

# Little Higgs Models

## Theory & Phenomenology

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- How to make a light Higgs (without SUSY)
- Minimal models: The Littlest Higgs and the Minimal Moose
- Phenomenology

## Electroweak symmetry breaking (beyond the SM):

- Dynamics responsible for the electroweak scale?

$$v = 246 \text{ GeV} \approx 10^{-17} M_{\text{Planck}}$$

- New idea (and very old): geometrical origin = **extra dimensions**
- Old idea: field-theoretical origin = **renormalization-group running and field condensation**

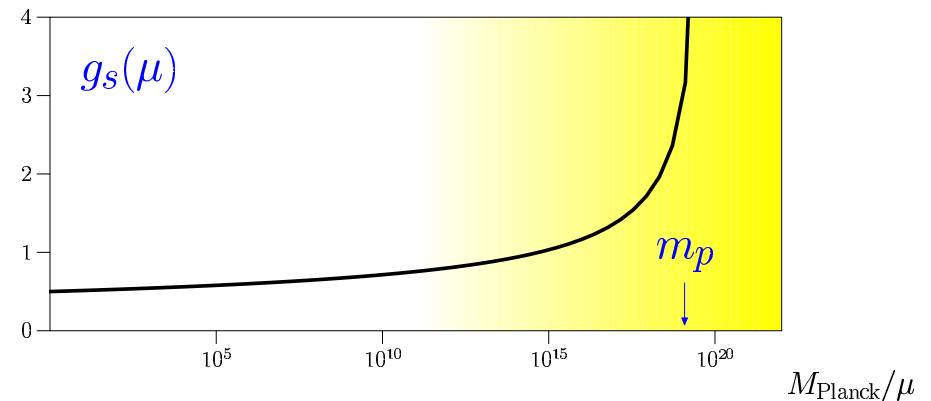
- Why would we like dynamical scale generation?

- We see it at work (QCD):

$$\langle GG \rangle \sim \langle qq \rangle \sim 1 \text{ GeV} \approx 10^{-19} M_{\text{Planck}}$$

$$\text{without GUT: } g_s(M_{\text{Planck}}) \approx 1/2$$

⇒ The proton mass is **natural**

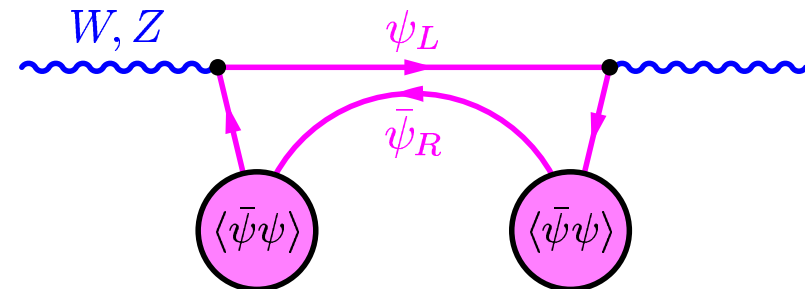


## Dynamical scale generation:

Why not simply copy QCD?

Susskind/Weinberg (1979)

- New QCD' (**Technicolor**),  $\Lambda \sim 3 \text{ TeV} = 4\pi v$
- ⇒ **Compositeness** (of  $W, Z, \dots$ )
- ⇒ **Resonances** in the TeV range

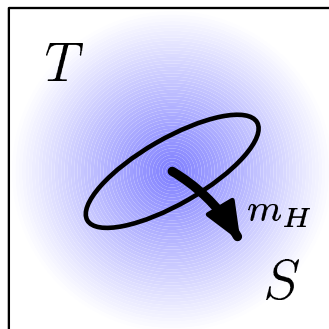


This is elegant, but ...

- **EW precision data**
- **Flavor Physics**

suggest:

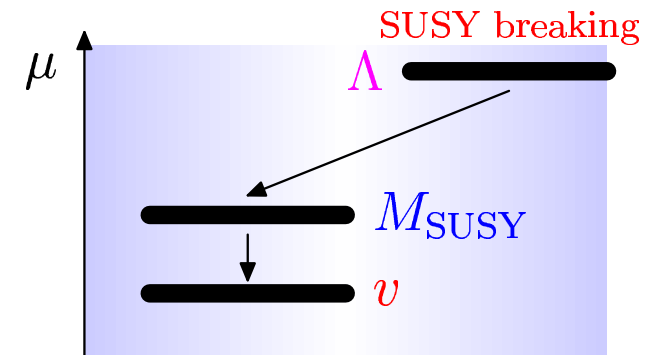
- ⇒ TeV-scale physics is **weakly interacting**
- ⇒ Strongly interacting physics (if any) is **beyond 10 TeV** and/or decoupled
- ⇒ There is a  $\sim 100 \text{ GeV}$  scalar:  
The **Higgs boson**



## Possible solution: Indirect (dynamical) scale generation

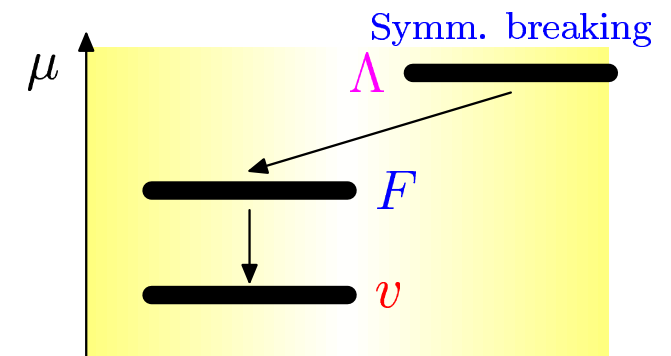
### Supersymmetric Models:

- Field condensation in new (hidden) sector with SUSY  
⇒ SUSY breaking ⇒ Scale generation
- Scalars are present because of SUSY
- Scalar potential by radiative corrections
- Top sector triggers EWSB



### Little Higgs Models:

- Field condensation in new sector with global symmetry  
⇒ spontaneous symmetry breaking ⇒ Scale generation
- Scalars are present because of Goldstone theorem
- Scalar potential by radiative corrections
- Top sector triggers EWSB



**Higgs = Pseudo-Goldstone boson:** Georgi, Pais (1974); Georgi, Dimopoulos, Kaplan (1984); ...

**Problem:** Without fine-tuning,  $v = F \sim \Lambda/4\pi$

⇒ **two-scale model** (like technicolor)

**Three-scale models:** Arkani-Hamed, Cohen, Georgi (2001); ...

$v \sim F/4\pi$  and  $F \sim \Lambda/4\pi \Rightarrow$  **little Higgs**

- Scale  $\Lambda$ :

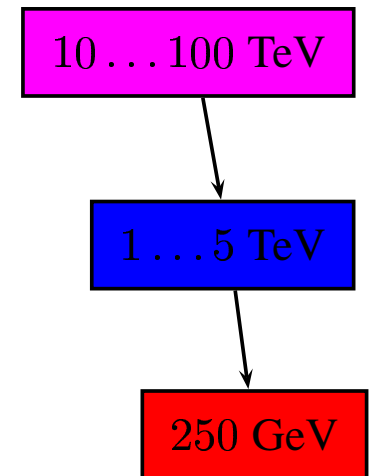
Compositeness, new dynamics

- Scale  $F$ :

Pseudo-Goldstone scalars, extra vector bosons, extra fermions

- Scale  $v$ :

Higgs bosons, known vector bosons and fermions



## How to make a Little Higgs Model:

- Extended gauge theory:  $SM' \otimes \mathcal{G}_{\text{new}} \rightarrow SM$

Compare: **GMSB-type SUSY,  $Z'$  models**

- Enlarged global symmetry:  $SM' \otimes \mathcal{G}_{\text{new}}$  embedded in  $\mathcal{H}$

Compare: **SM custodial symmetry**

- Extended top sector: New vector-like quark(s) coupled to both  $SM'$  and  $\mathcal{G}_{\text{new}}$

Compare: **See-saw topcolor models**

## LHM = Extended Gauge Theory:

Generic symmetry breaking pattern:

$$SM' \otimes \mathcal{G}_{\text{new}} \rightarrow SM$$

Fermions: Charged under SM and SM', neutral under  $\mathcal{G}_{\text{new}}$

Minimal version:

$$\begin{aligned} & (SU(2) \otimes U(1))_1 \otimes (SU(2) \otimes U(1))_2 \\ & \quad (g_1, g'_1) \quad \quad \quad (g_2, g'_2) \\ & \rightarrow (SU(2) \otimes U(1))_{SM} \\ & \quad \quad \quad (g, g') \end{aligned}$$

⇒ Four heavy vector bosons (at least):  $W^{\pm'}$ ,  $Z'$  and  $B'$

**Vector bosons:**  $W^{\pm'}$ ,  $Z'$  and  $B'$

Mass scale:  $F \sim 4\pi v \Rightarrow$  **TeV range**

Couplings:

$$\frac{1}{g^2} = \frac{1}{g_1^2} + \frac{1}{g_2^2} \quad \text{and} \quad \frac{1}{g'^2} = \frac{1}{g_1'^2} + \frac{1}{g_2'^2}$$

$\Rightarrow$  Both  $g_1$  and  $g_2$  are  $> g$

Limits:

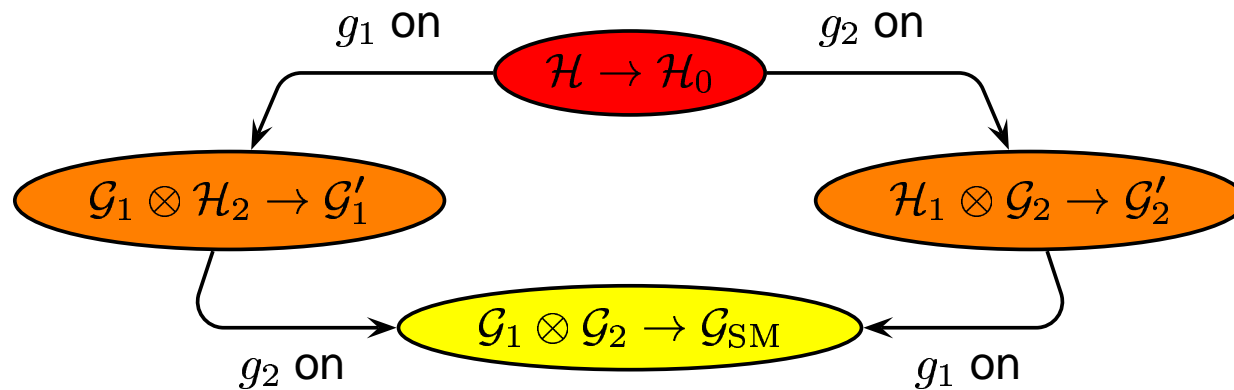
- $g_1 \gg g_2 \Rightarrow W', Z'$  **superheavy** and strongly coupled to fermions
- $g_1 \approx g_2 \Rightarrow W', Z'$  masses of **order  $F$** , standard fermion couplings
- $g_1 \ll g_2 \Rightarrow W', Z'$  **superheavy** and decouple from fermions (**fermiophobic**)

Larger gauge symmetry: Additional heavy vectors, e.g.

$$\mathcal{G}_2 = SU(3) \Rightarrow \text{complex } SU(2) \text{ doublet } V^{\pm}, V^0, \bar{V}^0$$

## Enlarged global symmetry

Embed the gauge group  $\mathcal{G}_1 \otimes \mathcal{G}_2$  in a larger **global symmetry**  $\mathcal{H}$



Switch on gauge couplings  $g_1, g_2$ : Symmetry is reduced

- Some Goldstone bosons of  $\mathcal{H} \rightarrow \mathcal{H}_0$  are eaten  $\Rightarrow W'_L, Z'_L, B'_L$
- Some become **heavy** (mass of order  $F \sim \Lambda/4\pi$ )  $\Rightarrow$  **TeV-scale scalars**
- Some become **light** (mass of order  $v = g_1 g_2 \Lambda / 16\pi^2$ )  $\Rightarrow$  **Higgs boson candidates**

Requirement:  $\mathcal{H}_1$  and  $\mathcal{H}_2$  don't commute

## Radiatively generated Higgs potential

Heavy scalars:  $H$

Light scalars:  $h$  (includes Goldstone bosons  $W_L, Z_L$ )

1-loop Coleman-Weinberg potential for  $H$  and  $h$ : determined by symmetry pattern

$$V(h, H) = aF^2 H^2 + bF(hHh) + ch^4 + \dots$$

Integrate out the  $H$  fields:

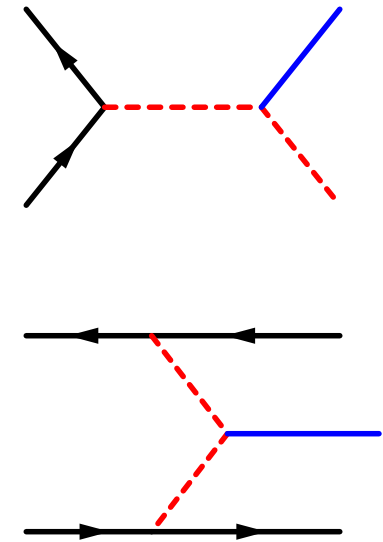
$$V(h) = c'h^4 + \dots$$

⇒  $h$  is massless to one-loop order

⇒ Positive mass<sup>2</sup> at two-loop order

⇒ Trilinear couplings  $Hhh$ :

production and decay channels for heavy scalars



## Extended top quark sector:

Extra massive singlet fermion(s) with top quantum numbers:  $t'_{L,R}$

Yukawa coupling (model dependent):

$$\lambda_1 F \text{Tr}[(\bar{t}_L \bar{b}_L \bar{t}'_L) \times \Sigma(H, h)] t_R + \lambda_2 F \bar{t}'_L t'_R$$

$$\Rightarrow \text{Top Yukawa coupling } \frac{1}{\lambda_t} = \frac{1}{\lambda_1} + \frac{1}{\lambda_2}$$

$t, t'$  loop correction  $\Rightarrow$  negative Higgs mass squared

$$V(h) = -a \frac{\lambda_1^2 F^2}{16\pi^2} h^2$$

$\Rightarrow$  Higgs VEV  $\langle h \rangle = v$

$\Rightarrow$  Top mass  $m_t = \lambda_t v$

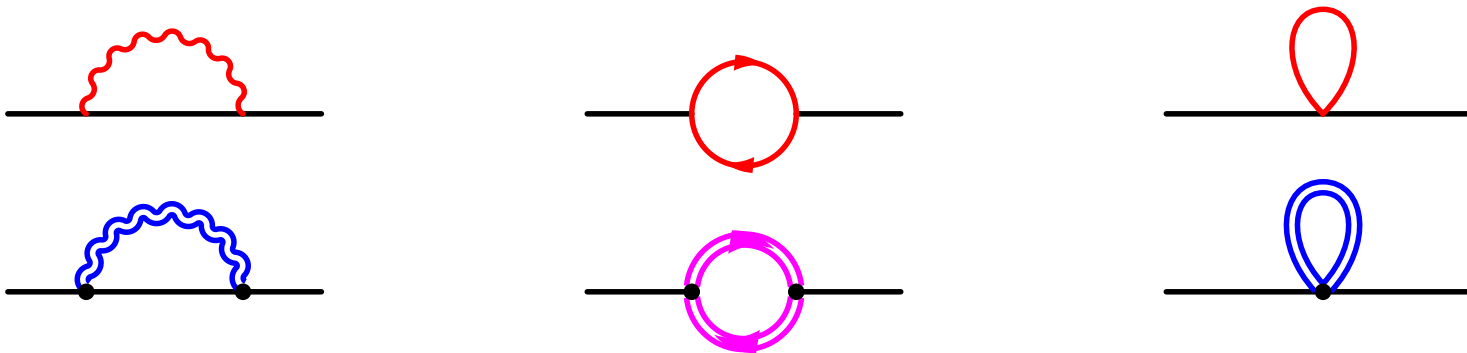
$\Rightarrow$  Top' mass  $m_{t'} = \lambda_2 F$

Yukawa couplings of heavy scalars  $H$ : suppressed by  $v/F$

## Summary: LHM particle content

- **Light Higgs doublet(s)**, possibly extra light scalar multiplets
- **Heavy Higgs multiplets**, coupled to Higgs/Goldstone pairs, decoupled from fermions, mixed with light Higgses (mixing angle  $\sim v/F$ )
- **Heavy vector bosons:**  
 $W^{\pm'}$ ,  $Z'$ ,  $B'$  coupled to SM fermions, mixed with  $W^{\pm}$ ,  $Z$  (mixing angle  $\sim v/F$ )
- **Heavy up-type quark(s):**  $t'$

Effective theory above  $F$ : cancellation of quadratic divergences in Higgs self-energy



Two classes of LHM have been proposed:

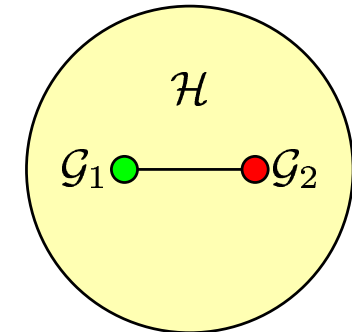
## Simple group:

Extended gauge symmetry  $\mathcal{G}_1 \otimes \mathcal{G}_2$  embedded in **simple group**  $\mathcal{H}$

⇒ exotic Higgs representations,

⇒ low-energy effective theory can be the **minimal SM**

**Minimal version:**  $(SU(2) \otimes U(1))^2$  embedded in  $SU(5)$



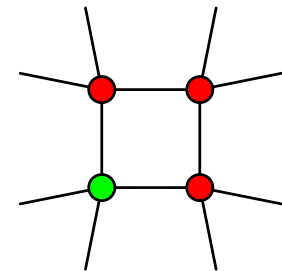
## Theory space (Moose Models):

Extended gauge symmetry  $(\mathcal{G}_{\text{SM}})^n$  embedded in **semisimple group**  $\mathcal{G}^{4n}$

⇒ many copies of scalar (and vector) multiplets

⇒ low-energy effective theory contains at least **two Higgs doublets**

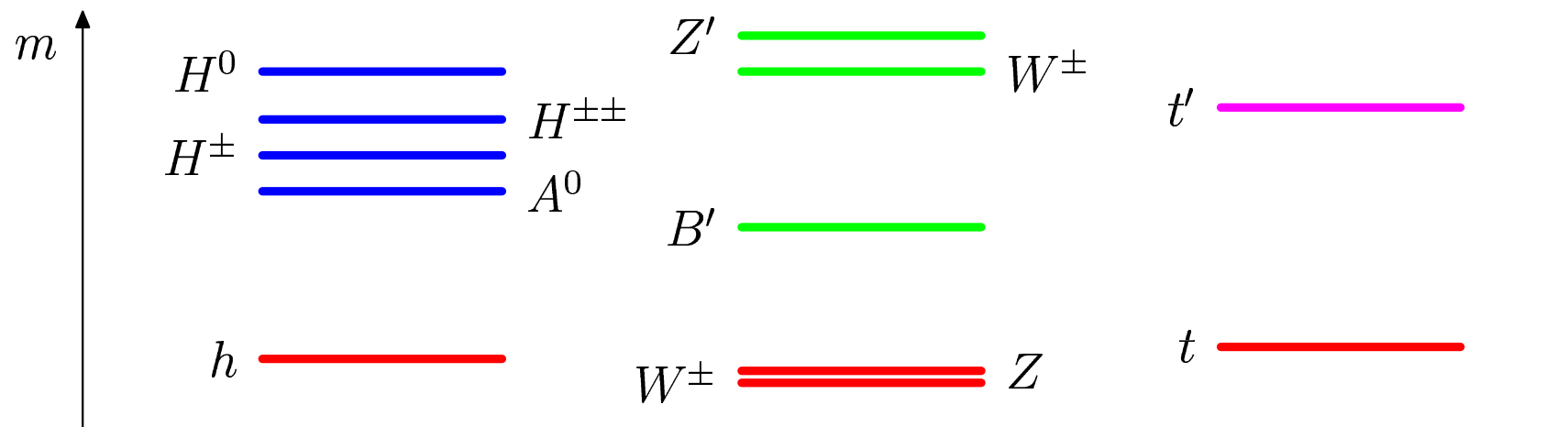
**Minimal version:**  $(SU(2) \otimes U(1))^2$  embedded in  $SU(3)^8$



## The Littlest Higgs: Arkani-Hamed/Cohen/Katz/Nelson (2002)

- Global symmetry breaking  $SU(5) \rightarrow SO(5) \Rightarrow 24 - 10 = 14$  Goldstone bosons
- Gauge symmetry breaking  $(SU(2) \otimes U(1))^2 \rightarrow SU(2) \otimes U(1) \Rightarrow W_L^{\pm'}, Z_L', B_L'$
- 10 scalars remain:
  - Complex doublet (light)  $\Rightarrow$  Higgs boson  $h$  and  $W_L^{\pm}, Z_L$  with  $\langle h \rangle = v$
  - Complex triplet (heavy)  $\Rightarrow H, A, H^{\pm}, H^{\pm\pm}$  with  $\langle H \rangle \sim v^2/F$

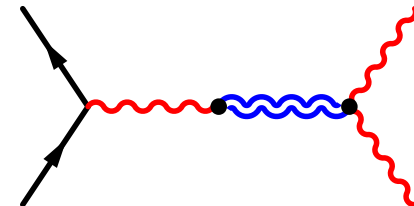
Spectrum (example):



## Heavy vectors

Vector bosons with TeV scale masses:

- \*  $B'$  significantly lighter than  $W', Z'$ , production at LC 800?
  - ⇒ Resonance in  $e^+e^- \rightarrow f\bar{f}$  and  $qq \rightarrow qq$  (→ Tevatron Run 2)
  - ⇒ Fermion couplings suppressed if  $g'_1 < g'_2$ , enhanced if  $g'_1 > g'_2$
  - ⇒ Hypercharge boson: coupling to  $W^+W^-$  suppressed
  - ⇒ Measurements of mass, width and BRs determine scale  $F$ , quantum numbers and gauge sector parameters
- \*  $W'^{\pm}$  and  $Z'$  are heavier
  - ⇒ Indirect constraints from contact interactions
  - ⇒ Effect on triple gauge couplings
  - ⇒ GigaZ



## Heavy scalars

Scalars with TeV scale masses, unsuppressed couplings to light Higgs and Goldstones:

\*  $A^0$  could be the lightest state

⇒ Associated production  $e^+e^- \rightarrow Ah$

⇒ Fermion couplings suppressed ⇒ decay into  $hZ_L$

\* Charged Higgs  $H^\pm$

⇒ Single production in  $e^+e^- \rightarrow H^+W_L^-$ , decay into  $W_L Z_L$

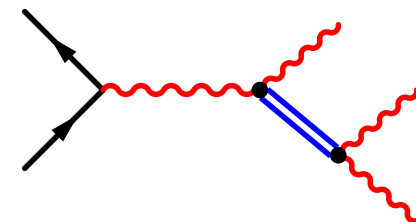
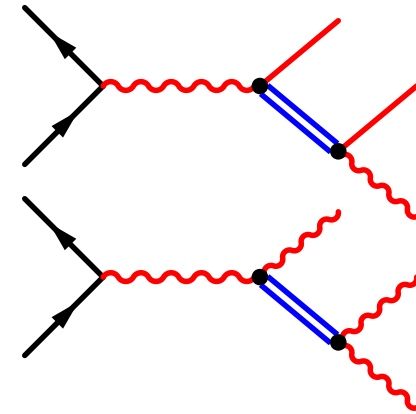
\* Doubly charged Higgs  $H^{\pm\pm}$

⇒ Single production in  $e^-e^- \rightarrow \nu\nu H^{--}$ , decay into  $W_L^- W_L^-$

\* Heavy Higgs  $H^0$

⇒ Higgs-strahlung  $e^+e^- \rightarrow Z_L H$

⇒ Decay into  $Z_L Z_L$  and  $hh$ , but not  $W^+W^-$



... but they may be too heavy ⇒ indirect constraints from Higgs,  $Z$  and  $W$  observables

## Light scalars

Higgs boson  $h$  mixes with heavy scalar  $H$ , mixing angle  $\sim v/F$ :

⇒ Light Higgs  $h$  mass significantly decreased

⇒ Yukawa couplings slightly **decreased**

⇒ **Anomalous**  $hVV$  and  $hhh$  **couplings**

Precise measurements of all Higgs couplings very important!

## Heavy quarks

Top quark  $t$  mixes with **heavy fermion**  $t'$

\*  $t'$  **production** at LHC?

\* Decay into  $bW_L$  and  $th$

Linear collider: Precise measurements of  $t$  quark couplings!

## Precision observables

Violation of the custodial  $SU(2)_C$  symmetry  $\Rightarrow$  contributions to the  $\rho$  parameter  $O(v^2/F^2)$

- \* extra hypercharge boson  $B'$
- \* Higgs triplet
- \* extended top quark sector
- \* [ $m_h > 100$  GeV]

## Light hypercharge boson $B'$

- \* Limits from Drell-Yan processes / contact terms at Tevatron, LEP2

## Extra $t'$

- \* Contributions to  $t, b$  couplings and  $R_b$

## Present limits:

Csaki/Hubisz/Kribs/Meade/Terning (2002)

$$M_{B'} > 650 \text{ GeV} \quad \text{and} \quad M_{W'} > 2.7 \text{ TeV}$$

Caveat: **Top sector neglected**

Hewett/Petriello/Rizzo (2002)

$$F > 3.5 \text{ TeV} \quad \text{and} \quad m_{t'} > 10 \text{ TeV}$$

Caveat: **Scalar sector neglected**

⇒ Need to take into account **all contributions** (cancellations!)

⇒ Favored: Models with

- Different **hypercharge embedding**:  $B'$  heavier
- Custodial  **$SU(2)$  symmetry**:  $SU(2) \otimes SU(2) \subset \mathcal{G}_2$ ?
- Fermiophobic vector bosons

Han/Logan/McElrath/Wang (2003)

**The Minimal Moose Model:** Arkani-Hamed/Cohen/Katz/Nelson/Gregoire/Wacker (2002)

- Global symmetry breaking:

$$(SU(3) \otimes SU(3))^4 \rightarrow SU(3)^4$$

with four independent scalar multiplets  $\Rightarrow$  32 Goldstone bosons

- Gauge symmetry breaking:

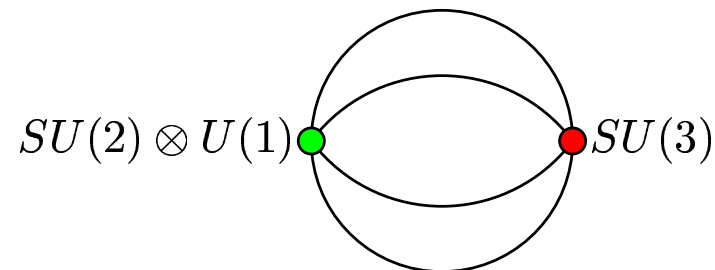
$$SU(3) \otimes SU(2) \otimes U(1) \rightarrow SU(2) \otimes U(1)$$

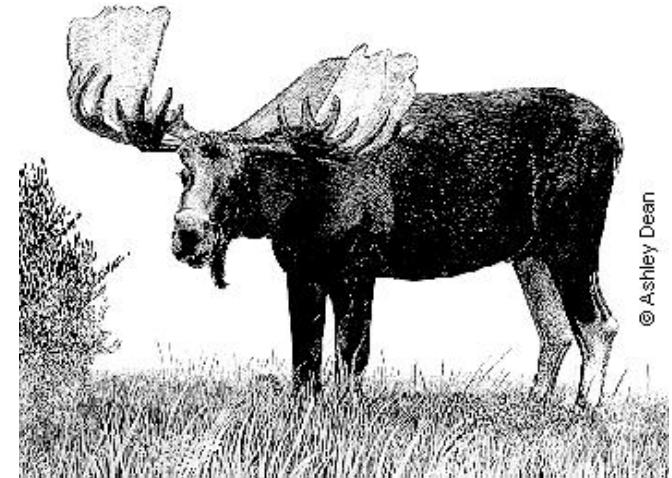
$\Rightarrow$  Heavy vectors  $W^{\pm'}$ ,  $Z'$ ,  $B'$  and  $V^{\pm}$ ,  $V^0$ ,  $\bar{V}^0$

- Light scalars: 2 Higgs doublets,  
2 real triplets, 2 real singlets

- Heavy scalars:  
1 singlet, 1 doublet, 1 triplet

- Top sector: Vector-like heavy fermion(s)





## Phenomenology of the Minimal Moose:

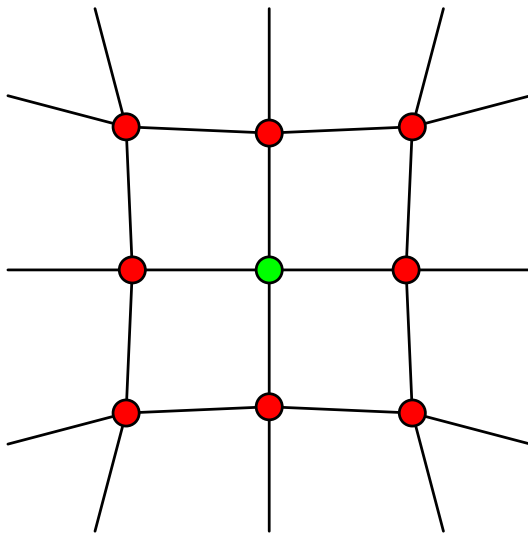
- \*  $B'$  is candidate for the **lightest new vector boson**
- \* Rich **low-energy Higgs phenomenology** (100...1000 GeV)
  - Two Higgs doublets:  $h, H, A^0, H^\pm$
  - Scalar triplet  $\phi^\pm, \phi^0$  and singlet  $\eta^0$  bosons (no VEV)
- \* Additional **heavy Higgs states** (mixing with light Higgs)
- \* Anomalous **top sector**

## Theory space

Multiple gauge groups: *sites*

Goldstone multiplets, charged under two gauge groups: *links*

Spontaneous symmetry breaking in the Goldstone sectors  $\Rightarrow$  breaking of  $\mathcal{G}^n \rightarrow \text{SM}$



## Little Higgs Models:

- Extra links  $\Rightarrow$  closed cycles
  - Two-dimensional lattice  $\Rightarrow$  unsuppressed quartic couplings  $\Rightarrow$  EWSB by Goldstone VEV
- $\Rightarrow$  Many multiplets of scalars and vector bosons
- $\Rightarrow$  Spectrum and interactions similar to 6D extra-dimension model  
(Scalars = 5-6 components of 6D gauge fields)

## What should we look for?

### Direct observation

- LHC/LC: Light Higgs boson?
- LHC: Heavy vector bosons, heavy colored fermions
- LHC/LC: Uncover Higgs sector completely (how many sub-TeV scalars are there?)

### Precision observables

- Tevatron/LHC/LC: Contact terms  $\Rightarrow$  constrain four-fermion operators
- LHC/LC: Triple gauge couplings (fermiophobic vector bosons)
- LHC/LC: Quartic gauge couplings (vector bosons and heavy scalars)  
 $\Rightarrow$  Analysis similar to strong-WW scattering, but in the presence of light Higgs
- LC: Higgs self-coupling(s), quartic Higgs-vector boson couplings

## More phenomenology:

- Grand Unification:

- The SM  $SU(2) \otimes U(1)$  symmetry is continued above the TeV scale. If the extra  $\mathcal{G}_2$  symmetry is strongly coupled, the shifts in the SM couplings are small
- ⇒ The SM mismatch in unification could be a consequence of  $\mathcal{G}_2$
- ⇒ Coupling sum rule: The GUT scale is lowered ... unless  $SU(3)_{\text{QCD}}$  is also affected.

- Dark Matter:

- If there is an unbroken discrete symmetry in the LHM, there will be a dark matter candidate. (Not in the minimal model.)
- ⇒ Heavy neutral scalar?

- Supersymmetry:

- LHM symmetry structure could come with SUSY breaking in the 10 TeV range
- ⇒ Superpartners much more heavy?

## Conclusions:

- Little Higgs Models = indirect EW scale generation: **symmetry breaking in new sector** (analogous to SUSY models)
- Special role of **top quark**
- ... but no solution to flavor problem Chivukula/Evans/Simmons (2002)
- Phenomenology:
  - **Light Higgs boson** (multiple light states possible)
  - **Heavy scalar states**
  - **Heavy vector bosons**: recurrences of  $W, Z, \gamma$  (and others)
  - **Heavy quarks** mixed with 3rd generation
- Significant sensitivity of precision observables, even if no new states accessible  
⇒ **get a complete and consistent picture of all Higgs and vector boson couplings!**