

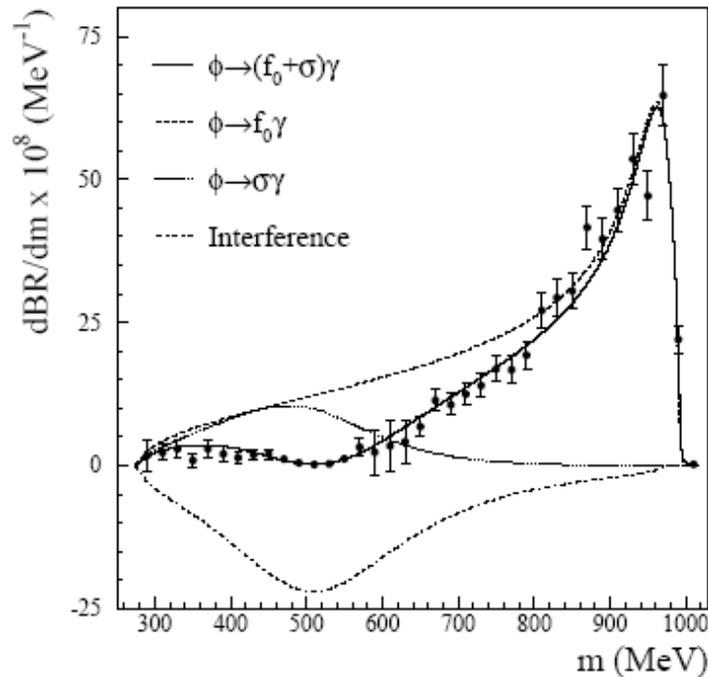
Diquark-antidiquark Mesons : a new spectroscopy?
*Report on work done with F. Piccinini, A. Polosa, V.
Riquer*

Luciano MAIANI,
Univ. di Roma1. Italia; INFN, Roma. Italia

Tetraquark Mesons

- An idea which got momentum in the 70s (R. Jaffe, H. Lipkin...);
- QCD encourages the speculation that such states are indeed possible;
- As we shall see, the light scalar mesons $a(980)$, $f(980)$ really look like
- $[qq]_{col=\bar{3}}[\bar{q}\bar{q}]_{col=3}$
- In alternative: $a(980)$, $f(980)$ could be K-Kbar “molecules”, bound by one- π exchange, i.e. in the configuration:
- $(q\bar{q})_{col=1}(q\bar{q})_{col=1}$
- the existence of lighter partners, σ and κ is crucial;
- Recent expts, at FNAL(E791), Frascati (KLOE) and BES, have seen again σ and κ : this would be against “molecules”;
- ...and there are states with hidden/open charm that do not look like charmonium states: X(3872), X(3940), (Belle, Babar in B decays), X(2632) (SELEX).

$\phi(1020) \rightarrow \pi^0 \pi^0 \gamma$



Study of the Decay $\phi(1020) \rightarrow \pi^0 \pi^0 \gamma$ with the KLOE Detector
 The KLOE Collaboration
arXiv:hep-ex/0204013 Apr 2002

Fit results using f_0 only, Fit (A), and including the σ , Fit (B).

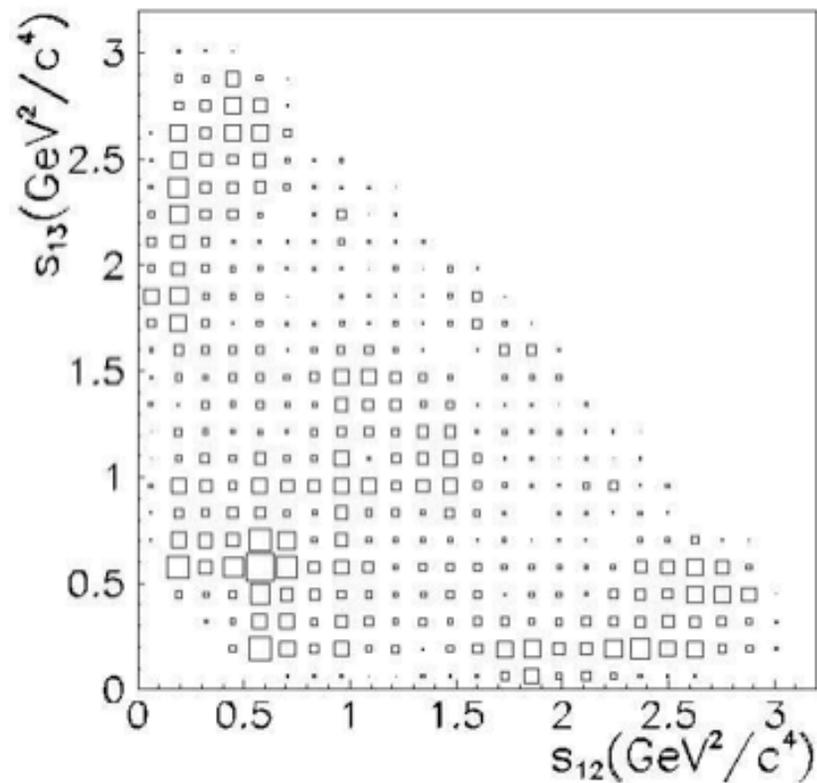
	Fit (A)	Fit (B)
χ^2/ndf	109.53/34	43.15/33
M_{f_0} (MeV)	962 ± 4	973 ± 1
$g_{f_0 K^+ K^-}^2 / (4\pi)$ (GeV^2)	1.29 ± 0.14	2.79 ± 0.12
$g_{f_0 K^+ K^-}^2 / g_{f_0 \pi^+ \pi^-}^2$	3.22 ± 0.29	4.00 ± 0.14
$g_{\phi \sigma \gamma}$	—	0.060 ± 0.008

Sigma parameters from E791

$$M_{\sigma} = (478_{-23}^{+24} \pm 17) \text{ MeV}$$

$$\Gamma_{\sigma} = (324_{-40}^{+42} \pm 21) \text{ MeV}$$

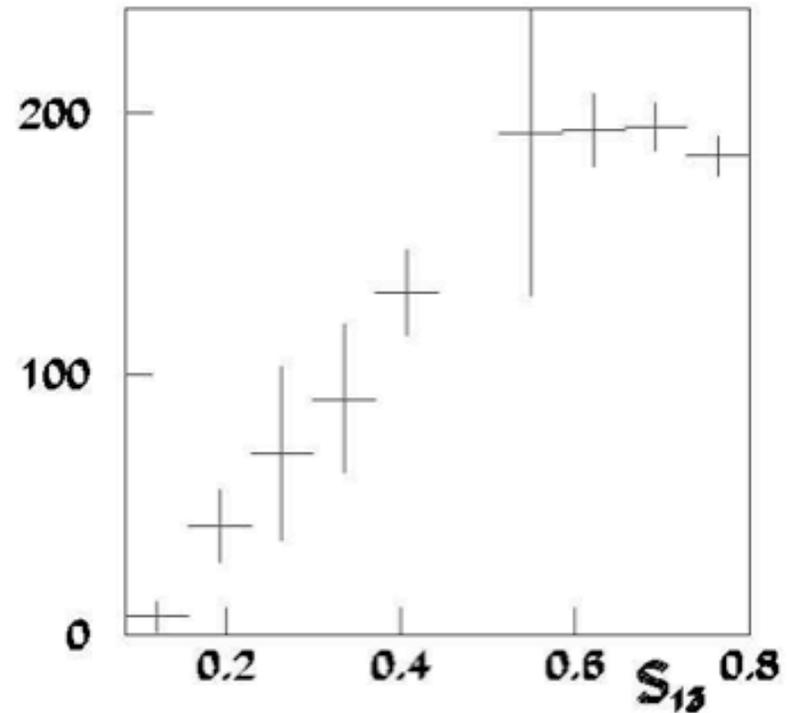
$\sigma(600)$ at E791



PRL 86, 770 (2001)

$$M = 478 \pm 23 \pm 17 \text{ MeV}$$

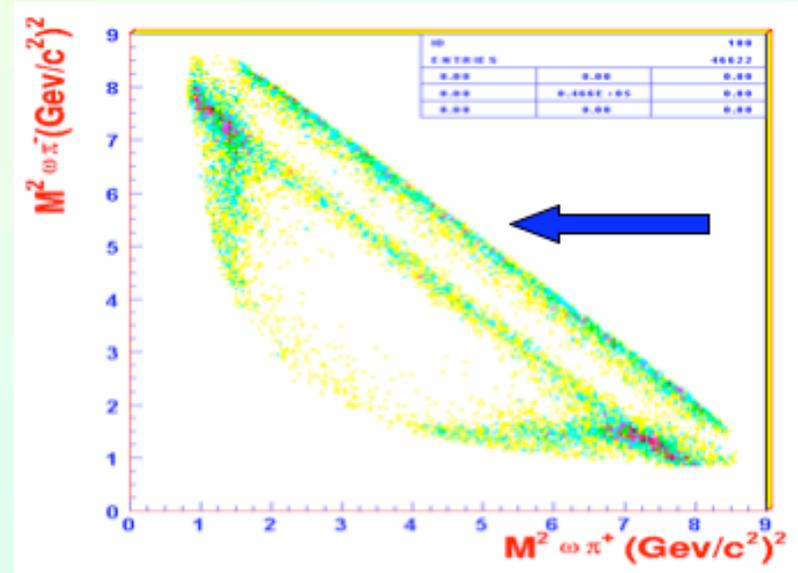
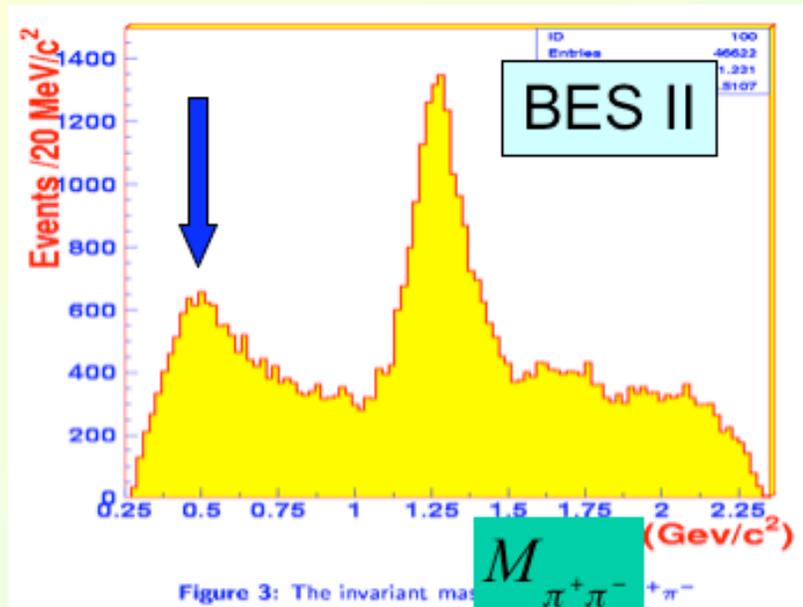
$$\Gamma = 324 \pm 40 \pm 21 \text{ MeV}$$



hep-ex/0307008 (2003)

$\sigma(600)$ @ BES II

- BES II: σ in $J/\psi \rightarrow \omega\pi^+\pi^-$.

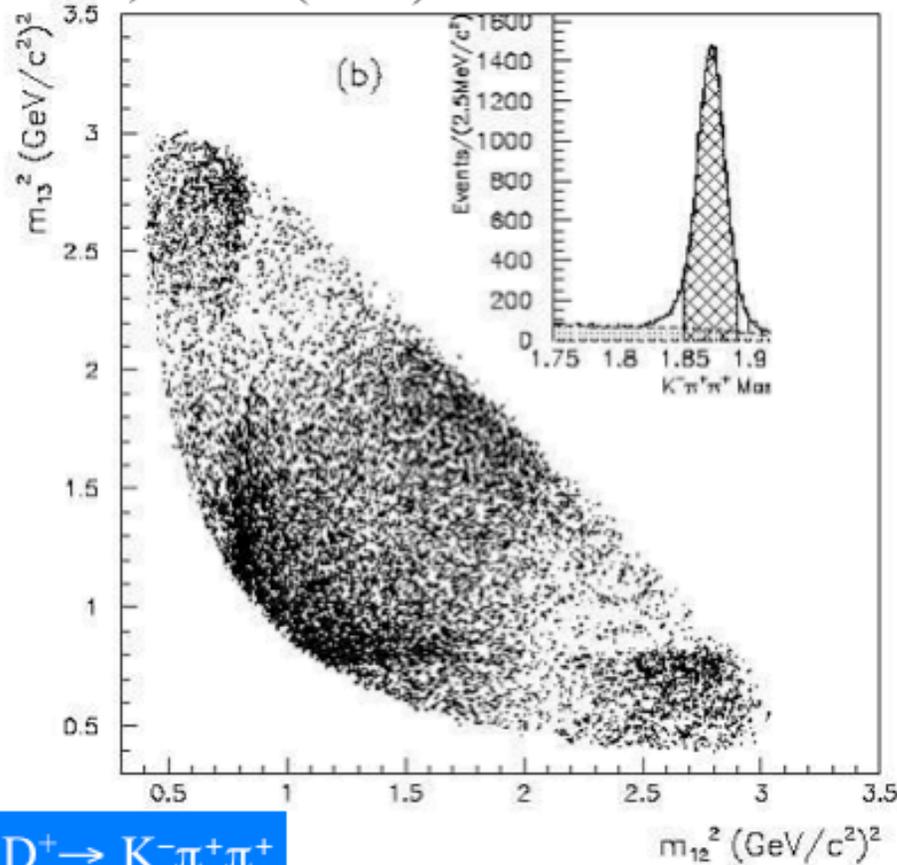


- Partial wave analysis: pole position

$$(541 \pm 39) - i(252 \pm 42) \text{ MeV}$$

$\kappa(800)$ @ E791

PRL 89, 12801 (2002)



$D^+ \rightarrow K^- \pi^+ \pi^+$

$M = 797 \pm 19 \pm 43 \text{ MeV}$

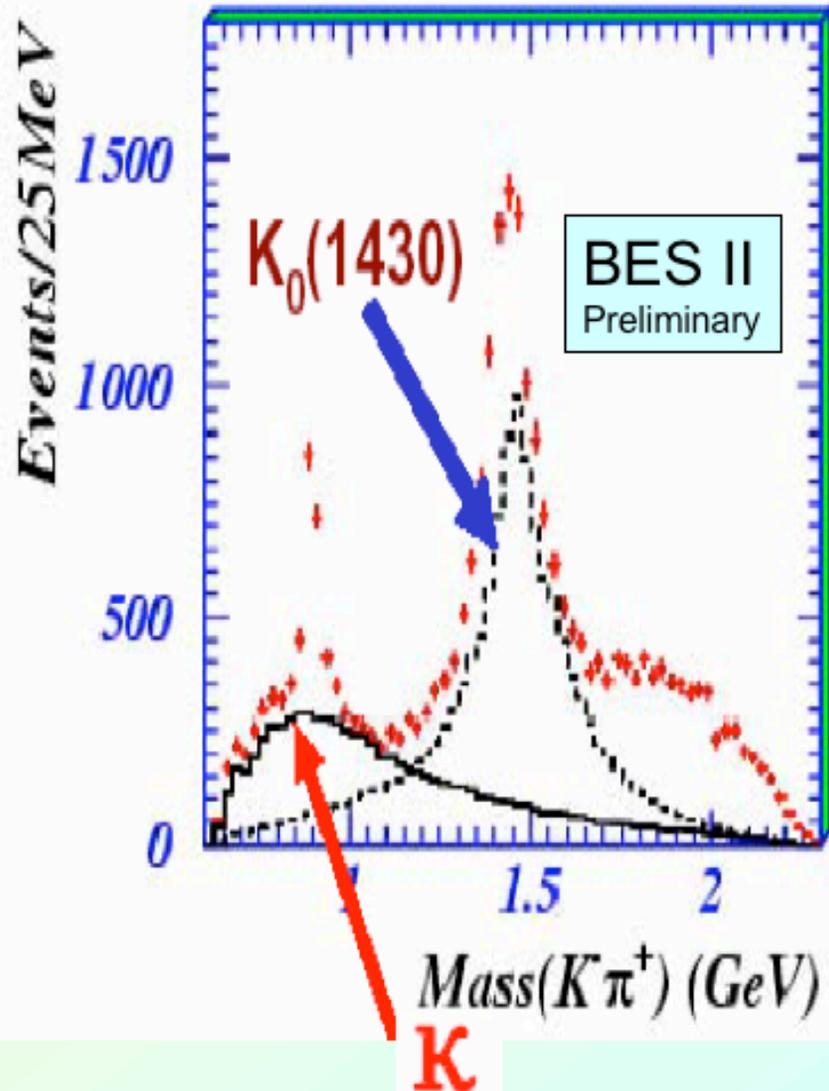
$\Gamma = 410 \pm 43 \pm 87 \text{ MeV}$

Non Res. Bkg. :90% (no k) \rightarrow 13% (k)

No k needed in Dalitz plot fit of
 $D^0 \rightarrow K^- \pi^+ \pi^0$ e $D^0 \rightarrow K_S \pi^+ \pi^-$
(CLEO)

No k needed in Dalitz plot fit of
 $D^0 \rightarrow K^0 K^- \pi^+$ e $D^0 \rightarrow K K^+ \pi^-$
(BABAR)

$\kappa(800)$ @ BES II



- BES II: κ in $J/\psi \rightarrow K^* K \pi \rightarrow K \pi K \pi$.

$$(760 \sim 840) - i(310 \sim 420) \text{ MeV}$$

The present work (1)

Recent evidence for σ at low energy led us to reconsider the case of sub-GeV scalar mesons. Many previous investigations (Joffe, Close&Tornqvist, Schechter and coll...).

We propose:

all scalars below 1 GeV are diquark-antidiquark bound states (1 nonet), the q - q bar scalar nonet ($L=1, S=1, J=0$) has to be above.

Results:

- Low energy states show inverted mass spectrum, consistent with “perfect mixing”;
- Strong decays are reasonably accounted for;
- Decays $\phi \rightarrow S \gamma$ still to be studied;
- Relations with earlier proposal by Rossi&Veneziano suggests connection to baryon-antibaryon, rather than meson-meson states (or molecule)

PRL 93, 212002 (2004), hep-ph/0407017

The present work (2)

- Heavy quark interactions are spin independent: new spin states?
 - We propose that X(3872) observed by Belle and by Babar *is a diquark-antidiquark bound state and estimate the spectrum of states of the spin multiplet with the same flavors*.
 - $X(3872) = (J=1^{++}) = (cq)_{col=\bar{3},S=1} (\bar{c}\bar{q})_{col=3,S=1}$
 - with the same parameters, we can accommodate the X(2632) observed by SELEX:
 - $X(2632) = (J=2^{++}) = (cq)_{col=\bar{3},S=1} (\bar{s}\bar{q})_{col=3,S=1}$
- *we predict X(3872) made by two states with $\Delta m = (5-8) \text{ MeV} \approx 2 (m_d - m_u)$*
- *one state only in the decay: $B^+ \rightarrow K^+ X(3872)$, the other appears in $B^0 \rightarrow K_S X(3872)$*
- *bounds to the production of the charged partner X^+ are close but not in contradiction with BaBar.*

PR D70 054009 (2004), hep-ph/0407028;
PR D to appear, hep-ph/0412098

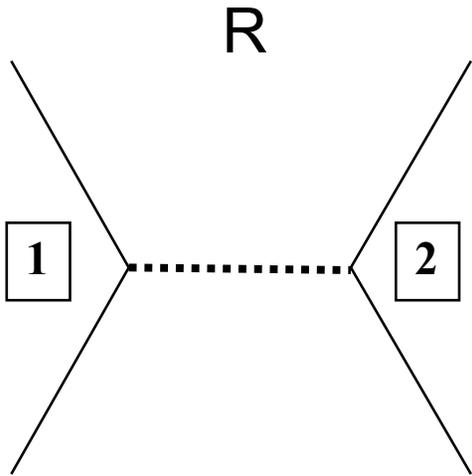
Summary

-
- Attractive and repulsive channels in QCD
- String structures: the “baryonium” model (Rossi & Veneziano, 1977)
- The light scalar mesons;
- Two-meson decays;
- Surprising charmonium states seen by Belle, Babar and Selex;
- S-wave Tetraquarks, the X(3872) and spectrum of related states;
- Selex particle, X(2632), and associated spectrum;
- Alignment to quark masses, isospin breaking;
- Conclusions

Attractive & repulsive channels in QCD

Interaction of two colored objects:

$$\propto g^2 \langle \vec{T}_1 \cdot \vec{T}_2 \rangle_R = \frac{g^2}{2} [\langle (\vec{T}_1 + \vec{T}_2)^2 \rangle_R - \langle \vec{T}_1^2 \rangle - \langle \vec{T}_2^2 \rangle] =$$



$$= \frac{g^2}{2} [C^{(2)}(R) - C^{(2)}(1) - C^{(2)}(1)]$$

$$\bar{q}q = \begin{cases} \text{octet} = +1/3 & \text{repulsion} \\ \text{singlet} = -8/3 & \text{attraction} \end{cases}$$

$$qq = \begin{cases} \text{"3bar"} = -4/3 & \text{attraction} \\ \text{"6"} = +2/3 & \text{repulsion} \end{cases}$$

Spin-spin interaction

$$\propto -g^2 \langle \sigma_1 \vec{T}_1 \cdot \sigma_2 \vec{T}_2 \rangle_R \propto -g^2 [C^{(2)}_{\text{eff}}(R)] [J(J+1) - 3/2]$$

$$\langle \bar{q}q \rangle_1 \text{ and } \langle qq \rangle_{\bar{3}} = \begin{cases} \text{spin } 1 = +1/2 & \text{repulsion} \\ \text{spin } 0 = -3/2 & \text{attraction} \end{cases}$$

Baryons in the octet:

$$\Lambda = ([ud]_{J=0} s); \Sigma^0 = (\{ud\}_{J=1} s) \rightarrow \Lambda \text{ is lighter than } \Sigma$$

With **antisymmetry** in **color** and **spin** and a common spatial configuration, Fermi statistics $\Rightarrow \bar{3}_f$

Good diquarks: $[qq]_{\bar{3}_c, 1_s, \bar{3}_f}$

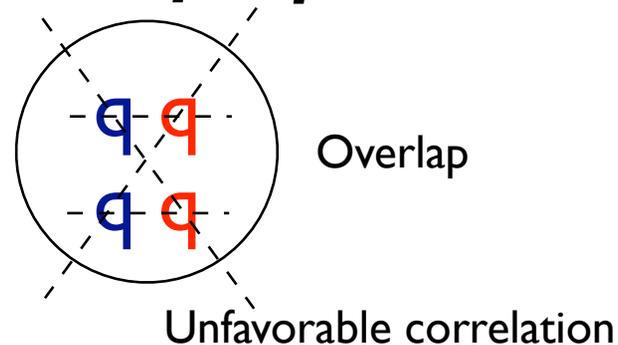
Bad diquarks: $(qq)_{\bar{3}_c, 3_s, 6_f}$

Since spin interaction is a relativistic effect we might expect it stronger for the lightest quarks....

$$\text{Splitting : } (ud) - [ud] > (us) - [us] > \underline{(uc) - [uc]} \approx 0$$

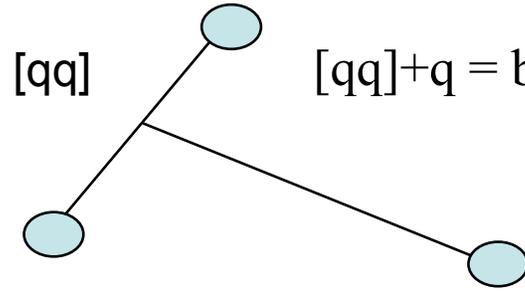
HQ Spin Symmetry

Repulsion of diquarks

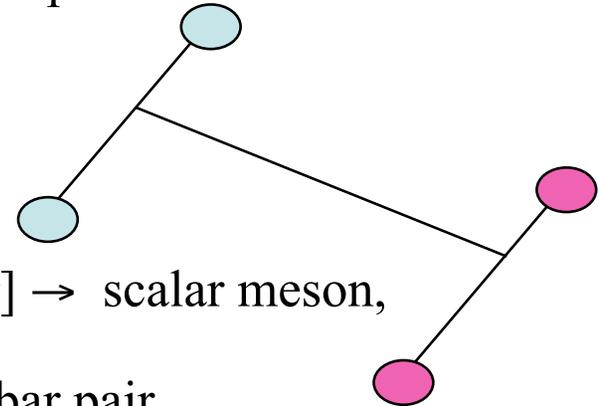


Diquark: $[qq]$ in color = $3\bar{}$, spin=0, SU3 flavour = $3\bar{}$ makes a simple unit to form color singlets (Jaffe..more recently Jaffe&Wilczek, Karliner & Lipkin for penta-quark)

Diquark needs to combine with other colored objects

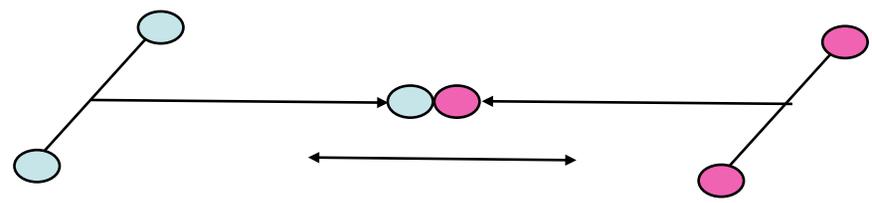


$[qq]+q =$ baryon (e.g. Λ), Y-shape



$[qq]+ [q\bar{q} q\bar{q}] \rightarrow$ scalar meson,

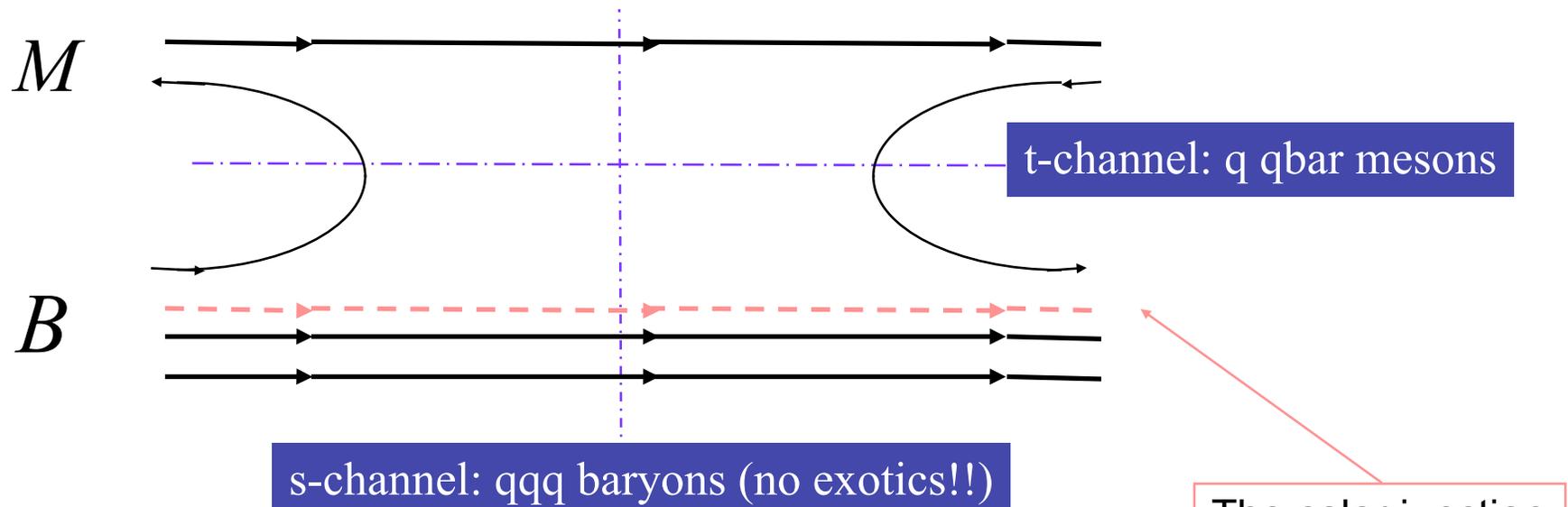
if you stretch the string, $[qq][q\bar{q} q\bar{q}] \rightarrow B B\bar{}$ pair
a new topology, related to B-B bar.



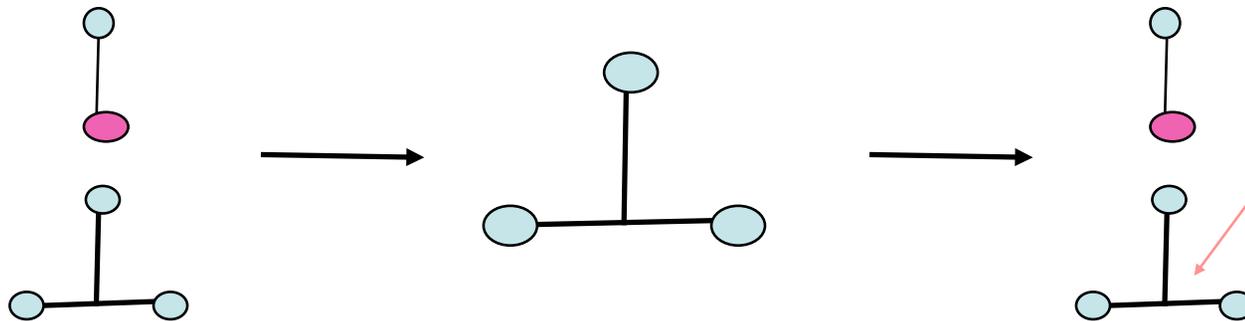
meson-meson molecules are in different color configuration. But: do “residual” forces bind?

Duality in Meson-Baryon scattering

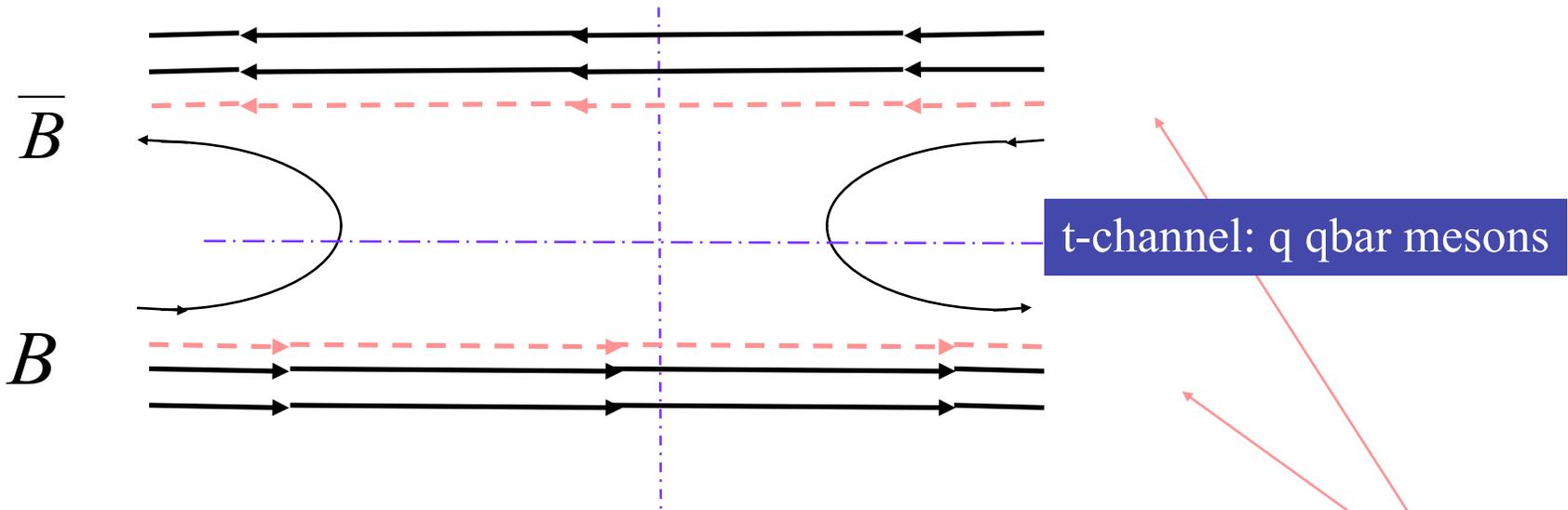
G.C. Rossi and G. Veneziano, Nucl. Phys. **B123** (1977) 507



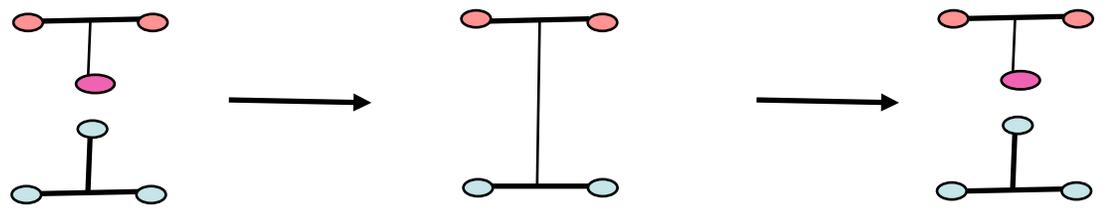
The color junction



What about Baryon-Antibaryon scattering ?



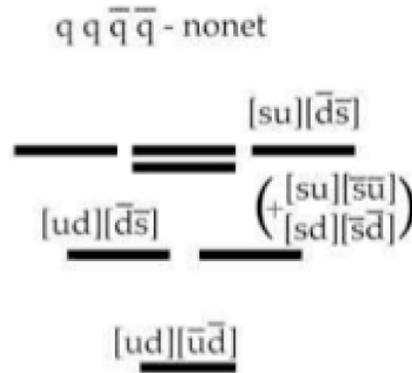
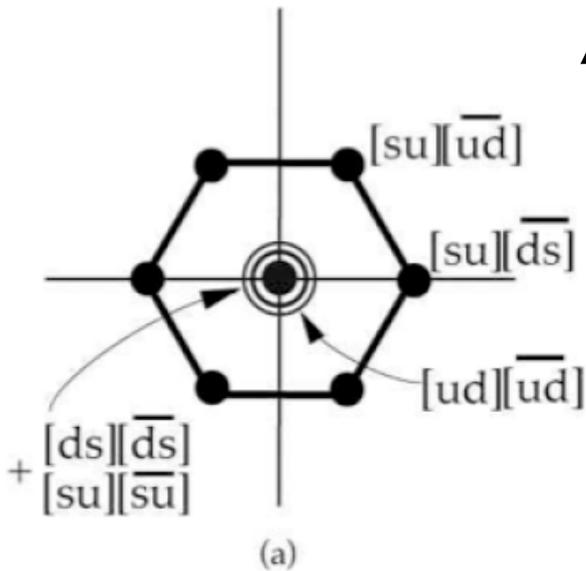
s-channel: [qq][qbar qbar] (Baryonium)



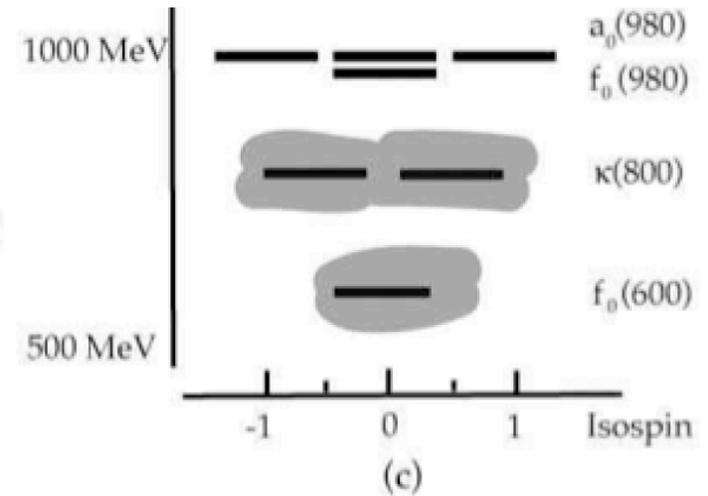
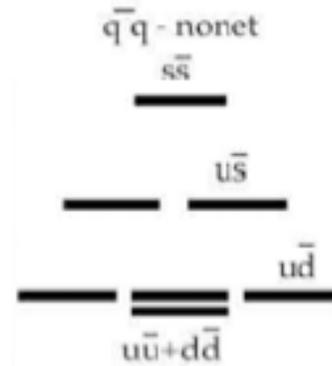
Color junctions

[qq][qbar qbar] mesons are dual to q qbar mesons in B-Bbar scattering. The relation to B-Bar persists in decays!!

Quantum numbers and mass formula



(b)



$$m = \begin{pmatrix} \alpha & & \\ & \alpha & \\ & & \beta \end{pmatrix}$$

= diquark masses

4 parameters, 4 masses+1 mixing, one overall relation:

$$\cos 2\phi + 2\sqrt{2} \sin 2\phi = 1 + 4 \frac{a + \sigma - 2\kappa}{a - \sigma}$$

$$f_o(I = 0) = \frac{1}{\sqrt{2}} ([su][\bar{s}\bar{u}] + [sd][\bar{s}\bar{d}])$$

$$\sigma_o(I = 0) = [ud][\bar{u}\bar{d}]$$

$$|f\rangle = \cos \phi |f_o\rangle + \sin \phi |\sigma_o\rangle$$

$$|\sigma\rangle = -\sin \phi |f_o\rangle + \cos \phi |\sigma_o\rangle.$$

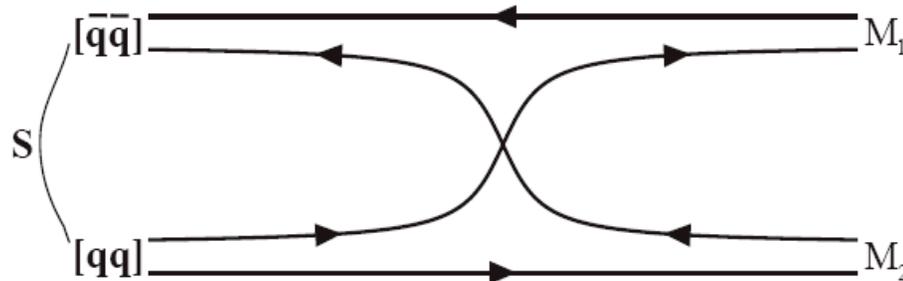
$$\tan 2\phi = -0.07, \quad \sigma = (570 \text{ MeV})^2$$

$$\tan 2\phi = -0.19, \quad \sigma = (470 \text{ MeV})^2$$

$$\tan 2\phi = -0.31, \quad \sigma = (370 \text{ MeV})^2$$

- Almost “ideal mixing”
- Linear mass formula gives very similar results
- With Linear m.f., parameters related to diquark masses: $\alpha=480$ MeV, $\beta = 250$ MeV
- Note: $\alpha-\beta=230$ MeV vs $m_s=150$ MeV.

Strong Decays



$$L_1 = A(S_i^l \varepsilon_{jlm} \varepsilon^{ikn}) M_k^l M_n^m$$

No derivative coupling!

FIG. 1: The decay of a scalar meson S made up of a diquark-antidiquark pair in two mesons $M_1 M_2$ made up of standard $(q\bar{q})$ pairs.

$$\Gamma(S \rightarrow i) = \frac{A^2}{8\pi} \frac{p}{M_s^2} x_{s \rightarrow i}, \quad (13)$$

where p is the decay momentum, M the mass of the scalar meson and $x_{s \rightarrow i}$ a factor which includes numerical coefficients in the individual amplitudes and isospin multiplicities.

	Theory	Experiment	Theory	Experiment
	$\pi\pi$		KK	
σ	345 MeV	324 ± 50 MeV	-	
f	$g_\pi < 0.02$	$g_\pi = 0.19 \pm 0.05$	$g_K = 0.28$	$g_K = 0.40 \pm 0.6$
	$\eta\pi$			
a	43 MeV	60 ± 13 MeV	23 MeV	12 ± 3 MeV
	$K\pi$			
κ	138 MeV	410 ± 100 MeV	-	

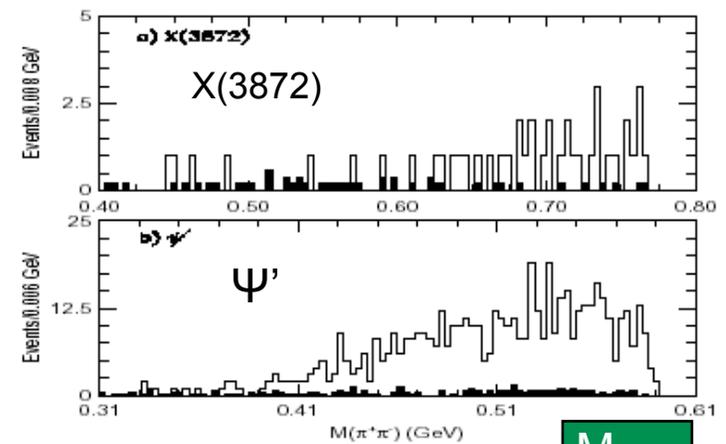
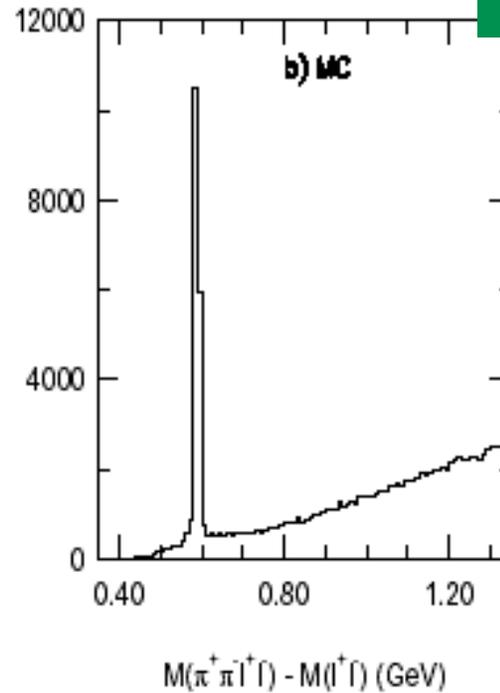
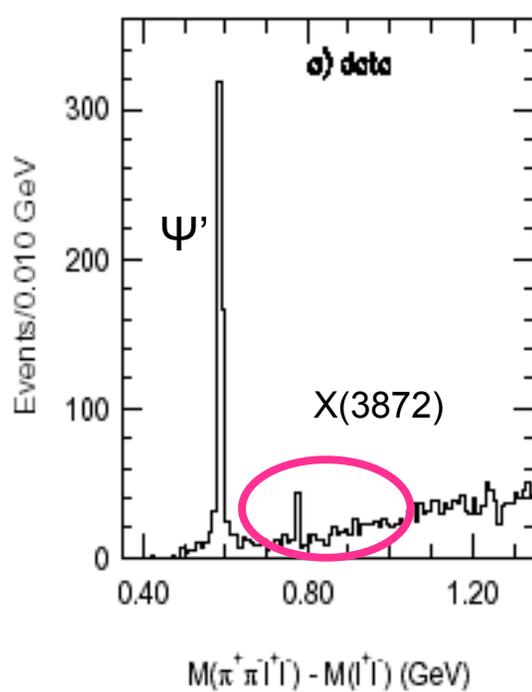
TABLE II: Fit with a single parameter $A = 2.6$ GeV. For g_π we have reported the upper limit to the decay rate obtained from the $f - \sigma$ mixing considered previously, see text.

Maybe $f \pi\pi$ comes from “one-loop”: $f \rightarrow K\bar{K} \rightarrow \pi\pi$,
or perhaps (!!) $f \rightarrow B\bar{B} \rightarrow \pi\pi$ (Baryonium ?, see later)

All in all we get quite a consistent picture, reconciles the large σ width with narrow a and f widths and reinforces $[qq][q\bar{q}q\bar{q}]$ assignement

Observation of a narrow charmonium-like state in exclusive $B^\pm \rightarrow K^\pm \pi^+ \pi^- J/\psi$ decays

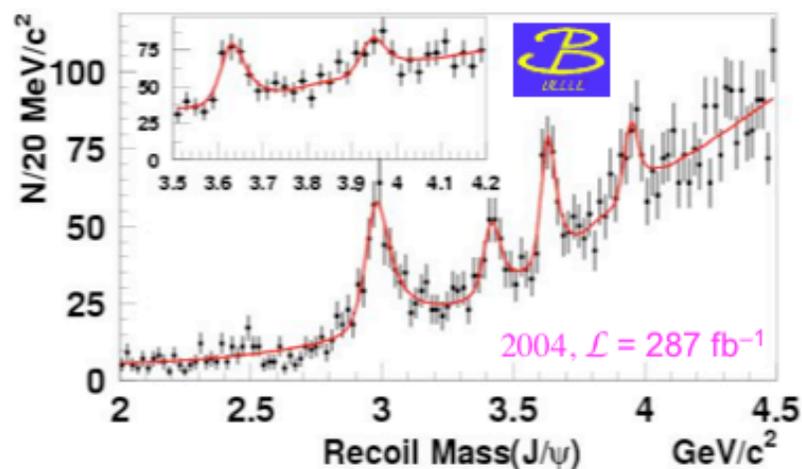
BELLE Collaboration



M_{TTTT}

A new peak in J/ψ recoil

- At ICHEP '04 Belle presented an update with full statistics, extending the mass range again to the high part of the spectrum (3.8 - 4.5 GeV/c^2)
- No evidence for $X(3872)$ on the recoil
- New peak observed at higher mass:
 - $m_Y = (3940 \pm 11) \text{ MeV}/c^2$;
 - $N_Y = 148 \pm 33$ (4.5σ);
 - width consistent with experimental resolution



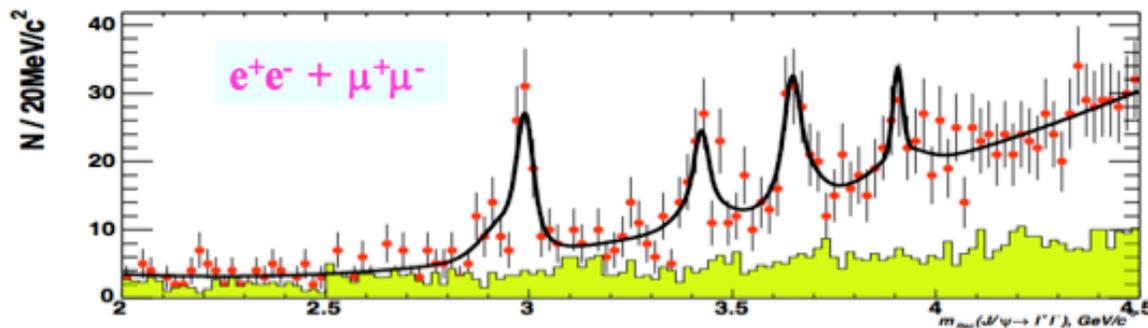
BABAR double $c\bar{c}$ analysis

Analysis similar to Belle's

– Run 1-2-3, $\mathcal{L} = 124.4 \text{ fb}^{-1}$

S. Ye, X. Lou, G. Williams

BAD 718



Similar peak observed at high mass. Very preliminarily:

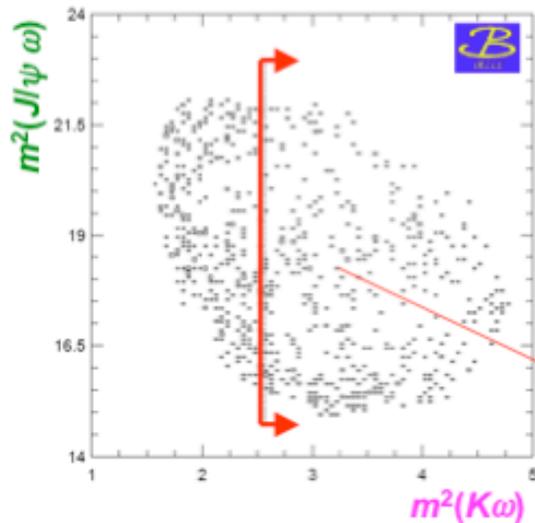
– $m_Y \approx 3910 \text{ MeV}/c^2$

– $N_Y = 51 \pm 27$

Missed ICHEP because of serious problem with lepton lists.

Problem fixed. A few remaining issues before preliminary result ready.

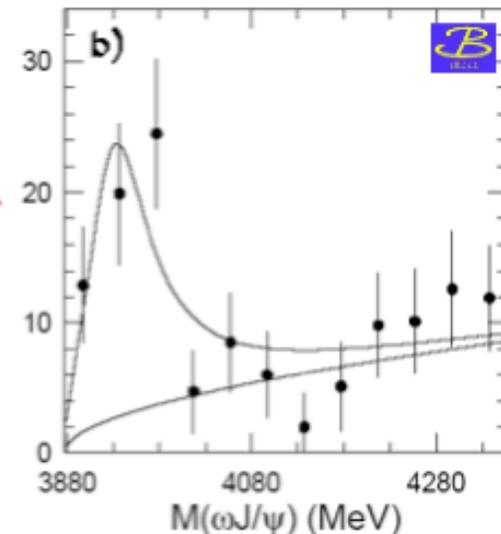
New status decaying to $J/\psi \omega$?



- Select events in $m(\pi^+\pi^-\pi^0) \sim m_\omega$ region
- Dalitz plot $m^2(J/\psi \omega)$ vs. $m^2(K\omega)$:
 - concentration at low $m^2(K\omega)$;
 - select events with $m^2(K\omega)$ to exclude $B \rightarrow J/\psi K^{**}$

- Plot $m(J/\psi \omega)$ distribution after combinatorial bkg. subtraction:

- enhancement just above threshold, not compatible with phase-space $B \rightarrow J/\psi \omega K$;
- Fit to S-wave B-W yields: $N_\gamma = 61 \pm 11 (>8\sigma)$;
 $m_\gamma = (3491 \pm 11) \text{ MeV}/c^2$; $\Gamma_\gamma = (92 \pm 24) \text{ MeV}$

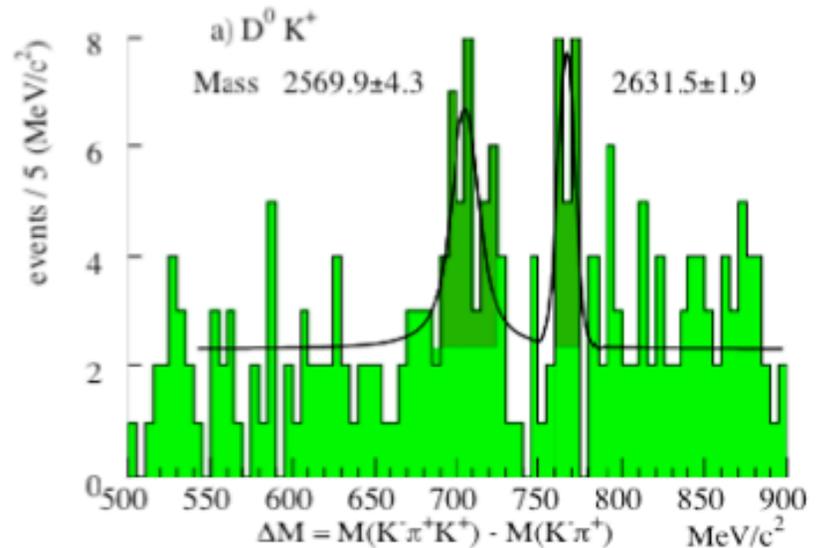
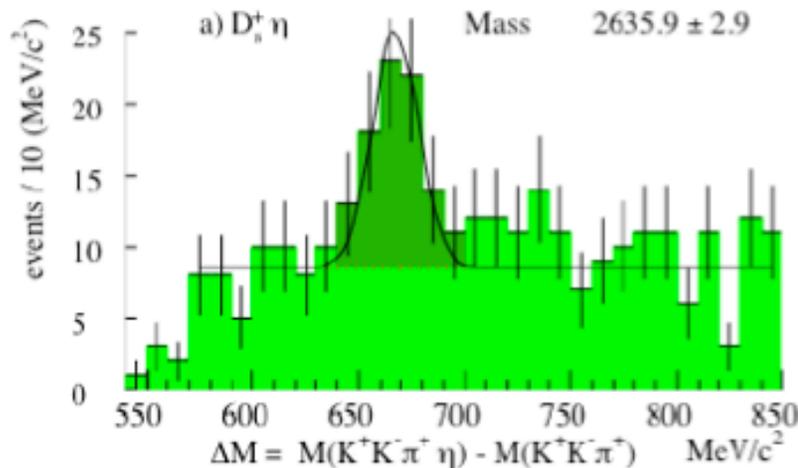


SELEX-Fermilab

hep-ex/0406045

'SELEX reports these peaks as the first observation of yet another high mass Ds state decaying strongly to a ground state charm plus a pseudoscalar meson. The mechanism which keeps this state narrow is unclear[...] The Ds η decay rate dominates the D0K $^+$ by a factor of ~ 6 despite having half the phase space.'

D_sj(2632) (?)

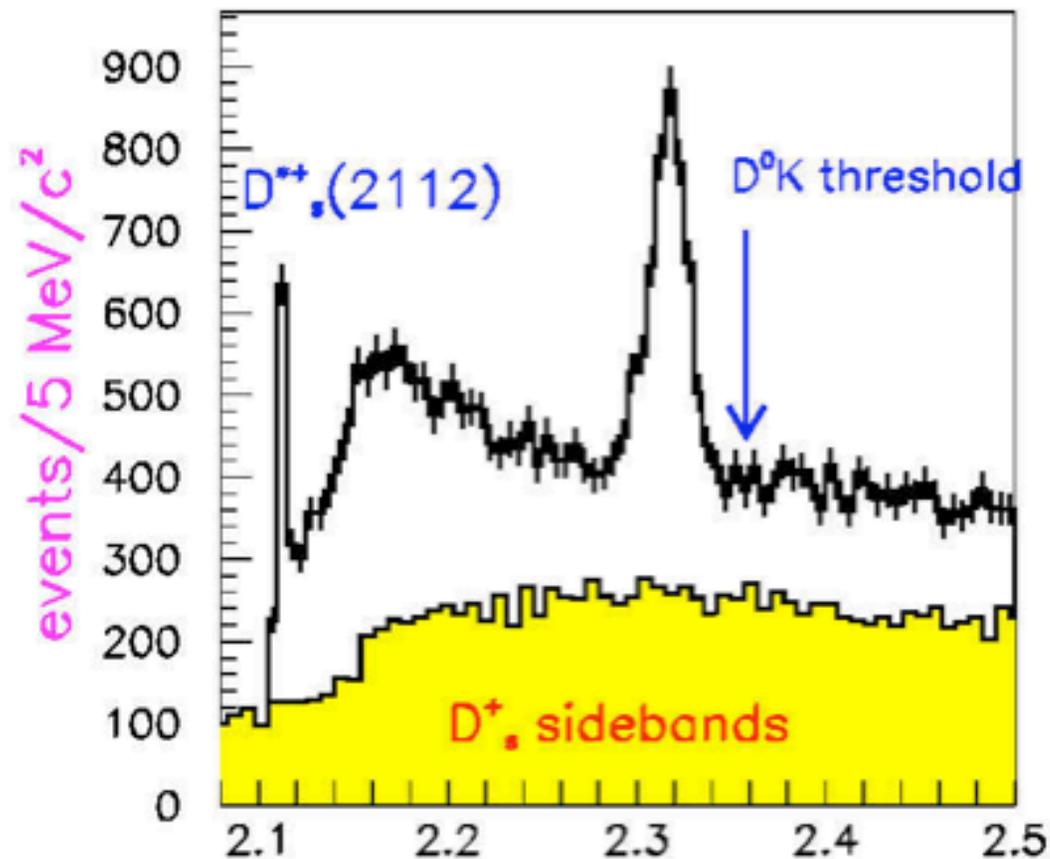


$D_s^*(2317)$ @ BABAR

Observed in 2003 in $D_s \pi^0$

Confirmed by BELLE e CLEO

Does not decay in photon



$M = 2317.4 \pm 0.9$ MeV/c²
 $\Gamma < 4.6$ MeV

HQET multiplets ($Qq\bar{q}$) light dof	$(0^-, 1^-)$	$(0^+, 1^+)$	$(1^+, 2^+)$
$q=u, d$	(D, D^*)	(D_0^*, D_1) broad	(D_1, D_2^*) narrow
$q=s$	(D_s, D_s^*)	$D_s^*(2317)$ $D_s(2460)$ BaBar '03	(D_{s1}, D_{s2}^*) narrow

Likely the missing ($c\bar{s}$) states.
Very narrow.

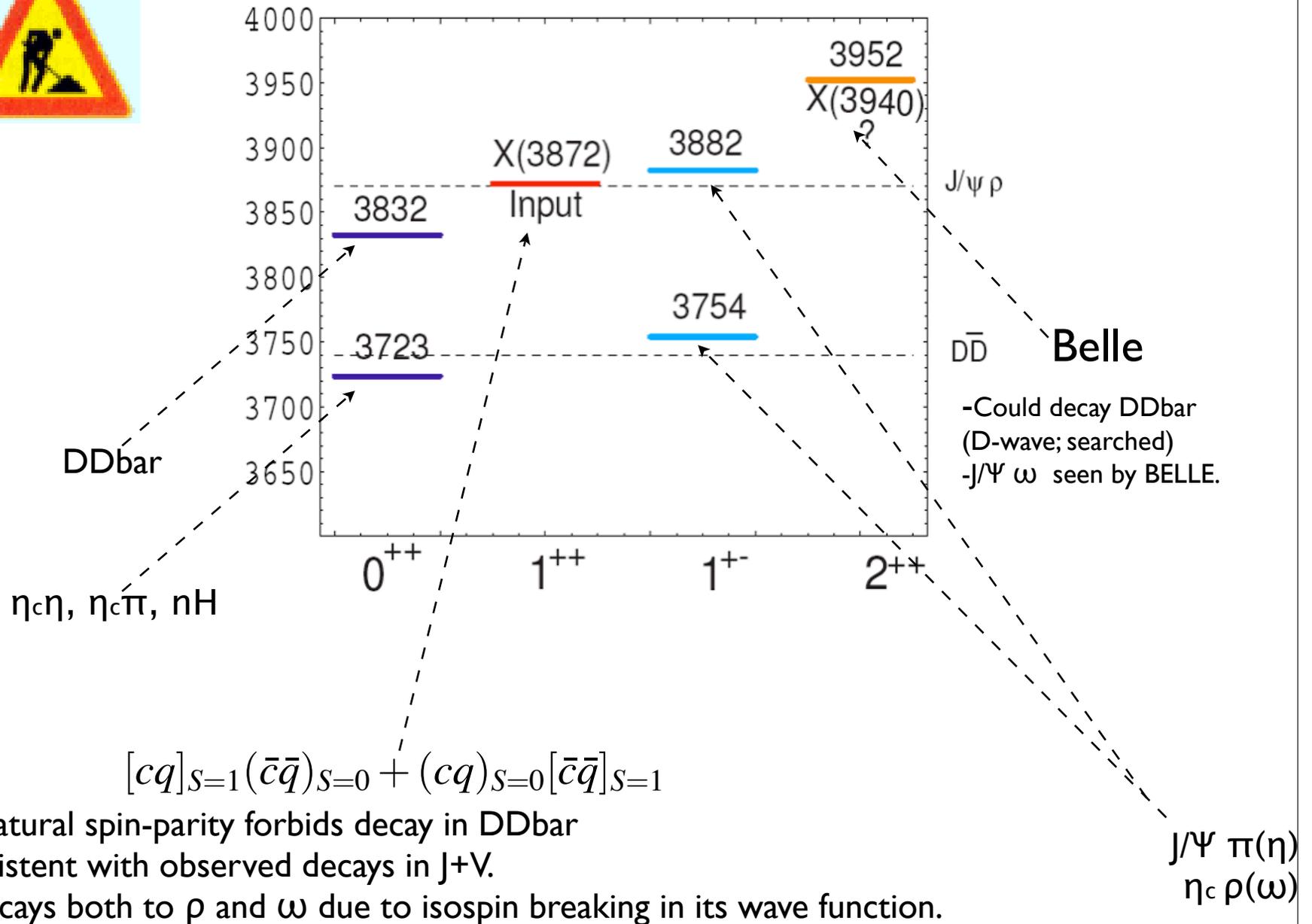
The $0^+(c\bar{s})$ state decays to $D_s\pi^0$ violating isospin because it is below (~ 40 MeV) threshold for DK decay. Relativistic potential models were predicting these states higher in mass \Rightarrow broader.

Tetraquarks with open and hidden charm (Phys.Rev. **D70**, 054009 (2004); hep-ph/0412098)

- The spin-spin interaction between heavy quarks is $O(1/M)$
 - If $S=0$ diquarks are bound, $S=1$ diquarks do
 - All states in the composition $(S=0 \oplus S=1) \otimes (S=0 \oplus S=1)$ must exist
 - not natural spin-parity only!
 - a large multiplet with composition:
$$2 (J^{PC}=0^{++}) + (J=1^{++}) + 2 (J=1^{+-}) + (J=2^{++}).$$
- Mass spectrum determined by:
 - constituent diquark masses
 - spin-spin interactions
 - the latter: from meson and baryon spectrum or from one gluon exchange



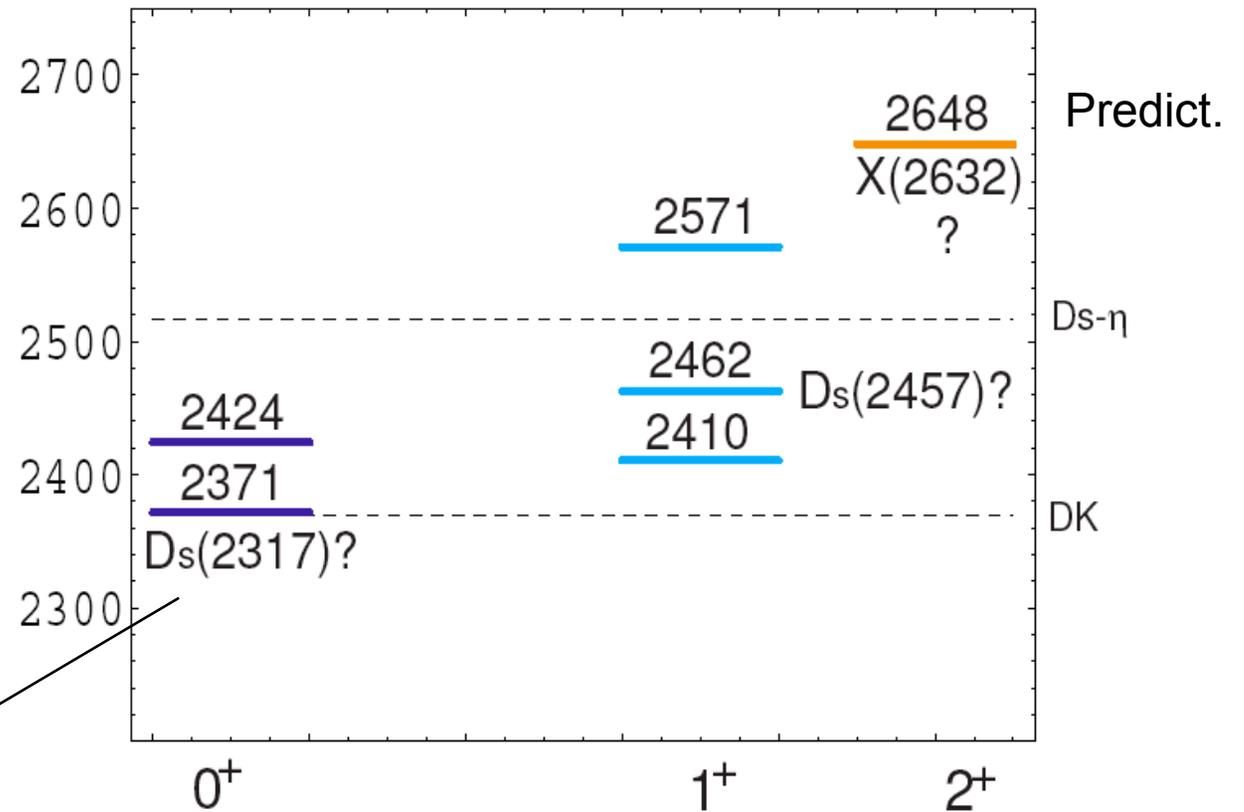
X-states





The states:

$$[cu]_{S=0,1} [\bar{s}\bar{u}]_{S=0,1}$$



Masses

- Contributions to the mass of bound states: two conflicting mechanisms:
 - *Annihilation*: $d\bar{d} \rightarrow u\bar{u}$
 - $u\bar{u} \rightarrow d\bar{d}$

gives a matrix with all equal elements, which is diagonal in the isospin basis;

 - *Quark Masses*: the eigenvectors are the states with quarks of definite flavor (e.g. ω/ϕ mixing)
 - TOTAL:
$$\begin{pmatrix} 2m_u + \delta & \delta \\ \delta & 2m_d + \delta \end{pmatrix}$$
 -
 - At charmonium scale, quark mass should dominate (Rossi -Veneziano;Maiani-Piccinini-Polosa-Riquer)
 - therefore the approximate mass eigenstates are $X_u = [cu][\bar{c}\bar{u}]$
 - $X_d = [cd][\bar{c}\bar{d}]$
 - rather than the $I=1,0$ states
- Belle sees both: $X \rightarrow J+\rho, J+\omega$ with similar B.R.!!!! A new phenomenon !!!!

Isospin breaking

We consider in this section the finer structure of the $X(3872)$. In particular, we consider the neutral states with the composition:

$$X_u = [cu][\bar{u}\bar{c}]; X_d = [cd][\bar{d}\bar{c}] \quad (32)$$

Physical states could be expected to fall in isospin multiplets with $I = 1, 0$:

$$\begin{aligned} a_{c\bar{c}} &= (X_u + X_d)/\sqrt{2}; \\ f_{c\bar{c}} &= (X_u - X_d)/\sqrt{2} \end{aligned} \quad (33)$$

$$\begin{aligned} X_{low} &= \cos\theta X_u + \sin\theta X_d; \\ X_{high} &= -\sin\theta X_u + \cos\theta X_d \end{aligned} \quad (34)$$

we get:

$$\begin{aligned} M(X_d) - M(X_u) &= \frac{2(m_{down} - m_{up})}{\cos(2\theta)} = \\ &= \frac{(6 - 8) \text{ MeV}}{\cos(2\theta)} \end{aligned} \quad (35)$$

$$\begin{aligned} \frac{\Gamma(3\pi)}{\Gamma(2\pi)}_{X_l} &= \frac{(\cos\theta + \sin\theta)^2}{(\cos\theta - \sin\theta)^2} \cdot \frac{\langle p_\omega \rangle}{\langle p_\rho \rangle} \\ \frac{\Gamma(3\pi)}{\Gamma(2\pi)}_{X_h} &= \frac{(\cos\theta - \sin\theta)^2}{(\cos\theta + \sin\theta)^2} \cdot \frac{\langle p_\omega \rangle}{\langle p_\rho \rangle} \end{aligned} \quad (44)$$

BELLE attributes all events with $\pi^+\pi^-\pi^0$ mass above 750 MeV to ω decay and divides by the total number of observed 2π events. They find:

$$\left(\frac{\Gamma(3\pi)}{\Gamma(2\pi)}\right)_{BELLE} = 0.8 \pm 0.3_{stat} \pm 0.1_{syst} \quad (45)$$

The central value is compatible with eq.(44) for:

$$\theta \simeq \pm 20^\circ \quad (46)$$

for X_l or X_h , respectively. Correspondingly, the mass difference of the two states is:

$$M(X_h) - M(X_l) \simeq 7 - 10 \text{ MeV} \quad (47)$$

Interference !

$$\begin{aligned} \frac{d\Gamma(X \rightarrow \psi + e^+e^-)}{ds} &= \\ &= \frac{|A|^2 B_{(\rho \rightarrow e^+e^-)} M_\rho \Gamma_V}{8\pi M_X^2 \pi} \cdot p(s) \cdot \\ &\cdot \left| \frac{1}{(s - M_\rho^2) + i(M_\rho \Gamma_\rho)} \pm \frac{1/3}{(s - M_\omega^2) + i(M_\omega \Gamma_\omega)} \right|^2 \end{aligned}$$

we have assumed the quark-model ratio for the leptonic amplitudes of ρ and ω and used the narrow width approximation. The sign \pm applies to X_u and X_d , respectively. Combining with eq.(43), with $\theta = 0$, we find:

$$\begin{aligned} B(X_u \rightarrow J/\Psi + e^+e^-) &= 0.8 \cdot 10^{-4} \\ B(X_d \rightarrow J/\Psi + e^+e^-) &= 0.3 \cdot 10^{-4} \end{aligned} \quad (49)$$

Decay widths

- The baryonium picture implies that the two-meson decays go via intermediate baryon-antibaryon states of high mass. This implies basically narrow widths.
- We describe the decay by a single switch amplitude, associated to the process (subscripts indicate color configuration): $[cu][\bar{c}\bar{u}] \rightarrow (c\bar{c})_{col=1}(u\bar{u})_{col=1}$

$$L_{X_u \Psi V} = g_V \epsilon^{\mu\nu\rho\sigma} P_\mu X_\nu \psi_\rho V_\sigma = g_V M_X (\mathbf{X} \wedge \boldsymbol{\psi}) \cdot \mathbf{V} \quad g_V M_X = \frac{A}{\sqrt{2}}$$

- a bold guess: $A=2.6 \text{ GeV}$
- $\Gamma(X_l \rightarrow J/\psi + \pi^+ \pi^-) = \frac{2x_{l,\rho}|A|^2}{8\pi M_X^2} \langle p \rangle_\rho =$
- $= 2x_{l,\rho} \cdot 2.3 \text{ MeV};$
- $\Gamma(X_l \rightarrow J/\psi + \pi^+ \pi^- \pi^0) = \frac{2x_{l,\omega}|A|^2}{8\pi M_X^2} \langle p \rangle_\omega =$
- $= 2x_{l,\omega} \cdot 0.4 \text{ MeV}$
- We anticipate small widths, comparable to the resolution of Belle and Babar

X particles in B decays

- Taking Belle data at face value, we conclude that only one of the two neutral states is produced appreciably in B^+ decay
- The the other has to appear in B^0 decay:
 - The X particles in B^+ and B^0 decays are not the same, and have a mass difference of 7 ± 2 MeV
- Bounds to the production of X^+ :

$$\begin{aligned}
 R^+ &= \\
 &= \frac{\mathcal{B}(B^+ \rightarrow K_S X^+) \cdot \mathcal{B}(X^+ \rightarrow J/\Psi + \pi^+ \pi^0)}{\mathcal{B}(B^+ \rightarrow K^+ X_{l/h}) \cdot \mathcal{B}(X_{l/h} \rightarrow J/\Psi + \pi^+ \pi^-)} > 0.2 \\
 R^0 &= \\
 &= \frac{\mathcal{B}(B^0 \rightarrow K^+ X^-) \cdot \mathcal{B}(X^- \rightarrow J/\Psi + \pi^- \pi^0)}{\mathcal{B}(B^0 \rightarrow K_S X_{h/l}) \cdot \mathcal{B}(X_{h/l} \rightarrow J/\Psi + \pi^+ \pi^-)} > 0.53
 \end{aligned}$$

to be compared with the upper limit given by BaBar [27]:

$$R^+ < 0.8$$

with large errors.

Conclusions

- A convincing picture of light scalars as $[qq][\bar{q}\bar{q}]$ states:
 - Masses
 - Ideal mixing
 - Decays reasonably described (exact SU3!) but for OZI violating (??)
 - Note: $\Delta m(f-a) \sim 10\text{MeV}$, $\Delta m(\text{up-down}) \sim 5\text{MeV}$: are $f(980)$ and $a(980)$ pure I-spin eigenstates?
- New phenomena
 - States $[cq][\bar{c}\bar{q}]$ and $[cq][\bar{c}\bar{s}]$ should exist, with both natural and unnatural spin parity;
 - I-spin breaking expected maximal in certain decay: was the SELEX particle just the first case?
 - X(3872) a good candidate, X(3940) predicted

Study of Φ decays in KLOE very important

WERE ARE THE EXOTIC STATES???

APPENDIX

E791-PRL86(2000) claims that almost the 46% of

$$D \rightarrow 3\pi$$

could resonate on a *scalar bump* with

$$m_\sigma = 478 \pm 24 \text{ MeV}$$

$$\Gamma_\sigma = 324 \pm 41 \text{ MeV}$$

Has this something to do with the pole on the second sheet of the isoscalar S-wave in $\pi\pi$ scattering found, e.g., by Colangelo-Gasser-Leutwyler NPB603(2001) and many others?

$$\sqrt{s} = (479 \pm 30) - i(295 \pm 20)$$

Their reply is '*there is no harm in calling this an unusually broad resonance...*'. Is it a broad enhancement whose dynamical origin is in the strong pionic FSI for these quantum numbers?

E791 made the data analysis describing this scalar bump as a *Breit-Wigner resonance!*

The full SU(3) effective lagrangian...

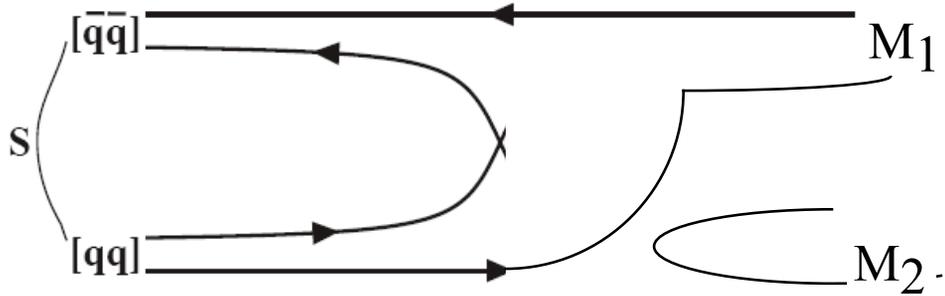
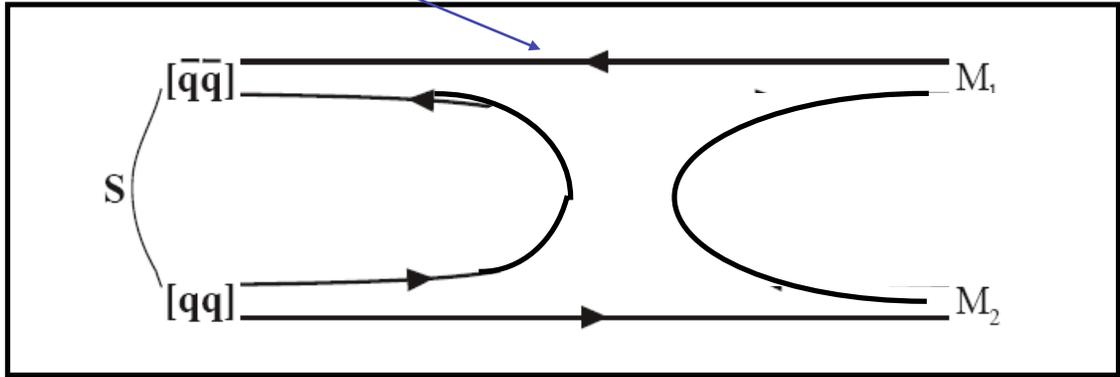
has three couplings:

$$\mathcal{L} = (S_i^j) \epsilon_{jlm} \epsilon^{ikn} [a M_k^l M_n^m + b \delta_k^l (M^2)_n^m + c \delta_k^l (M)_n^m \text{Tr} M],$$

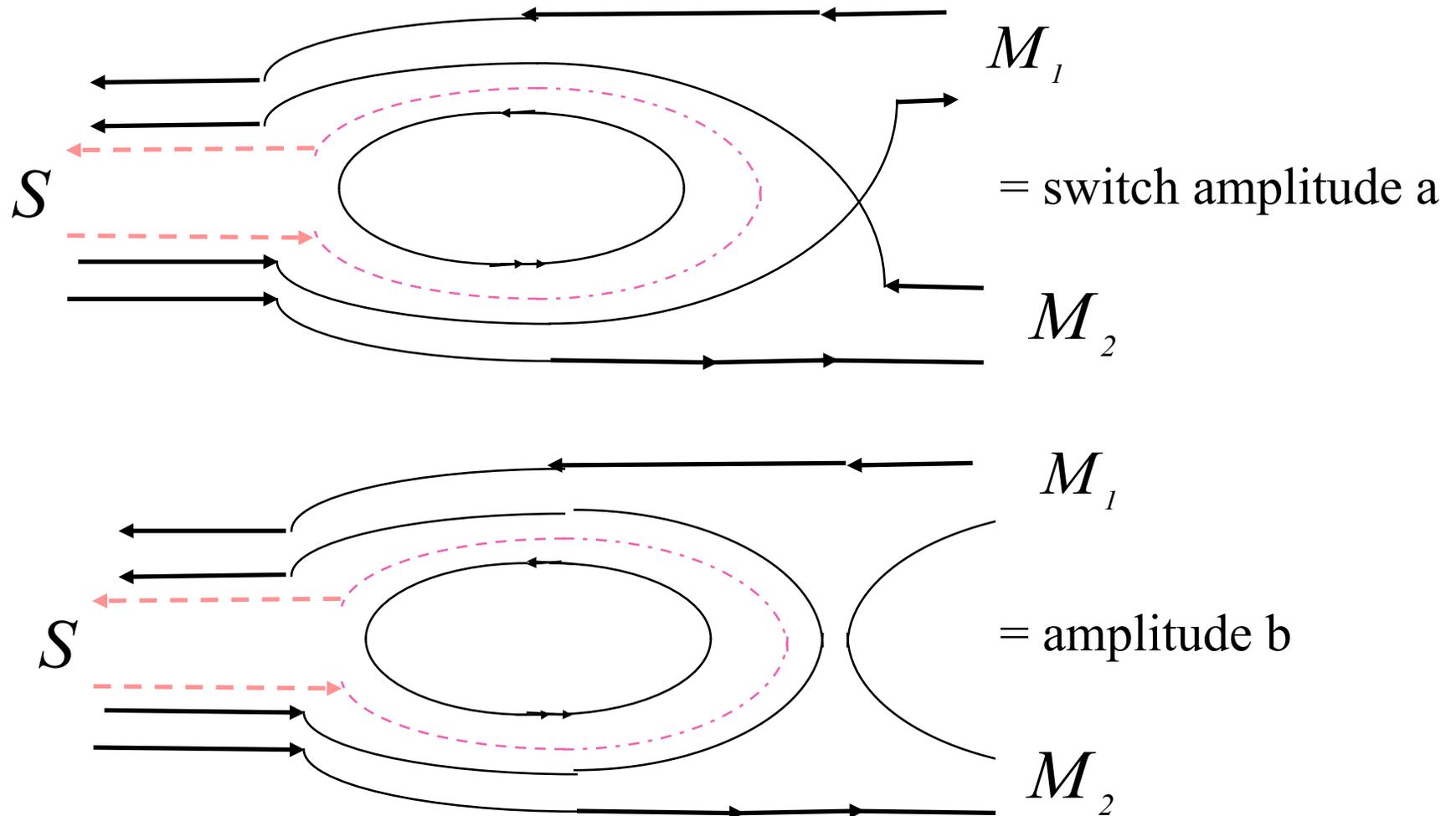
$$\begin{aligned} \mathcal{L} = & f_0 \left[b\sqrt{2} \frac{\pi \cdot \pi}{2} - (a - 3b) \frac{\bar{K} K}{\sqrt{2}} + \dots \right] \\ & + \sigma_0 \left[-(a - 2b) \frac{\pi \cdot \pi}{2} + b \bar{K} K + \dots \right] \\ & + a^0 \left[(a - b) \frac{\bar{K} \tau_3 K}{\sqrt{2}} - (a - c) \eta_s \pi^0 \right. \\ & \left. - \sqrt{2} (b - c) \eta_q \pi^0 + \dots \right] \\ & + (a - b) \left(\frac{\bar{K}^+ \pi^0}{\sqrt{2}} + \bar{K}^0 \pi^- \right) \kappa^+ + \dots \end{aligned}$$

Quark rearrang.

$$\mathcal{L} = (S_i^j) \epsilon_{jlm} \epsilon^{ikn} [a M_k^l M_n^m + b \delta_k^l (M^2)_n^m + c \delta_k^l (M)_n^m \text{Tr} M],$$



Diquark-antidiquark decays



Decay proceed via B- Bbar virtual states! Should we try the same for $\phi \rightarrow S\gamma$?

$$\begin{aligned}
|a - 2b| &= 2.6 \text{ GeV} (= A) \text{ (from } \sigma \rightarrow \pi\pi), \\
|a - 3b| &= 3.1 \text{ GeV (from } f_0 \rightarrow K\bar{K}), \\
|a - b| &= 1.8 \text{ GeV (from } a_0 \rightarrow K\bar{K}), \quad (19)
\end{aligned}$$

indeed equally spaced with: $b = -0.7 \text{ GeV}$.

We find further: $\Gamma(\kappa) = 66 \text{ MeV}$; $g_\pi = 0.06$.

c should be small: with $c=0$ we get $\Gamma(a \rightarrow \eta\pi) = 30 \text{ MeV}$ vs $60 \pm 13 \text{ MeV}$

Maybe $f \rightarrow \pi\pi$ comes from “one-loop”:

$$f \rightarrow K\bar{K} \rightarrow \pi\pi,$$

or perhaps (!!) $f \rightarrow B\bar{B} \rightarrow \pi\pi$ (Baryonium picture, see later)

Still too low.

Overall we get quite a consistent picture, reconciles σ with a and f widths, reinforces $[qq][q\bar{q}q\bar{q}]$ assignment

A=2.6 GeV

1. $\Gamma_{tot}(a_0) = 72 \pm 16$ MeV and the KK branching ratio $\mathcal{B}(a_0 \rightarrow K\bar{K}) = 0.17 \pm 0.03$ thus obtaining:

$$\Gamma(a_0 \rightarrow \eta\pi) = 60 \pm 13 \text{ MeV}, \quad (14)$$

$$\Gamma(a_0 \rightarrow K\bar{K}) = 12 \pm 3 \text{ MeV}. \quad (15)$$

2. We compute the decay momentum with the central values of the parent mass, with the exception of the decay $a \rightarrow K\bar{K}$, which is below threshold at the central mass value. In this case we have averaged the decay momentum over a Breit-Wigner, using the $\Gamma_{tot}(a)$ given above, and find: $\langle p(a \rightarrow K\bar{K}) \rangle \approx 84$ MeV, which gives the value of the partial width reported in Table II.

3. In the case of $f \rightarrow K\bar{K}$ or $\pi\pi$, the authors of ref. [11] define:

$$\Gamma(S \rightarrow i) = g_i p(M) \quad (16)$$

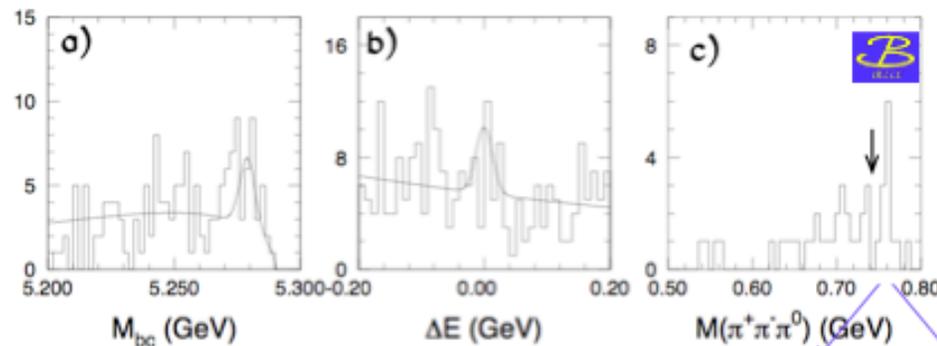
and fit the data to a Breit-Wigner formula with mass-dependent width, thus giving directly the values of g_i that we report in the Table II.

D. Barberis et al., Phys. Lett. **B453** (1999) 325.

By the way...

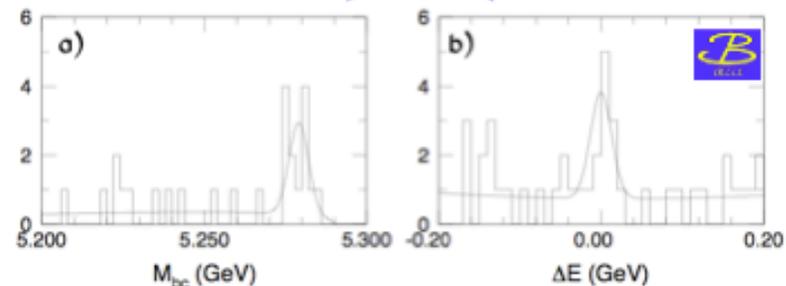
□ Events selected within 2σ ($12 \text{ MeV}/c^2$) of $3872 \text{ MeV}/c^2$ show hints of a B signal.

– After cutting at $m(\pi^+\pi^-\pi^0) > 750 \text{ MeV}/c^2$, a cleaner signal is present over a low background. Fit yields 10.0 ± 3.6 events (5.8σ)



□ Assuming all signal is ω (and $X(3872)$!):

$$\frac{\Gamma(X \rightarrow J/\psi \omega)}{\Gamma(X \rightarrow J/\psi \pi^+ \pi^-)} = 0.8 \pm 0.3 \pm 0.1$$



Diquark-Antidiquarks with Hidden or Open Charm and the Nature of $X(3872)$

L. Maiani*

Università di Roma 'La Sapienza' and I.N.F.N., Roma, Italy

F. Piccinini†

I.N.F.N. Sezione di Pavia and Dipartimento di Fisica Nucleare e Teorica, Pavia, Italy

A.D. Polosa‡

Dipartimento di Fisica, Università di Bari and I.N.F.N., Bari, Italy

V. Riquer§

CERN Theory Department, CH-1211, Switzerland

(Dated: December 22, 2004)

Heavy-light diquarks can be the building blocks of a rich spectrum of states which can accommodate some of the newly observed charmonium-like resonances not fitting a pure $c\bar{c}$ assignment. We examine this possibility for hidden and open charm diquark-antidiquark states deducing spectra from constituent quark masses and spin-spin interactions. Taking the $X(3872)$ as input we predict the existence of a 2^{++} state that can be associated to the $X(3940)$ observed by Belle and re-examine the state claimed by SELEX, $X(2632)$. The possible assignment of the previously discovered states $D_s(2317)$ and $D_s(2457)$ is discussed. We predict $X(3872)$ to be made of two components with a mass difference related to $m_u - m_d$ and discuss the production of $X(3872)$ and of its charged partner X^\pm in the weak decays of $B^{+,0}$.

ROMA1-1396/2004, FNT/T-2004/20, BA-TH/502/04, CERN-PH-TH/2004-239

PACS numbers: 12.40.Yx, 12.39.-x, 14.40.Lb

I. INTRODUCTION

It is an old idea that the light scalar mesons $a(980)$ and $f(980)$ may be 4-quark bound states [1]. The idea was more or less accepted in the mid-seventies but then it lost momentum, due to contradictory results that led the lowest-lying candidate members of a diquark-antidiquark nonet, σ and κ , to disappear from the Particle Data Tables.

From this point of view, the recent observations of $\sigma(480)$ and $\kappa(800)$ in D non-leptonic decays at Fermilab [3] and in the $\pi\pi$ spectrum in $\phi \rightarrow \pi^0\pi^0\gamma$ at Frascati [4], have considerably reinforced the case of a full nonet with inverted spectrum, as expected for $[qq][\bar{q}\bar{q}]$ states and fully antisymmetric diquark ($[qq]$: color = $\bar{3}$, flavor = $\bar{3}$, spin = 0). The isolated $I = 0$ state is the lightest and it likes to decay in $\pi\pi$; the heaviest particles have $I = 1, 0$ and like to decay in states containing strange quark pairs.

iv:hep-ph/0412098 v2 21 Dec 2004