

INFN/AE-73/10  
18 Dicembre 1973

M. Brini-Penzo and A. Pullia: FOCUSING DEVICES FOR A NARROW-BAND NEUTRINO BEAM AT HIGH ENERGY. -

## INTRODUCTION -

In the new range of high energies now becoming available at the large accelerators (NAL and future SPS), a narrow-band neutrino beam can be conceived as a facility exploiting the forward peaked particle production and involving important advantages from the experimental point of view. There would be substantial simplifications in the evaluation of the  $\nu$ -spectrum; furthermore the knowledge of the incoming neutrino energy (within errors of the order of  $\pm 10\%$ ) would be a valuable information to complement the measurement of the "visible" energy of the event.

A narrow-band neutrino beam can be in principle obtained by the decay of a monochromatic beam of neutrino parents. Several designs based on quadrupole systems have been proposed<sup>(1-4)</sup> and realized<sup>(5)</sup>. We studied a different system<sup>(6, 7, 8)</sup>, based on focusing elements similar to the horns<sup>(9)</sup> used up to now in the neutrino experiments at CERN.

## 1. - DESCRIPTION OF THE SYSTEM -

A system realized with one horn only was the object of a first study<sup>(7, 8)</sup> on a device for focusing parallel to the beam axis only the parents with the right momentum: the others, coming out of

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the horn not parallel to the axis, were eliminated by means of a tunnel with a small radius or, realistically, by a system of slits. In an optical comparison the focusing system is represented by a point source (target) of  $\pi$  and K placed in the focus of a lens and by a collimator (tunnel), (see fig. 1 a)).

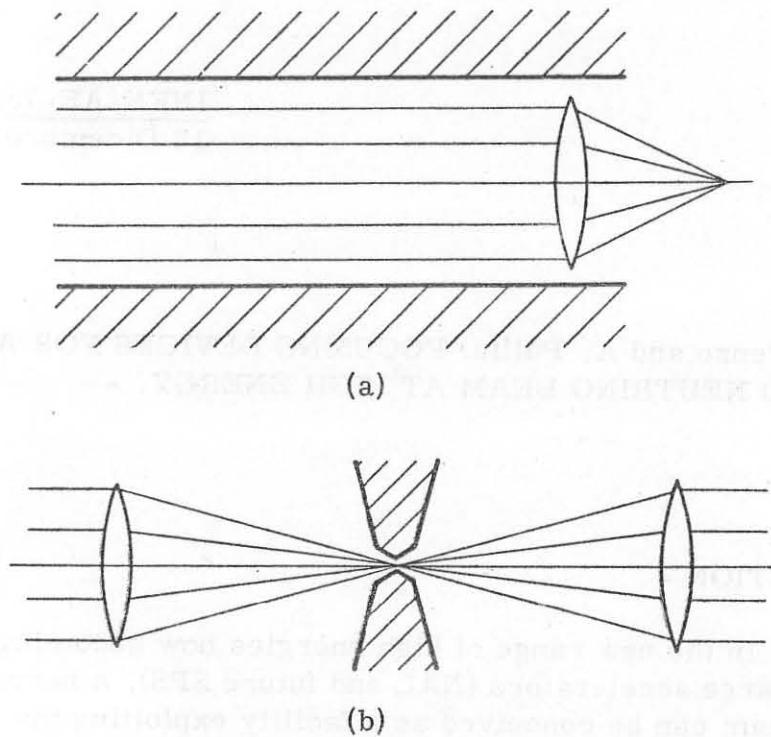


FIG. 1 - a) Optical representation of the focusing system with one horn; b) Optical representation of the focusing system with two horns.

The results were encouraging, because the resulting spectrum has a characteristic dichromatic shape, with two separated peaks corresponding to neutrinos  $\nu_{\pi}$  (from  $\pi \rightarrow \mu_{\nu}$ ) and  $\nu_K$  (from  $K \rightarrow \mu_{\nu}$ ) with the maximum momentum allowed by the kinematics ( $0.43 p_{\pi}$  and  $0.95 p_K$  respectively).

The  $\nu_K$  peak, which is the most interesting one, is however, only a factor 6 above the background. Therefore, as it is unlikely to obtain substantial improvements with one horn only, we use in the system studied now<sup>(6)</sup> two horns, designed in the following way: the first one focuses into a slit only the parents of a given momentum, which are then aligned in a beam parallel to the axis by a second horn. All the other parents, of a different momentum, do not cross the slit. In the optical comparison the focusing system is represented by a first lens (first horn) giving the image of the source of  $\pi$  and K (target) in a definite point (slit) which constituted the focus of a second

lens (second horn) (see fig. 1 b)).

## 2. - DESCRIPTION OF THE LAYOUT. -

As an example of the proposed method<sup>(6)</sup> we describe in detail the neutrino narrow-band beam which can be set up at CERN SPS with a primary beam momentum of 200 GeV/c, selecting  $\pi$ 's and K's of 100 GeV/c momentum. A schematic layout of the beam is shown in fig. 2 a); the longitudinal sections of the two horns, indicated as  $R_1$  and  $R_2$ , and the slit are shown in detail in figs. 2 b) to d).

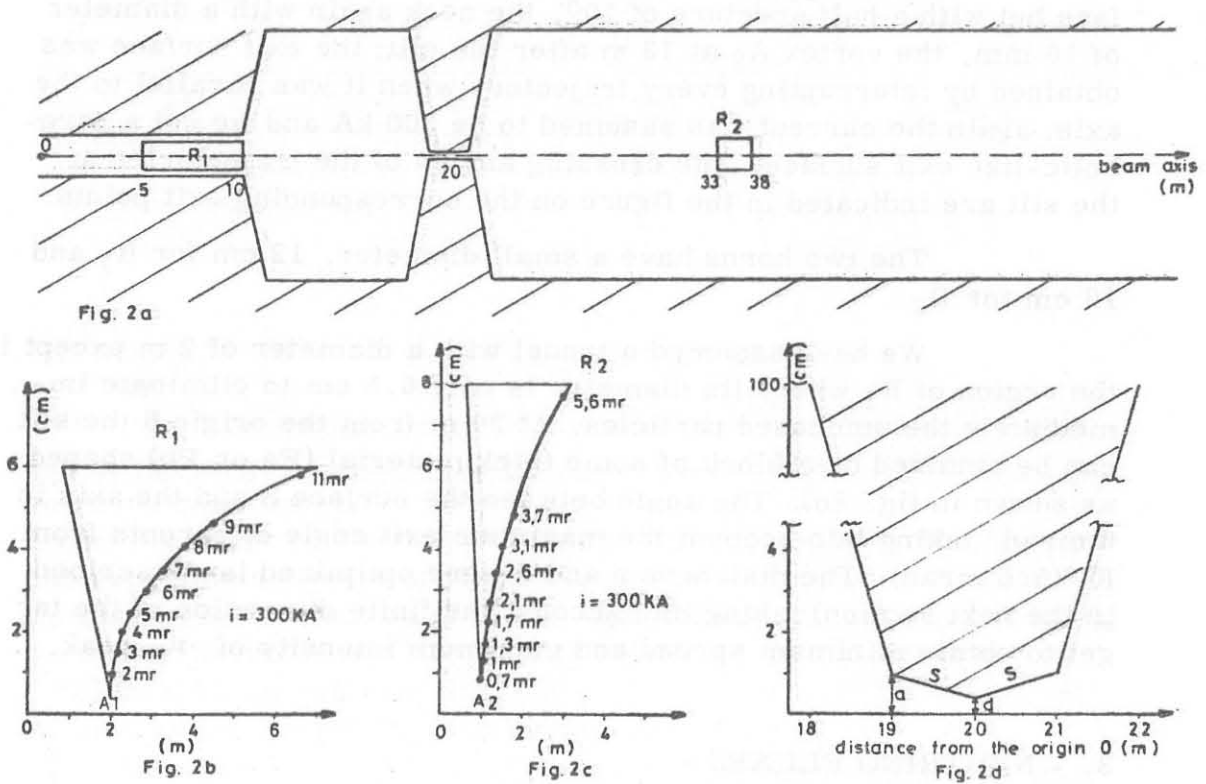


FIG. 2 - a) Schematic layout of the focusing section of the neutrino beam; b) Longitudinal section of  $R_1$ ; c) Longitudinal section of  $R_2$ ; d) Longitudinal section of the slit.

The procedure used to design  $R_1$  was the following: the vertex  $A_1$  was fixed at 5 m from the origin 0 (assumed point-like) of the parents trajectories of 100 GeV/c momentum, the neck had a diameter of 16 mm, the entry surface was chosen conical with a half aperture of  $2.5^\circ$ ; then the exit surface was defined by the points obtained interrupting every trajectory when its direction pointed to a focus at 20 m from the origin 0; the current of the horn was assumed to be 300 kA. The resulting exit surface is parabolic-like. In the figure the

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exit points corresponding to the trajectories leaving from the origin at different angles starting from 2 mrad (the minimum angle focused) are shown. This value of the minimum angle is due to the necessity of keeping the neck of the horn wide enough in diameter ( $\geq 16$  mm) to stand the mechanical and thermic stress<sup>(10)</sup> and the requirement of limiting this focusing section of the beam to a reasonable length. Considering the K production spectra<sup>(11)</sup> at 100 GeV/c for a primary momentum of 200 GeV/c, one can see that the fraction of K flux corresponding to the angles less than 2 mrad, which is lost, is less than 10% of the total.

Then  $R_2$  was designed assuming again a conical entry surface but with a half aperture of  $10^\circ$ , the neck again with a diameter of 16 mm, the vertex  $A_2$  at 13 m after the slit; the exit surface was obtained by interrupting every trajectory when it was parallel to the axis; again the current was assumed to be 300 kA and we get a parabolic-like exit surface. The crossing angles of the trajectories at the slit are indicated in the figure on the corresponding exit points.

The two horns have a small diameter, 12 cm for  $R_1$  and 16 cm for  $R_2$ .

We have assumed a tunnel with a diameter of 2 m except in the region of  $R_1$  where its diameter is only 6.5 cm to eliminate immediately the unfocused particles. At 20 m from the origin 0 the slit can be obtained by a block of some thick material (Fe or Pb) shaped as shown in fig. 2d). The angle between the surface S and the axis is 6 mrad, taking into account the maximum exit angle of parents from  $R_1$  (5.6 mrad). The distances a and d were optimized (as described in the next section) taking into account the finite dimension of the target to obtain minimum spread and maximum intensity of  $\nu_K$  peak.

### 3. - NEUTRINO FLUXES -

In fig. 3 the general layout of the neutrino beam is shown. We have calculated the  $\nu$ -fluxes for this layout by a Monte-Carlo method, with the following data:

- the  $\pi^+$  and  $K^+$  production spectra have been calculated by the Hagerdon thermo-dynamical model<sup>(11)</sup>
- the decay tunnel is 510 m long
- the distance target-detector T-D is 830 m (corresponding about to the position of BEBC as detector) or 920 m (corresponding to Gargamelle)
- the fluxes are averaged over a circular surface of 0.6 m radius
- the  $\pi$  and K absorptions in the walls of the horns are taken into account assuming a thickness of 3 mm

- the target is made of a thick material ( $W$ ) and short (one interaction length) to approximate the ideal conditions of a point-source; the interaction length in  $W$  has been assumed  $\lambda_p = 13$  cm for protons and  $\lambda_{\pi, K} = 17$  cm for pions and kaons; the target is in such a position that the origin 0 of the beam axis coincides with the average interaction point; a beam stopper three interaction lengths long, just before  $R_1$ , with a radius  $r = 8.5$  mm, absorbs the protons which have not interacted and the parents emitted at angles less than 2 mrad.

In fig. 4 we show the  $\nu$ -spectra in Gargamelle calculated with the above data and for different dimensions of the slit, that is for different values of the parameters  $a$  and  $d$  (see Fig. 2 d).

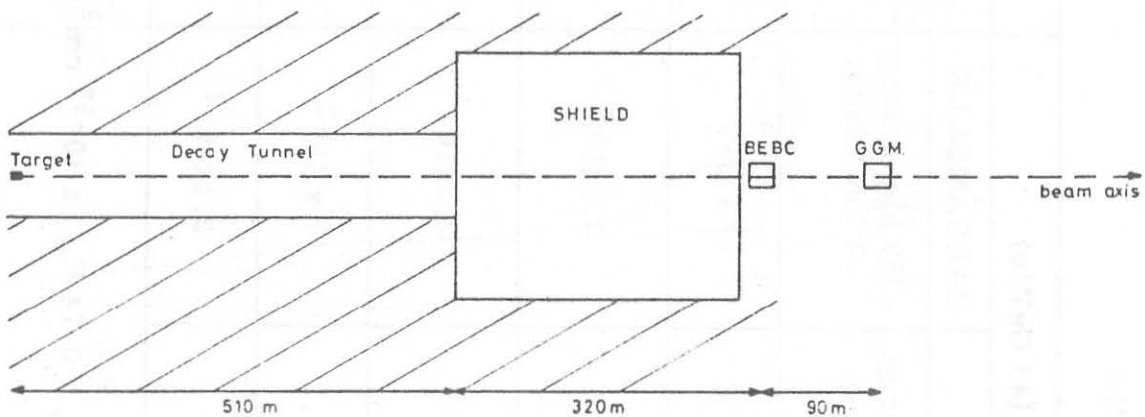


FIG. 3 - General layout of the neutrino beam.

It can be seen that the best dimensions for the slit are  $a = 1$  cm and  $d = 4$  mm (with larger dimensions the peak becomes too large).

In fig. 4 also the contribution of the beam stopper as a meson source (hence a neutrino one), is shown for the configuration with  $d = 4$  mm: it is negligible, mainly in the  $\nu_K$  energy region.

Fig. 5 shows the energy distributions of parents  $\pi$  and  $K$  after the slit with  $d = 4$  mm: the momentum spread of the parents (F. W.H.H.) is  $\sim 10$  GeV/c.

In fig. 6 the  $\nu$ -spectra in Gargamelle and in BEBC are shown together and compared with the ideal spectra: it can be seen that the  $\nu_K$  peaks are about a factor 3.5 under the ideal fluxes.

In table I the principal parameters of these spectra are reported.

We have also evaluated the expected events rates at 300 and 400 GeV/c primary momentum by multiplying the calculated ra-

TABLE I

	$\nu_\pi$ peak (43 GeV/c)		$\nu_K$ peak (95 GeV/c)	
	BEBC	GARGAMELLE	BEBC	GARGAMELLE
Ratio $I_{\text{peak}}/I_{\text{background}}$ : $\left\{ \begin{array}{l} 200 \text{ GeV/c} \\ \text{primary} \\ \text{momentum} \end{array} \right.$	$\sim 250$ high energy $\sim 7$ low energy	$\sim 250$ high energy $\sim 7$ low energy	$\sim 100$ high energy 25 low energy	$\sim 100$ high energy 22 low energy
F. W. H. H. : $\left\{ \begin{array}{l} 200 \text{ GeV/c} \\ \text{primary mo} \\ \text{mentum} \end{array} \right.$	$\sim 15$ GeV	$\sim 15$ GeV	$\sim 12$ GeV	$\sim 12$ GeV
Neutrinos in the detector/ $10^{13}$ protons/ $\text{m}^2$ : $\left\{ \begin{array}{l} 200 \text{ GeV/c} \\ \text{primary} \\ \text{momentum} \end{array} \right.$	$2.8 \times 10^8$	$2.3 \times 10^8$	$2.4 \times 10^7$	$1.5 \times 10^7$
Event/rate/ton. $10^{13}$ protons <sup>(1)</sup> : $\left\{ \begin{array}{l} 200 \text{ GeV/c} \\ \text{primary} \\ \text{momentum} \end{array} \right.$	$3.8 \times 10^{-3}$	$3.3 \times 10^{-3}$	$8.5 \times 10^{-4}$	$5.3 \times 10^{-4}$
Extrapolated event rate for 300 GeV/c primary momentum	$1.1 \times 10^{-2}$	$1 \times 10^{-2}$	$5 \times 10^{-3}$	$3.1 \times 10^{-3}$
Extrapolated event rate for 400 GeV/c primary momentum	$2.5 \times 10^{-2}$	$2.1 \times 10^{-2}$	$7 \times 10^{-3}$	$4.4 \times 10^{-3}$

(1) - The event rates are calculated assuming  $\sigma_\nu = 0.74 E_\nu \times 10^{-38} \text{ cm}^2/\text{nucleon}$ .

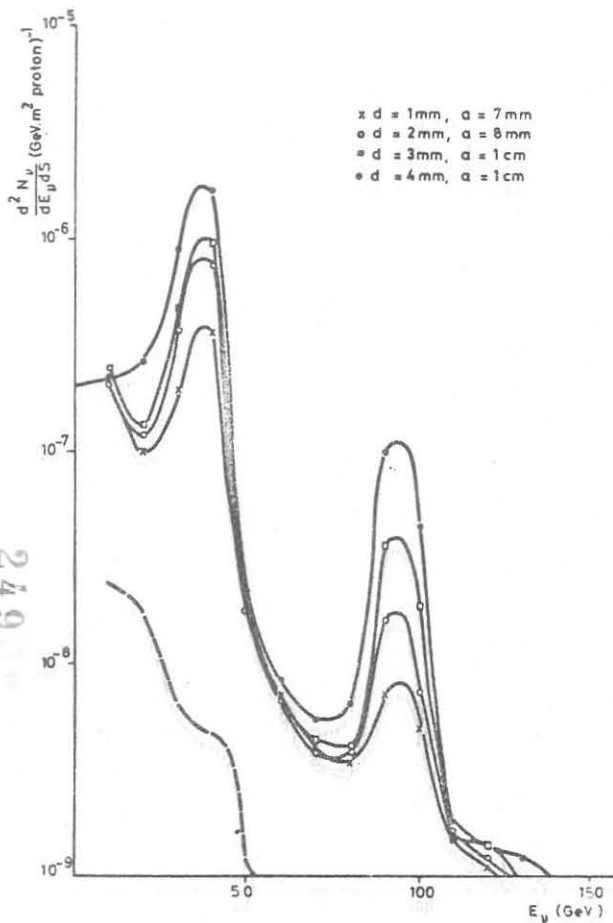


FIG. 4 - Neutrino narrow-band spectra in Gargamelle for different sizes of the slit and contribution of the beam stopper to the neutrino spectrum.

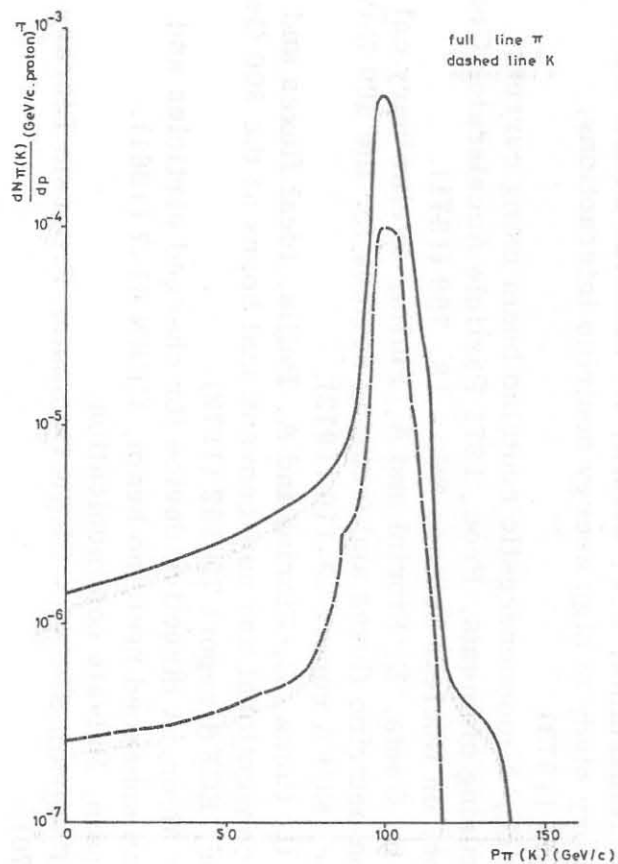


FIG. 5 - Energy distribution of the neutrino parents after the slit.

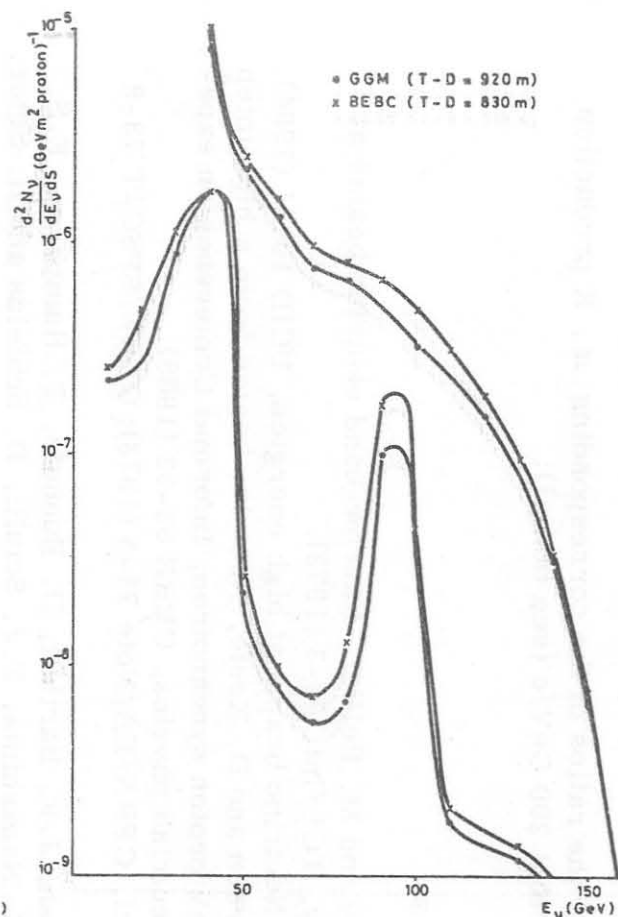


FIG. 6 - Neutrino narrow-band and ideal spectra in Gargamelle and in BEBC for the best slit size.

te at 200 GeV/c by the ratios of the corresponding  $\pi$ , K production spectra to the ones at 200 GeV/c (see table I).

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