



# Fit to the gluonium content using PDG 2008 and new KLOE measurement of $\text{Br}(\omega \rightarrow \pi^0 \gamma)$

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# Outline

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- ★ **Model description;**
- ★ **Published KLOE result respect to other estimates;**
- ★ **Re-fit to all relevant measurements;**
- ★ **Re-fit using PDG-2008 and new KLOE measurement of  $\text{Br}(\omega \rightarrow \pi^0 \gamma)$ .**



# $\eta, \eta'$ : mixing and gluonium

The  $\eta, \eta'$  mesons wave function can be decomposed in the quark mixing base as in the following (J. L. Rosner, Phys. Rev. D 27 (1983) 1101. ).

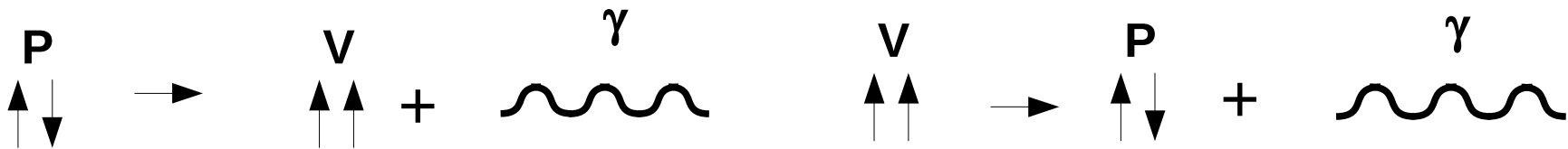
$$|\eta'\rangle = X_{\eta'} |q\bar{q}\rangle + Y_{\eta'} |s\bar{s}\rangle + Z_{\eta'} |G\rangle \quad |\eta\rangle = \cos\varphi_P |q\bar{q}\rangle - \sin\varphi_P |s\bar{s}\rangle \quad |q\bar{q}\rangle = \frac{|u\bar{u}\rangle + |d\bar{d}\rangle}{\sqrt{2}}$$

$$X_{\eta'} = \sin\varphi_P \cos\varphi_G$$

$$Y_{\eta'} = \cos\varphi_P \cos\varphi_G$$

$$Z_{\eta'} = \sin\varphi_G$$

The  $\phi \rightarrow \eta, \eta' \gamma$  transition is modelled according a spin flip transition



$$\Gamma(P \rightarrow V \gamma) = \frac{g^2}{4\pi} |p_y|^3$$

$$\Gamma(V \rightarrow P \gamma) = \frac{1}{3} \frac{g^2}{4\pi} |p_y|^3$$

Only quarks participate to the electromagnetic transition, gluonium is a spectator. It appears in the  $\eta'$  decay amplitudes only through the normalisation to 1 ( $Y_{\eta'} \sim \cos\varphi_G$ )



# Magnetic dipole transition

Decay width:

$$\Gamma(V \rightarrow P\gamma) = \frac{2}{3}\alpha\omega^3 \left(\frac{E_P}{m_V}\right) \sum \left| \left\langle V \left| \frac{\mu_q e_q \sigma_q}{e} \right| P \right\rangle \right|^2$$

Quark charge

Pauli matrices

example: Matrix element for  $\rho \rightarrow \eta\gamma$  decay

$$|\rho\rangle = \frac{|u\bar{u}\rangle - |d\bar{d}\rangle}{\sqrt{2}}$$

$$\left\langle \rho \left| \frac{\mu_q e_q \sigma_q}{e} \right| \eta \right\rangle = \mu \cos(\varphi_P) \frac{2}{3} \langle u\bar{u}_\rho | u\bar{u}_\eta \rangle + \frac{1}{3} \langle d\bar{d}_\rho | d\bar{d}_\eta \rangle = \mu \cos(\varphi_P) \langle q\bar{q}_\rho | q\bar{q}_\eta \rangle$$

Isospin conservation

$e/m_q$  (effective quark mass)



# Magnetic dipole transition

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$$\Gamma(V \rightarrow P\gamma) = \frac{2}{3}\alpha\omega^3 \left(\frac{E_P}{m_V}\right) \sum \left| \left\langle V \left| \frac{\mu_q e_q \sigma_q}{e} \right| P \right\rangle \right|^2$$

Quark charge

Pauli matrices

example: Matrix element for  $\rho \rightarrow \eta\gamma$  decay

$$|\rho\rangle = \frac{|u\bar{u}\rangle - |d\bar{d}\rangle}{\sqrt{2}}$$

$$\left\langle \rho \left| \frac{\mu_q e_q \sigma_q}{e} \right| \eta \right\rangle = \mu \cos(\varphi_P) \frac{2}{3} \langle u\bar{u}_\rho | u\bar{u}_\eta \rangle + \frac{1}{3} \langle d\bar{d}_\rho | d\bar{d}_\eta \rangle = \mu \cos(\varphi_P) \langle q\bar{q}_\rho | q\bar{q}_\eta \rangle$$

$e/m_q$  (effective quark mass)

wave function overlapping

$$C_q = \langle \eta_q | \omega_q \rangle = \langle \eta_q | \rho \rangle \quad C_s = \langle \eta_s | \phi_s \rangle \quad C_\pi = \langle \pi | \omega_q \rangle = \langle \pi | \rho \rangle$$

In the formulas only the ratios appear:  $Z_{NS} = C_q/C_\pi$      $Z_S = C_s/C_\pi$

QCD effects reside in mixing parameters, overlapping parameters and effective quark masses.



# V P $\gamma$ and P $\gamma\gamma$ transitions

**KLOE [Phys. Lett. B648 (2007) 267] has fitted:**

$$\frac{\Gamma(\eta' \rightarrow \rho \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{z_q^2}{\cos^2 \phi_V} \cdot 3 \left( \frac{m_{\eta'}^2 - m_\rho^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 X_{\eta'}$$

$$\frac{\Gamma(\eta' \rightarrow \omega \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{1}{3} \left( \frac{m_{\eta'}^2 - m_\omega^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 \left[ z_q X_{\eta'} + 2 \frac{m_s}{\bar{m}} z_s \cdot \tan \phi_V \cdot Y_{\eta'} \right]^2$$

**together with the measured branching ratio:**

$$R_\phi = (4.77 \pm 0.09 \pm 0.19) \times 10^{-3}$$

$$R_\phi = \frac{Br(\phi \rightarrow \eta' \gamma)}{Br(\phi \rightarrow \eta \gamma)} = \cot^2 \phi_P \cdot \cos^2 \phi_G \left( 1 - \frac{m_s}{\bar{m}} \frac{z_q}{z_s} \cdot \frac{\tan \phi_V}{\sin 2\phi_P} \right)^2 \cdot \left( \frac{p_{\eta'}}{p_\eta} \right)^3$$

**and the ratio:**  $\frac{\Gamma(\eta' \rightarrow \gamma \gamma)}{\Gamma(\pi^0 \rightarrow \gamma \gamma)} = \frac{1}{9} \left( \frac{m_{\eta'}}{m_\pi} \right)^3 \left( 5X_{\eta'} + \sqrt{2} \frac{f_q}{f_s} Y_{\eta'} \right)^2$  **E. Kou, Phys. Rev. D 63 (2001) 54027**



# V P $\gamma$ and P $\gamma\gamma$ transitions

**KLOE [Phys. Lett. B648 (2007) 267] has fitted:**

Were taken from a global fit without gluonium:

$$\frac{\Gamma(\eta' \rightarrow \rho \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{z_q^2}{\cos \phi_V} \cdot 3 \left( \frac{m_{\eta'}^2 - m_\rho^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 X_{\eta'}^2$$

A. Bramon, R. Escribano,  
M.D. Scadron  
Phys. Lett. B503 (2001) 271

$$\frac{\Gamma(\eta' \rightarrow \omega \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{1}{3} \left( \frac{m_{\eta'}^2 - m_\omega^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 \left[ z_q^2 X_{\eta'} + 2 \frac{m_s}{\bar{m}} z_s \tan \phi_V Y_{\eta'} \right]^2$$

together with the measured branching ratio:

$$R_\phi = (4.77 \pm 0.09 \pm 0.19) \times 10^{-3}$$

$$R_\phi = \frac{Br(\phi \rightarrow \eta' \gamma)}{Br(\phi \rightarrow \eta \gamma)} = \cot^2 \phi_P \cdot \cos^2 \phi_G \left( 1 - \frac{m}{\bar{m}} \frac{z_q \tan \phi_V}{z_s \sin 2\phi_P} \right)^2 \cdot \left( \frac{p_{\eta'}}{p_\eta} \right)^3$$

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# V P $\gamma$ and P $\gamma\gamma$ transitions

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$$\frac{\Gamma(\eta' \rightarrow \rho \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{z_q^2}{\cos^2 \phi_V} \cdot 3 \left( \frac{m_{\eta'}^2 - m_\rho^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 X_{\eta'}^2$$

T. Feldmann, Int. J. Mod. Phys. A 15 (2000) 159

$$\frac{\Gamma(\eta' \rightarrow \omega \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{1}{3} \left( \frac{m_{\eta'}^2 - m_\omega^2 m_\omega}{m_\omega^2 - m_\pi^2 m_{\eta'}} \right)^3 \left[ z_q X_{\eta'} + 2 \frac{m_s}{\bar{m}} z_s \cdot \tan \phi_V \cdot Y_{\eta'} \right]^2$$

**together with the measured branching ratio:**

$$R_\phi = (4.77 \pm 0.09 \pm 0.19) \times 10^{-3}$$

$$R_\phi = \frac{Br(\phi \rightarrow \eta' \gamma)}{Br(\phi \rightarrow \eta \gamma)} = \cot^2 \phi_P \cdot \cos^2 \phi_G \left( 1 - \frac{m_s}{\bar{m}} \frac{z_q}{z_s} \cdot \frac{\tan \phi_V}{\sin 2\phi_P} \right)^2 \cdot \left( \frac{p_{\eta'}}{p_\eta} \right)^3$$

**and the ratio:**  $\frac{\Gamma(\eta' \rightarrow \gamma \gamma)}{\Gamma(\pi^0 \rightarrow \gamma \gamma)} = \frac{1}{9} \left( \frac{m_{\eta'}}{m_\pi} \right)^3 \left( 5X_{\eta'} + \sqrt{2} \frac{f_q}{f_s} Y_{\eta'} \right)^2$  E. Kou, Phys. Rev. D 63 (2001) 54027





# Fit redone leaving free all parameters

- 1) Leave the z's parameter free;
- 2) Add more constraints (needed to perform the fit with larger number of parameters);
- 3) Check the contribution from  $\eta' \rightarrow \gamma\gamma$  /  $\pi^0 \rightarrow \gamma\gamma$

$$\frac{\Gamma(\omega \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{1}{9} \left[ z_q \cos(\phi_p) - 2 \frac{m_s}{\bar{m}} z_s \tan(\phi_V) \sin(\phi_p) \right]^2 (1 - z_G^2) \left( \frac{m_\omega^2 - m_\eta^2}{m_\omega^2 - m_{\pi^0}^2} \right)^3$$

$$\frac{\Gamma(\rho \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = z_q^2 \frac{\cos^2(\phi_p)}{\cos^2(\phi_V)} \left( \frac{m_\rho^2 - m_\eta^2}{m_\omega^2 - m_\pi^2} \frac{m_\omega}{m_\rho} \right)^3$$

$$\frac{\Gamma(\phi \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{1}{9} \left[ z_q \tan(\phi_V) \cos(\phi_p) + 2 \frac{\bar{m}}{m_s} z_s \sin(\phi_p) \right]^2 \left( \frac{m_\phi^2 - m_\eta^2}{m_\omega^2 - m_\pi^2} \frac{m_\omega}{m_\phi} \right)^3$$

$$\frac{\Gamma(\phi \rightarrow \pi^0 \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \tan^2 \phi_V \cdot \left( \frac{m_\phi^2 - m_{\pi^0}^2}{m_\omega^2 - m_{\pi^0}^2} \cdot \frac{m_\omega}{m_\phi} \right)^3, \quad \frac{\Gamma(K^{*+} \rightarrow K^+ \gamma)}{\Gamma(K^{*0} \rightarrow K^0 \gamma)} = \left( \frac{2 \frac{m_s}{\bar{m}} - 1}{1 + \frac{m_s}{\bar{m}}} \right)^2 \cdot \left( \frac{m_{K^{*+}}^2 - m_{K^0}^2}{m_{K^{*0}}^2 - m_{K^0}^2} \cdot \frac{m_{K^0}}{m_{K^{*+}}} \right)^3$$



The experimental covariance matrix **B** contains correlation among common used quantities in the fitted relations:

$$\frac{\Gamma(\omega \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)}, \frac{\Gamma(\rho \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)}, \frac{\Gamma(\phi \rightarrow \eta \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} \longrightarrow \text{Introduces a correlation in the fitted quantities}$$

$$\frac{\Gamma(\eta' \rightarrow \gamma \gamma)}{\Gamma(\pi^0 \rightarrow \gamma \gamma)} = \frac{Br(\eta' \rightarrow \gamma \gamma) \Gamma_{\eta'}}{\Gamma(\pi^0 \rightarrow \gamma \gamma)} \quad \frac{\Gamma(\eta' \rightarrow \rho \gamma)}{\Gamma(\omega \rightarrow \pi^0 \gamma)} = \frac{Br(\eta' \rightarrow \rho \gamma) \Gamma_{\eta'}}{\Gamma(\omega \rightarrow \pi^0 \gamma)}$$

$x_2$	-34					
$x_3$	-78	-29				
$x_4$	-35	-24	32			
$x_5$	-26	-12	26	8		
$x_6$	-28	-11	35	11	9	
$\Gamma$	32	-2	-24	-5	-88	-8
	$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$

**Br and  $\Gamma$  strongly correlated (above all  $\Gamma(\eta' \rightarrow \gamma \gamma)$ )**

**the  $\Gamma$  is measured using:**

$$e^+e^- \rightarrow \eta' e^+e^-$$

**An independent measurement of the  $\eta'$  total width is welcome.**

Mode	Rate (MeV)	Scale factor
$\Gamma_1$ $\pi^+ \pi^- \eta$	0.090 $\pm$ 0.008	1.2
$\Gamma_2$ $\rho^0 \gamma$ (including non-resonant $\pi^+ \pi^- \gamma$ )	0.060 $\pm$ 0.005	1.2
$\Gamma_3$ $\pi^0 \pi^0 \eta$	0.042 $\pm$ 0.004	1.6
$\Gamma_4$ $\omega \gamma$	0.0062 $\pm$ 0.0008	1.2
$\Gamma_5$ $\gamma \gamma$	0.00430 $\pm$ 0.00015	1.1
$\Gamma_6$ $3\pi^0$	(3.2 $\pm$ 0.6) $\times 10^{-4}$	1.1



# Fit results without $\eta' \rightarrow \gamma\gamma / \pi^0 \rightarrow \gamma\gamma$

$$\Gamma(P \rightarrow V \gamma) = \frac{g^2}{4\pi} |p_\gamma|^3$$

Fit redone a la Escribano

(using couplings and not taking into account correlations)

very similar results

	Fit with width ratios	Escribano <i>et al.</i> , JHEP 0705:006 (2007)	Fit with couplings
$\chi^2 / n.d.f(Prob)$	1.8/2 (41 %)	4.2/4 (38 %)	4.7/4 (32 %)
$Z_G^2$	$0.03 \pm 0.06$	$0.04 \pm 0.09$	$0.04 \pm 0.07$
$\varphi_G$	$(10 \pm 10)^\circ$	$(12 \pm 13)^\circ$	$(11 \pm 11)^\circ$
$\varphi_P$	$(41.6 \pm 0.8)^\circ$	$(41.4 \pm 1.3)^\circ$	$(41.5 \pm 1.1)^\circ$
$Z_{NS}$	$0.85 \pm 0.03$	$0.86 \pm 0.03$	$0.86 \pm 0.03$
$Z_S$	$0.78 \pm 0.05$	$0.79 \pm 0.05$	$0.78 \pm 0.05$
$\varphi_V$	$(3.16 \pm 0.10)^\circ$	$(3.2 \pm 0.1)^\circ$	$(3.18 \pm 0.10)^\circ$
$m_s/\bar{m}$	$1.24 \pm 0.07$	$1.24 \pm 0.07$	$1.24 \pm 0.07$
disappear in the ratio $Z_K$		$0.89 \pm 0.03$	$0.89 \pm 0.03$
$g$		$0.72 \pm 0.01$	$0.72 \pm 0.01$

Table 3: Comparison among the fit results without the  $\eta' \rightarrow \gamma\gamma / \pi^0 \rightarrow \gamma\gamma$  measurement and the Escribano *et al.* results.



# Adding the $\eta' \rightarrow \gamma\gamma / \pi^0 \rightarrow \gamma\gamma$ constraint

	Glueonium allowed	Glueonium at zero
$\chi^2/n.d.f(Prob)$	5/3 (17.5 %)	13/4 (1.1 %)
$Z_G^2$	$0.105 \pm 0.037$	0 fixed
$\varphi_P$	$(40.7 \pm 0.7)^\circ$	$(41.6 \pm 0.5)^\circ$
$Z_{NS}$	$0.866 \pm 0.025$	$0.863 \pm 0.024$
$Z_S$	$0.79 \pm 0.05$	$0.78 \pm 0.05$
$\varphi_V$	$(3.15 \pm 0.10)^\circ$	$(3.17 \pm 0.10)^\circ$
$m_s/\bar{m}$	$1.24 \pm 0.07$	$1.24 \pm 0.07$

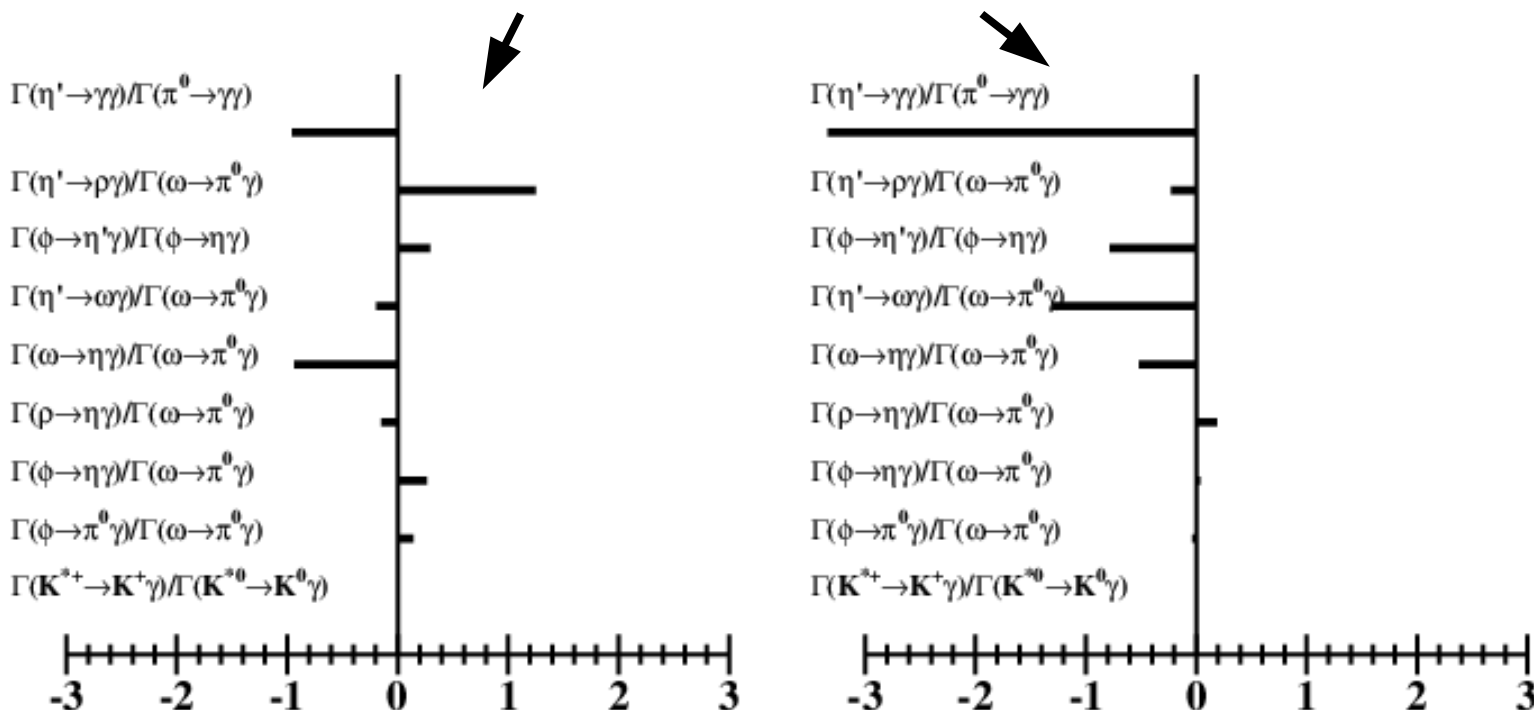
KLOE  
(Phys. Lett. B648 (2007) 267)

$$\phi_P = (39.7 \pm 0.7_{tot})^\circ$$

$$|\phi_G| = (22 \pm 3)^\circ$$

$$\sin^2\phi_G = Z^2 = 0.14 \pm 0.04$$

Fit pulls





# Update using PDG 2008



# Result of the fit

	Glunium allowed	Glunium at zero
$\chi^2/n.d.f(Prob)$	7.9/3 (5 %)	15/4 ( $5 \times 10^{-3}$ )
$Z_G^2$	$0.097 \pm 0.037$	0 fixed
$\varphi_P$	$(41.0 \pm 0.7)^\circ$	$(41.7 \pm 0.5)^\circ$
$Z_{NS}$	$0.86 \pm 0.02$	$0.858 \pm 0.021$
$Z_S$	$0.79 \pm 0.05$	$0.78 \pm 0.05$
$\varphi_V$	$(3.17 \pm 0.09)^\circ$	$(3.19 \pm 0.09)^\circ$
$m_s/\bar{m}$	$1.24 \pm 0.07$	$1.24 \pm 0.07$

**PDG08**

**The same glunium content but unsatisfying fit quality.**

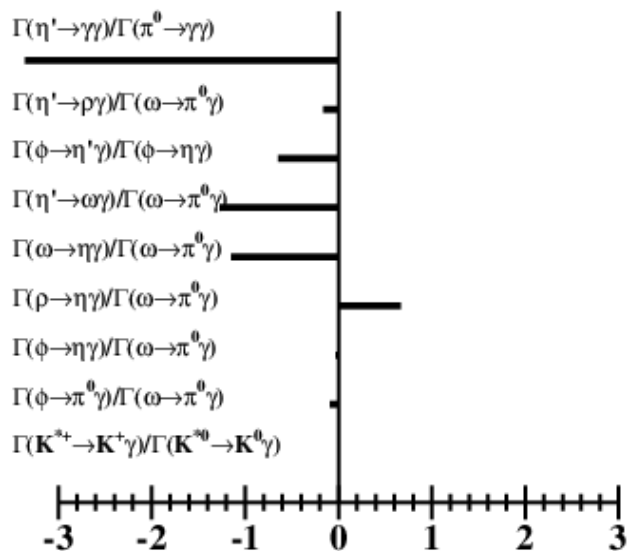
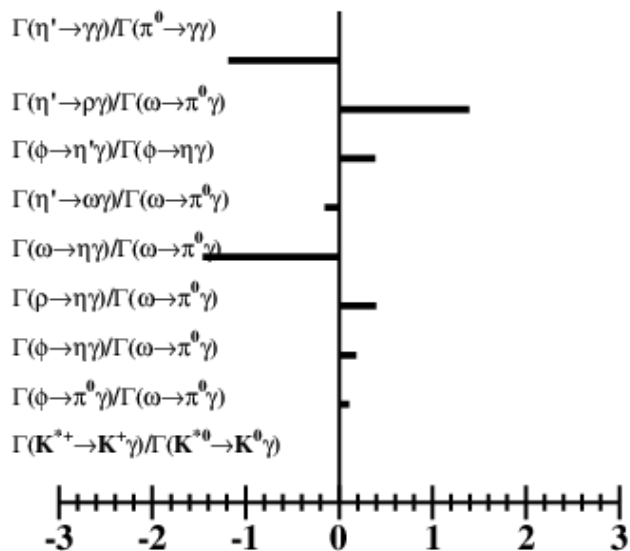
	Glunium allowed	Glunium at zero
$\chi^2/n.d.f(Prob)$	5/3 (17.5 %)	13/4 (1.1 %)
$Z_G^2$	$0.105 \pm 0.037$	0 fixed
$\varphi_P$	$(40.7 \pm 0.7)^\circ$	$(41.6 \pm 0.5)^\circ$
$Z_{NS}$	$0.866 \pm 0.025$	$0.863 \pm 0.024$
$Z_S$	$0.79 \pm 0.05$	$0.78 \pm 0.05$
$\varphi_V$	$(3.15 \pm 0.10)^\circ$	$(3.17 \pm 0.10)^\circ$
$m_s/\bar{m}$	$1.24 \pm 0.07$	$1.24 \pm 0.07$

**PDG06**

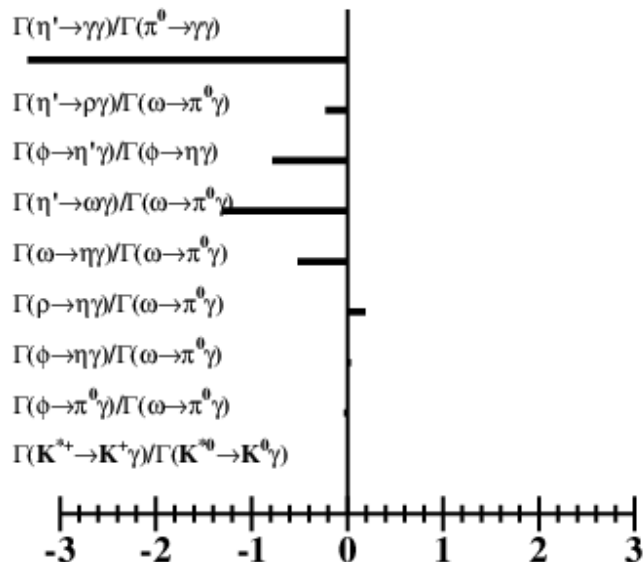
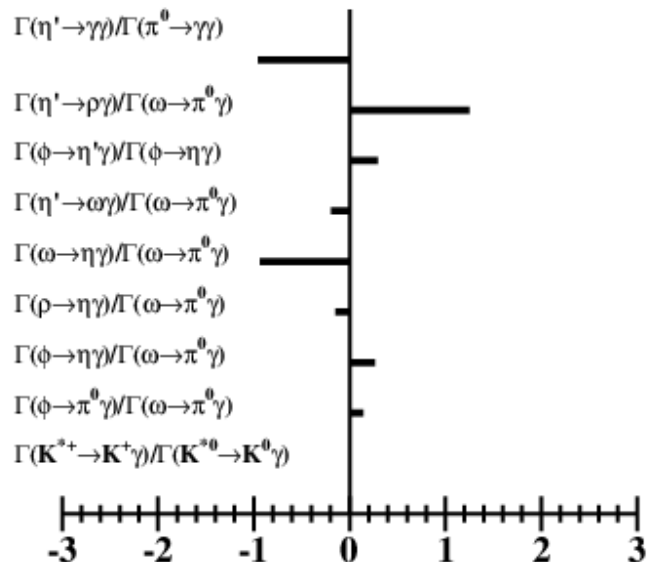
Table 2: Fit results.



# Pulls of the fit



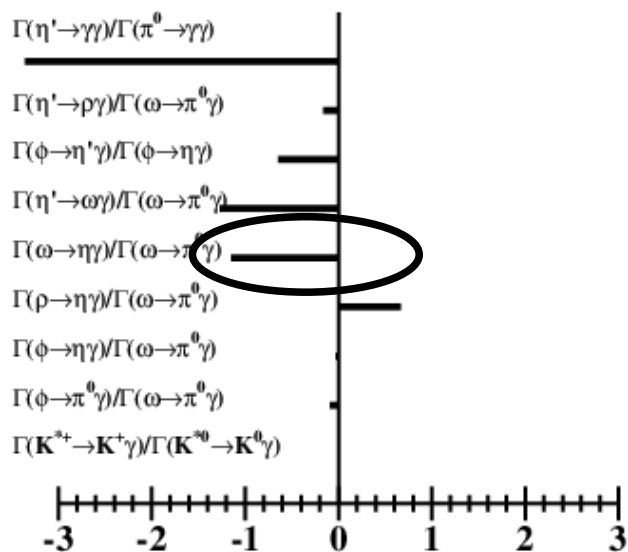
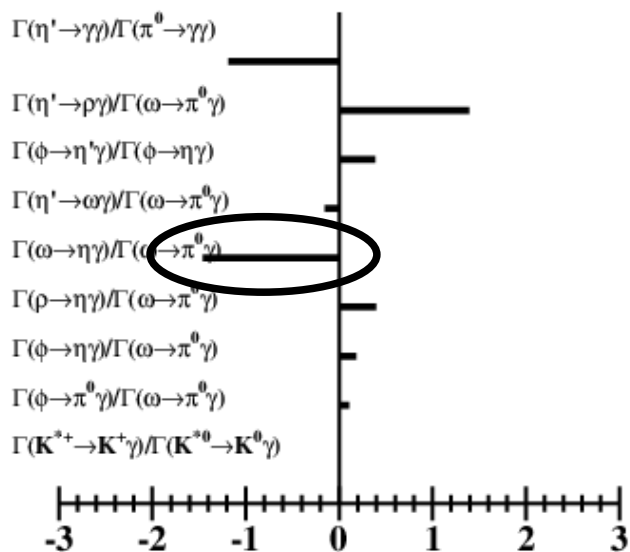
PDG08



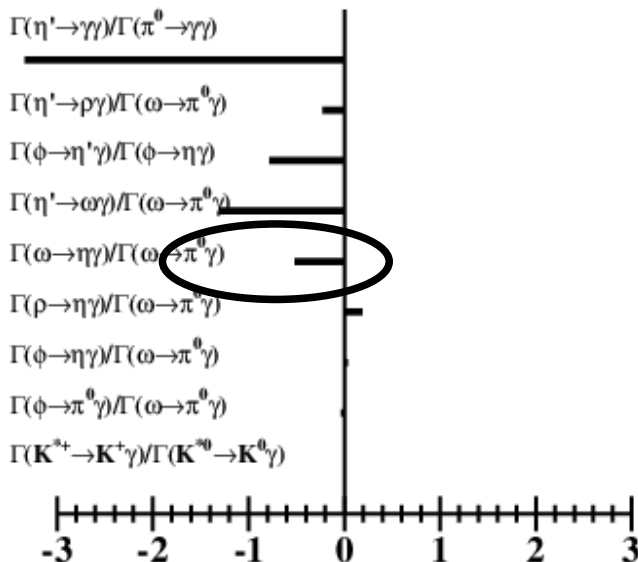
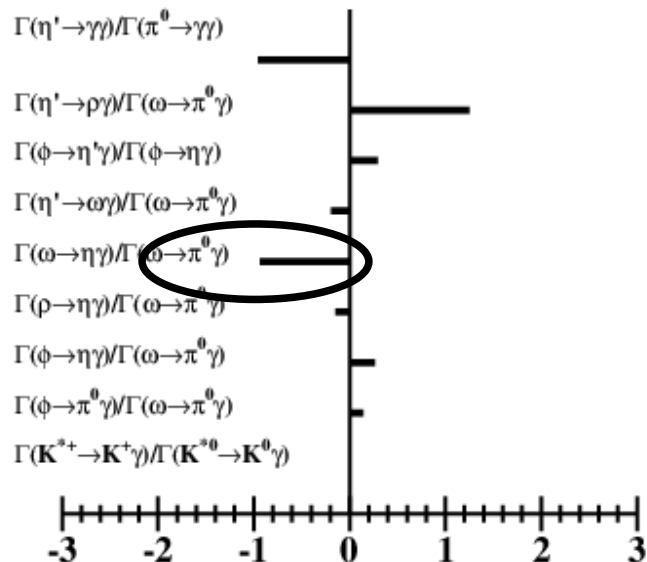
PDG06



# Pulls of the fit



PDG08



PDG06

$\omega \rightarrow \eta\gamma$  pull has increased in both gluonium hypothesis





# $\omega \rightarrow \eta\gamma$ measurement from PDG

The  $\omega \rightarrow \eta\gamma$  partial width changed from

$$(4.9 \pm 0.5) \times 10^{-4} \text{ to } (4.6 \pm 0.4) \times 10^{-4}$$

This value is determined by the global PDG fit, and it is mainly determined by:

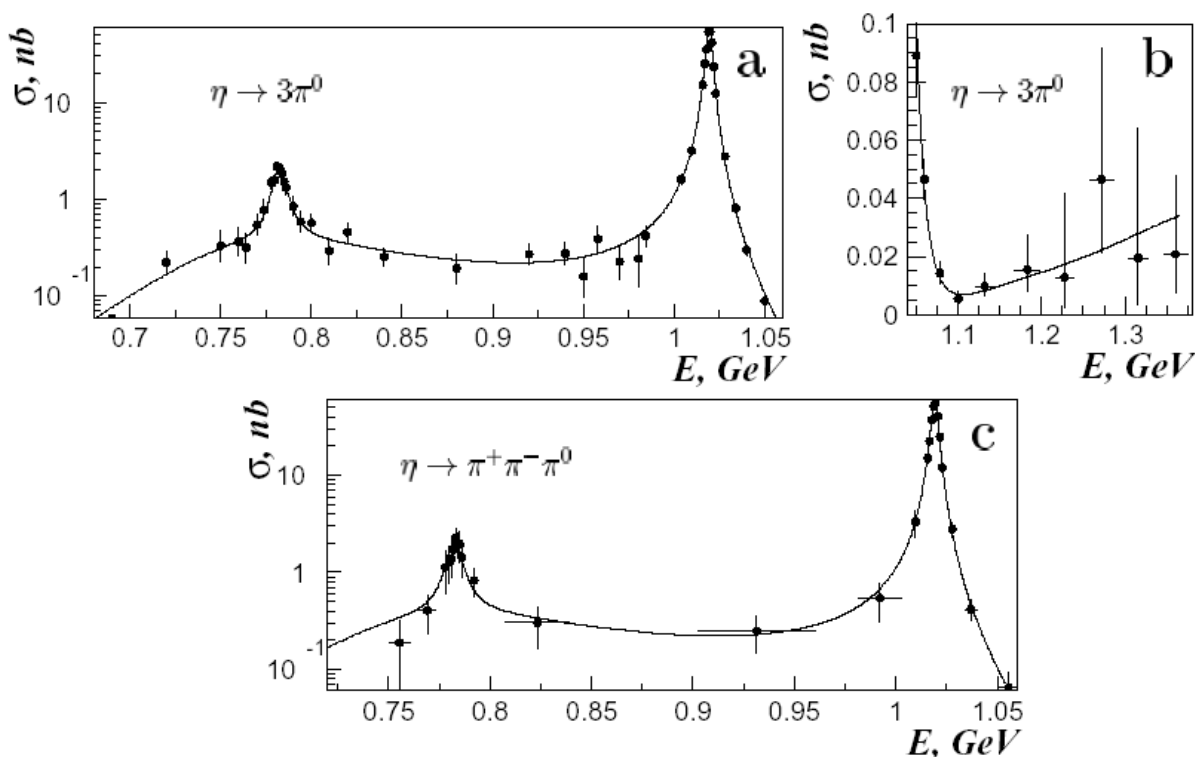
$\Gamma(e^+e^-) \times \Gamma(\eta\gamma) / \Gamma_{\text{total}}^2$					$\Gamma_9\Gamma_5 / \Gamma^2$
<u>VALUE (units <math>10^{-8}</math>)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b><math>3.31 \pm 0.28</math> OUR FIT</b>	Error includes scale factor of 1.1.				
<b><math>3.18 \pm 0.28</math> OUR AVERAGE</b>					
$3.10 \pm 0.31 \pm 0.11$	33k	<sup>24</sup> ACHASOV 07B	SND	0.6-1.38 $e^+e^- \rightarrow \eta\gamma$	↓
$3.17^{+1.85}_{-1.31} \pm 0.21$	17.4k	<sup>25</sup> AKHMETSHIN 05	CMD2	0.60-1.38 $e^+e^- \rightarrow \eta\gamma$	
$3.41 \pm 0.52 \pm 0.21$	23k	<sup>26,27</sup> AKHMETSHIN 01B	CMD2	$e^+e^- \rightarrow \eta\gamma$	

**ACHASOV 07B: Phys. Rev. D76 (2007) 077101**



# $\omega \rightarrow \eta\gamma$ branching ratio measurement from SND

The branching ratio is extracted with a global fit to the  $e^+e^- \rightarrow \eta\gamma$  with a VMD model with  $\rho, \omega, \phi, \rho'$  included ( $\rho'$  parameters varied to compute systematics and constrained from  $e^+e^- \rightarrow \eta\rho$ ).



$\omega$  contribution overwhelmed by the  $\rho \rightarrow \eta\gamma$  contribution  
no correlation matrix is given in the paper



# Direct $\omega \rightarrow \eta \gamma$ branching ratio measurement

The fit is dominated by the SND measurement

$$\Gamma(\eta\gamma)/\Gamma_{\text{total}}$$

VALUE (units $10^{-4}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>4.6 <math>\pm</math> 0.4</b>	<b>OUR FIT</b>	Error includes scale factor of 1.1.		
<b>6.3 <math>\pm</math> 1.3</b>	<b>OUR AVERAGE</b>	Error includes scale factor of 1.2.		
6.6 $\pm$ 1.7		<sup>53</sup> ABELE	97E CBAR	0.0 $\bar{p}p \rightarrow 5\gamma$
8.3 $\pm$ 2.1		ALDE	93 GAM2	38 $\pi^- p \rightarrow \omega n$
3.0 $^{+2.5}_{-1.8}$		<sup>54</sup> ANDREWS	77 CNTR	6.7–10 $\gamma\text{Cu}$

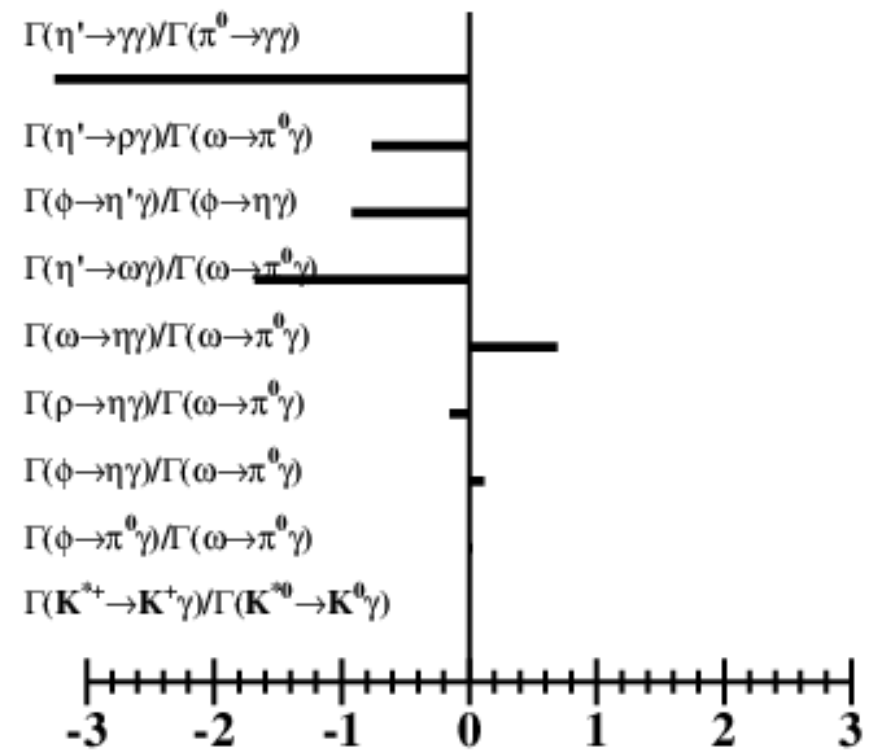
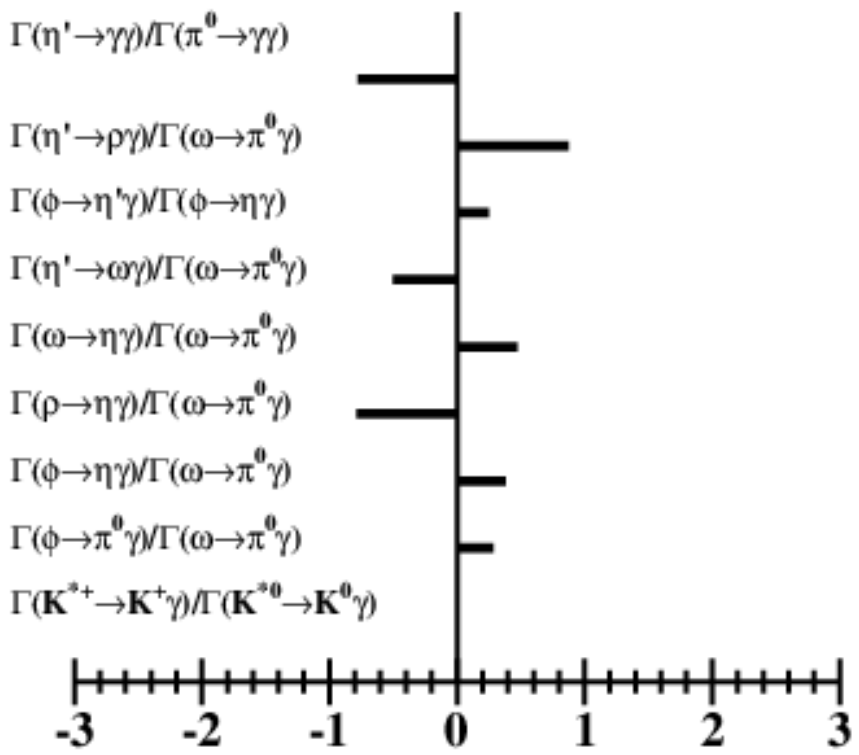
In Crystal Barrel the channel  $\bar{p}p \rightarrow \eta \omega$  is used that is 6 times larger than  $\bar{p}p \rightarrow \eta \rho$ , the  $\omega \rightarrow \eta \gamma$  Br was normalized to the  $\omega \rightarrow \pi^0 \gamma$  Br

using  $\omega \rightarrow \eta \gamma$  from PDG average

	Gluonium allowed	Gluonium at zero
$\chi^2/n.d.f(Prob)$	3.9/3 (27.5 %)	13/4 (1.1 %)
$Z_G^2$	0.111 $\pm$ 0.036	0 fixed
$\varphi_P$	(40.6 $\pm$ 0.7) $^\circ$	(41.5 $\pm$ 0.5) $^\circ$
$Z_{NS}$	0.890 $\pm$ 0.025	0.882 $\pm$ 0.023
$Z_S$	0.79 $\pm$ 0.05	0.78 $\pm$ 0.05
$\varphi_V$	(3.15 $\pm$ 0.10) $^\circ$	(3.18 $\pm$ 0.09) $^\circ$
$m_s/\bar{m}$	1.24 $\pm$ 0.07	1.24 $\pm$ 0.07



# Pulls





# Results using KLOE $\text{Br}(\omega \rightarrow \pi^0 \gamma)$

**KLOE  $\text{Br}(\omega \rightarrow \pi^0 \gamma) = 8.09 \pm 0.14 \%$**   
**PDG08  $= 8.92 \pm 0.24 \%$** 
 **$3\sigma$  away**

	Gluonium allowed	Gluonium at zero
$\chi^2/n.d.f(Prob)$	4.16/3 (25 %)	13/4 (1.1 %)
$Z_G^2$	$0.109 \pm 0.036$	0 fixed
$\varphi_P$	$(40.5 \pm 0.7)^\circ$	$(41.4 \pm 0.5)^\circ$
$Z_{NS}$	$0.935 \pm 0.025$	$0.926 \pm 0.023$
$Z_S$	$0.83 \pm 0.05$	$0.82 \pm 0.05$
$\varphi_V$	$(3.3 \pm 0.09)^\circ$	$(3.3 \pm 0.09)^\circ$
$m_s/\bar{m}$	$1.24 \pm 0.07$	$1.24 \pm 0.07$

**No big effect on  $Z_G$  but important contribution to  $Z_{S,NS}$  and  $\omega$ - $\phi$  mixing angle**



## Conclusion and Outlook (KLOE 2)

Memo almost ready, will be internally published next week for referee review (S. Giovanella & P. Gauzzi).

- 1) Measure  $\eta' \rightarrow \gamma\gamma$  branching ratio that is the only one pointing for the gluonium;
- 2) Improve  $\phi \rightarrow \eta' \gamma$  through the measurement of  $\eta' \rightarrow \pi\pi\eta$ ,  $\pi^+\pi^-\gamma$ ,  $\omega\gamma$ ,  $\gamma\gamma$ ,  $\pi^+\pi^-\pi^0$  (upper limit),  $n^0\pi$ ;



# Fit procedure.

The  $\chi^2$  is defined as follows:

$$\chi^2 = \sum_{i,j=1,3} (y_i - y_i^{th}) \times V_{ij}^{-1} (y_j - y_j^{th})$$

$V_{ij}$  is the error matrix which is a function of theoretical uncertainties, as well as the experimental

$$V_{ij} = [B_{ij} + (A_{ik} \times C_{kl} \times A_{lj}^T)]$$

Experimental  
covariance matrix

Theoretical parameters  
covariance matrix

$B_{ij}$  Full covariance matrix  
(correlation comes  
from the constrained  
fit to  $\eta'$  Br)

$$; A_{ik} = \begin{pmatrix} \frac{\partial y_1^{th}}{\partial f_s} & \frac{\partial y_1^{th}}{\partial f_q} & \frac{\partial y_1^{th}}{\partial C_{NS}} & \frac{\partial y_1^{th}}{\partial C_S} & \frac{\partial y_1^{th}}{\partial \frac{m_s}{m}} \\ \frac{\partial y_2^{th}}{\partial f_s} & \frac{\partial y_2^{th}}{\partial f_q} & \frac{\partial y_2^{th}}{\partial C_{NS}} & \frac{\partial y_2^{th}}{\partial C_S} & \frac{\partial y_2^{th}}{\partial \frac{m_s}{m}} \\ \frac{\partial y_3^{th}}{\partial f_s} & \frac{\partial y_3^{th}}{\partial f_q} & \frac{\partial y_3^{th}}{\partial C_{NS}} & \frac{\partial y_3^{th}}{\partial C_S} & \frac{\partial y_3^{th}}{\partial \frac{m_s}{m}} \end{pmatrix}$$

$$C_{kl} = \begin{pmatrix} \sigma_{f_q}^2 & 0 \\ 0 & \sigma_{f_s}^2 \end{pmatrix}$$

Re-evaluated at  
each minimization  
step