

Laboratori Nazionali di Frascati

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Three-Pion Form Factor from Electron-Positron Experiments.

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A recent work ⁽¹⁾ has shown the possibility of a direct measurement of the photon-(n-pion) vertices through processes of the kind $e^+ + e^- \rightarrow n\pi$. Such processes will be studied with the storage rings presently under development in Frascati ⁽²⁾.

Particularly the reaction

$$(1) \quad e^+ + e^- \rightarrow \pi^+ + \pi^- + \pi^0,$$

can be used to obtain information on the vertex $\gamma - 3\pi$, which is important in the theory of the isoscalar nucleon structure. Following reference ⁽¹⁾ the differential cross-section in the C.M.S. for the process (1) is given by

$$(2) \quad \frac{d^2\sigma}{d\omega^{(+)} d\omega^{(-)}} = \frac{\alpha}{(2\pi)^2} \frac{|H|^2}{64 E^2} |(\mathbf{P}^{(+)} \times \mathbf{P}^{(-)})|^2 \sin^2 \theta d(\cos \theta),$$

where $\mathbf{P}^{(+)}$, $\omega^{(+)}$, $\mathbf{P}^{(-)}$ and $\omega^{(-)}$ are the momenta and the energies of the π^+ and π^- , E is the energy of the electrons, θ is the angle between the normal to the plane in which the three pions are produced and the line of collision, and H is the form factor of the vertex $\gamma \rightarrow \pi^+ + \pi^- + \pi^0$. Recently BLANKENBECLER and TARSKI ⁽³⁾ have calculated this vertex essentially in terms of the $T=1$, $J=1$ 2-pion resonance.

From their work we obtain

$$(3) \quad |H|^2 = (2\pi)^6 |F_0|^2 |\exp [\Delta_{12} + \Delta_{23} + \Delta_{31}]|^2 D^{-2}(t)|_{t=4E^2},$$

where F_0 is a constant that must be determined experimentally; $|\exp [\Delta_{ij}]|^2$ are

⁽¹⁾ N. CABIBBO and R. GATTO: *Phys. Rev. Lett.*, **4**, 313 (1960).

⁽²⁾ See the reports by B. TOUSCHEK and by R. GATTO in the *Proceedings of the International Conference on the Theoretical Aspects of very high energy phenomena* CERN, June 1961 (CERN Report 61-22, pag. 75), and R. GATTO, *Proceedings of the Conference on the elementary particles* Aix-en-Provence, 1961 (to be published).

⁽³⁾ R. BLANKENBECLER and J. TARSKI: *Phys. Rev.* (to be published).

the pion-pion phase functions; t is the square of the four momentum transfer; and $D(t)$ is the function that sums the connected three pions diagrams (see ⁽³⁾).

For the total cross-section we have obtained the expression

$$(4) \quad \sigma(E) = \frac{\alpha}{3} \frac{(\vec{2}\pi)^7}{(4E)^4} |F_0|^2 D(4E^2) - 2 \int_{4E^2}^{(2E-\mu)^2} dx \int_{y_1}^{y_2} dy \cdot \\ \cdot |\exp [A(x) + A(y) + A(4E^2 - x - y + 3\mu^2)]^2 | \mathbf{P}^{(+)} \times \mathbf{P}^{(-)} |^2 .$$

Where μ is the pion mass, and

$$(5) \quad \omega^{(+)} = \frac{1}{4E} (4E^2 + \mu^2 - y) ,$$

$$(6) \quad \omega^{(-)} = \frac{1}{4E} (4E^2 + \mu^2 - x) ,$$

$$(7) \quad \cos (\widehat{\mathbf{P}^{(+)} \mathbf{P}^{(-)}}) = \frac{1}{2 |\mathbf{P}^{(+)}| |\mathbf{P}^{(-)}|} \cdot \\ \cdot [x - 2\omega^{(+)}(2E - \omega^{(-)})] ,$$

$$(8) \quad y_{1,2} = 2E\omega \mathbf{P}^{(-)} \mp |\mathbf{P}^{(-)}| \sqrt{\frac{x-4}{x}} .$$

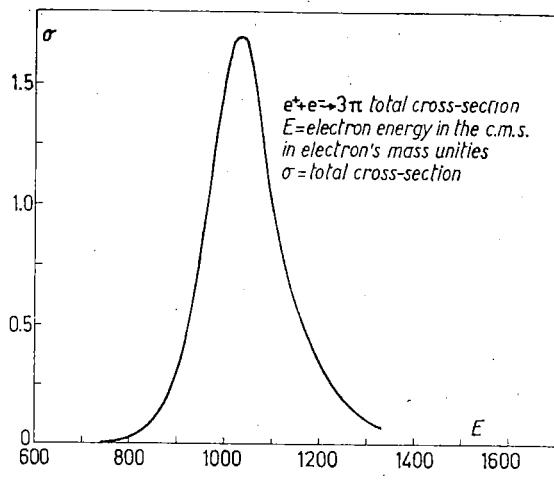


Fig. 1.

We have integrated (4) for the value $t_R = 20\mu^2$ of the $\pi\pi$ resonance, in the approximation $D(t) = 1$.

The results are reported in Fig. 1, that gives $\sigma(E)$ in arbitrary units.

However recently experimental evidence ⁽⁴⁾ has been found for a narrow three pion resonance around $\sqrt{t} = 787$ MeV, with a half-width less than 15 MeV. The rescattering function $D^{-1}(t)$ should therefore show a strong resonance in this zone, and the cross-section $\sigma(E)$ should have a peak around $E = 395$ MeV.

Our results give another peak at $E = 525$ MeV. It seems difficult at present to evaluate the relative importance of the two resonances.

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⁽⁴⁾ B. C. MAGLIC, L. W. ALVAREZ, A. H. ROSENFIELD and M. L. STEVENSON: *Phys. Rev. Lett.*, 7, 178 (1961).