

COSMOLOGIA

In che Universo viviamo?

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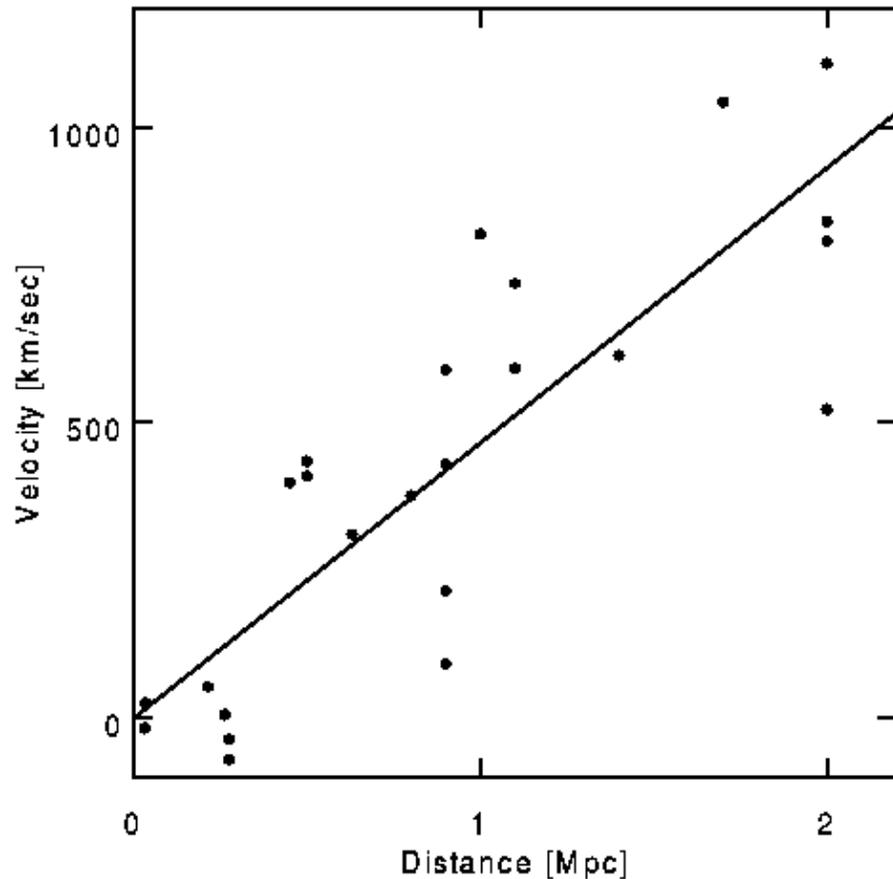


La legge di Hubble

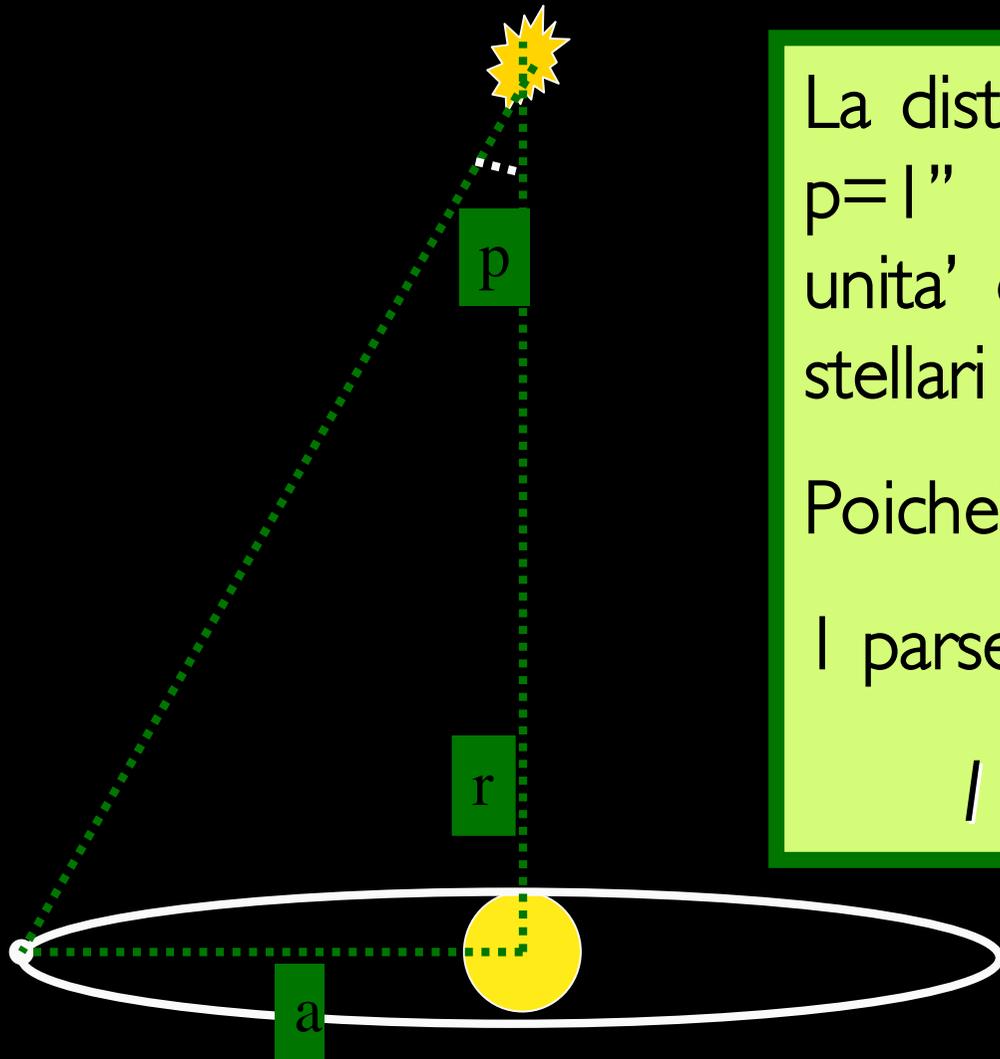
$$V = H_0 r$$



$$H_0 = 500 \frac{\text{km/s}}{\text{Mpc}}$$



Il Megaparsec



La distanza r corrispondente a $p=1''$ e' stata assunta come unita' di misura delle distanze stellari

Poiche' $a = 1.5 \cdot 10^{13}$ cm

1 parsec $= 3 \cdot 10^{18}$ cm

1 Mpc $= 3 \cdot 10^{24}$ cm

The Hubble Key Project

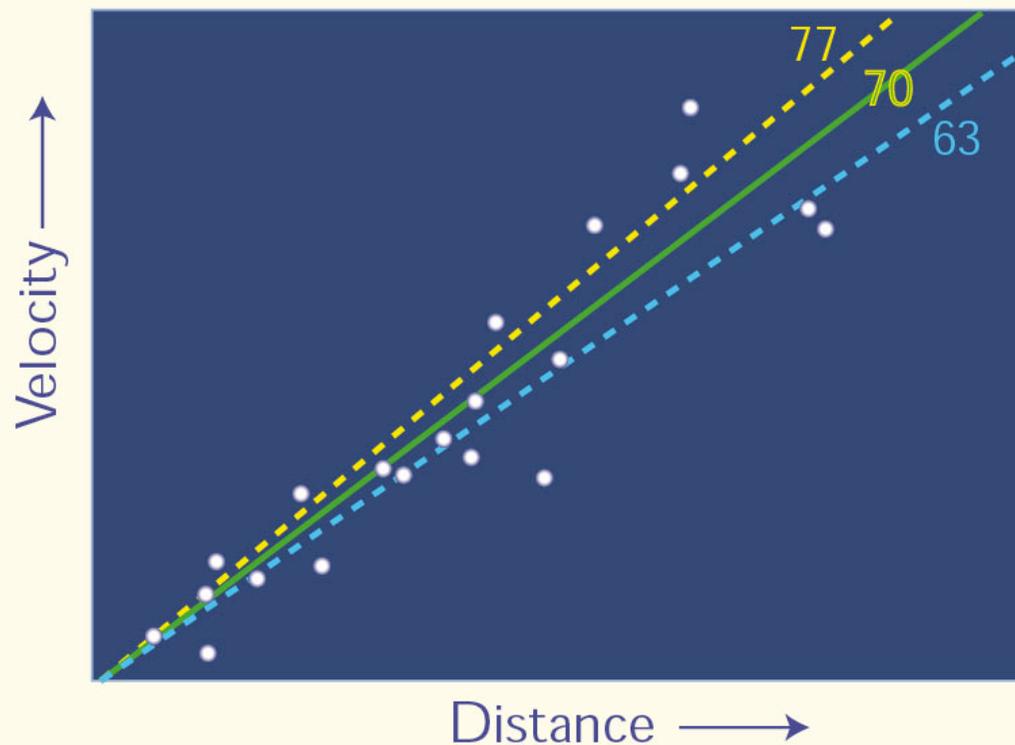
...to measure H_0 to 10% accuracy.....

HUBBLE CONSTANT

$$H_0 = (70 \pm 7) \frac{\text{km/s}}{\text{Mpc}}$$

W. Freedman *et al.* (1999)

Hubble Diagram for Cepheids



Un universo in espansione



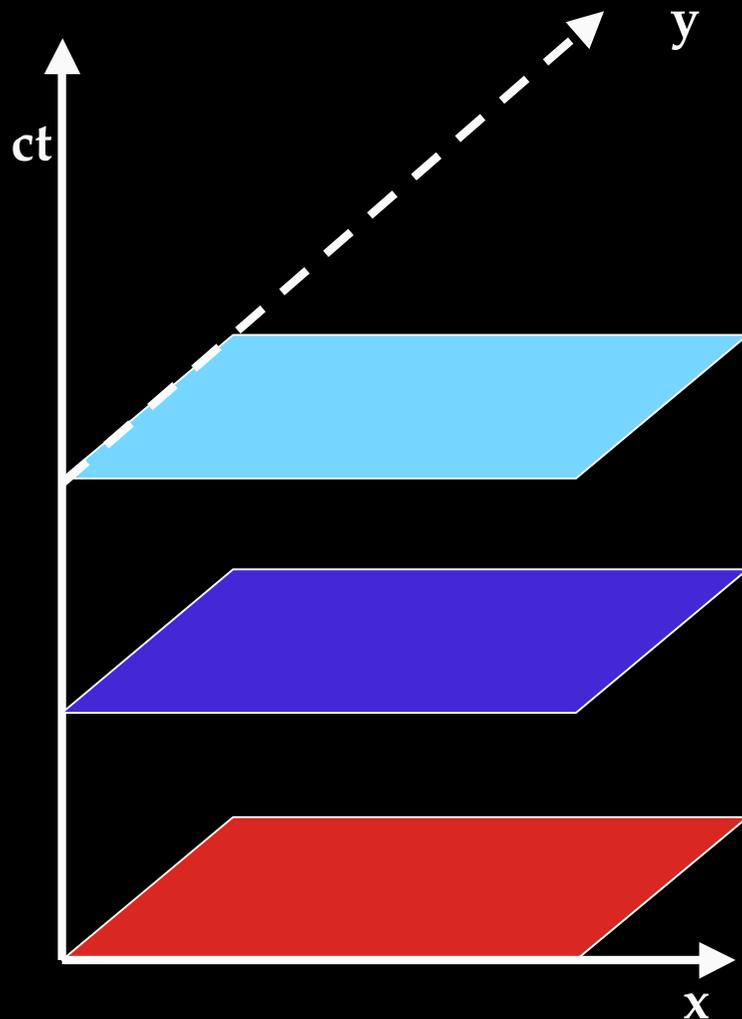
On the curvature of space

A.Friedmann, 1922

“The purpose of this note is ... to demonstrate the possibility of a world in which the curvature of space is independent of the three spatial coordinates, but does depend on time”

Aleksandr Aleksandrovich Friedman: 1888- 1925

Modelli cosmologici



Coordinate gaussiane

Gauge sincrona

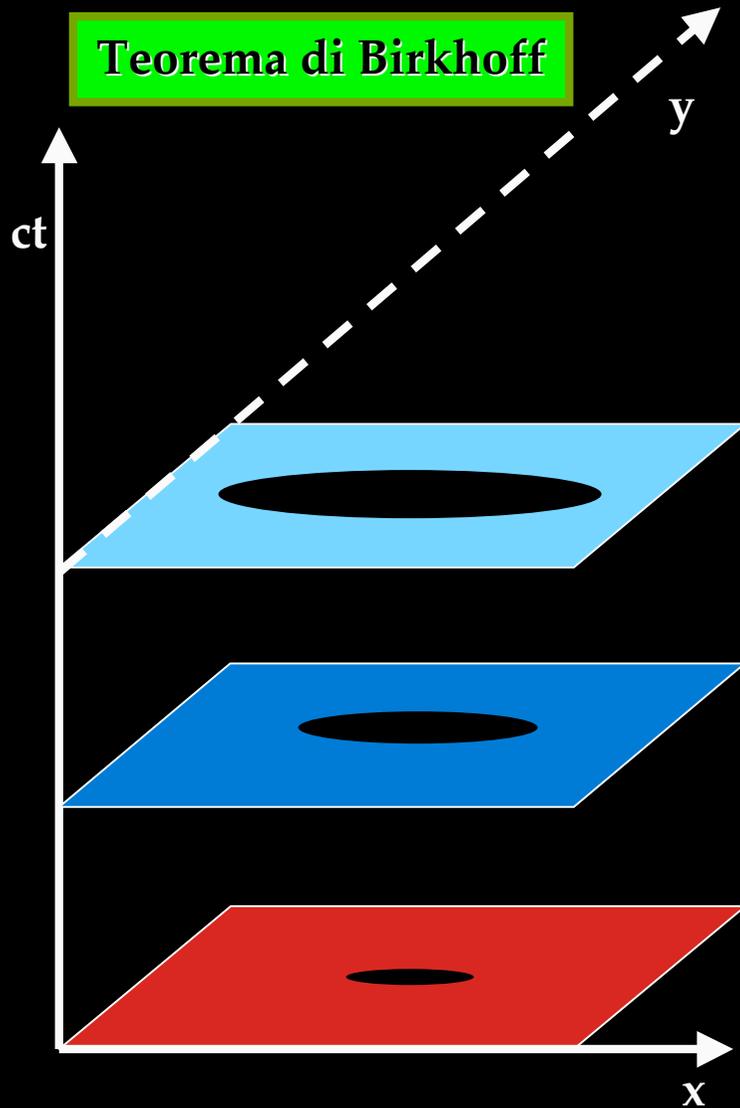
Coordinate comobili

$$g_{00} = 1; g_{0k} = 0;$$

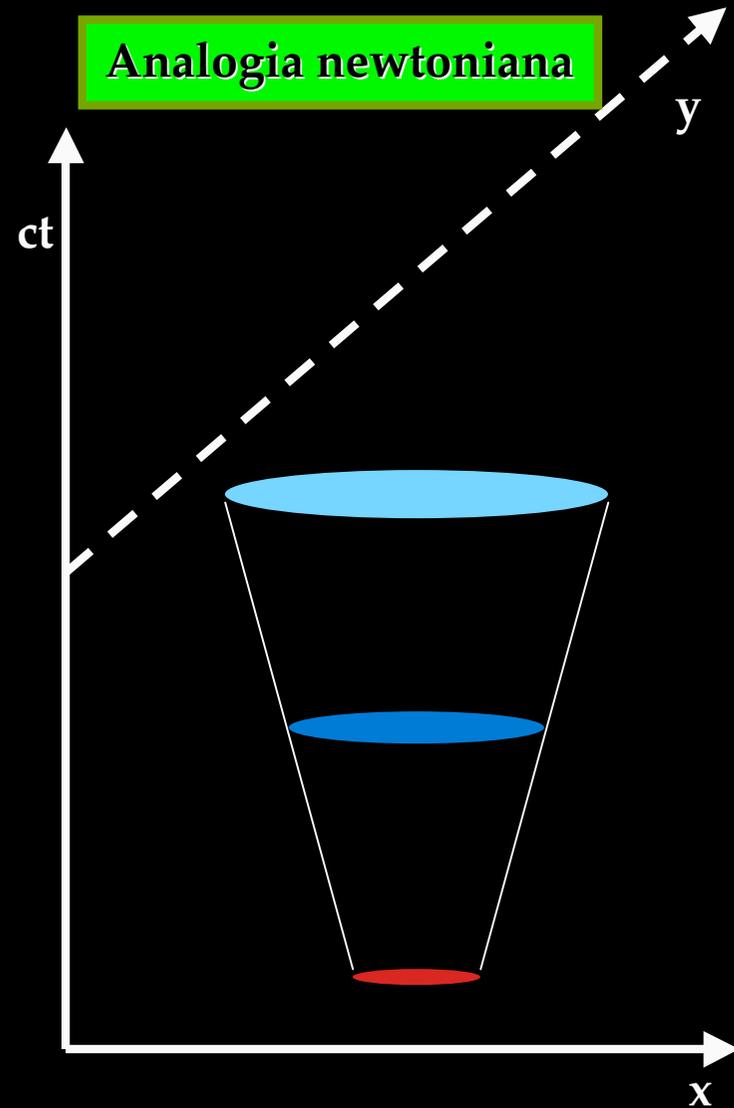
$$ds^2 = dx^{02} \square R^2(t) \left[\frac{dr^2}{\square kr^2} + r^2 d\square^2 \right] \quad k = \begin{matrix} +1 \\ 0 \\ -1 \end{matrix}$$

Modelli Cosmologici

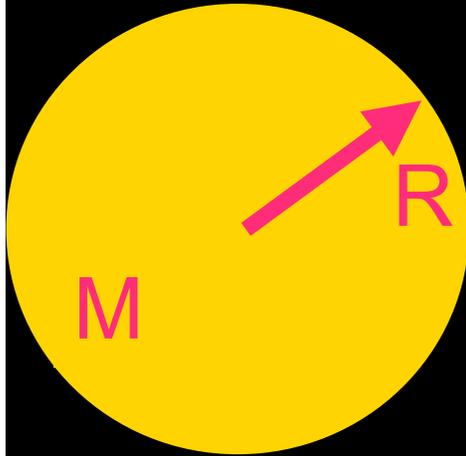
Teorema di Birkhoff



Analoga newtoniana



Modelli cosmologici



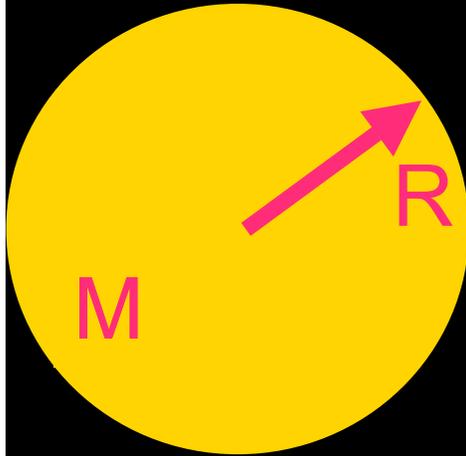
$$\ddot{R} = -\frac{GM}{R^2} \quad \square \quad \frac{1}{2}\dot{R}^2 - \frac{GM}{R} = e$$

$$M = \frac{4}{3}\pi R^3$$

$$\frac{\ddot{R}}{R} = -\frac{4}{3}\pi G \rho \quad \square \quad \frac{\frac{1}{2}\dot{R}^2}{R} - \frac{2e}{R^2} = \frac{8}{3}\pi G \rho$$

$$\frac{\ddot{R}}{R} = -\frac{4}{3}\pi G \rho \quad \square \quad \frac{\frac{1}{2}\dot{R}^2}{R} + \frac{kc^2}{R^2} = \frac{8}{3}\pi G \rho$$

Modelli cosmologici



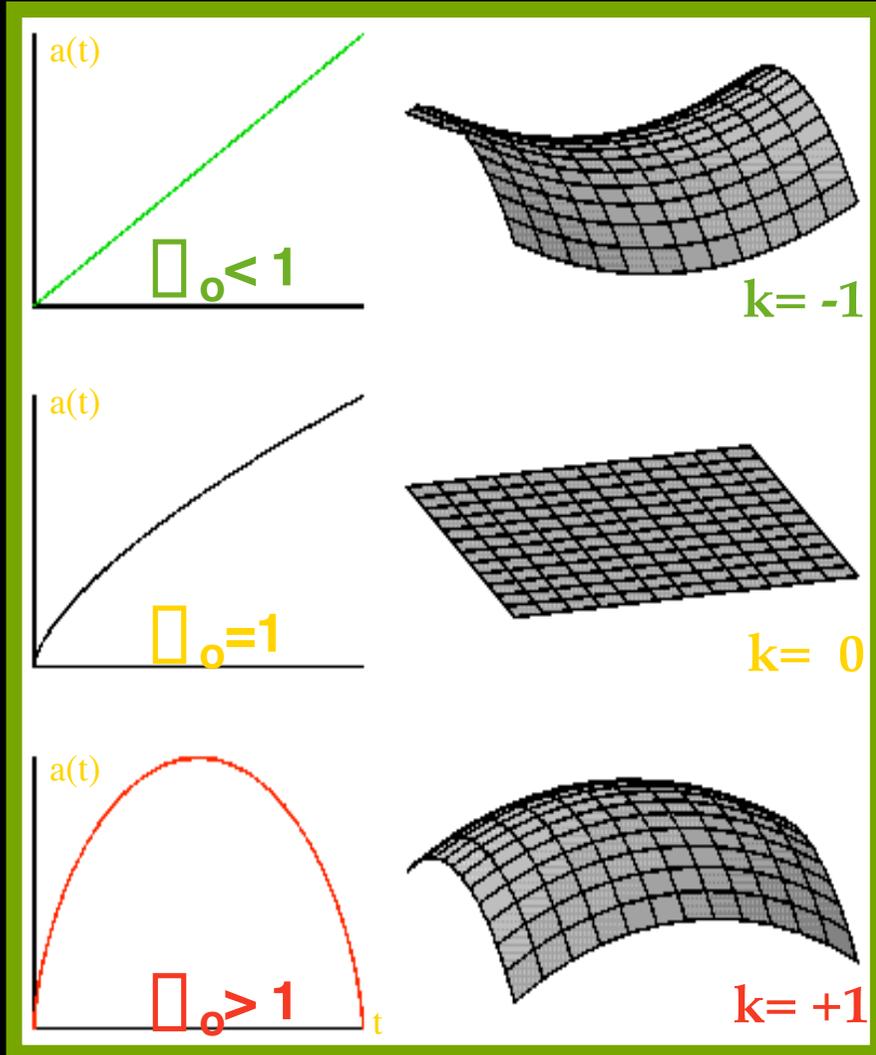
$$\left. \frac{|U|}{T} \right|_0 = \left. \frac{GM/R}{\dot{R}^2/2} \right|_0 = \frac{8\pi G \rho_0 R^2}{3 \dot{R}^2} \bigg|_0 = \frac{8\pi G \rho_0}{3H_0^2} = \frac{\rho_0}{\rho_{crit}} = \Omega_0$$

$$\rho_{crit} = 2 \times 10^{29} h^2 \frac{g}{cm^3}; \quad h = \frac{H_0}{100 \frac{km s^{-1}}{Mpc}}$$

$$\Omega_0 = 2e = kc^2$$

k	Ω_0	E
+1	>1	<0
0	=1	=0
-1	<1	>0

Cosmologia relativistica



$$a(t) = \frac{\Omega_o/2}{1 - \Omega_o} [\cosh \sqrt{\Omega_o} t]$$

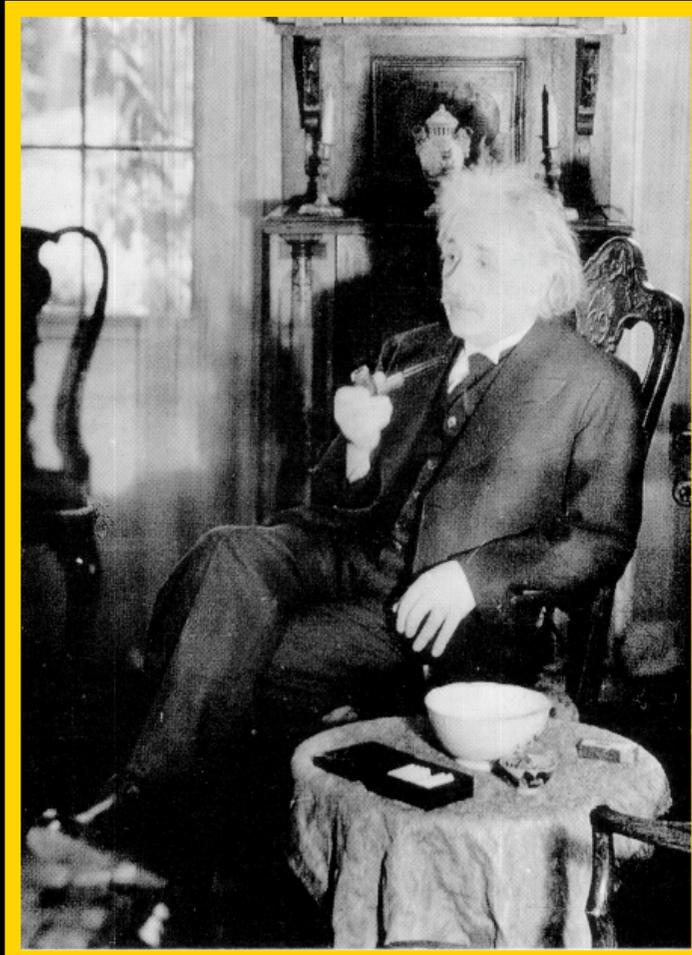
$$H_o t = \frac{\Omega_o/2}{(1 - \Omega_o)^{3/2}} [\sinh \sqrt{\Omega_o} t]$$

$$a(t) = \left[\frac{t - t_o}{t_o} \right]^{2/3}$$

$$a(t) = \frac{\Omega_o/2}{1 - \Omega_o} [1 - \cos \sqrt{\Omega_o} t]$$

$$H_o t = \frac{\Omega_o/2}{(1 - \Omega_o)^{3/2}} [\sqrt{\Omega_o} \sin \sqrt{\Omega_o} t]$$

Costante cosmologica: Λ



“Cosmological considerations
on the general theory of
relativity

A.Einstein, 1917

“In order to arrive at this consistent view, we admittedly had to introduce **an extension** of the field equations of gravitation **which is not justified**”

Costante cosmologica: Λ

- ◆ Modifica delle equazioni di campo della relatività generale

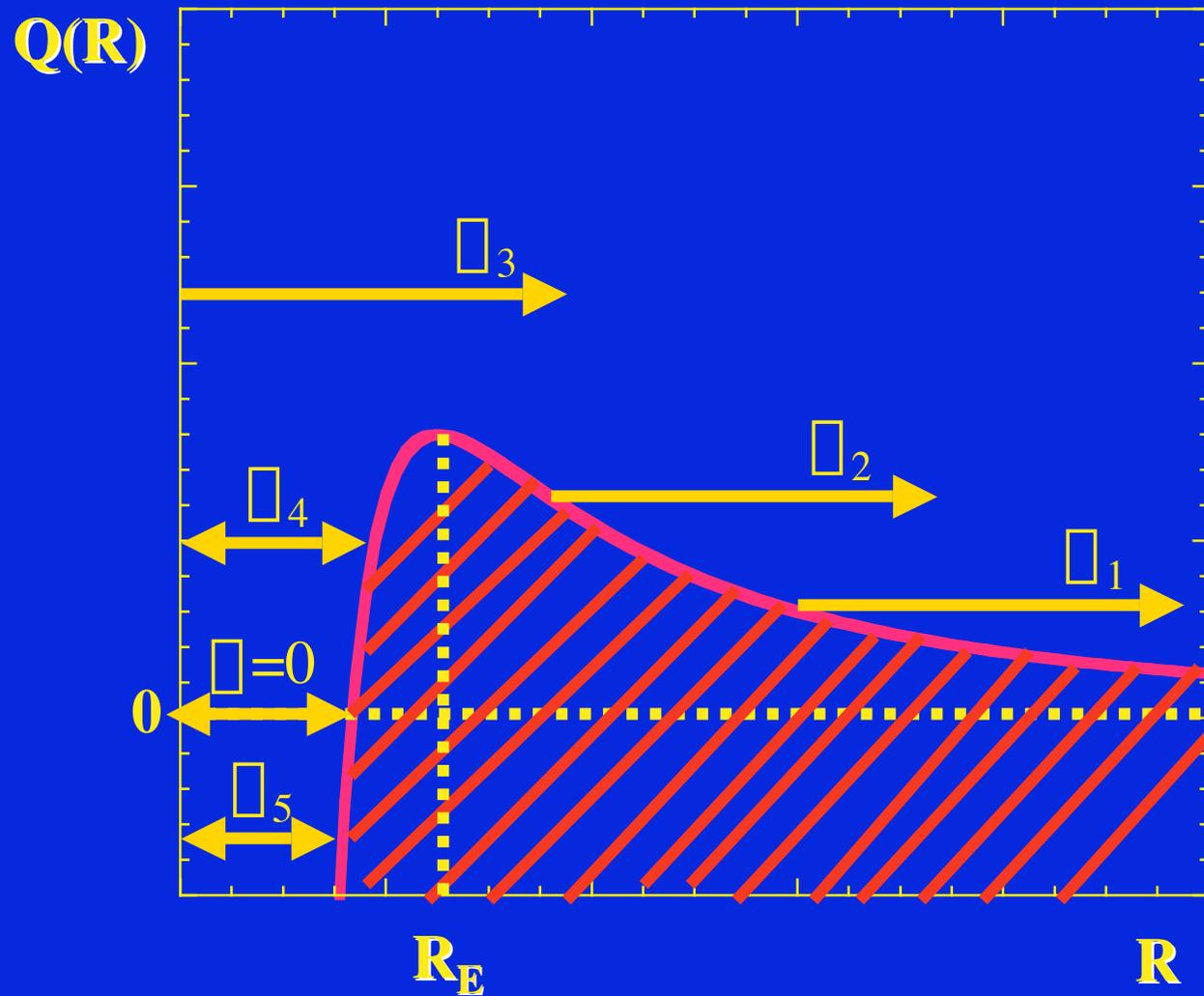
$$\frac{\dot{R}^2}{R^2} + \frac{kc^2}{R^2} = \frac{8}{3}\rho G + \frac{1}{3}\rho c^2 \quad \frac{\dot{R}^2}{R^2} = \frac{c^2}{3}[\rho - Q(R)]$$

$$\rho \geq Q(R) \equiv \frac{3k}{R^2} \left[\frac{8\rho G}{c^2} \right] +$$

- ◆ E' come avere un fluido con

$$\tilde{\rho} = \rho + \frac{c^2\rho}{8\rho G}; \quad \tilde{p} = p - \frac{c^2\rho}{8\rho G};$$

Costante Cosmologica: $k=+1$



\square & $k=+1$

Consideriamo un universo di polvere, $p = 0$. Abbiamo visto che

$$Q(R) = \frac{3}{R^2} \square \frac{8\square G}{c^2} \square$$

ha un massimo in corrispondenza di

$$\frac{1}{R_E^2} = \frac{4\square G}{c^2} \square_0$$

Scegliamo

$$\square_E = Q(R_E) = \frac{3}{R_E^2} \square \frac{8\square G}{c^2} \square_0 = \frac{1}{R_E^2}$$

Otteniamo la soluzione statica di Einstein

$$\square \frac{\dot{R}}{R} \square_E^2 = \frac{c^2}{3} [\square_E \square Q(R_E)] = 0;$$

$$\square \frac{\ddot{R}}{R} \square_E^2 = \square \frac{4\square G}{3} \square_0 = \square \frac{4\square G}{3} \square_0 + \frac{1}{3} \square_E c^2 = 0$$

ρ & $k=+1$

Assumiamo come ordine di grandezza $\rho_0 \approx 10^{30} \text{ g cm}^3$; allora

$$R_E^2 = \frac{c^2}{4\rho G \rho_0} \approx 3 \cdot 10^{28} \text{ cm} \quad \rho_E = \frac{1}{R_E^2} \approx 10^{57} \text{ cm}^{-2}$$

La massa dell'universo di Einstein e'

$$M = 2\rho^2 \rho_0 R_E^3 \approx 5 \cdot 10^{56} \text{ g}$$

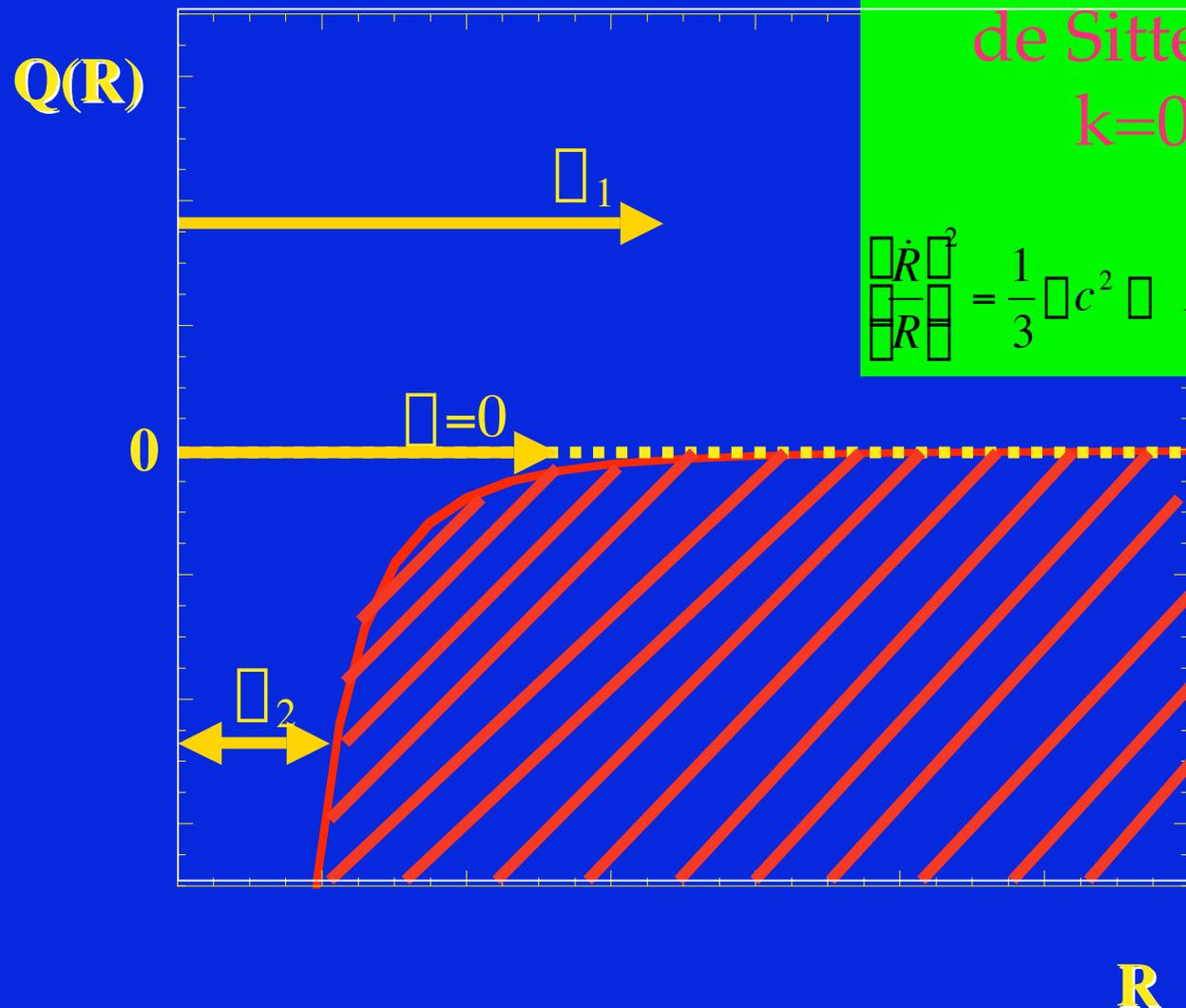
La densita' media del sistema solare (SS) e'

$$\rho_{SS} \approx \frac{3M_{Sole}}{4\rho R_{Plutone}^3} = \frac{3}{4\rho} \frac{2 \cdot 10^{33} \text{ g}}{(40 \cdot 1.5 \cdot 10^{13} \text{ cm})^3} \approx 2.2 \cdot 10^{12} \text{ g/cm}^3$$

L'influenza di ρ sulla dinamica del sistema solare e' trascurabile :

$$\frac{\rho_0}{\rho_{SS}} \approx 5 \cdot 10^{19}$$

Costante Cosmologica: $k=0, -1$



Equazione di Friedman

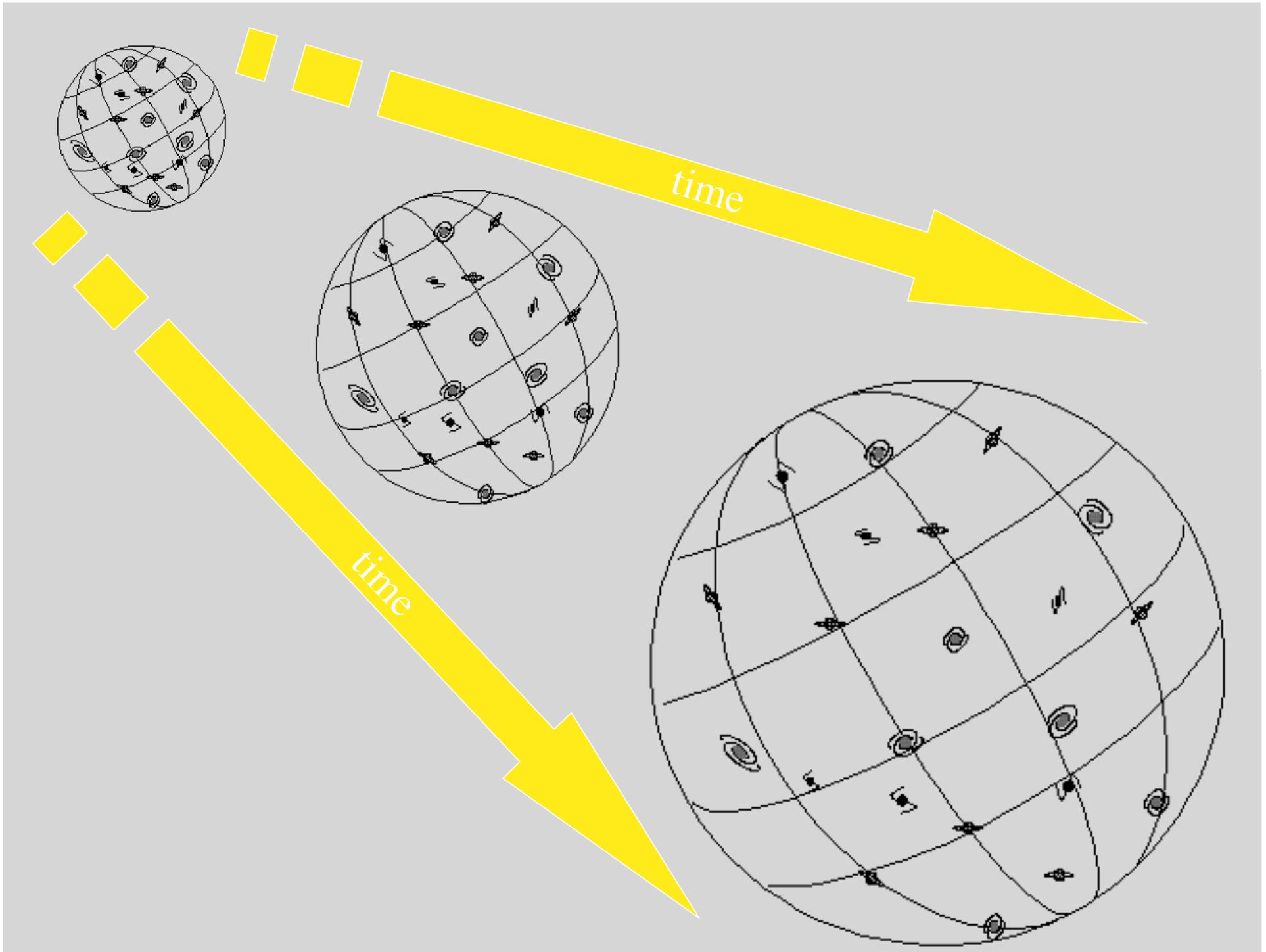
$$\left(\frac{\dot{R}}{R}\right)^2 + \frac{kc^2}{R^2} = \frac{8\pi G}{3} \frac{\rho_0}{R^3} + \frac{1}{3} \Lambda c^2$$

Fase accelerata:
domina la costante
cosmologica

Fase cinematica:
domina la curvatura

Tasso
di
espansione

Fase decelerata:
domina la materia



Problema della piattezza

Per definizione

$$\Omega(t) \equiv \frac{8\pi G}{3H^2(t)} \rho(t)$$

Dall'equazione di Friedman, con $\Lambda = 0$ e $a(t) = R(t)/R(t_0)$, si ottiene

$$H^2(t) \equiv \left(\frac{\dot{R}}{R}\right)^2 = H_0^2 \frac{\Omega_0}{a^3(t)} + \frac{1 - \Omega_0}{a^2(t)}$$

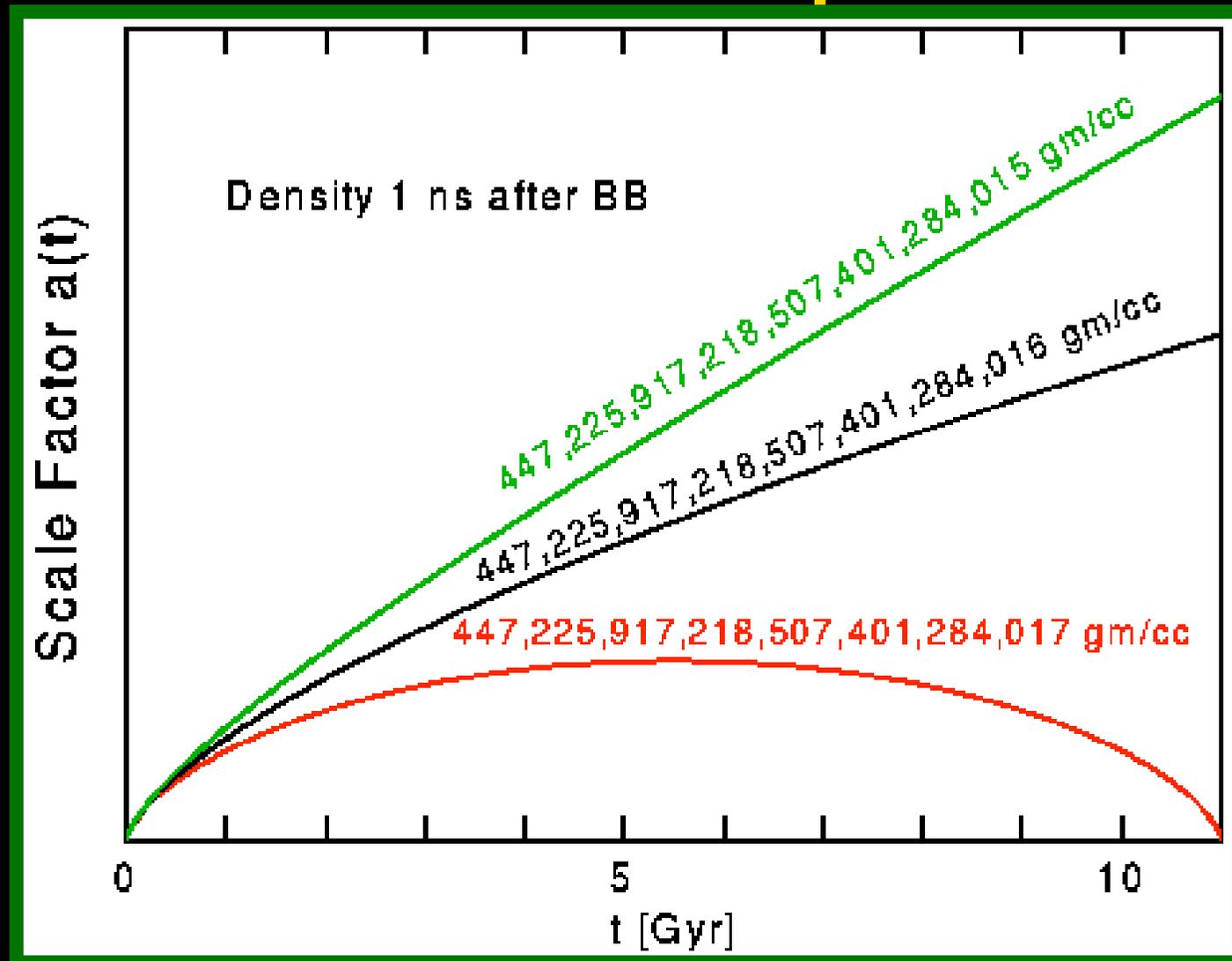
si ricava

$$\frac{1}{\Omega(t)} - 1 = \frac{1}{\Omega_0} - 1 \frac{1}{a(t)}$$

Quindi in un universo di Friedman

$$\lim_{t \rightarrow \infty} \Omega(t) = 1$$

Problema della piattezza



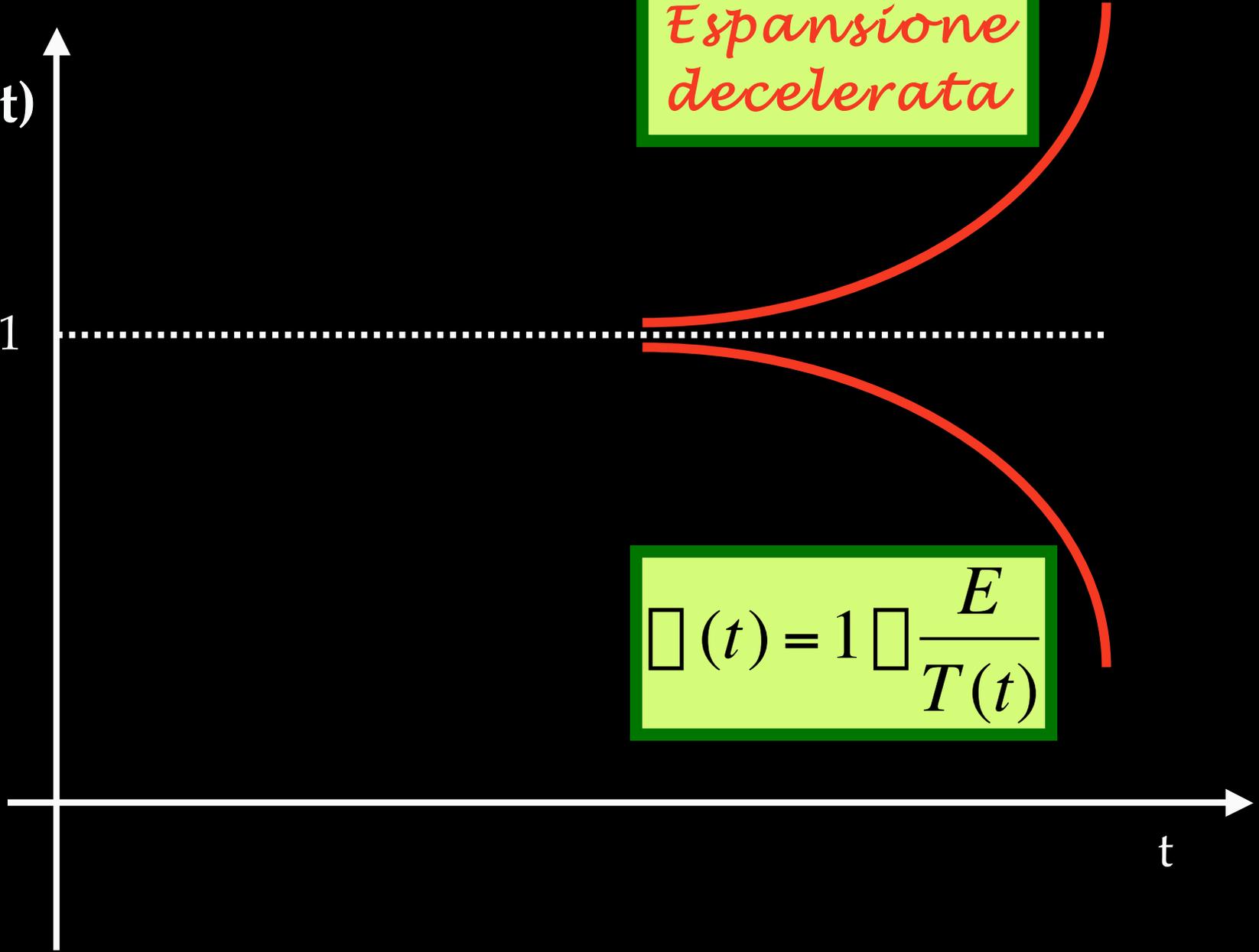
$\square(t)$

1

*Espansione
decelerata*

$$\square(t) = 1 \square \frac{E}{T(t)}$$

t



$\square(t)$

1

*Espansione
decelerata*

*Espansione
accelerata*

$$\square(t) = 1 \square \frac{E}{T(t)}$$

t

