

Evaluation of systematics for the analysis of the $\pi^0\pi^0\gamma$ final state

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- ✓ Introduction and analysis scheme
- ✓ Acceptance issues
- ✓ Photon pairing
- ✓ Kinematic fitting and ISR
- ✓ Evaluation of the background
- ✓ Fit procedure
- ✓ Prospects and plans

ϕ decays meeting – 9 May 2005

Composition of the $\pi^0\pi^0\gamma$ final state

Two main contributions to $\pi^0\pi^0\gamma$ final state @ M_ϕ :

1. $e^+e^- \rightarrow \omega\pi^0 \rightarrow \pi^0\pi^0\gamma$

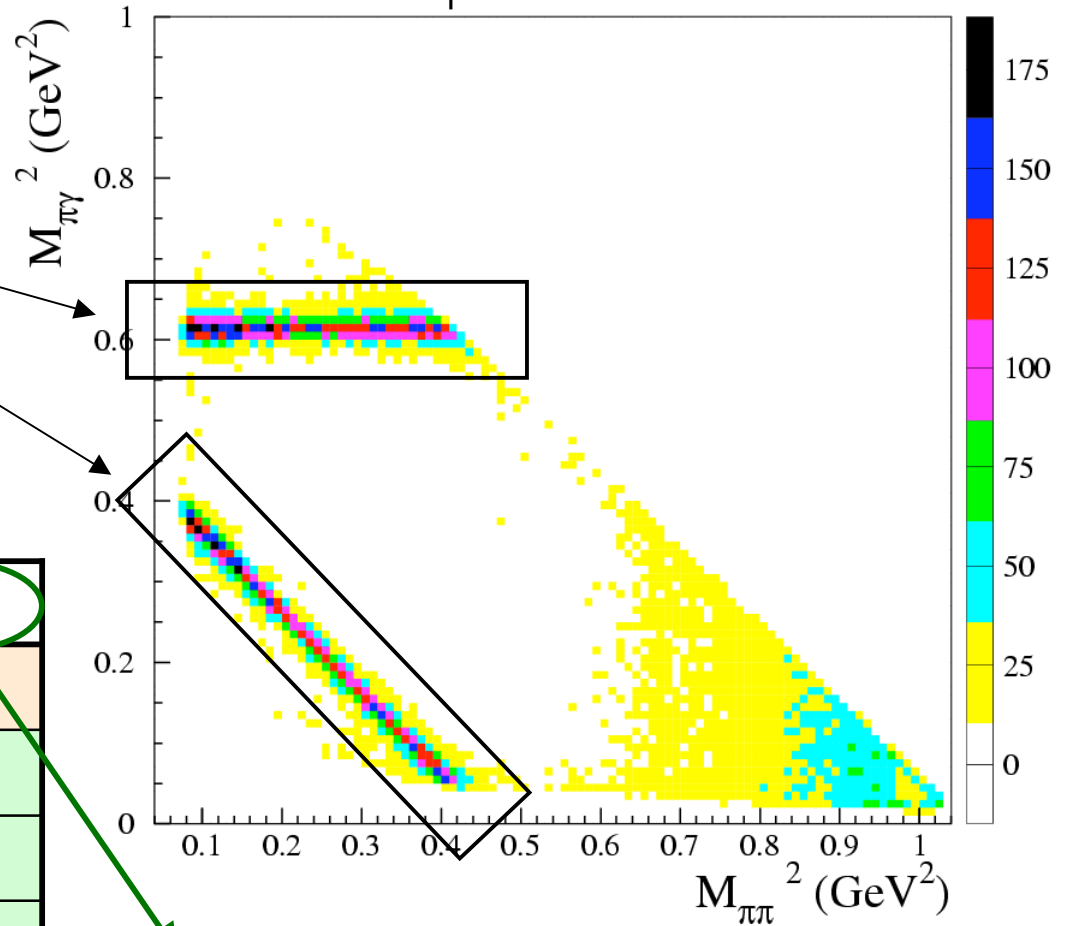
$\sigma_{\text{vis}}(M_\phi) \sim 0.5 \text{ nb}$

2. $\phi \rightarrow S\gamma \rightarrow \pi^0\pi^0\gamma$

$\sigma_{\text{vis}}(M_\phi) \sim 0.3 \text{ nb}$

Backgrounds:

Process	S/B
$\phi \rightarrow a_0\gamma \rightarrow \eta\pi^0\gamma \rightarrow \gamma\gamma\pi^0\gamma$	8.51
$\phi \rightarrow \eta\gamma \rightarrow \pi^0\pi^0\pi^0\gamma$	0.06
$\phi \rightarrow \eta\gamma \rightarrow \gamma\gamma\gamma$	0.05
$\phi \rightarrow \pi^0\gamma$	0.21
$e^+e^- \rightarrow \gamma\gamma(\gamma)$	0.002



$S = \omega\pi + S\gamma$

Data and Monte Carlo samples

DATA

2001+2002 data : $L_{\text{int}} = 450 \text{ pb}^{-1}$

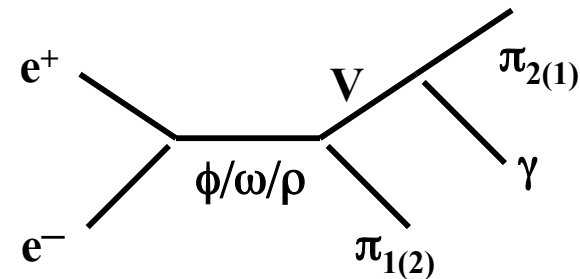
Data have some spread around the ϕ peak + two dedicated off-peak runs @ 1017 and 1022 MeV \Rightarrow **Data divided in 100 keV bins of \sqrt{s}**

Today we will discuss only about the 145 pb^{-1} at 1019.6 MeV

MC

RAD04 MC production: $5 \times L_{\text{int}}$

GG04 MC production: $1 \times L_{\text{int}}$



Improved $e^+e^- \rightarrow \omega\pi^0 \rightarrow \pi^0\pi^0\gamma$ generator Three body phase space according to VDM from NPB 569 (2000), 158

Sample preselection and kinematic fit

1. **Acceptance** cut:

5 neutral clusters in TW with $E > 7$ MeV and $|\cos\theta| < 0.92$

[TW: $|T_{cl} - R_{cl}/c| < \text{MIN}(5\sigma_T, 2 \text{ ns})$]

2. **Kinematic fit** requiring 4-momentum conservation and the “promptness” of γ 's ($T_{cl} - R_{cl}/c = 0$)

3. **Pairing**: best γ 's comb. for the $\pi^0\pi^0\gamma$ hypothesis

4. **Kinematic fit** for both γ 's pairing, requiring also constraints on π masses of the assigned $\gamma\gamma$ pairs

Study on Trigger, FILFO and ECL

1. Trigger and Cosmic Veto efficiency

Calorimeter trigger fully efficient on the signal. Cosmic Veto losses evaluated with prescaled events.

$$\varepsilon_{\text{CV}} = (99.54 \pm 0.08) \%$$

2. MC evaluation of FILFO and ECL losses

$$\varepsilon_{\text{FLF}} = (99.95 \pm 0.01) \% \quad \varepsilon_{\text{ECL}} = (96.5 \pm 0.1) \%$$

3. DATA evaluation of FILFO and ECL losses

The minimum bias sample streamed by C.DiDonato was used to evaluate with data ε_{FLF} and ε_{ECL} Only data with $\sqrt{s} = 1019.6$ used.

$$\varepsilon_{\text{FLF}} = (99.90 \pm 0.05) \% \quad \varepsilon_{\text{ECL}} = (99.2 \pm 0.1) \% !!$$

The large difference on ε_{ECL} mainly due to the a wrong parametrization of time res in neurad code! An overall correction $R=1.02 \pm 0.01$ applied to $\varepsilon(\text{MC})$

Study on the acceptance (I)

1. Effect of the wrong energy scale in MC

As for other analyses, the MC energy scale is shifted of $\sim 1.4\%$. We have corrected the MC energy of this amount and counted again the number of prompt photons. The **related systematics is ~ 0.3 of the statistical error** on the efficiency.

2. Systematics on cluster efficiency

We have varied the data-MC efficiency curve following the 3 search cones used in the calibration with $\phi \rightarrow \pi^+\pi^-\pi^0$.

The **related systematics is ~ 0.5 of the statistical error** on the efficiency.

WARNING:

The new “softer” data/MC efficiency curves obtained by Tommaso and Matteo may have some impact in the spectrum at low energy (below 70 MeV). We are planning to evaluate the systematics applying these new curves!

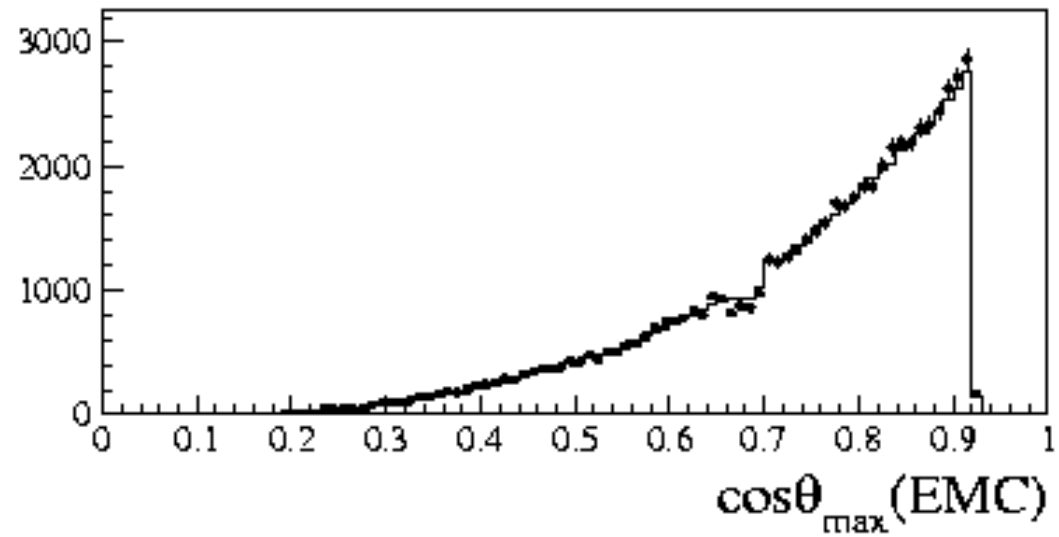
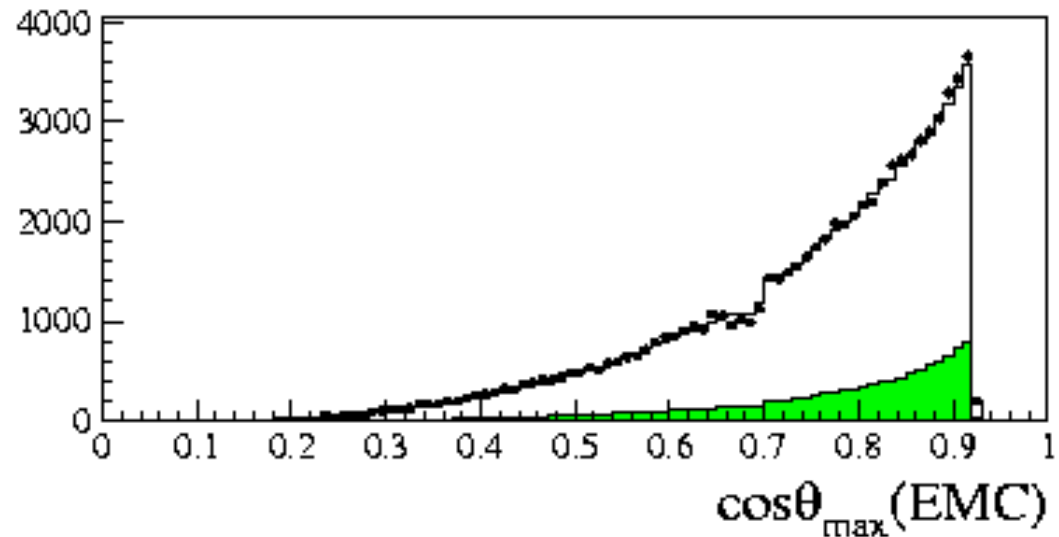
Study on the acceptance (II)

3. Angular acceptance

Calorimeter resolution on angular position has very high precision.

To see how well we define our acceptance we compare data and MC.

Agreement is excellent.



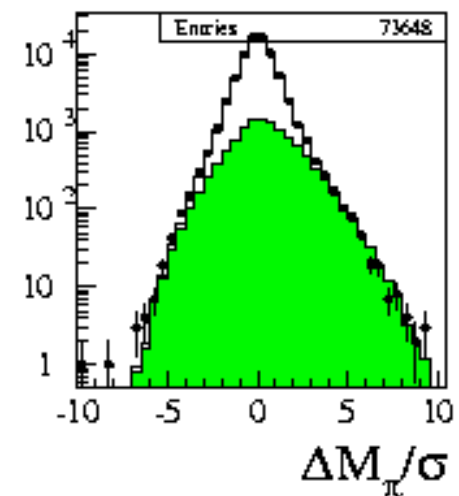
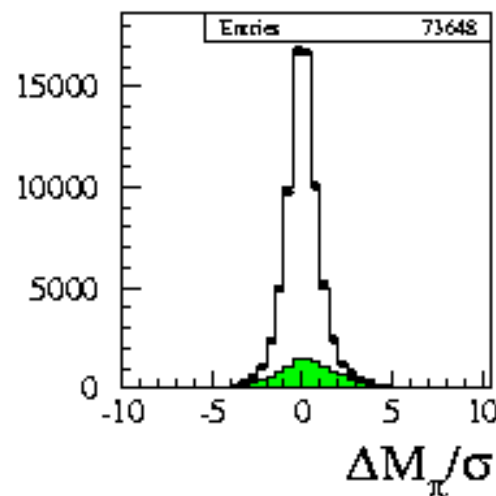
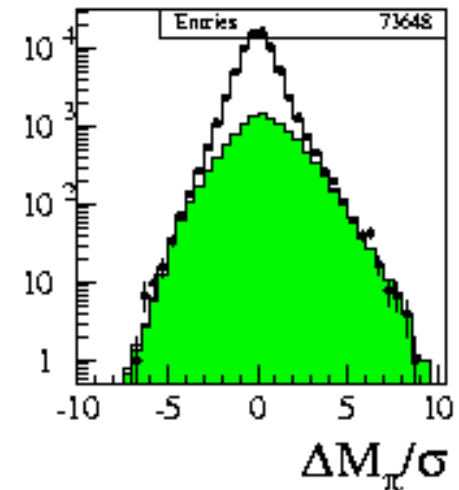
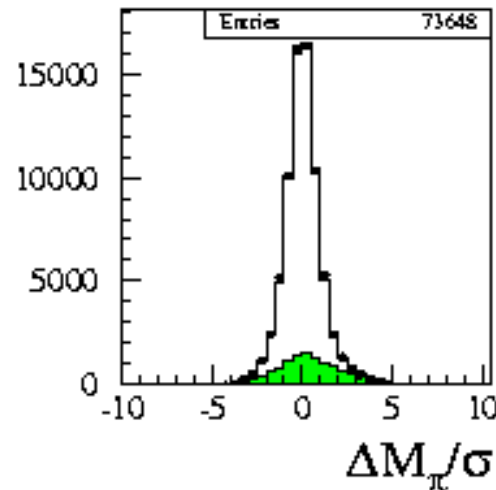
Photon pairing efficiency (I)

After the first kinematic fit we pair photons into π^0 depending upon the expected energy resolution and minimizing a χ^2_{SEL} estimator. Good agreement is observed between data and MC in $\Delta M/\sigma$

$$\Delta M/s = (M_{\gamma\gamma} - M_{\pi^0})/\sigma$$

Resolution is evaluated with same parametrization of energy resolution (after first kinematic fit) used in pairing procedure.

We keep only events with $|\Delta M/\sigma| < 5$

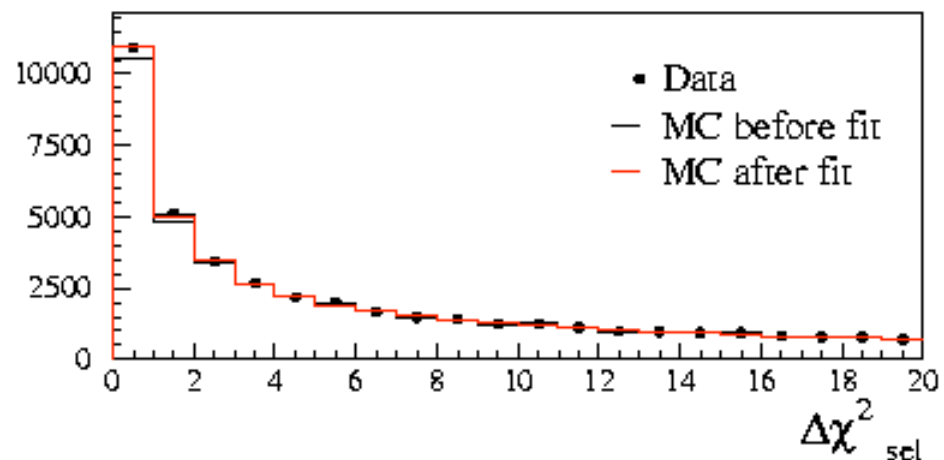
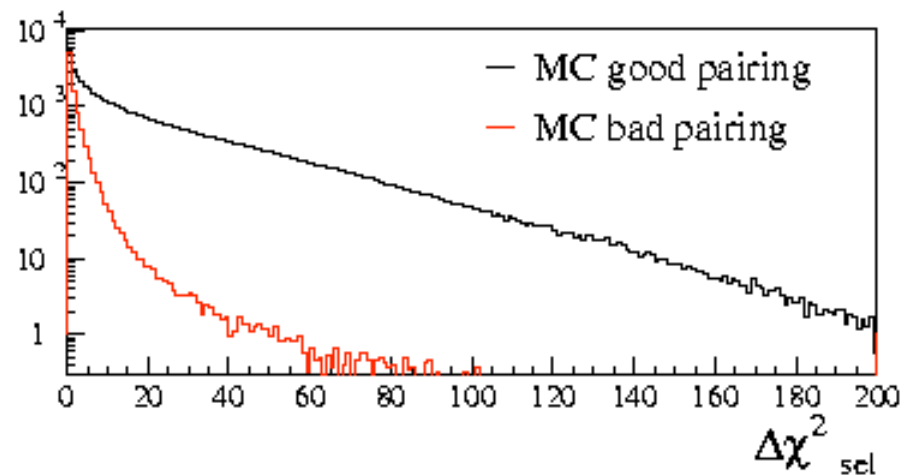


Photon pairing efficiency (II)

- ✓ To assign a systematic to the pairing procedure we studied the difference between χ^2_{SEL} for the best and second best choice of photons: $\Delta\chi^2_{\text{SEL}} = \chi^2_{\text{SEL}}(\text{Best}) - \chi^2_{\text{SEL}}(\text{SecondBest})$

- ✓ We then fit the $\Delta\chi^2_{\text{SEL}}$ distribution in data with a linear combination of MC spectra for the right and wrong choice of paired photons (by MC truth).

$$R_{\text{pair}} = 1.08 \pm 0.02$$



Second Kinematic fit and related problems

- ✓ After pairing photons a **second kinematic fit is performed constraining on M_π (11C)**.
- ✓ The data-MC comparison of the χ^2 distribution is excellent (cumulative distributions differ less than 1% at the analysis-cut value) ... but

we have problems in fitting the Dalitz when applying a tight χ^2 cut



Indeed, while testing the fit on the Dalitz we realized that :

- when leaving free Γ_ω and VDM couplings the width tends to grow with a correlated increase of the couplings.
- The χ^2 of the fit nicely improves!

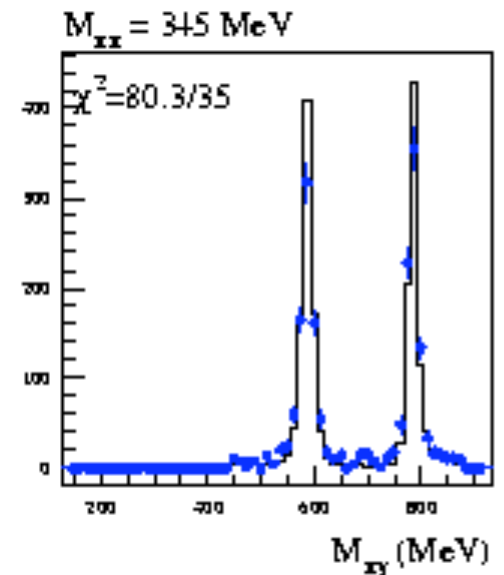
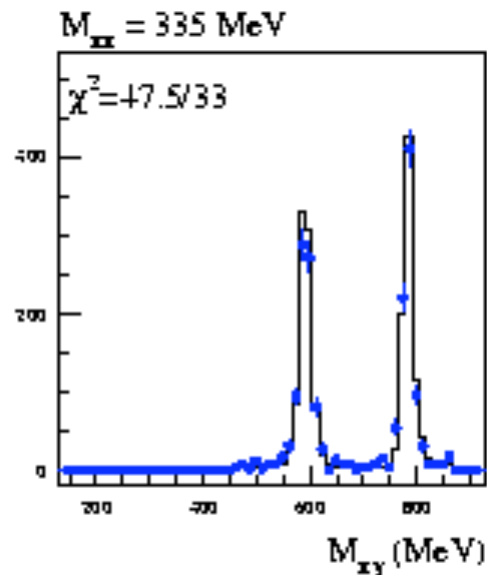
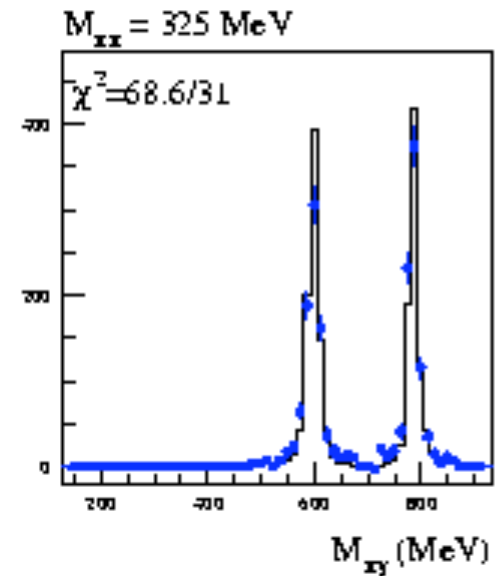
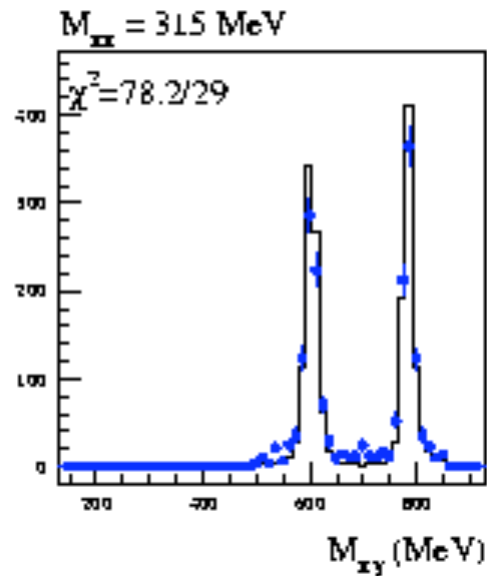
The fit problem: a growing Γ_ω ??

A flavour of the Dalitz regions where the fit fails and the χ^2 jumps up.



A too large Γ_ω cannot be a realistic origin/solution of the problem:

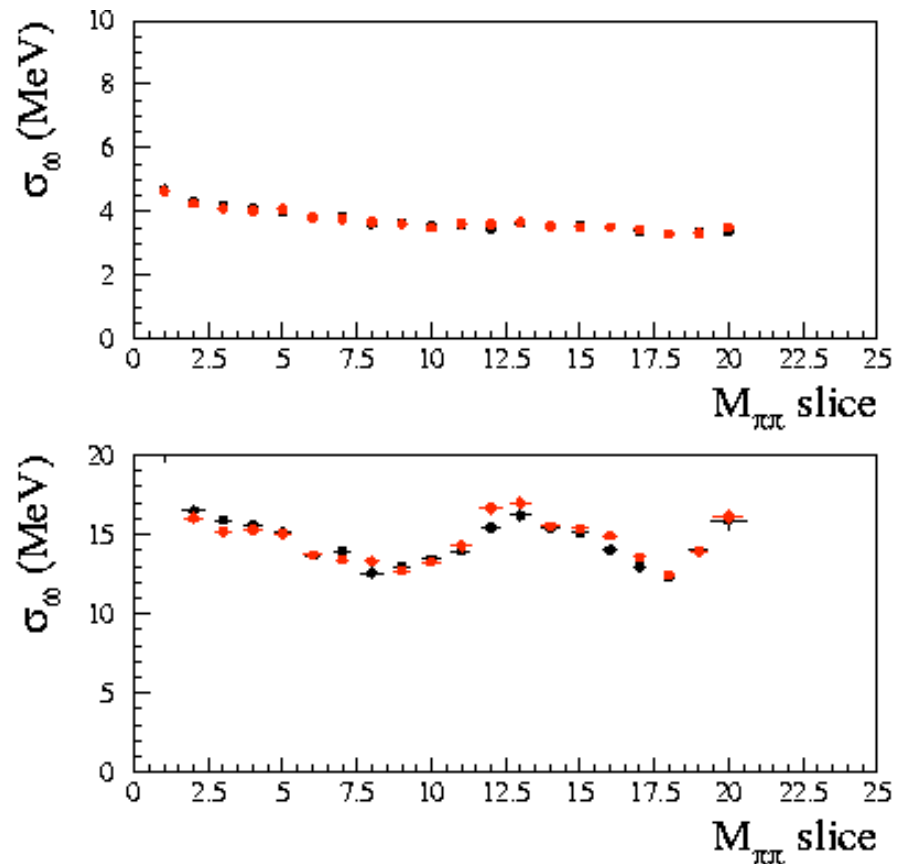
- Γ_ω is well known
- our own measurement of $\omega \rightarrow \pi\pi\pi$ agrees with PDG.
- even inserting in the propagator a $\Gamma_\omega(s)$ does not help.



The fit problem: different data-MC resolution??

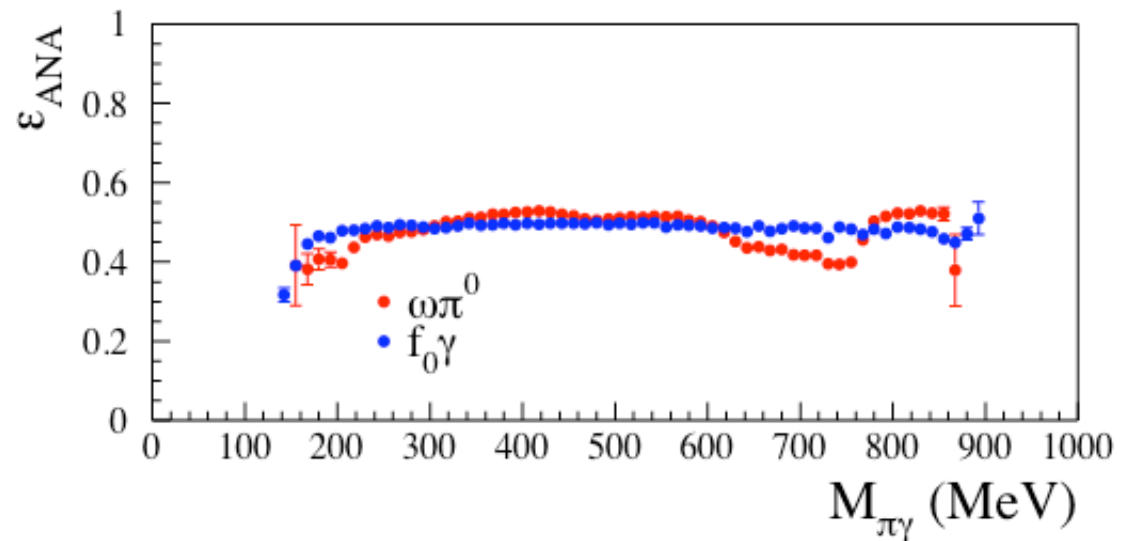
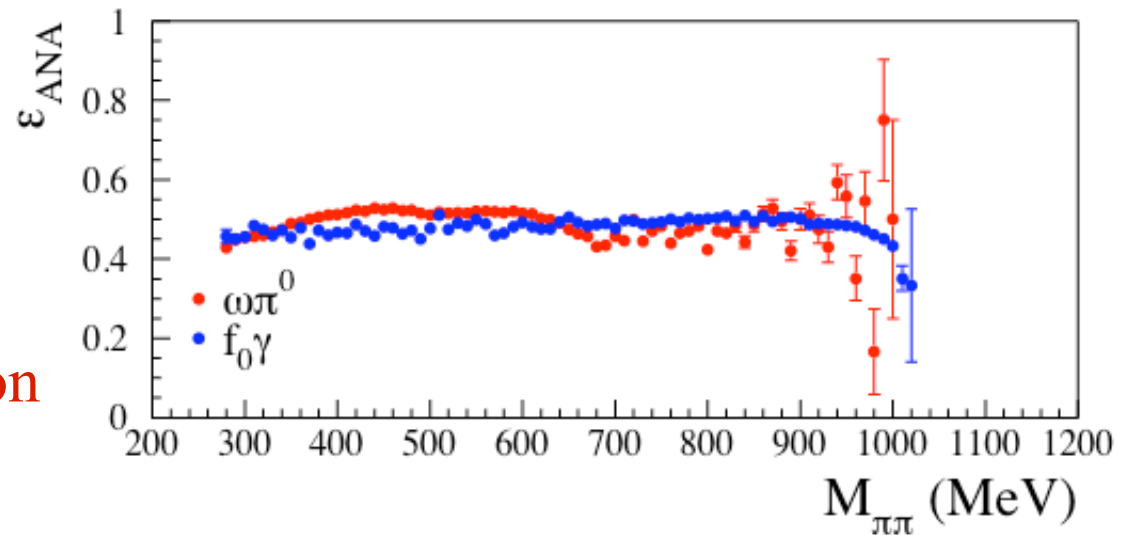
The mass resolution after the second kinematic fit cannot be the origin of such an apparent enlargement on Γ_ω . To simulate the effect of $\Gamma_\omega = 11$ MeV we should add in the MC an additional source of resolution able to contribute for ≈ 7 MeV. The MC resolution on $M_{\pi\gamma}$ shows a core of 3-4 MeV (F=0.66) and tails of 12-15 MeV (F=0.33).

From the agreement on ΔM
we consider this effect negligible



The fit problem: dependence of ϵ_{ana} on the dalitz

- By looking at the behaviour of ϵ_{ana} in the Dalitz we recall that:
 - The dependence of ϵ_{ana} for the $S\gamma$ process is flat
 - ϵ_{ana} for the VDM process shows instead fast variation along the plane.



Relation between ISR and ϵ_{ana} (I)

We searched for the origin of the difference on ϵ_{ana} dependence

- We found that this is due to the different spectrum of ISR photons between the $S\gamma$ and VDM processes.
- The resonant behaviour of $S\gamma$ originates large (25%) radiative corrections with an ISR energy spectrum constrained by $\Gamma\phi$ to be below 10 MeV.
- The not-resonant behaviour of VDM makes small the radiative correction while the ISR energy spectrum shows large tails.

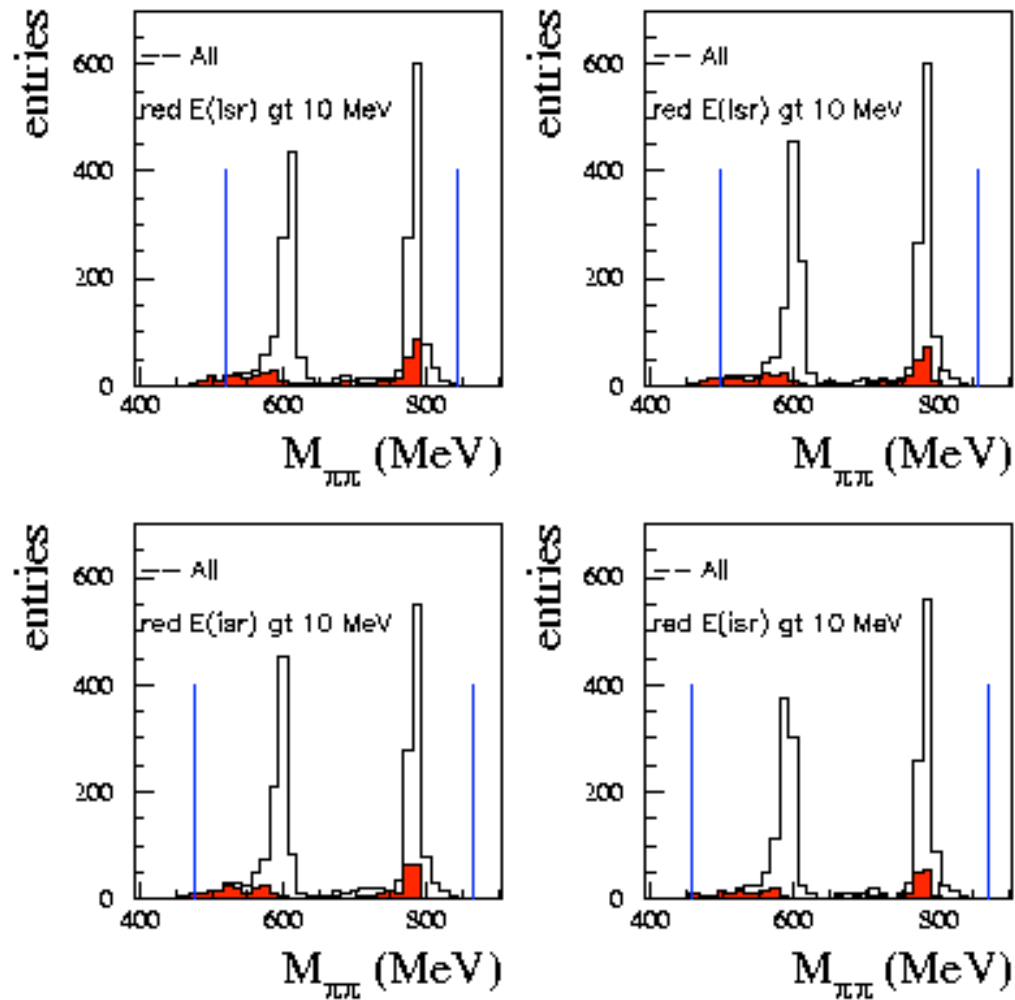
Our kinematic fit constrains the energy sum to the beam energy with an error of 300 keV (BES).

Require a tight χ^2_{FIT} cut implies rejection of the VDM events with large ISR tails.

Relation between ISR and ϵ_{ana}

(II)

The MC distribution of VDM events with an ISR photon >10 MeV is concentrated on the low- $M_{\pi\gamma}$ tails of the omega peaks.



Trying to handle the ISR problem

Two different roads to attack / understand this point :

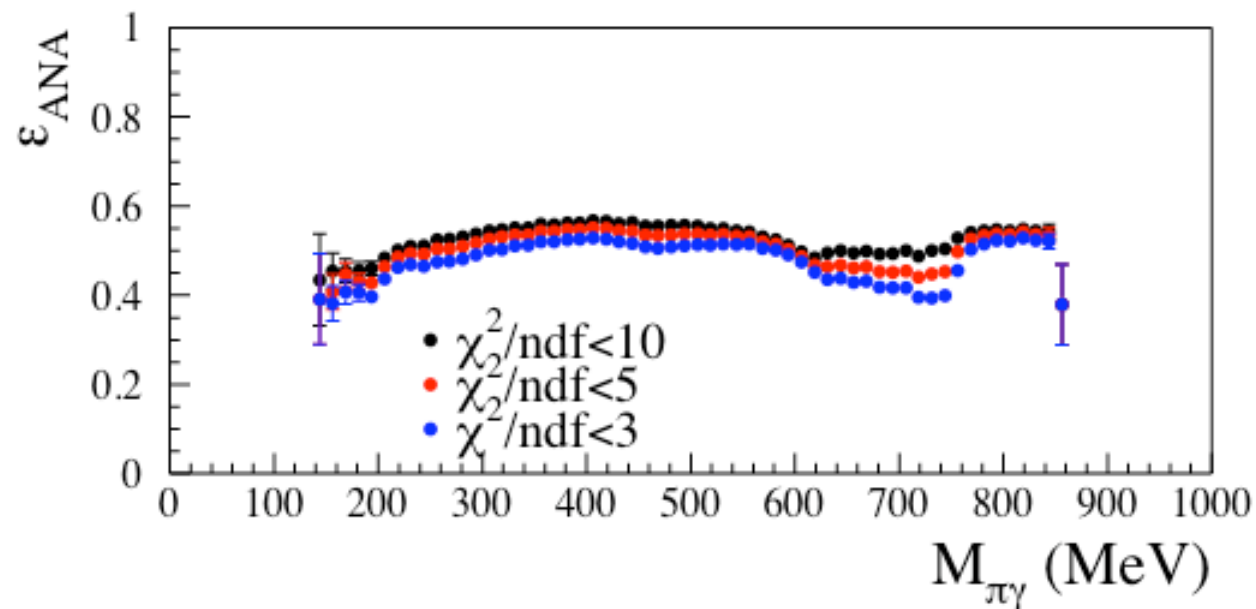
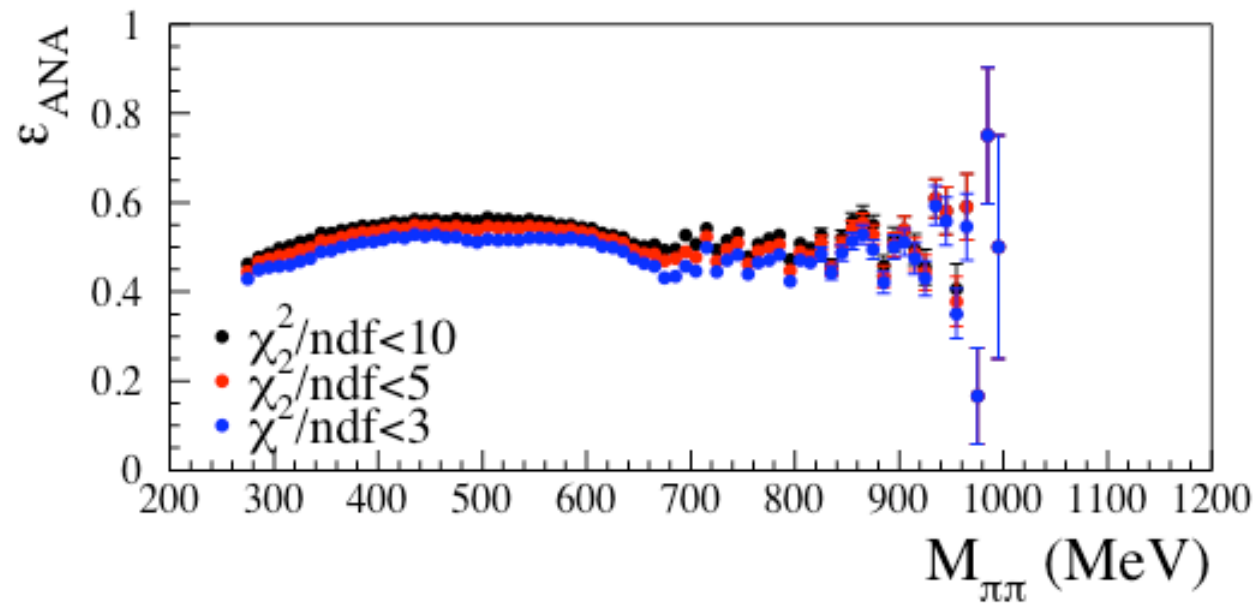
1) Relax the cut on $\chi^2_{\text{FIT}}/\text{ndof}$. Moving it from 3 to 5, 10

the efficiency gets flatter but the background goes up

Process	ϵ_{ana} ($\chi^2 < 3$)	S/B ($\chi^2 < 3$)	S/B ($\chi^2 < 5$)	S/B ($\chi^2 < 10$)
$e^+e^- \rightarrow \omega\pi^0 \rightarrow \pi^0\pi^0\gamma$	50.2 %			—
$\phi \rightarrow S\gamma \rightarrow \pi^0\pi^0\gamma$	48.7 %			—
$\phi \rightarrow a_0\gamma \rightarrow \eta\pi^0\gamma \rightarrow \gamma\gamma\pi^0\gamma$	18.7 %	22.6	18.9	16.5
$\phi \rightarrow \eta\gamma \rightarrow \pi^0\pi^0\pi^0\gamma$	3.6×10^{-3}	8.7	3.9	2.0
$\phi \rightarrow \eta\gamma \rightarrow \gamma\gamma\gamma$	6.0×10^{-3}	43.4	32.3	30.9
$\phi \rightarrow \pi^0\gamma$	0.2×10^{-3}	448.0	346.1	324.6
$e^+e^- \rightarrow \gamma\gamma(\gamma)$	0.9×10^{-6}	804.2	407.5	267.9

2) Study/check ISR in MC and learn how to assign a syst. error to it

ϵ_{ana} vs χ^2 cut for VDM processes



Background composition vs χ^2 cut

The systematic study on BKG repeated for the different χ^2 cuts.

In this table we report the data-MC weight factors found by fitting a set of Background-enriched distributions.

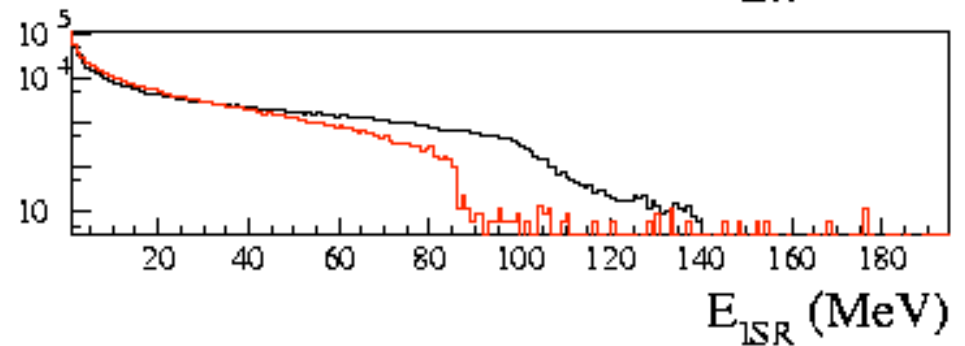
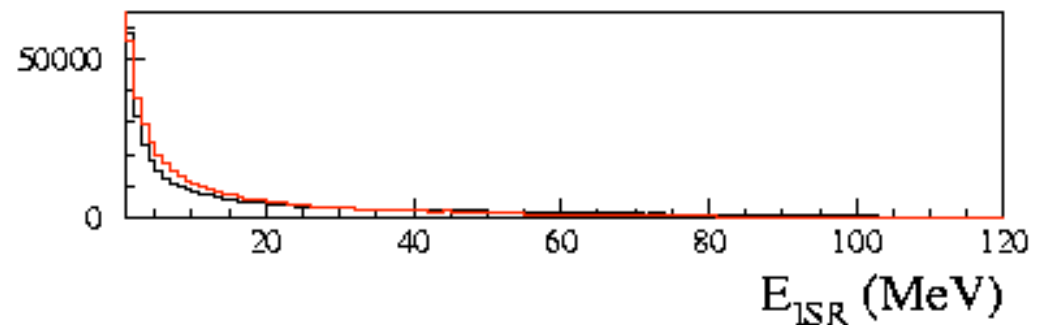
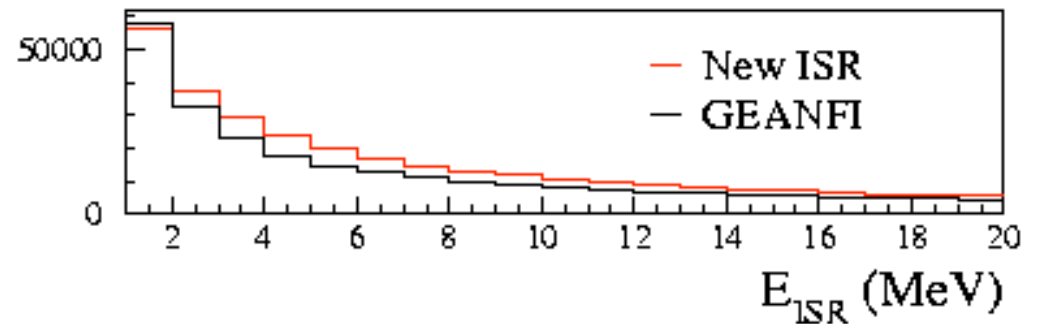
Process	Wb ($\chi^2 < 3$)	Wb ($\chi^2 < 5$)	Wb ($\chi^2 < 10$)
$\phi \rightarrow a_0\gamma \rightarrow \eta\pi^0\gamma \rightarrow \gamma\gamma\pi^0\gamma$	0.787 ± 0.003	0.787 ± 0.005	0.760 ± 0.020
$\phi \rightarrow \eta\gamma \rightarrow \pi^0\pi^0\pi^0\gamma$	1.064 ± 0.002	1.064 ± 0.002	1.064 ± 0.002
$\phi \rightarrow \eta\gamma \rightarrow \gamma\gamma\gamma$	0.892 ± 0.005	0.820 ± 0.040	0.802 ± 0.003
$\phi \rightarrow \pi^0\gamma$	1.78 ± 0.33	1.70 ± 0.24	1.61 ± 0.03
$e^+e^- \rightarrow \gamma\gamma(\gamma)$	1.85 ± 0.03	1.85 ± 0.03	1.85 ± 0.03

Situation is under control. For the main $\eta \rightarrow 3\pi^0$ background the related systematics have been checked with a χ^2 region between 4-20

ISR in Geanfi vs fast generator

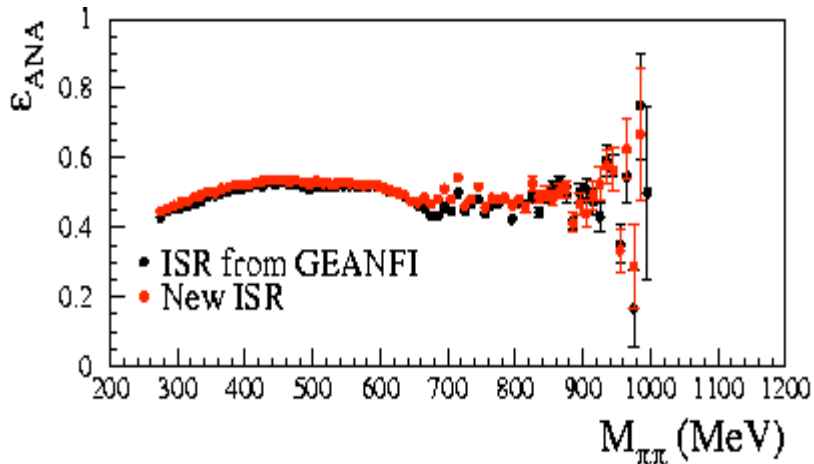
Another way is to test if the ISR in GEANFI is correct.

- We applied the Greco-Nicosini ISR to our standalone VDM generator
- we apply a threshold cutoff at $4M\pi^2$
- a correct \sqrt{s} dependence on the xsec with a threshold behaviour around $M\omega + M\pi$

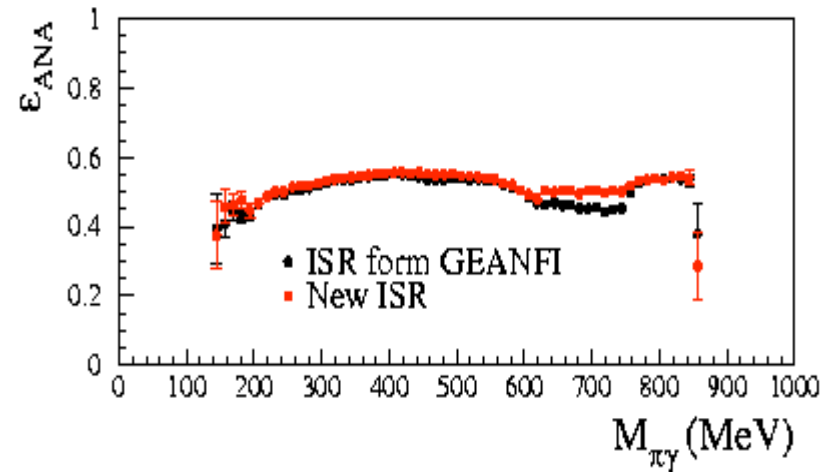
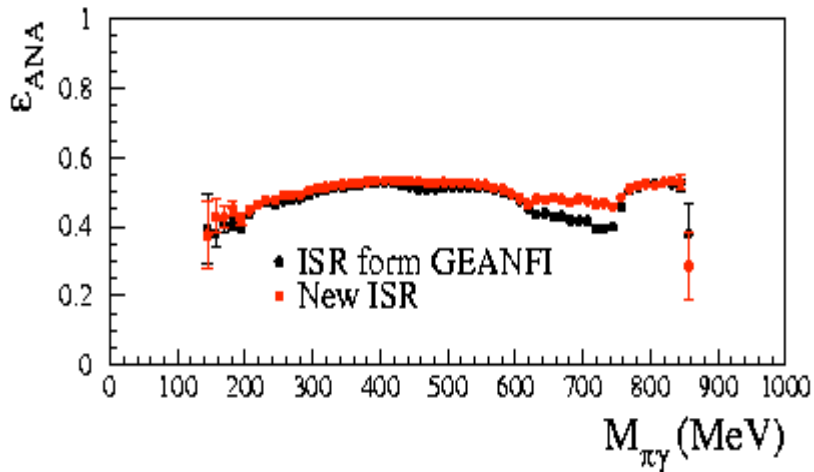
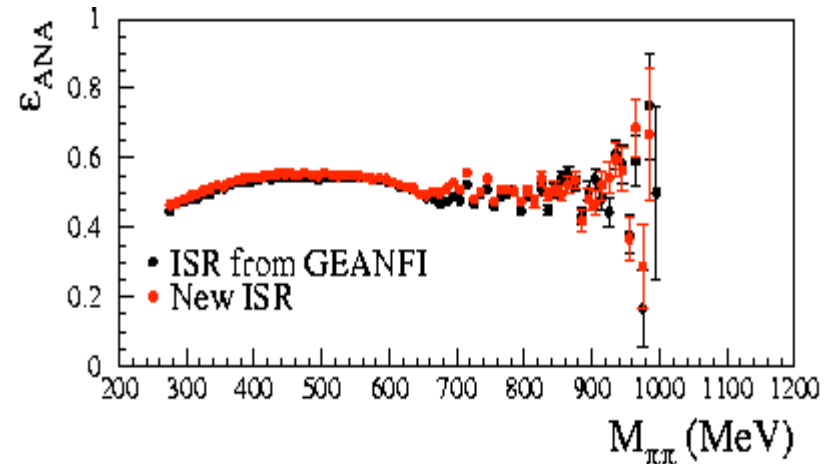


ϵ_{ana} (VDM) vs simulation of ISR

$\chi^2 < 3$



$\chi^2 < 10$



Fit function: the Achasov parametrization

$$\frac{d\sigma(e^+e^- \rightarrow \pi^0\pi^0\gamma)}{dm dm_{\pi\gamma}} = \frac{\alpha m_{\pi\gamma} m}{3(4\pi)^2 s^3} \left\{ \frac{2g_{\phi\gamma}^2}{|D_\phi(s)|^2} |g(m)|^2 \left| \frac{g_{f_0K+K-} - g_{f_0\pi^0\pi^0}}{D_{f_0}(m)} \right|^2 + \right.$$

$$\frac{1}{16} F_1(m^2, m_{\pi\gamma}^2) \left| \left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} - C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(m_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(m_{\pi\gamma}^2)} \right|^2 +$$

$$\frac{1}{16} F_1(m^2, \tilde{m}_{\pi\gamma}^2) \left| \left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} + C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(\tilde{m}_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(\tilde{m}_{\pi\gamma}^2)} \right|^2 +$$

$$\frac{1}{8} F_2(m^2, m_{\pi\gamma}^2) \text{Re} \left[\left(\left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} + C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(m_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(m_{\pi\gamma}^2)} \right) \times \right.$$

$$\left. \left(\left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} + C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(\tilde{m}_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(\tilde{m}_{\pi\gamma}^2)} \right)^* \right] \ominus$$

$$\frac{1}{\sqrt{2}} \text{Re} \left[g(m) e^{i\delta_B(m)} \frac{g_{f_0K+K-} - g_{f_0\pi^0\pi^0}}{D_{f_0}(m)} \frac{g_{\phi\gamma}}{D_\phi(s)} \left(\right.$$

$$F_3(m^2, m_{\pi\gamma}^2) \left(\left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} + C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(m_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(m_{\pi\gamma}^2)} \right)^* +$$

$$F_3(m^2, \tilde{m}_{\pi\gamma}^2) \left(\left(\frac{e^{i\phi_{\omega\phi}(m_\phi^2)} g_{\phi\gamma} g_{\phi\rho\pi} g_{\rho\pi\gamma}}{D_\phi(s)} + C_{\rho\pi} \right) \frac{e^{i\delta_{b\rho}}}{D_\rho(\tilde{m}_{\pi\gamma}^2)} + \frac{C_{\omega\pi^0}}{D_\omega(\tilde{m}_{\pi\gamma}^2)} \right)^* \left. \right] \left. \right\},$$

$f_0\gamma$

Model dependent term

$\omega\pi/\rho\pi$

It was modified in + cos ϕ (???)

NO MORE

$f_0\gamma/VP$ interf

[N.N.Achasov, A.V.Kiselev, private communication]

VDM parametrization: C_{VP} fixed – K_{VDM} (norm factor), $\delta_{b\rho}$, M_V , G_V free

Fit function: the Isidori-Maiani parametrization

- Point-like $\phi S\gamma$ coupling. Corrections to a “standard” BW-like f_0 (fixed Γ_S) described by the a_0 , a_1 parameters
[Isidori-Maiani, private communication]

$$A_1^{\text{scal}} = \frac{e}{4F_\Phi} \frac{sM_\Phi^2}{D_\Phi(s)} \left[\frac{g_{12}^f g_{f\gamma}^\Phi}{D_S[(1-x)s]} + \frac{a_0}{M_\Phi^2} + a_1 \frac{(1-x)s - M_S^2}{M_\Phi^4} \right]$$

In the interference term, a global phase due to the $\pi\pi$ re-scattering is included. No other free phases used.

A more refined parametrization received today:

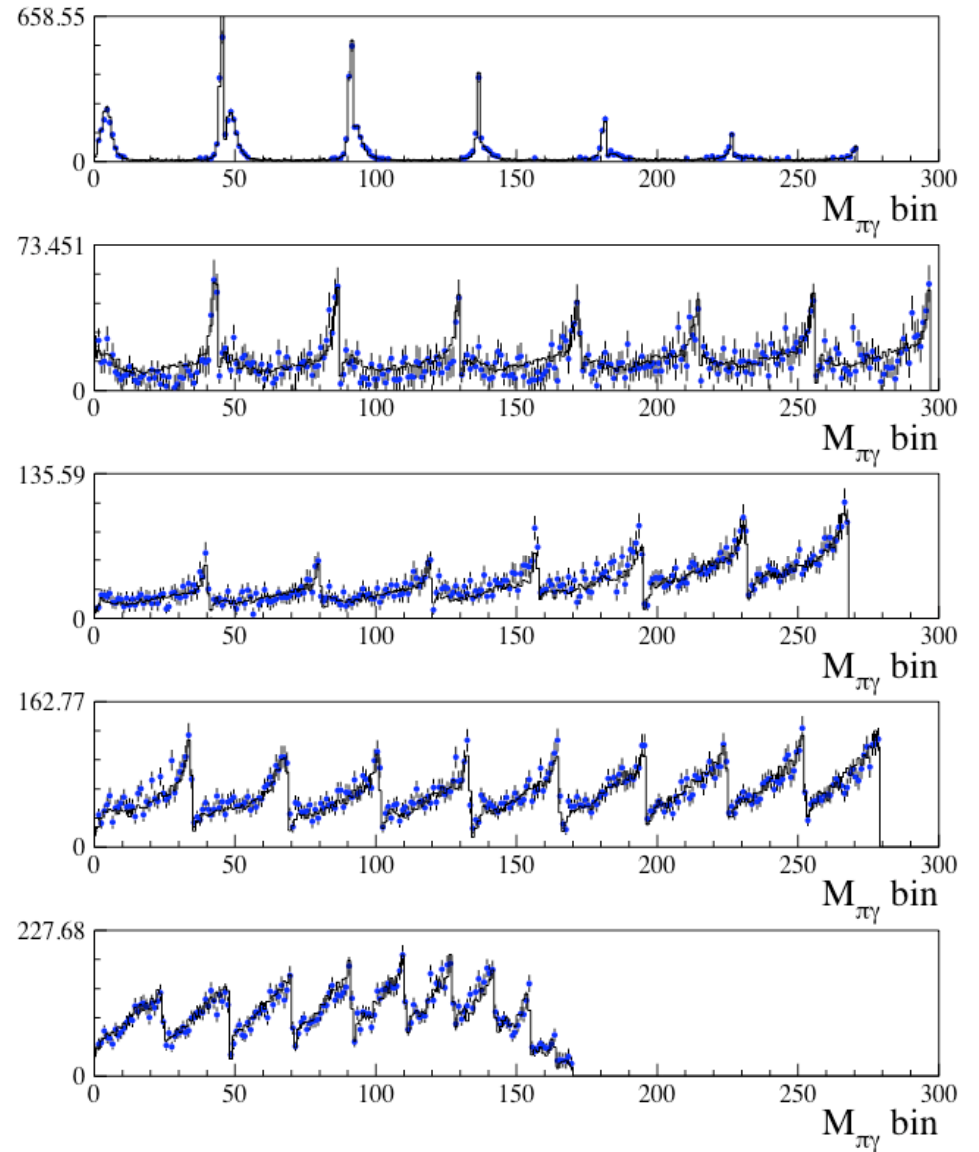
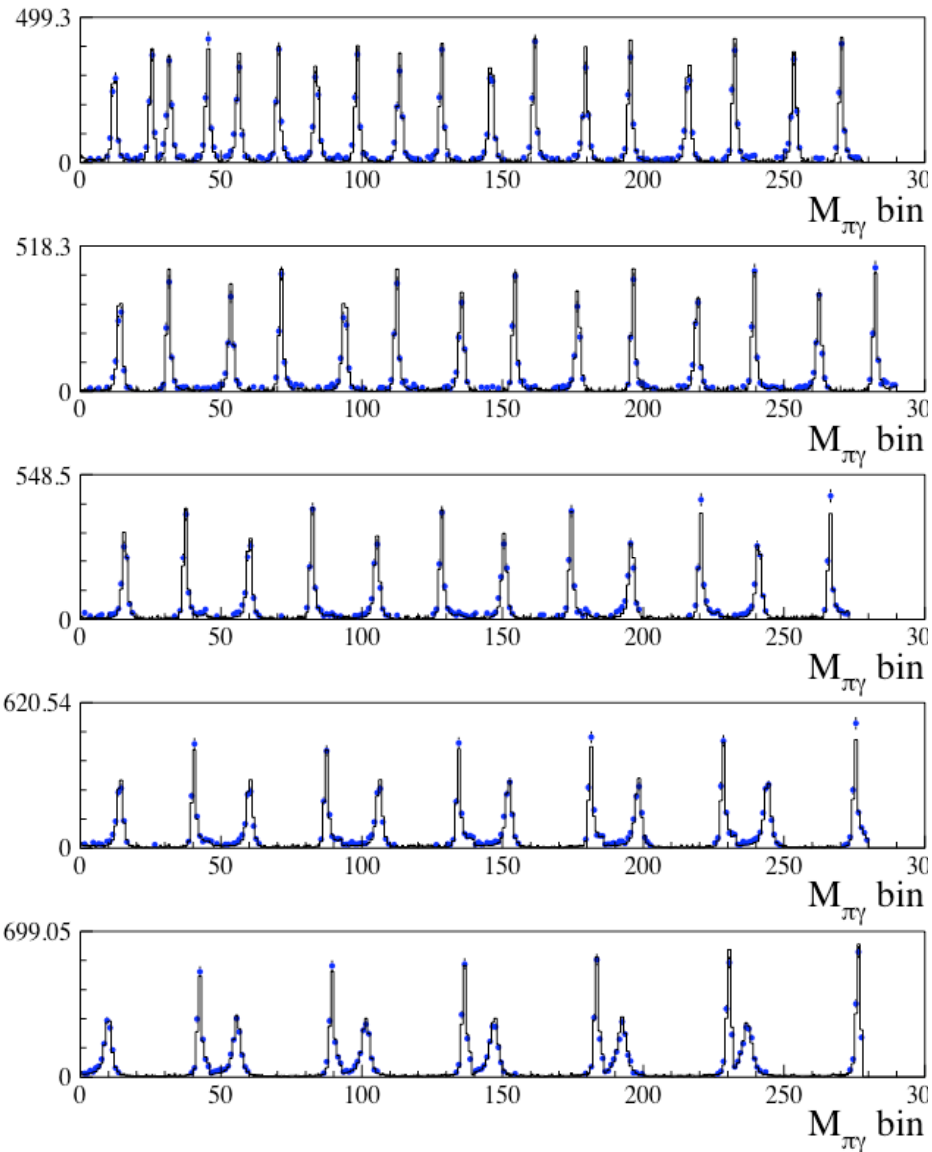
- Propagator with correct threshold behaviour (Flatte-like)
- Different $\pi\pi$ re-scattering phases for the a_0 and a_1 terms in order to reproduced behaviour of δ^0 at low $m_{\pi\pi}$ values

Stability of Fit results for ACH vs χ^2 (new/old ISR)

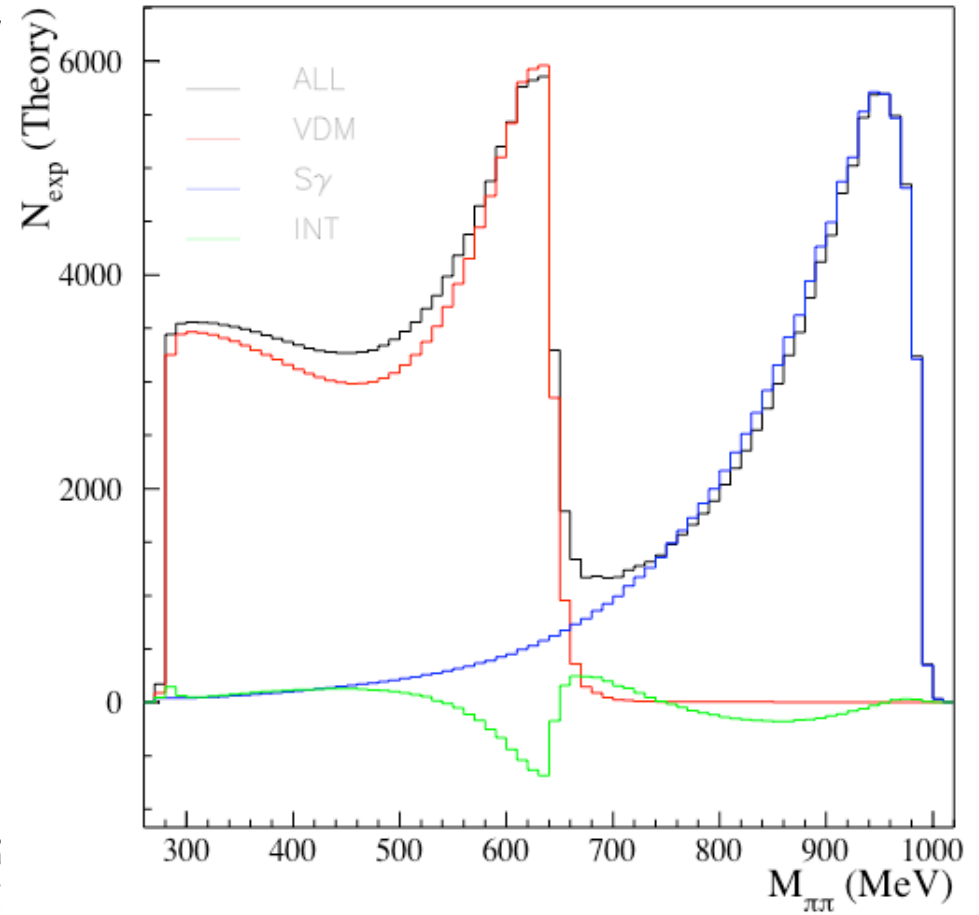
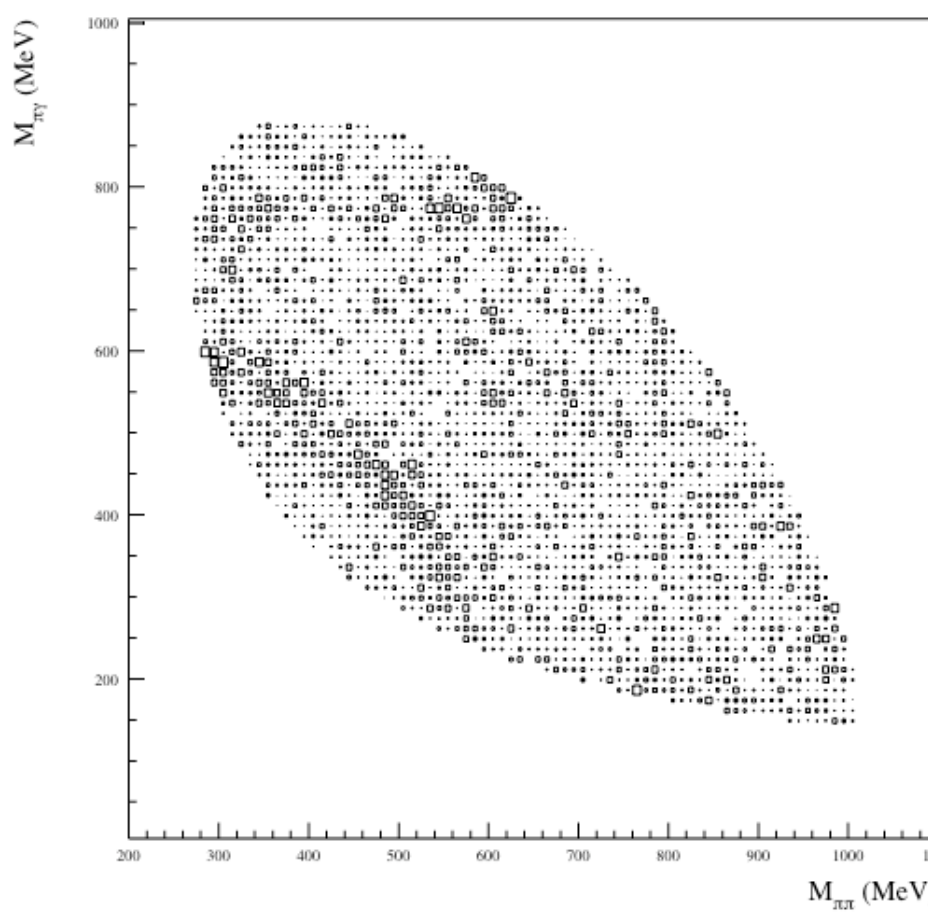
ISR χ^2 (KineFit) cut	NEW			OLD
	3	5	10	5
M_{f_0} (MeV)	959.0 . 2.5	959.0 \pm 1.0	959.5 2.0	961.4 0.8
g (f K^+K^-)(GeV)	4.75 0.12	4.83 \pm 0.06	4.84 \pm 0.11	5.00 \pm 0.06
g (f $\pi^+\pi^-$)(GeV)	1.64 0.09	1.62 \pm 0.05	1.66 \pm 0.09	1.71 \pm 0.03
g (σ $\pi^+\pi^-$)(GeV)	4.56 0.09	4.5 \pm 0.1	4.2 \pm 0.1	4.5 \pm 0.1
Cf_0s (GeV ⁻²)	0.107 0.020	0.111 \pm 0.008	0.119 \pm 0.020	0.130 \pm 0.008
Pf_0s	0.74 0.05	0.79 \pm 0.02	0.83 \pm 0.06	0.73 \pm 0.04
K_{VDM}	0.848 0.005	0.857 \pm 0.005	0.847 \pm 0.005	0.862 \pm 0.005
M_ω (MeV)	782.4 0.1	782.6 \pm 0.1	783.3 \pm 0.1	777.0 \pm 0.2
δ_{b_0} (degree)	100.5 1.4	102.7 \pm 1.2	102.5 \pm 1.7	102.4 \pm 1.3
χ^2/ndf	3310.4/2675	3051.2/2675	2684.0/2675	3061.5/2675
Prob(χ^2)	0.3x10-15	0.4x10-6	0.442	0.2x10-6

VDM-Scalar interference +, $M(\text{Sigma}) = 540$ MeV (Bes)

Dalitz-fit with ACH for $\chi^2(\text{KineFit}) < 5$: projections



Dalitz-fit with ACH for $\chi^2(\text{KineFit}) < 5$: contributions

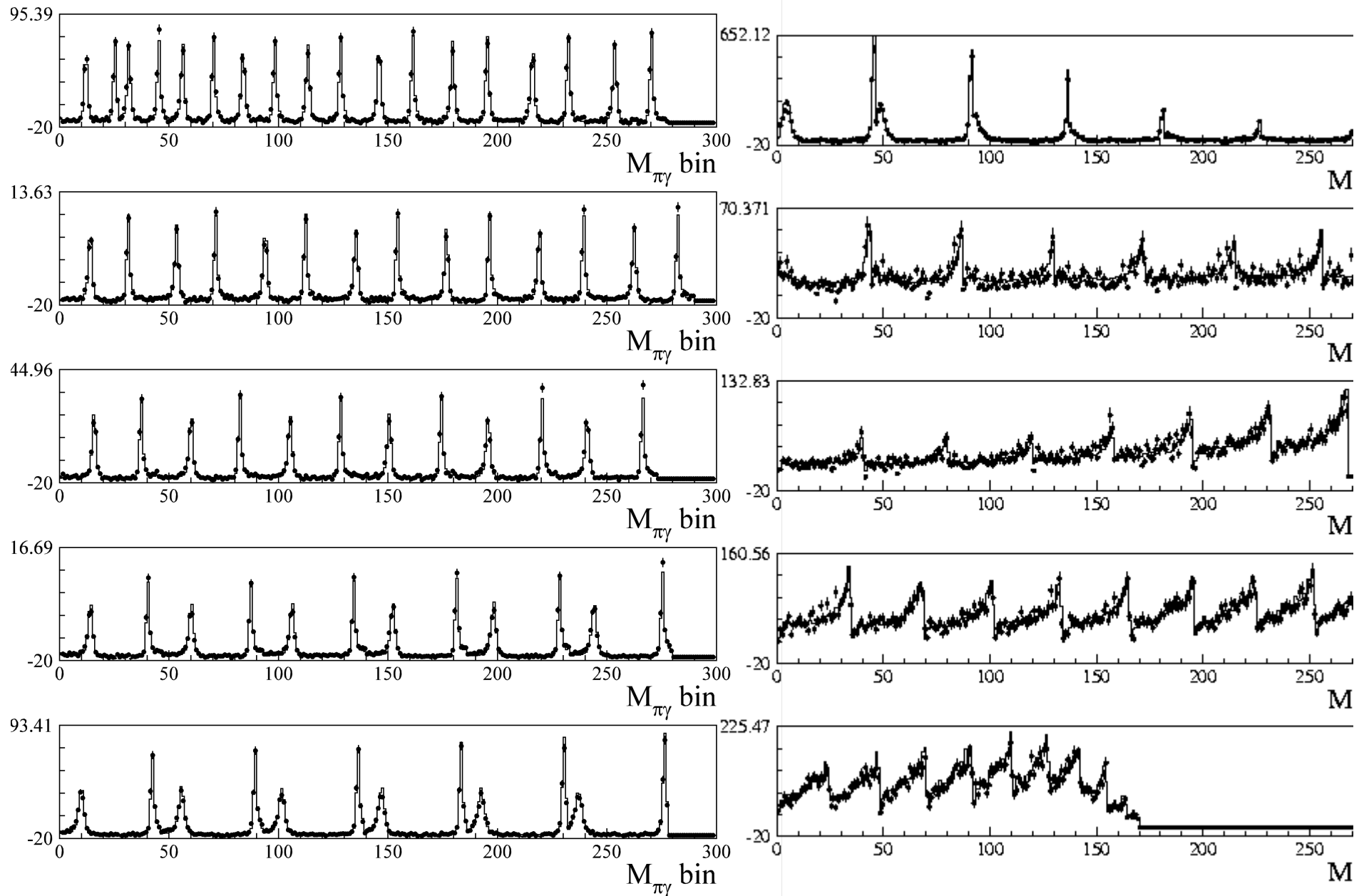


Stability of Fit results for IM vs χ^2 cut (new/old ISR)

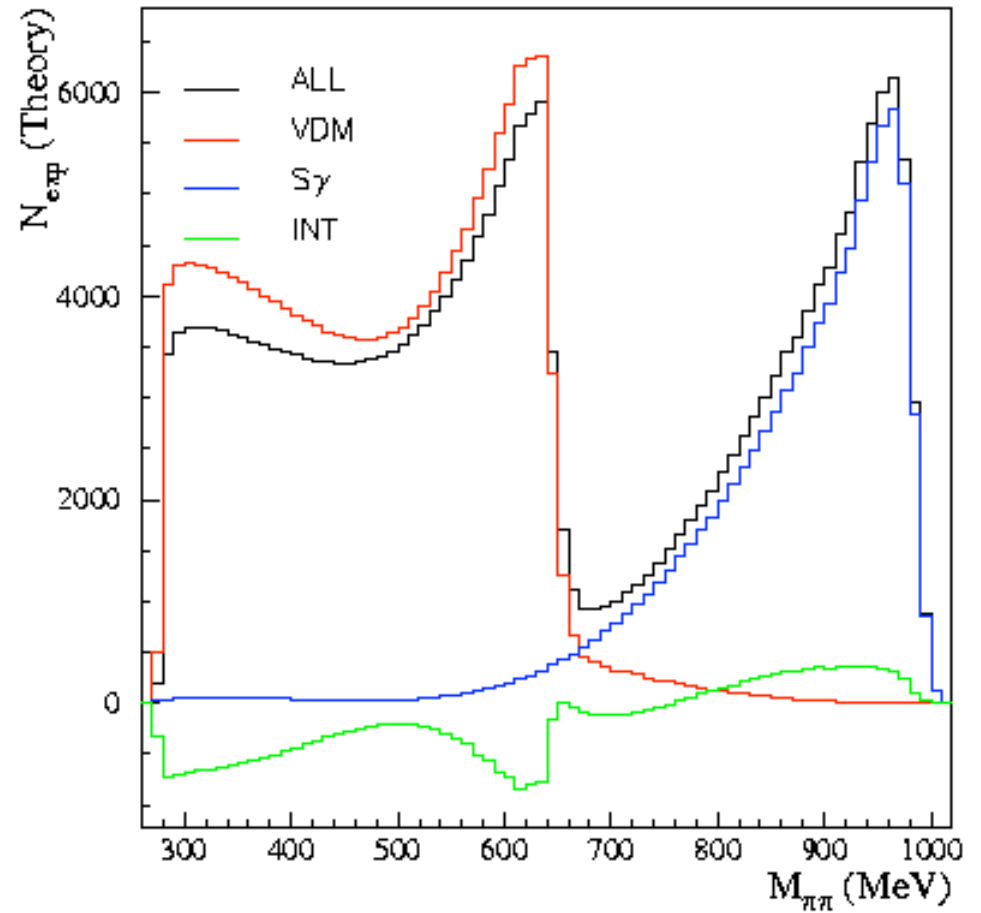
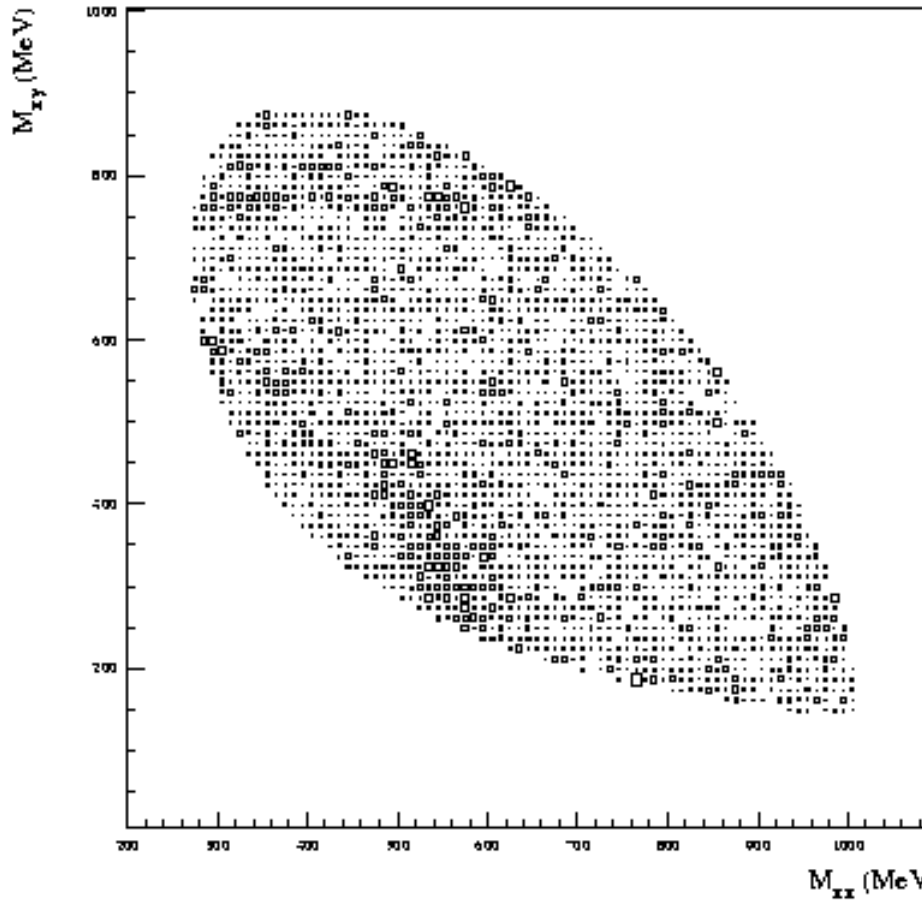
ISR χ^2 (KineFit) cut	NEW			OLD
	3	5	10	5
M_{f_0} (MeV)	981.7 ± 0.6	983.0 ± 0.7	982.8 ± 0.6	982.7 ± 0.2
Γ_{f_0} (MeV)	43.3 ± 0.7	43.9 ± 0.6	43.9 ± 0.7	43.8 ± 0.6
$g_{\phi f\gamma} \times g_{f\pi\pi}$	2.02 ± 0.02	2.10 ± 0.03	2.11 ± 0.02	2.08 ± 0.01
a_0	3.42 ± 0.07	3.74 ± 0.07	3.80 ± 0.09	3.64 ± 0.05
a_1	0.79 ± 0.07	1.08 ± 0.06	1.11 ± 0.10	1.03 ± 0.06
K_{VDM}	0.683 ± 0.004	0.683 ± 0.004	0.675 ± 0.004	0.686 ± 0.004
M_ω (MeV)	782.03 ± 0.07	782.09 ± 0.07	782.85 ± 0.07	782.13 ± 0.05
$\delta_{b\rho}$ (degree)	2.4 ± 1.3	2.2 ± 1.3	0.0 ± 2.0	1.7 ± 1.2
χ^2/ndf	3020.8/2675	2766.2/2675	2380.4/2675	2776.0/2675
Prob(χ^2)	0.27×10^{-3}	0.107	1.00	0.085

To be tried with the new parametrization and fixing $\delta_{b\rho}$ to grant a correct VDM dip vs \sqrt{s}

Dalitz-fit with IM for $\chi^2(\text{KineFit}) < 5$: projections



DALITZ-fit with IM for $\chi^2(\text{KineFit}) < 5$: contributions



Summary of Fit results (I)

- We have \approx finished to look for systematic effects ..we will incorporate their error on the χ^2 evaluation for the Dalitz-fit.
No large improvements foreseen.
- The ISR for VDM processes introduces a variation of the analysis efficiency along the plane. **Easy to overcome enlarging the cut ... but the background increases of a large factor.**
Intermediate case chosen: $\chi^2 < 5$.
- Due to the large statistics the $\text{prob}(\chi^2)$ reaches very low values also for this case (much better values obtained for the adjacent bins in \sqrt{s} ...)

Summary of Fit results (II)

- HOWEVER ... we have stable results for a “given” parametrization!
- We are unable to fit the dalitz with the ach-f0 parametrization without including the SIGMA meson.
- The inclusion of sigma seems to be fine but different results are obtained if leaving free or not the value of the mass.
- For both models we are not yet convinced of the proper δ^0_0 behaviour of the found solution.

This is still work in progress.

Conclusions/prospects

The two positive sides of the story:

- 1) above 700 MeV the situation is stable at 5-10 % level for whatever parametrization we use. Also at low masses we have similar shapes.
- 2) We are not anymore in disagreement with the $\pi^+\pi^-$ case

IF/WHEN we use an identical parametrization.

The two negative sides of the story:

- 1) We have to complete the description of δ_0^0 to be sure these fits make any sense!
- 2) although many efforts done on this item and we consider the EXPERIMENTAL SIDE “CONCLUDED” **we have not yet come to a solid conclusion.**

PLAN is however to start writing a note, incorporate latest changes of IM and go for a paper even with large $\text{prob}(\chi^2) \dots$

Spare Slides
Old Slides

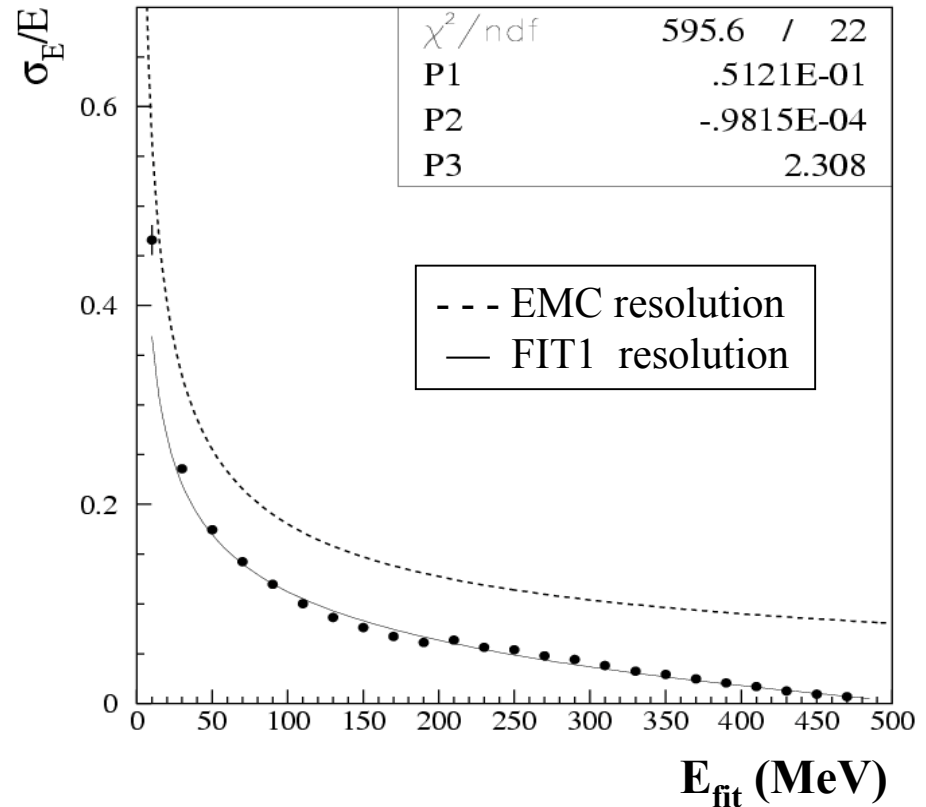
γ 's pairing

π^0 mass resolution parametrized as a function of the γ 's energy resolution after kinematic fit:

$$\sigma_M/M = 0.5 (\sigma_{E_1}/E_1 \oplus \sigma_{E_2}/E_2)$$

Fit function for energy resolution:

$$\sigma_E/E = (P_1 + P_2 E) / E[\text{GeV}]^{P_3}$$



The photon combination that minimize the following χ^2 is chosen:

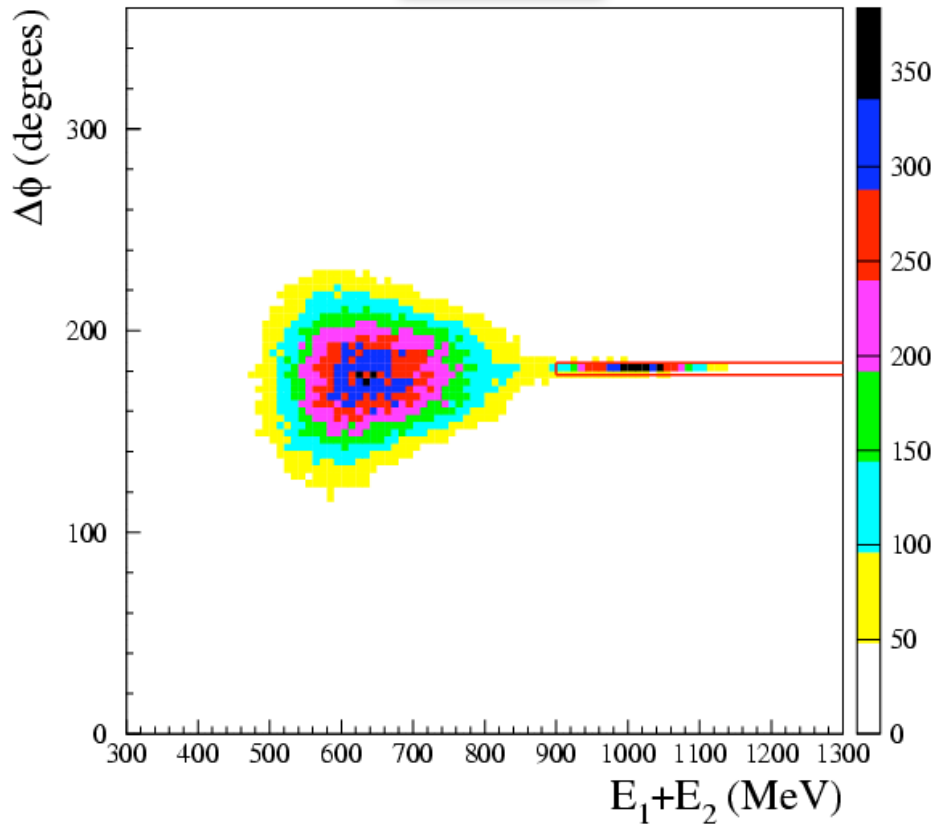
$$\chi^2 = (M_{\gamma_i \gamma_j} - M_\pi) / \sigma_{M_{ij}} + (M_{\gamma_k \gamma_l} - M_\pi) / \sigma_{M_{kl}}$$

$e^+e^- \rightarrow \gamma\gamma$ rejection

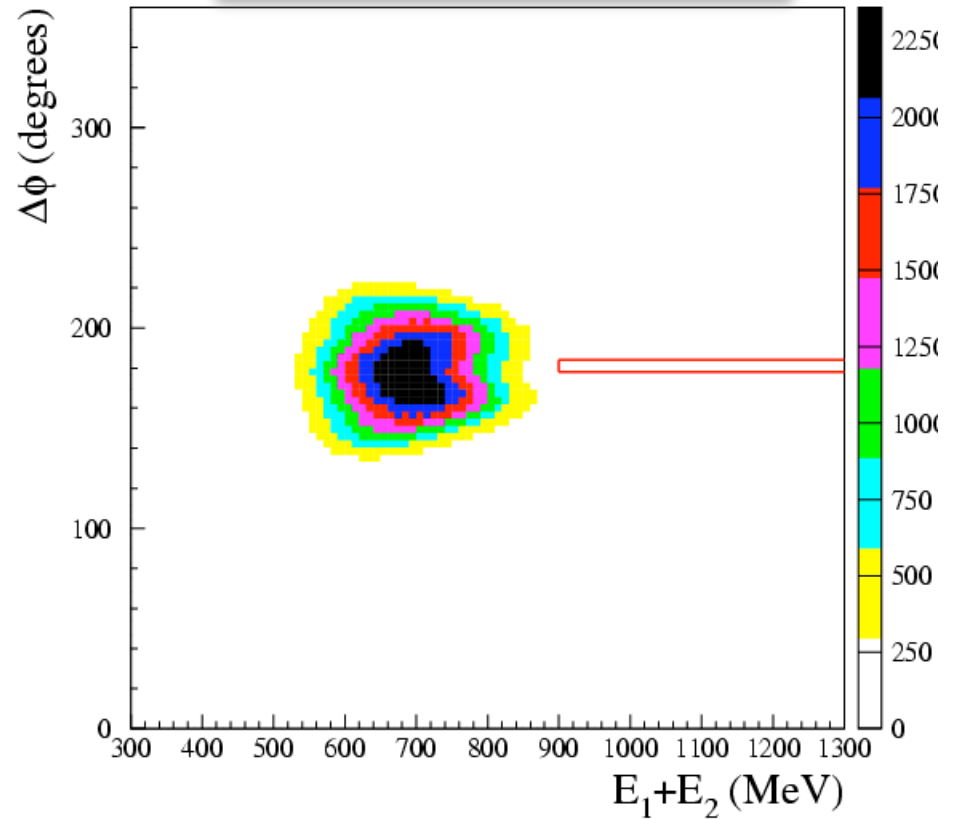
$e^+e^- \rightarrow \gamma\gamma$ rejection done using the two most energetic clusters of the event:

$$E_1 + E_2 > 900 \text{ MeV}$$

Data

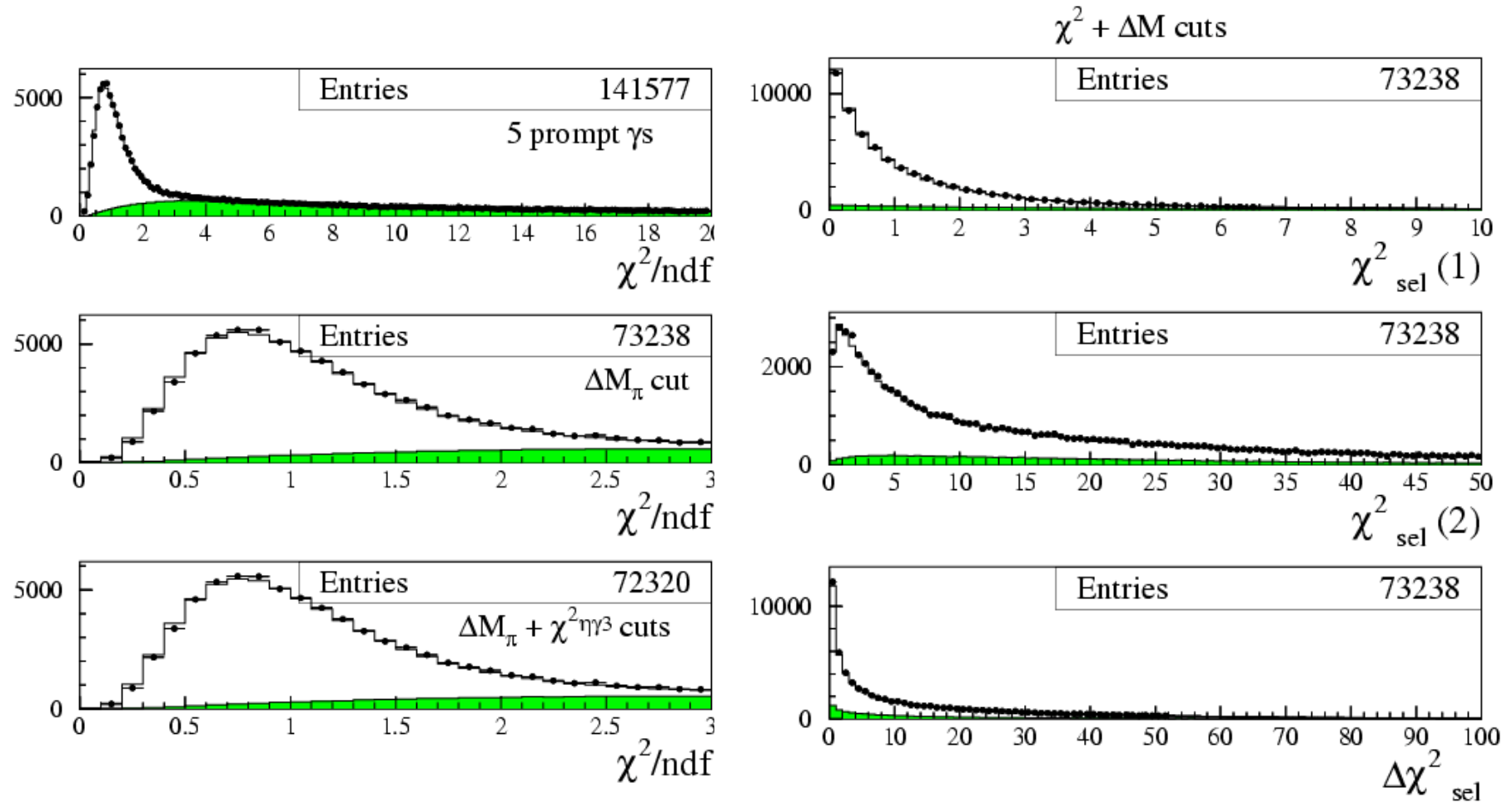


MC $\pi^0\pi^0\gamma$ events



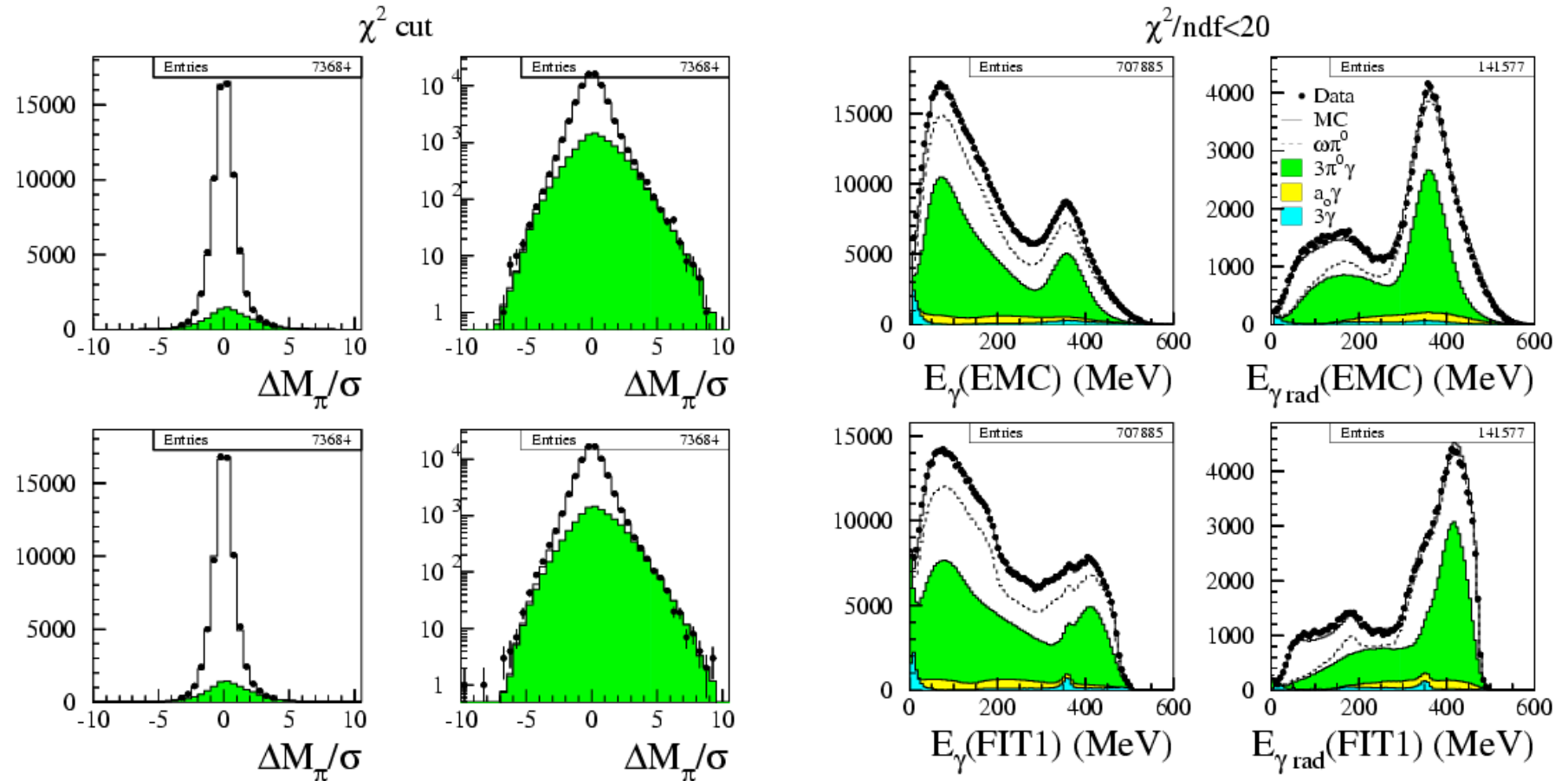
Dalitz plot analysis: data-MC comparison (I)

Analysis @ $\sqrt{s} = 1019.6$ MeV ($L_{\text{int}} = 145$ pb $^{-1}$)



Dalitz plot analysis: data-MC comparison (II)

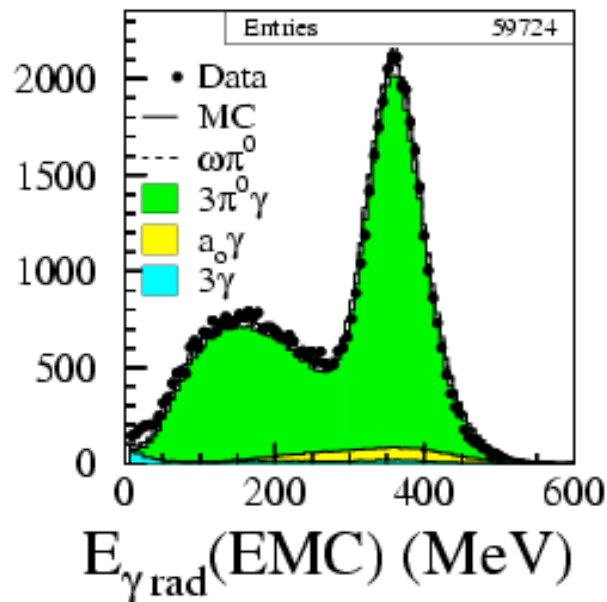
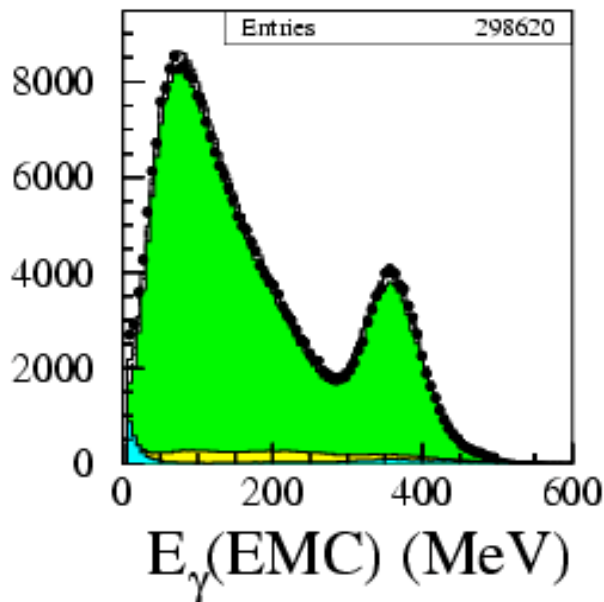
Analysis @ $\sqrt{s} = 1019.6 \text{ MeV}$ ($L_{\text{int}} = 145 \text{ pb}^{-1}$)



Background study for Dalitz plot analysis (I)

In order to study the systematics connected to the background subtraction we found for each category a distribution “background dominated” to be fitted

- $\phi \rightarrow \eta \gamma \rightarrow \pi^0 \pi^0 \pi^0 \gamma$ (most relevant bckg contribution)
 - Background enriched sample : $4 < \chi^2/\text{ndf} < 20$

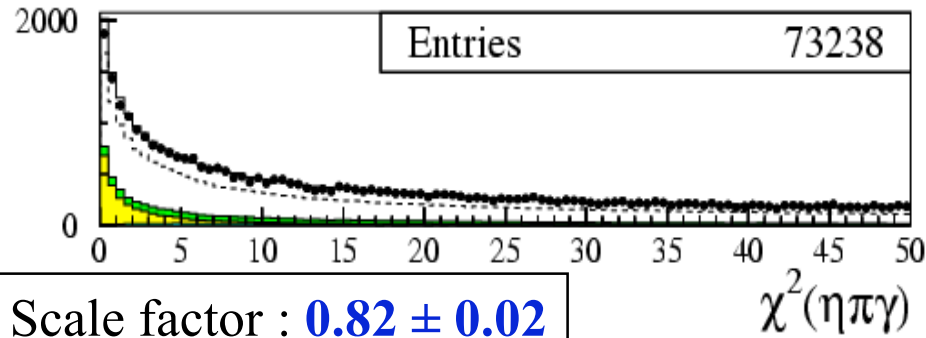
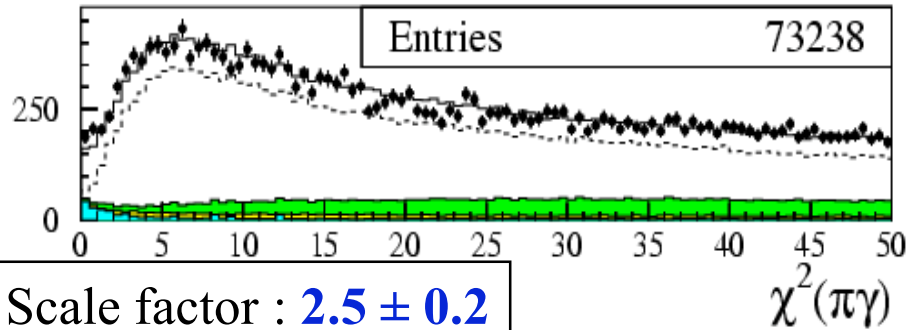
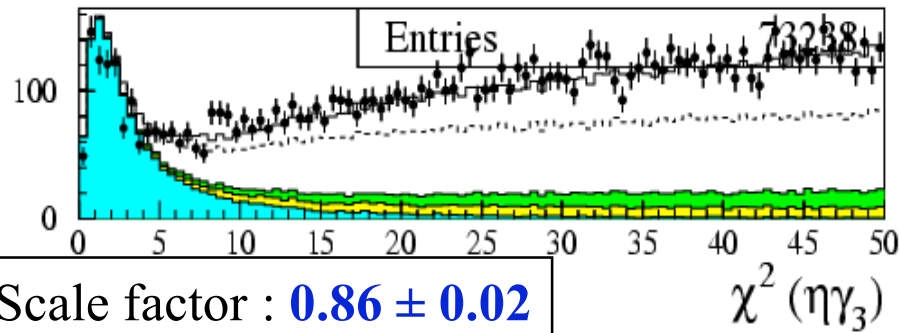


Scale factor :
 1.0156 ± 0.0002

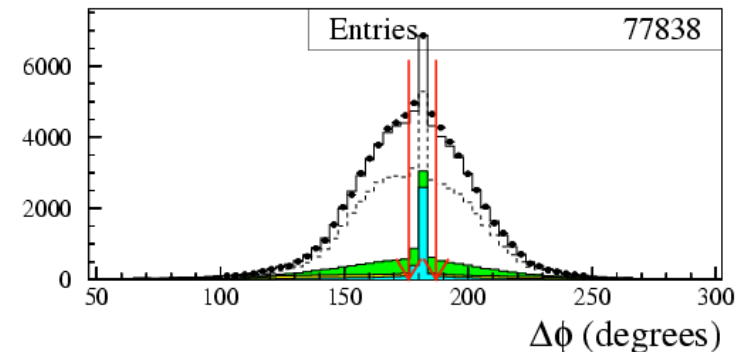
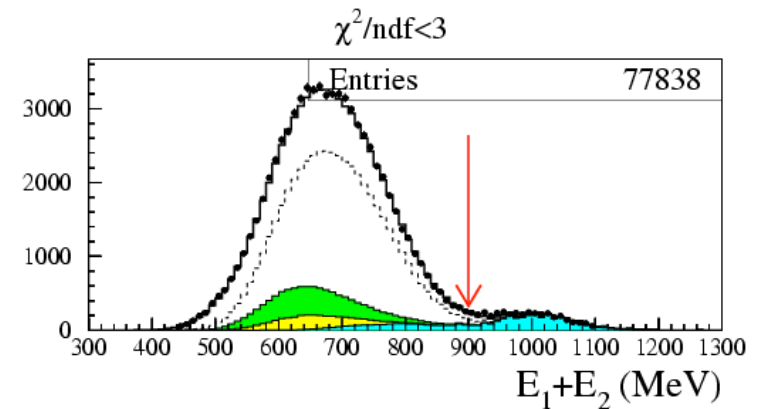
All of this fit results are used to evaluate the systematics on the background counting : **half of the difference (1 – scale factor)** is used

Background study for Dalitz plot analysis (II)

For $\phi \rightarrow \eta\gamma \rightarrow \gamma\gamma\gamma$, $\phi \rightarrow \pi^0\gamma$, $\phi \rightarrow a_0\gamma$
we calculate a χ^2 in the mass hypothesis



For $e^+e^- \rightarrow \gamma\gamma$, we fit the $\Delta\phi$ distribution for $\chi^2/\text{ndf} < 3$
(and no $\gamma\gamma$ rejection cut)



Scale factor : 1.85 ± 0.03

Dalitz plot @ $\sqrt{s}=1019.6$ MeV

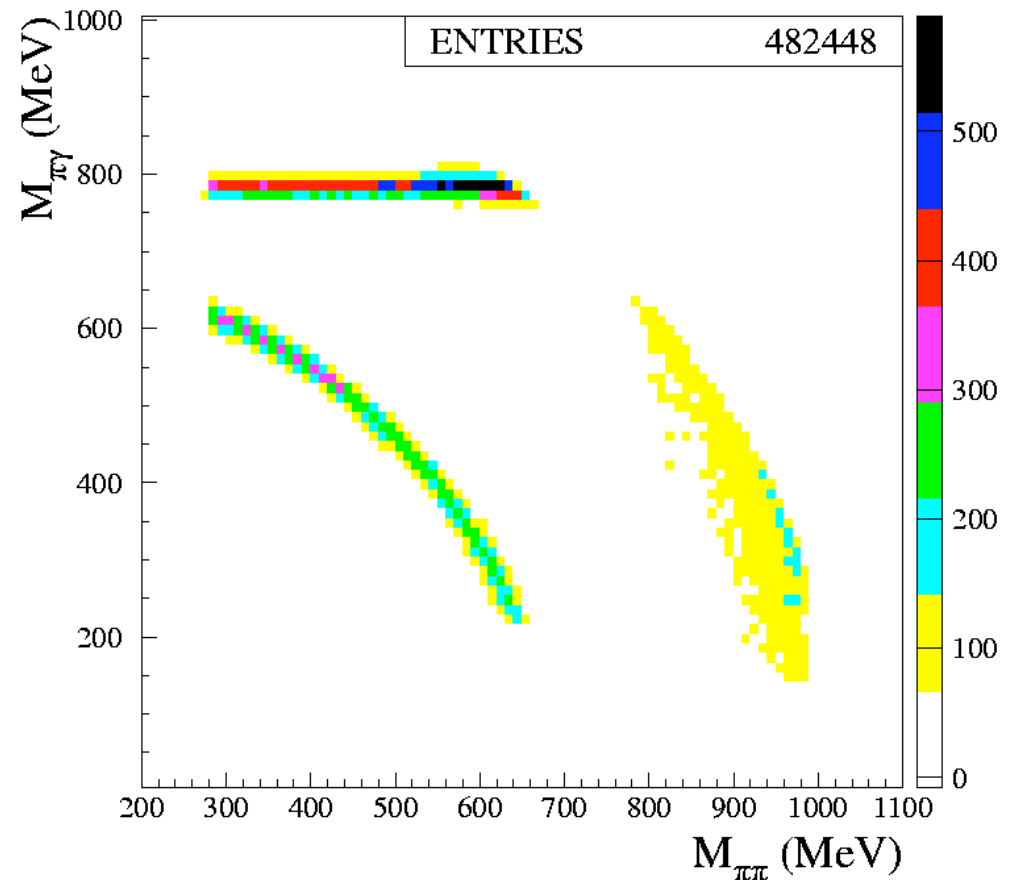
Fit to the Dalitz plot with the VDM and scalar term, including also interference

Binning: 10 MeV in $M_{\pi\pi}$, 12.5 MeV in $M_{\pi\gamma}$

What is needed:

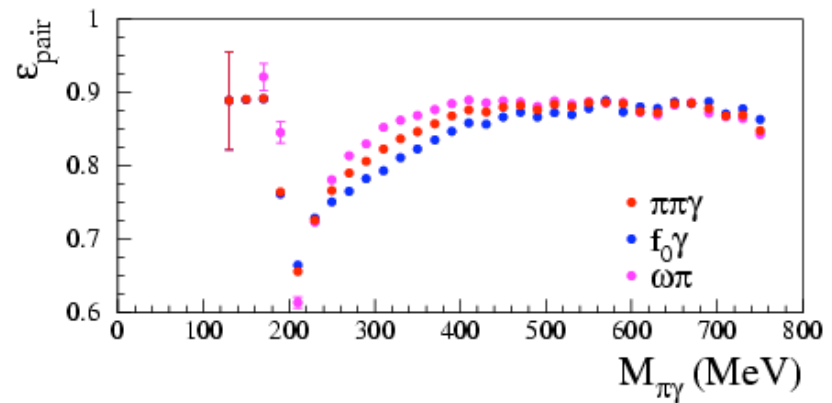
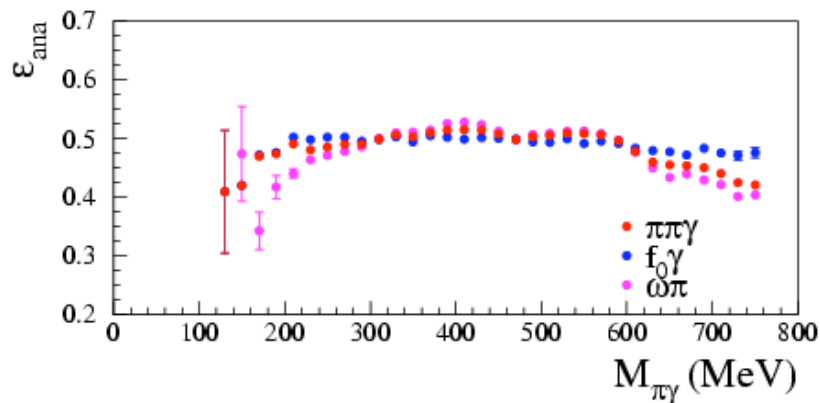
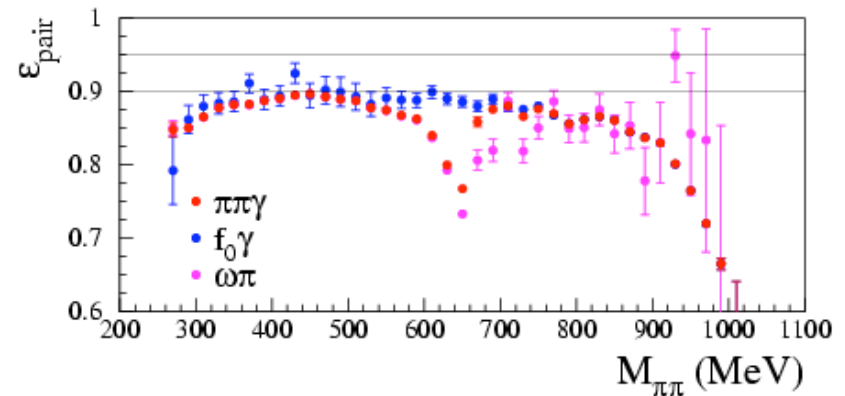
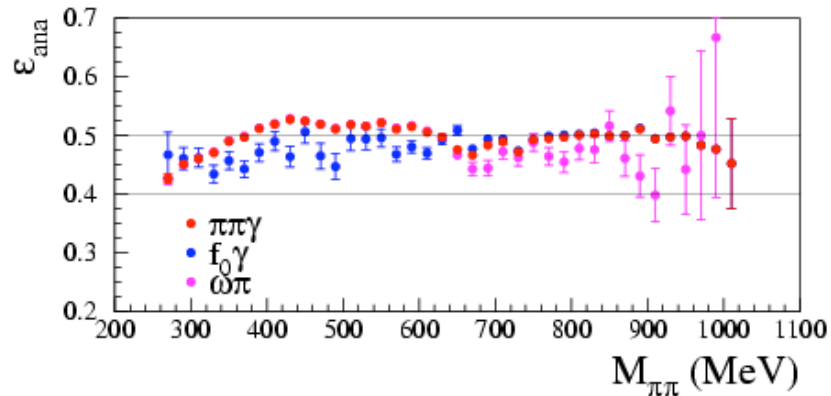
- Analysis efficiency
- Smearing matrix
- Theoretical functions
- ISR

Only statistical error and systematics on background considered for the moment



Analysis and pairing efficiencies ν s $M_{\pi\pi}$, $M_{\pi\gamma}$

Analysis efficiency and smearing matrix evaluated from MC for each bin of the $M_{\pi\pi}$ - $M_{\pi\gamma}$ plane

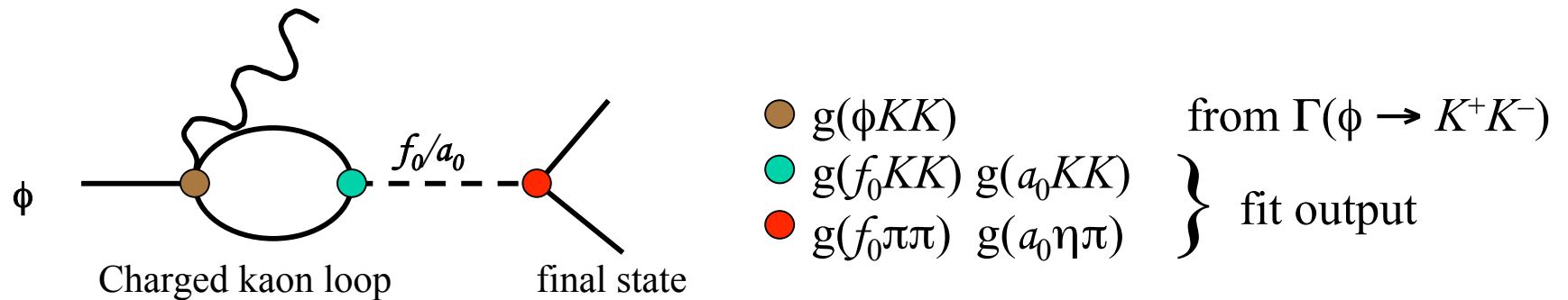


Different for the two processes!

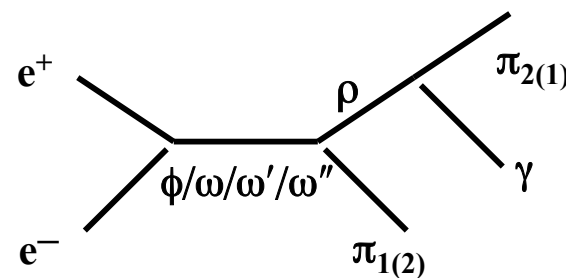
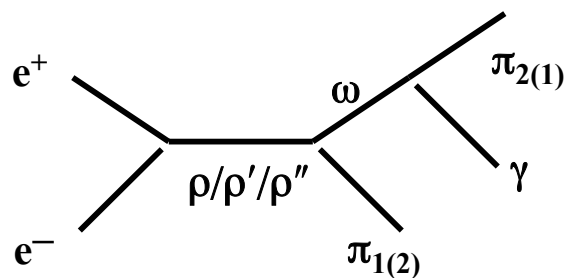
In the fit of the Dalitz different ϵ_{ana} and smearing used for the VDM and scalar contributions. For the moment the VDM results are used also for the interference term

Fit function: the Achasov parametrization (I)

➤ Scalar produced through a kaon loop



➤ VDM contribution from the following diagrams :



➤ All interferences considered

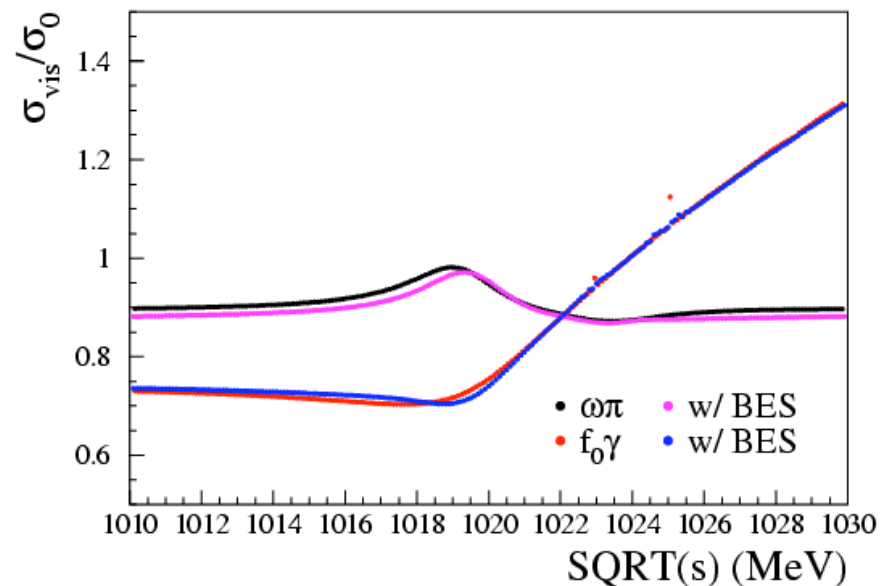
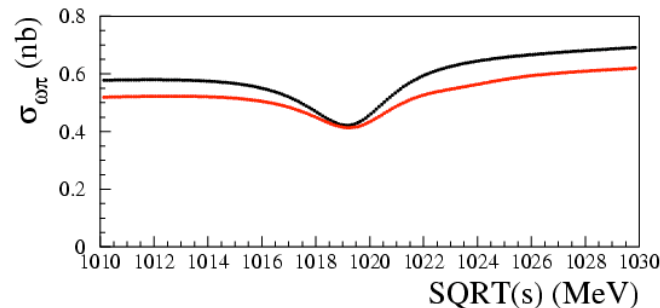
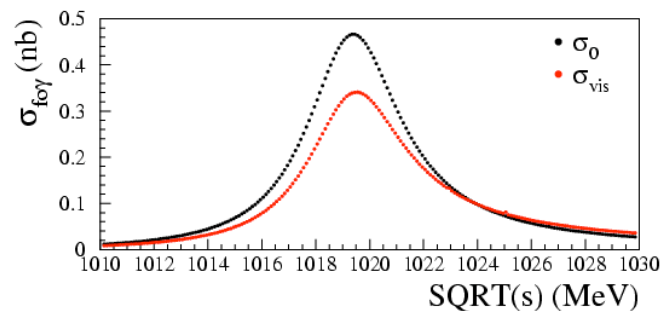
Calculation of the radiative corrections

ISR evaluated starting from the following σ_0 :

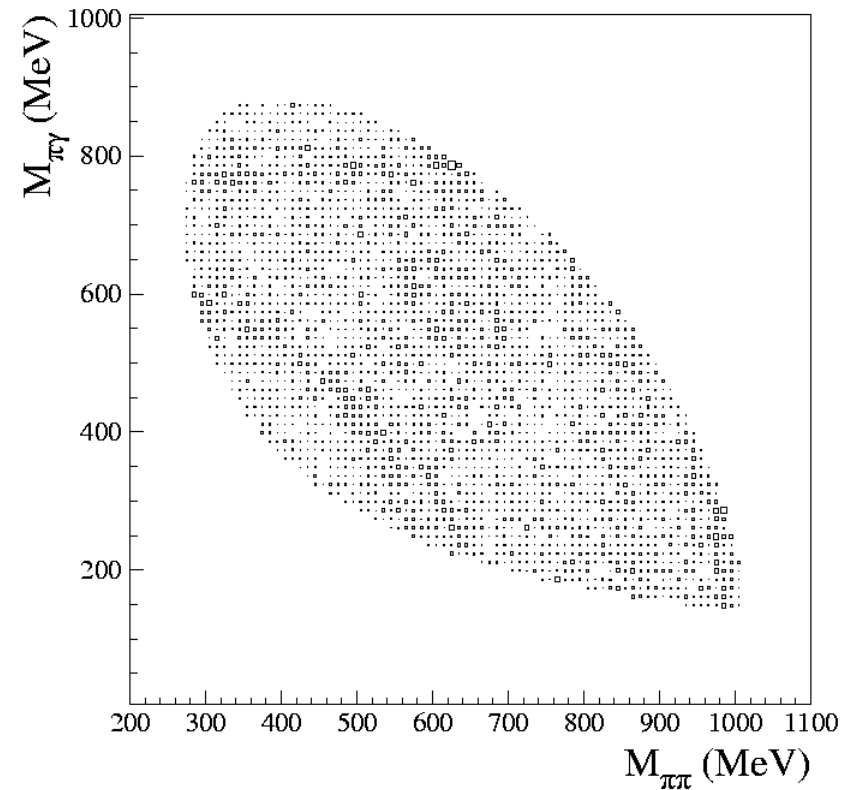
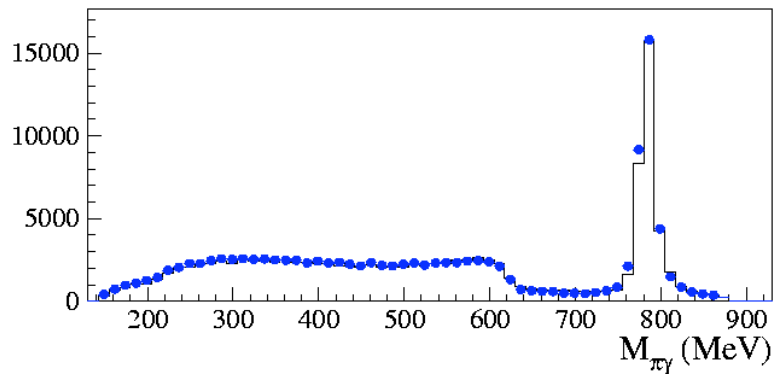
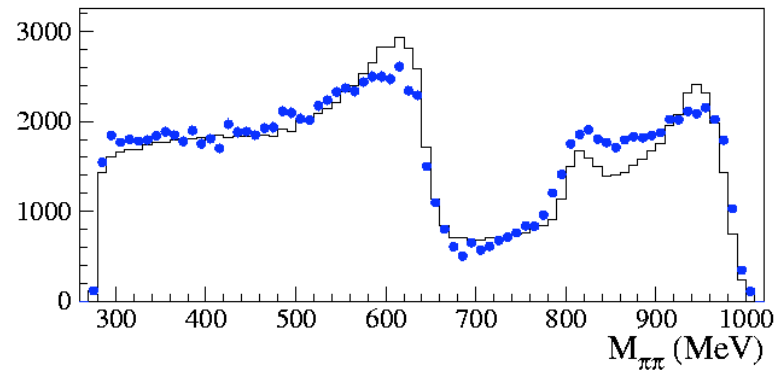
f_0 = “simple” BW (by integrating the Achasov scalar term)

$\omega\pi$ = SND parametrization from JETP-90 6 (2000) 927, obtained by fitting over a large \sqrt{s} range ... Proper threshold behaviour

$$\sigma_{vis} = \int_0^{4m_\pi^2} \sigma_0 [(1-x)s] H(s, x) \quad H(s, x) \text{ from Antonelli, Dreucci}$$



Fit results: the Achasov parametrization



The parametrization with the σ meson (I)

The σ is introduced in the scalar term as in ref. PRD 56 (1997) 4084.

- The two resonances are not described by the sum of two BW but with the matrix of the inverse propagators G_{R1R2} .
- Non diagonal terms on G_{R1R2} are the transitions caused by the resonance mixing due to the final state interaction which occurred in the same decay channels $R1 \rightarrow ab \rightarrow R2$

$$\frac{g_{f_0 K^+ K^-} g_{f_0 \pi^+ \pi^-}}{D_{f_0}(M_{\pi\pi})} \rightarrow \sum g_{R_1 kk} G_{RR}^{-1} g_{R_2 \pi\pi}$$

Where

$$G_{R1R2} = \begin{pmatrix} D_{f_0} & -\Pi_{f_0\sigma} \\ -\Pi_{\sigma f_0} & D_{\sigma} \end{pmatrix}$$

$$\Pi_{R1R2} = \sum_{ab} g_{R2ab} P_{R1}^{ab}(m) + C_{R1R2}$$

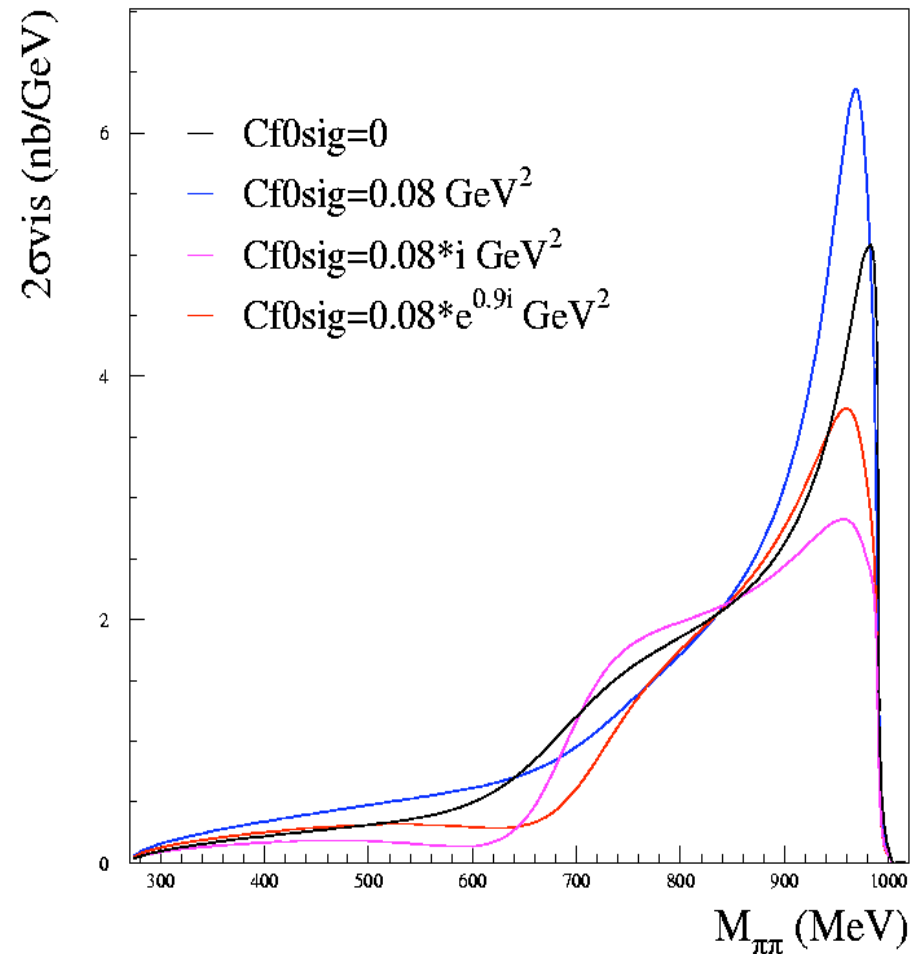
$C_{R1R2} = C_{f_0\sigma}$ takes into account the contributions of VV, 4 pseudoscalar mesons and other intermediate states. In the 4q,2q models there are free parameters

The parametrization with the σ meson (II)

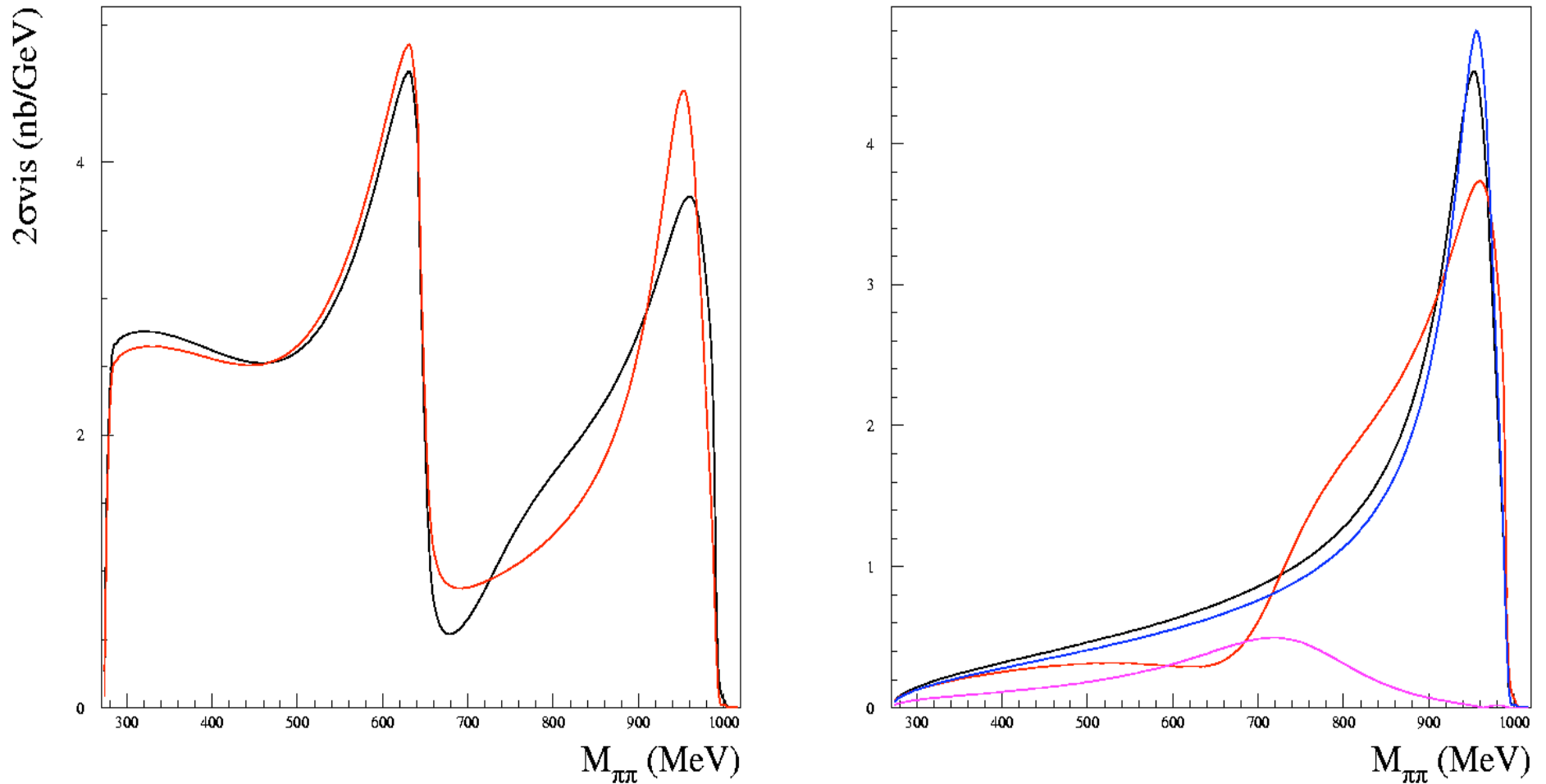
Extensive tests have been done on the formula used.

- Good agreement found between our coding and the one of Cesare we agreed that there is a mistype in the PRD
- We have asked also the help of G.Isidori-S.Pacetti to check this

The effect of the free term $C_{f0\sigma}$ and of its phase is large

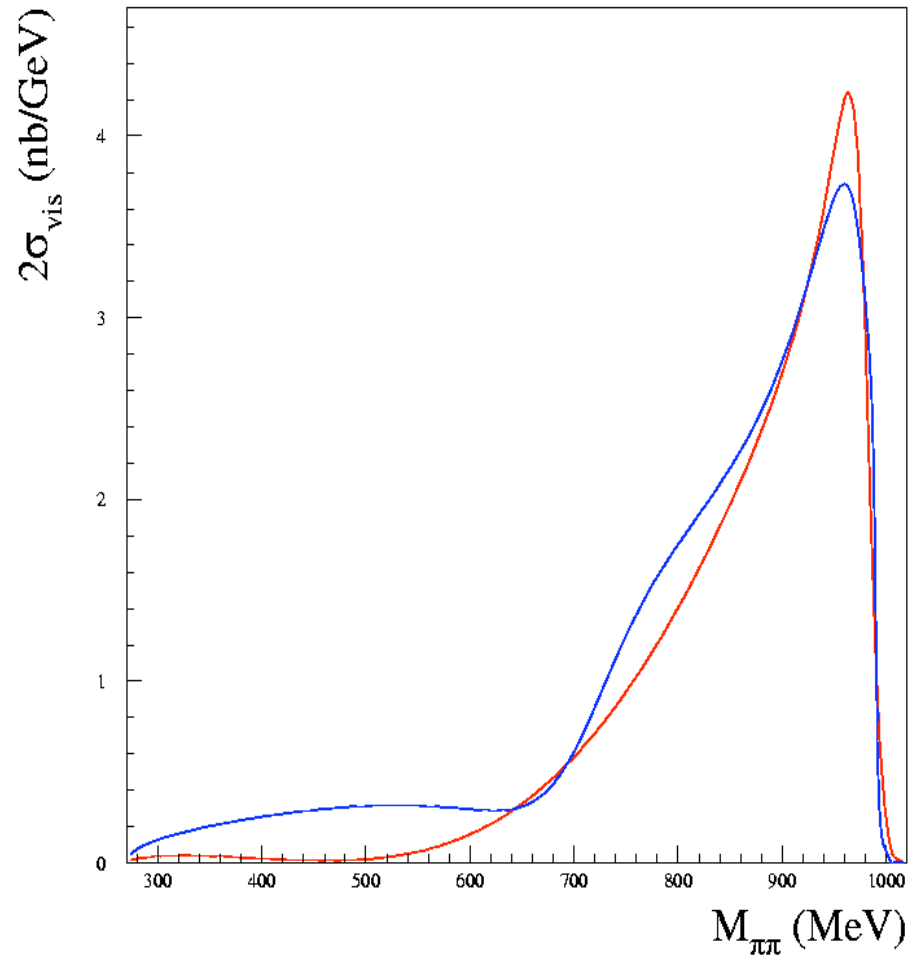


Fit results: the Achasov parametrization with σ (II)



- Black (red) curve are ACH model with (without) the inclusion of the σ meson
- Blue (purple) curve are the contribution due to the f_0 (σ) meson only with the ACH model when including the σ meson

Comparison between ACH-IM for the scalar term



Without the inclusion of the σ meson the agreement between ACH model and IM is not excellent although the integrals do not differ more than 20% above 700 MeV . Including the σ the agreement is better!