ± 2.8	8.4 ± 8.4
3.8 ± 1.1	2.7 ± 0.9	2.0 ± 0.8	1.7 ± 0.8	1.5 ± 0.7	3.2 ± 1.1	3.3 ± 1.2	5.2 ± 1.7	5.5 ± 2.1	6.4 ± 2.9	7.6 ± 5.4
3.7 ± 1.8	0.9 ± 0.9	0.9 ± 0.9	0.9 ± 0.9	2.9 ± 1.7	2.2 ± 1.5	5.0 ± 2.5	3.1 ± 2.2	4.2 ± 3.0	3.4 ± 3.4	0

The events were fitted kinematically with three constraints. When the spectator proton was not visible ($p_s \leq 80$ MeV/c) in the chamber, the standard assumption was made: $p_x = p_y = p_z = 0$ MeV/c and $\Delta p_x = \Delta p_y = 30$ MeV/c, $\Delta p_z = 41$ MeV/c.

7469 events were found which satisfied all the imposed selection criteria ($\chi^2 \leq 10$, $p_s \leq 250$ MeV/c, fiducial volume, at least one visible proton besides the π^- , etc.).

The total and differential cross-sections are reported in Table I for 20 intervals of E^{*2} , the c.m. energy squared of the π^-p system.

These results are in substantial agreement with those of the ABBHMM collaboration obtained at DESY with the same technique^(*).

Indeed, if we calculate the photon energy with the same procedure of ref. (2), the total cross-sections coincide within errors. We remark, however, that this procedure, for the same c.m. total energy, gives a photon energy systematically greater than that for the reaction on free nucleon.

In the computation of the total and differential cross-sections the loss of events with two invisible protons, as well as the effects due to the possible wrong choice of the spectator proton, were corrected by means of a Monte Carlo calculation. The data at small angles, where the corrections are larger than 25% and the statistics is poor, are quoted without error evaluation and should be considered only as indicative.

No corrections have been made to take into account the effects due to the deuterium target, as the theoretical and experimental information on this subject is still rather poor.

In Fig. 1 the total cross-section as a function of E^{*2} is shown. The superimposed curves represent the total cross-sections for the reactions $\gamma + p \rightarrow \pi^+ + n$ and $\gamma + p \rightarrow \pi^0 + p$. The comparison is made for the same total center-of-mass energy, while the energy scale is also given, for convenience, in terms of E_γ , the corresponding incident-photon energy on free nucleon. The quoted errors are purely statistical; a possible systematic error of the order of 5% on the measured photon flux is not included.

Figure 2 a and b shows the differential cross-sections for various intervals of E^{*2} as a function of θ^* , the angle between the π^- and the incident γ -ray in the π^-p c.m. system.

In Fig. 3 the cross-section at $\theta^* = 90^\circ$ as a function of E^{*2} is reported. Its behaviour is very similar to that of the total cross-section. The values were obtained by fitting the experimental angular distributions with $\cos \theta^*$ polynomials. The available data from other experiments are also reported in the Figure.

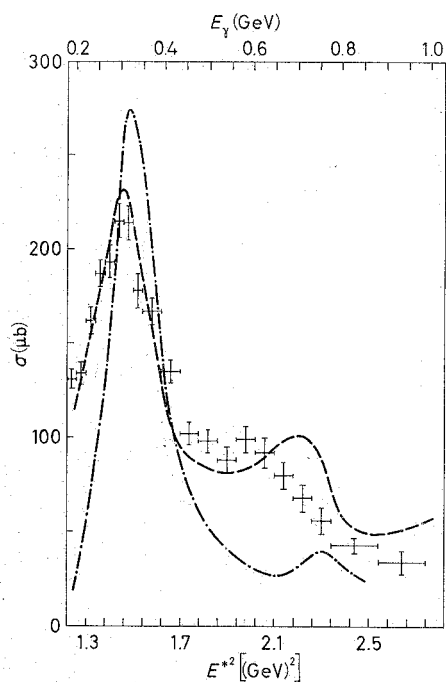


Fig. 1. - Total cross-section for the reaction $\gamma + n \rightarrow p + \pi^-$ as a function of E^{*2} (c.m. energy squared for the π^-p system): + present work, --- $\gamma + p \rightarrow n + \pi^+$, - - - $\gamma + p \rightarrow p + \pi^0$.

(*) AACHEN-BERLIN-BONN-HAMBURG-HEIDELBERG-MÜNCHEN COLLABORATION: *Nucl. Phys.*, 8 B, 535 (1968).

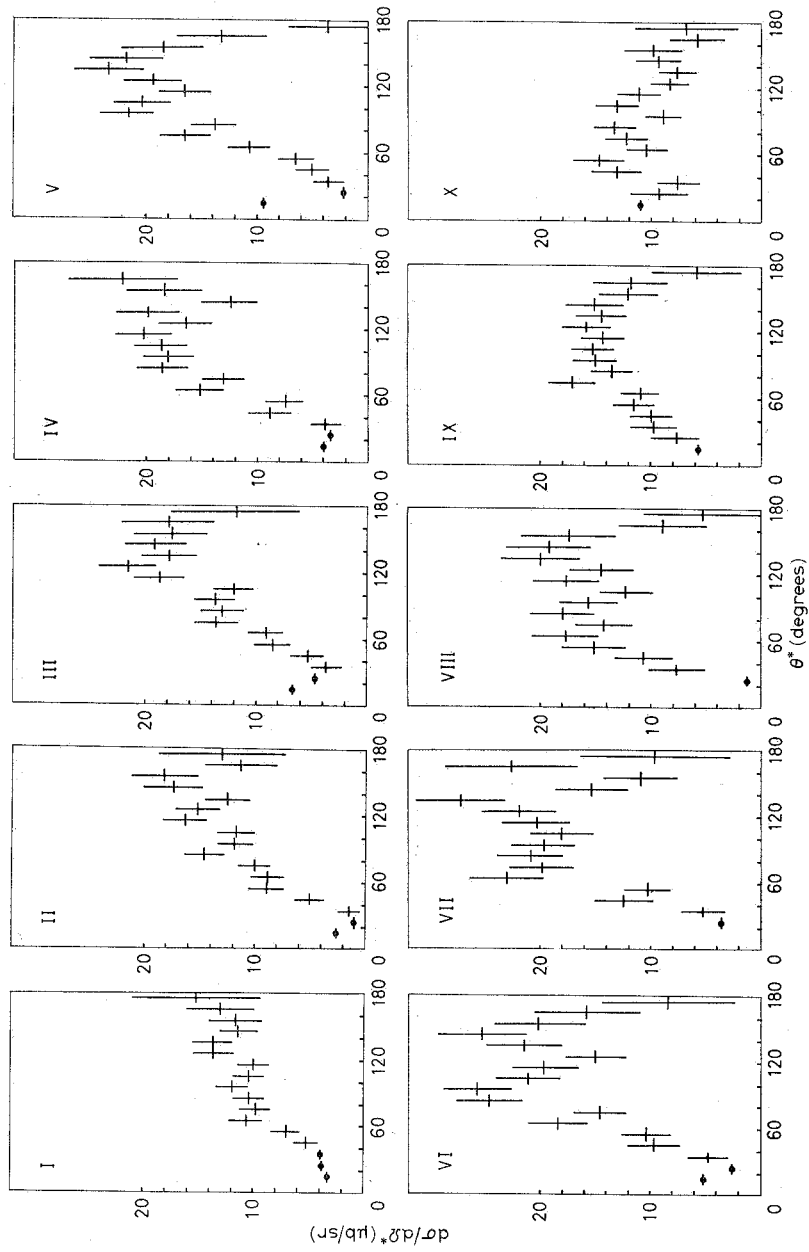


Fig. 2 a. - Differential cross-sections for the I to X intervals of E^{**} as a function of θ^* , the angle between the π^- and the incident photon in the π^-p c.m.s.

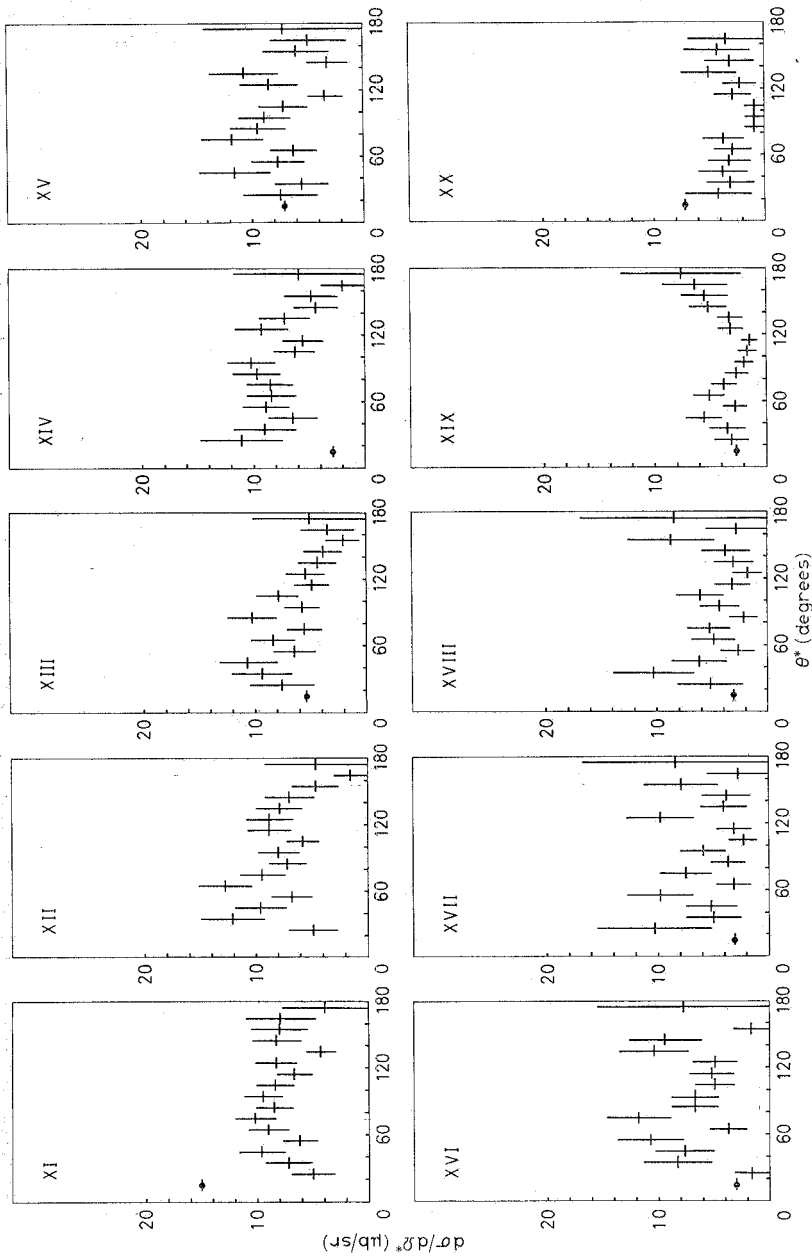


Fig. 2b. - Differential cross-sections for the XI to XX intervals of E^* as a function of θ^* , the angle between the π^- and the incident photon in the π^-p c.m.s.

At this point the controversial question of the existence of the P_{11} resonance in π^- photoproduction off neutrons has again to be considered. The experimental angular

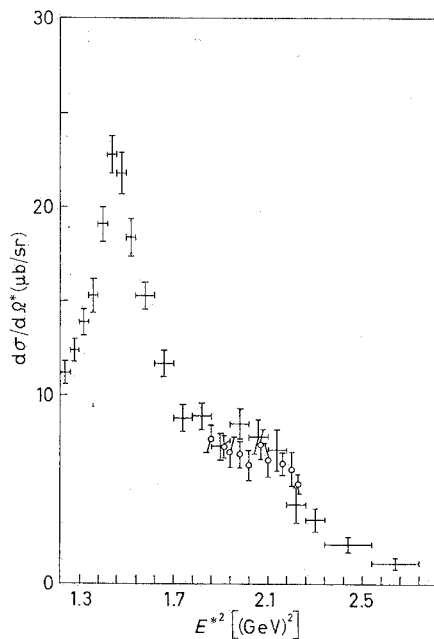


Fig. 3. - Differential cross-section at $\theta^* = 90^\circ$ as a function of E^{*2} : + present work, o BENEVENTANO *et al.* (ref. (*)).

distributions do not agree with those calculated by DONNACHIE⁽⁴⁾ assuming the presence of the P_{11} . Therefore, we tried to interpret the bump at about 1400 MeV in the total cross-section as due to the D_{13} resonance and to explain the energy shift by means of the interference between this resonance and the Born terms. The calculated effect however is not sufficient to explain the displacement found, so that the cross-section bump could be indicative of a relevant P_{11} contribution, in spite of the above-mentioned disagreement.

(*) M. BENEVENTANO, F. DE NOTARISTEFANI, P. MONACELLI, L. PAOLUZI, F. SEBASTIANI and M. SEVERI: *Lett. Nuovo Cimento*, **1**, 113 (1969).

(4) A. DONNACHIE: *Phys. Lett.*, **24 B**, 420 (1967).