



# **The average phase factor from chiral perturbation theory**

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**Niels Bohr Institute**

**Carlsberg fellow**





**What** Analytic computation of  $\langle e^{2i\theta} \rangle$        $\mu < m_\pi/2$

**Why** Understand numerical QCD at  $\mu \neq 0$

**How** Chiral Perturbation Theory       $T < T_c$

Toussaint NPPS 17 (1990) 248      de Forcrand Laliena PRD 61 (2000) 034502

Sasai Nakamura Takaishi NPPS 129 (2004) 539      Ejiri PRD 73 (2006) 054502

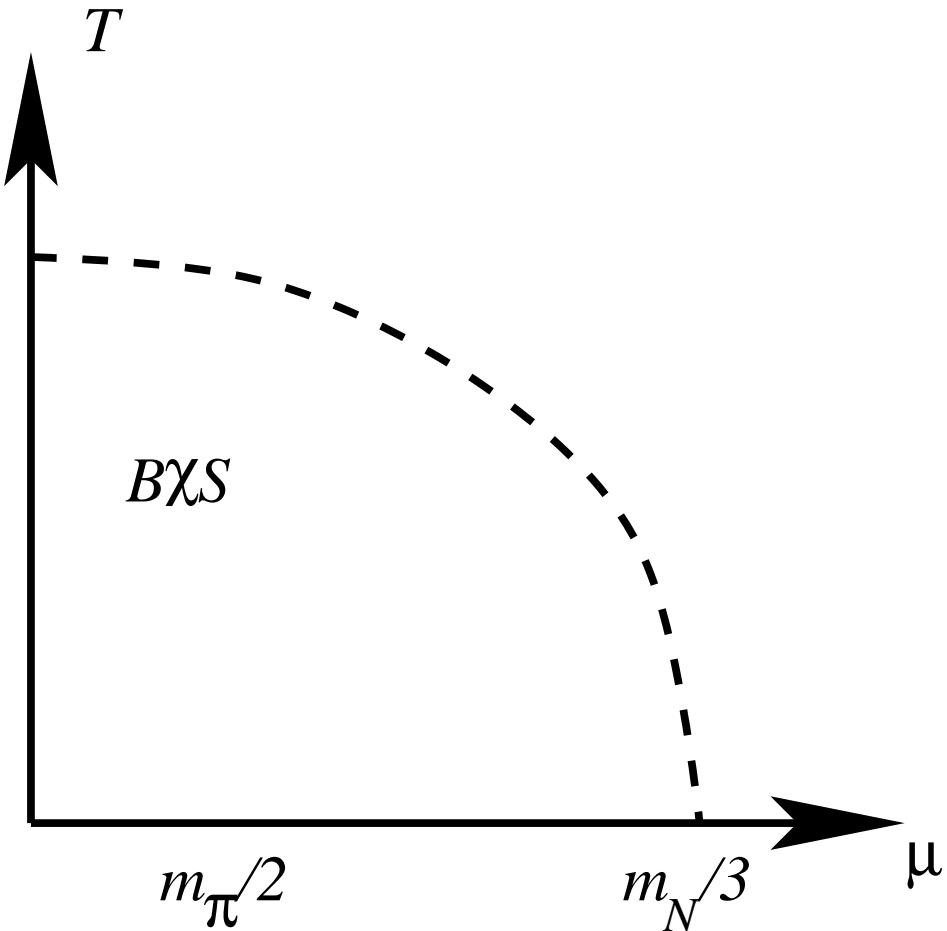
Allton Ejiri Hands Kaczmarek Karsch Laermann Schmidt Scorzato PRD 66 (2002) 074507

Schafer PRD 57 (1998) 3950      Fodor Katz Schmidt JHEP 03 (2007) 121

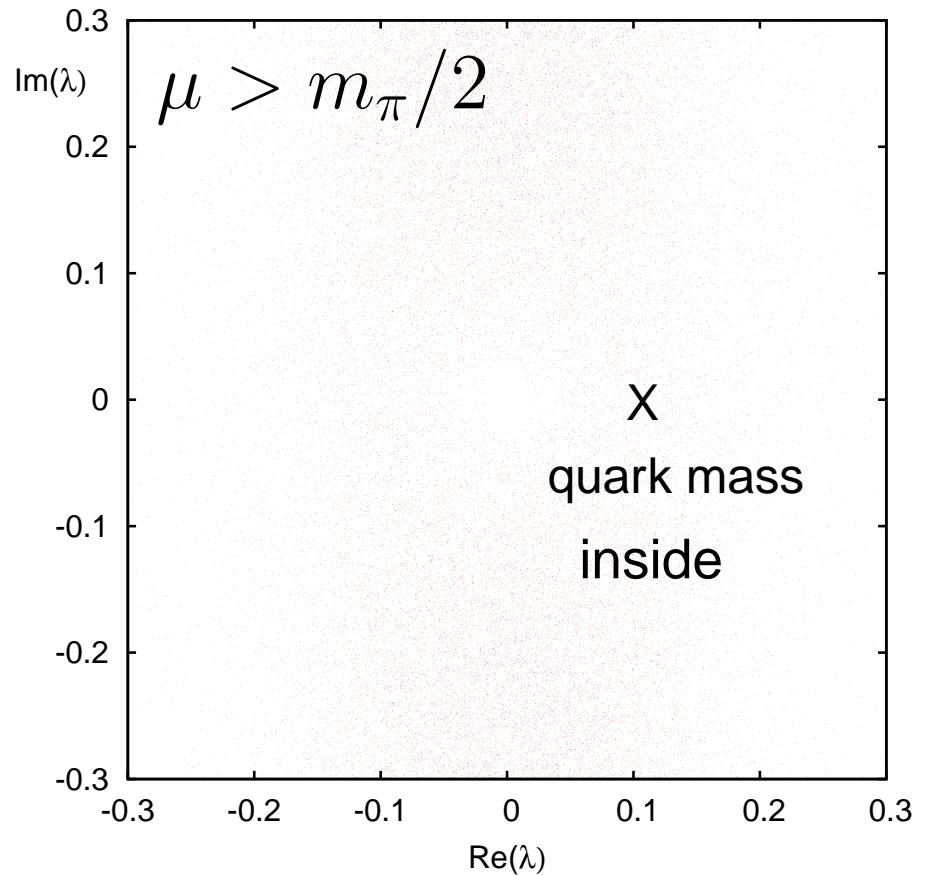
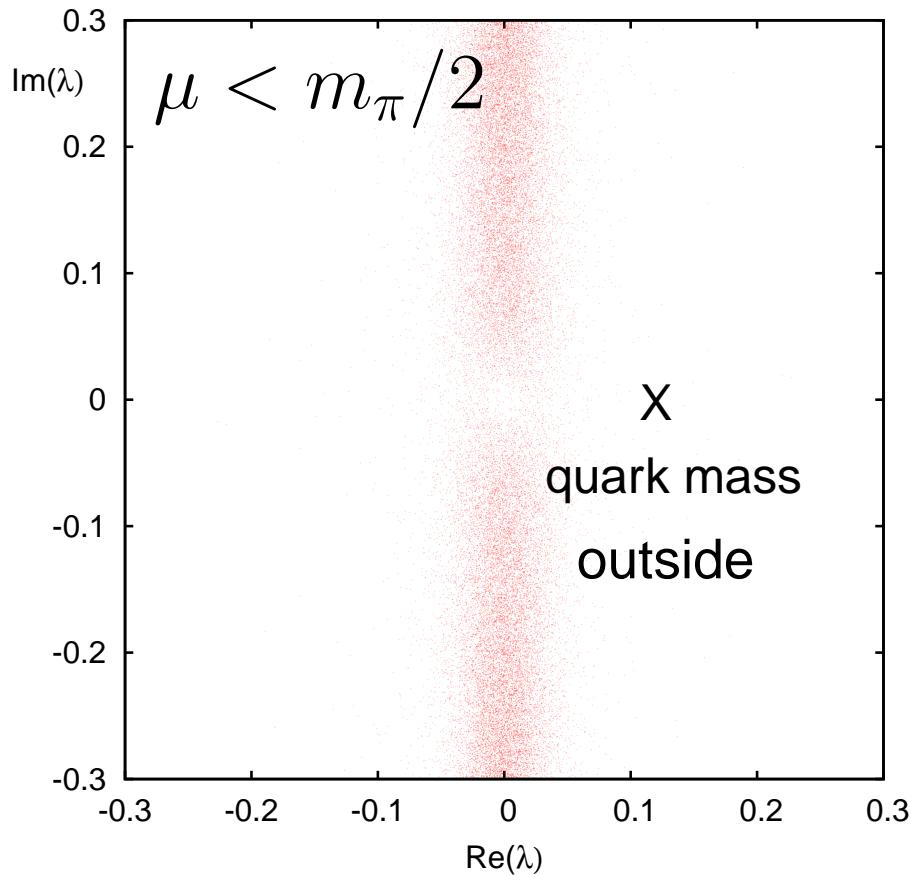
D'Elia Renzo Lombardo arXiv:0705.3814      Conradi D'Elia arXiv:0707.1987



# The Big Picture



# Quark mass & the eigenvalue distribution



$$(D + \mu\gamma_0)\psi_k = \lambda_k\psi_k$$

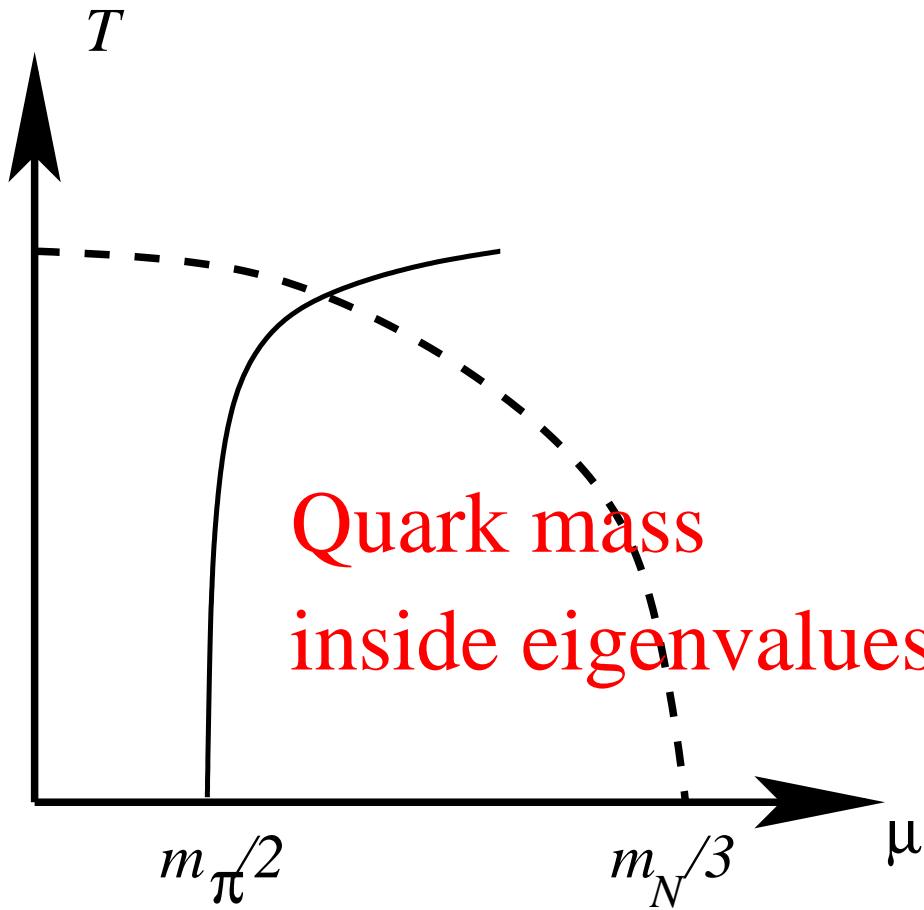


Bloch Wettig Lattice 2006

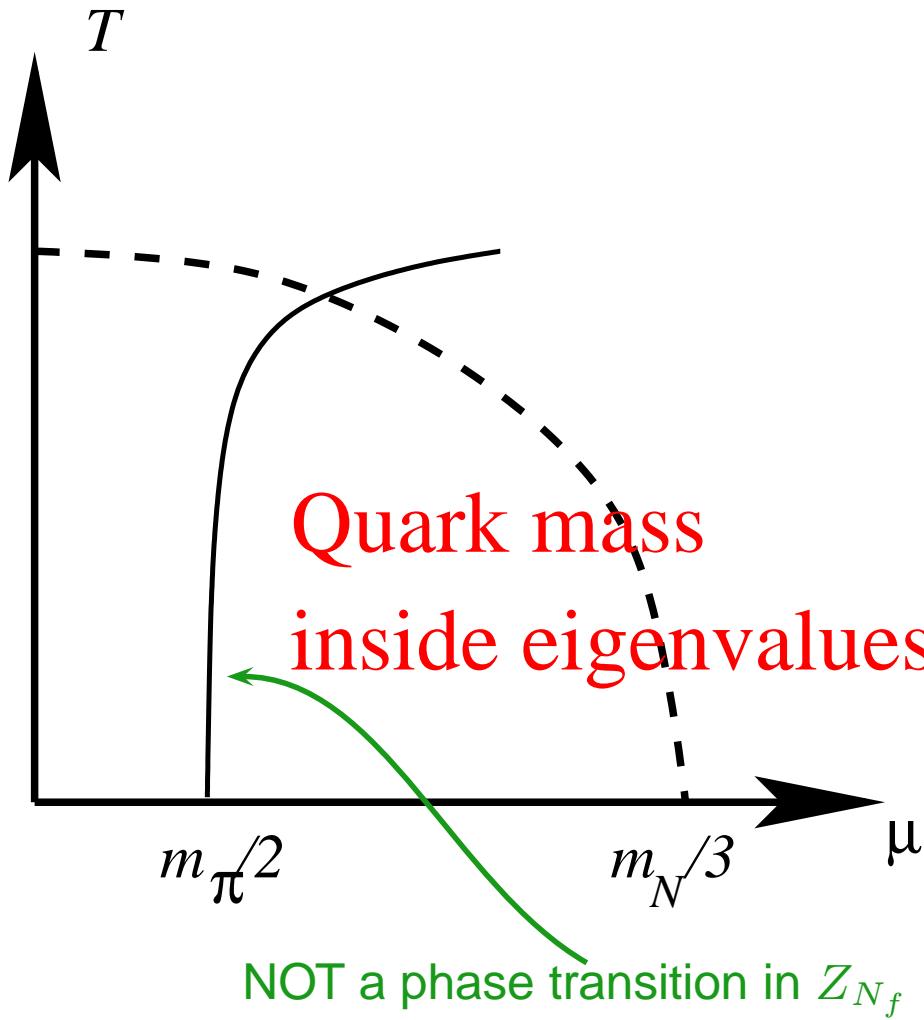
Gibbs PRINT-86-0389

Davies Klepfish PLB 256 (1991) 68 Lombardo Kogut Sinclair PRD 54 (1996) 2303

# The big picture



# The big picture



# The average phase factor



$$\langle e^{2i\theta} \rangle_{N_f} \equiv \left\langle \frac{\det(D + \mu\gamma_0 + m)}{\det(D + \mu\gamma_0 + m)^*} \right\rangle_{N_f}$$

is a ratio of two partition functions

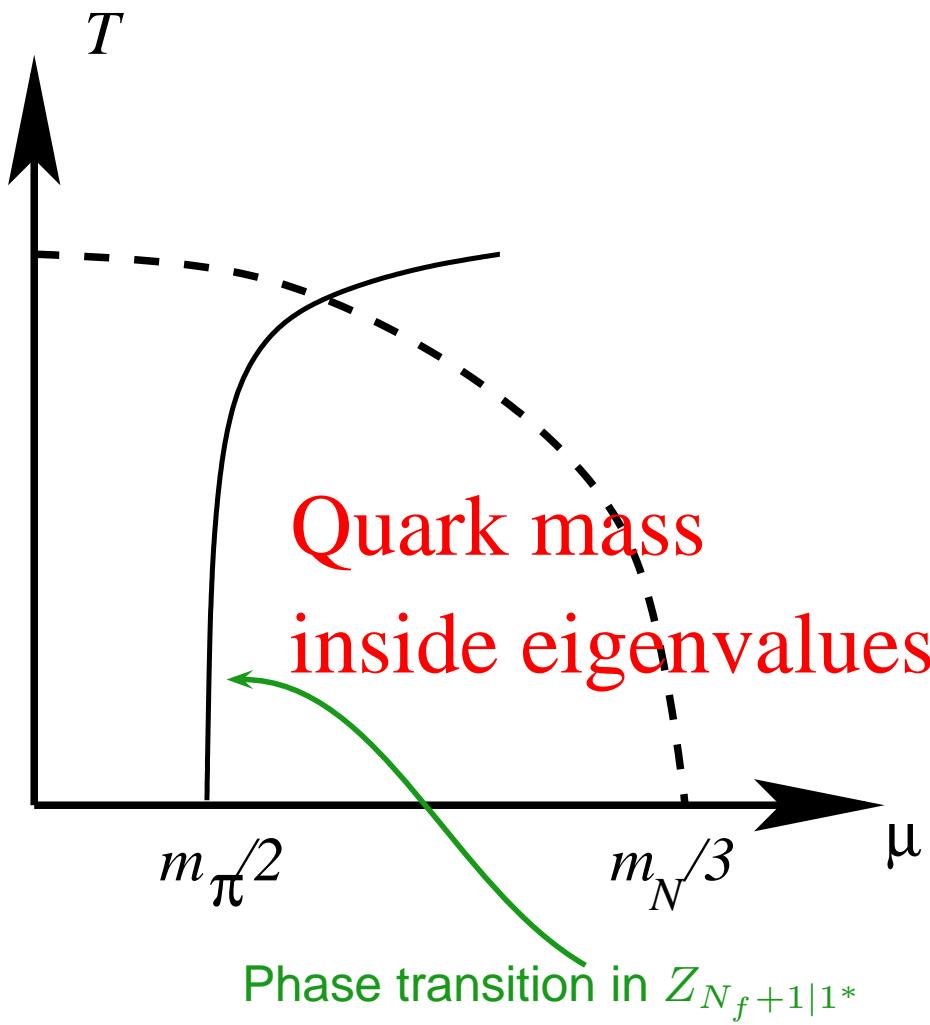
Phase transition at  $\mu = m_\pi/2$



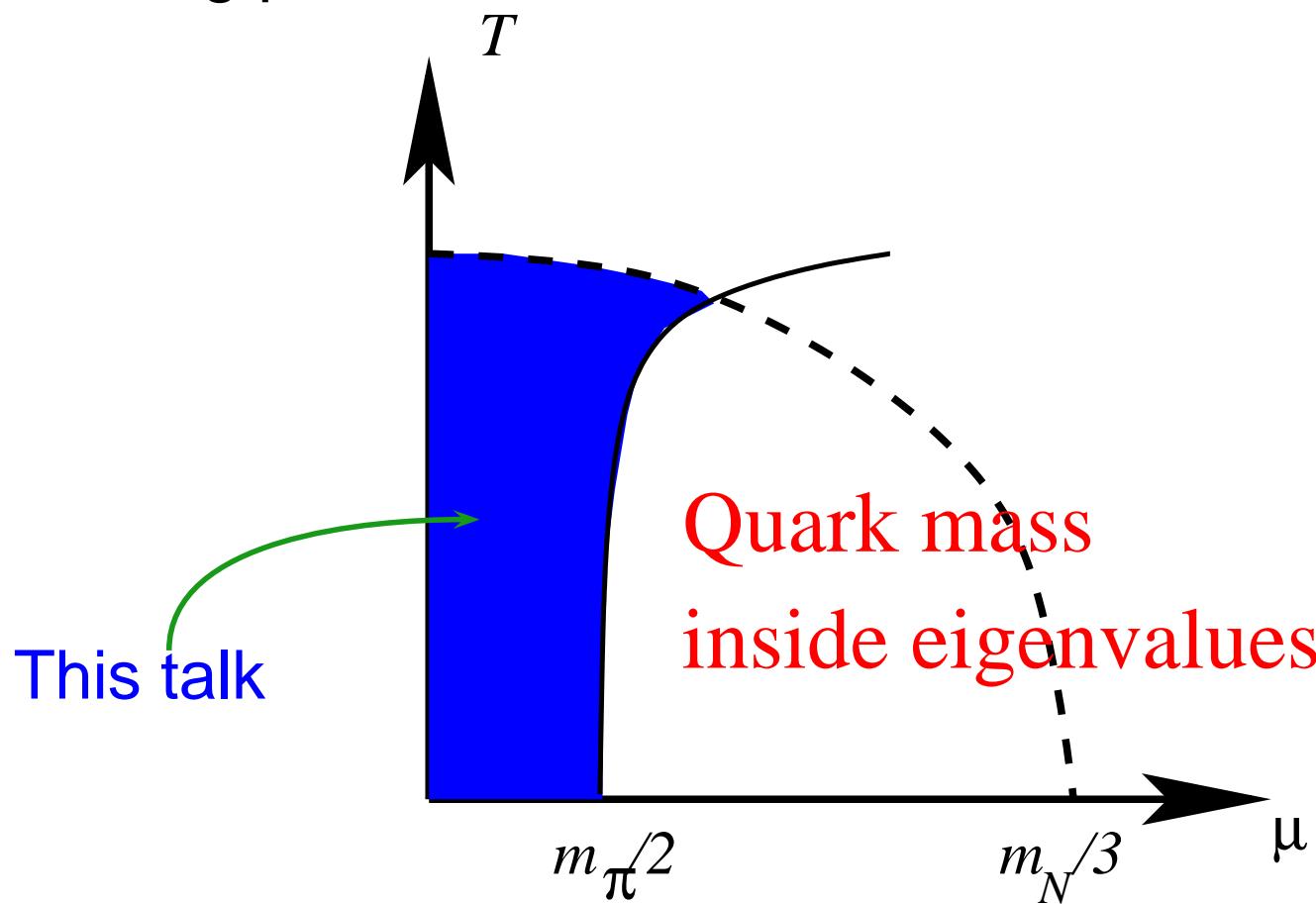
$$\langle e^{2i\theta} \rangle_{N_f} = \frac{Z_{N_f+1|1^*}}{Z_{N_f}} = e^{-V\Delta\Omega}$$



# The big picture



# The big picture





# Why the average phase factor can be computed in CPT

In CPT:

$Z_{N_f}$  independent of  $\mu$

(ordinary pions have baryon charge 0)

$Z_{N_f+1|1^*}$  depends on  $\mu$

(pions with a conjugate quark have baryon charge)



Kogut Stephanov Toublan PLB 464 (1999) 183

Kogut + .. NPB 582 (2000) 477

Hasenfratz Leutwyler NPB 343 (1990) 241

Splittoff Verbaarschot NPB 57 (2006) 259



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$\mu > m_\pi/2$  **Bose condensate of charged pions**



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# The average phase factor in CPT

$$\langle e^{2i\theta} \rangle_{N_f} = \frac{Z_{N_f+1|\mathbf{1}^*}}{Z_{N_f}} = e^{-(N_f+1)V\Delta G_0}$$

$\Delta G_0$  is the difference between charged and neutral pions

$\Delta G$  to 1-loop order in a box with dimensions  $V = L_i^3 L_0$

$\Delta G_0$  is independent of the cutoff

Splittorff Verbaarschot PRL 98 (2007) 031601

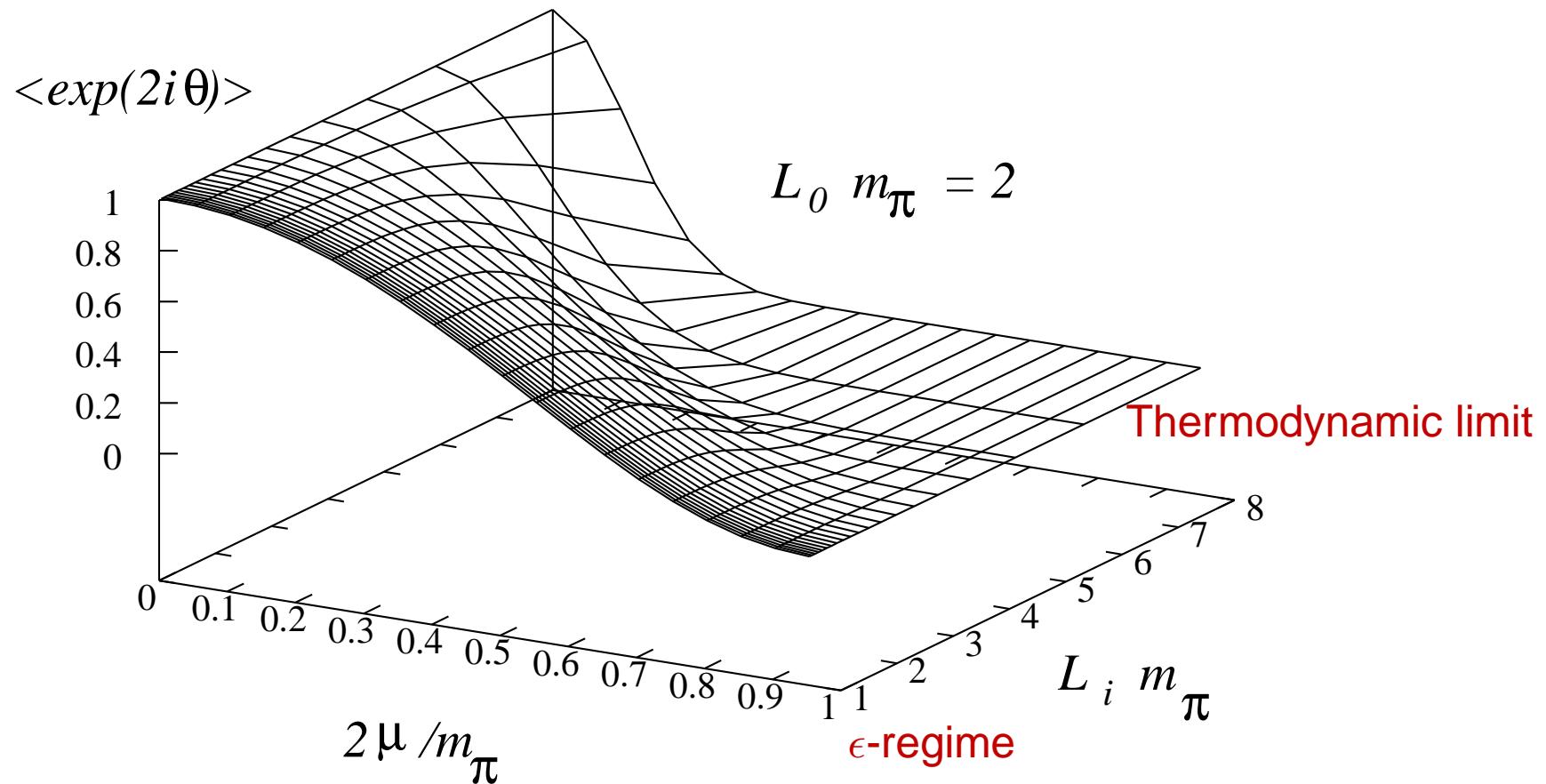
Splittorff Svetitsky hep-lat/0703004



# Average phase factor $N_f = 2$ from One-Loop CPT



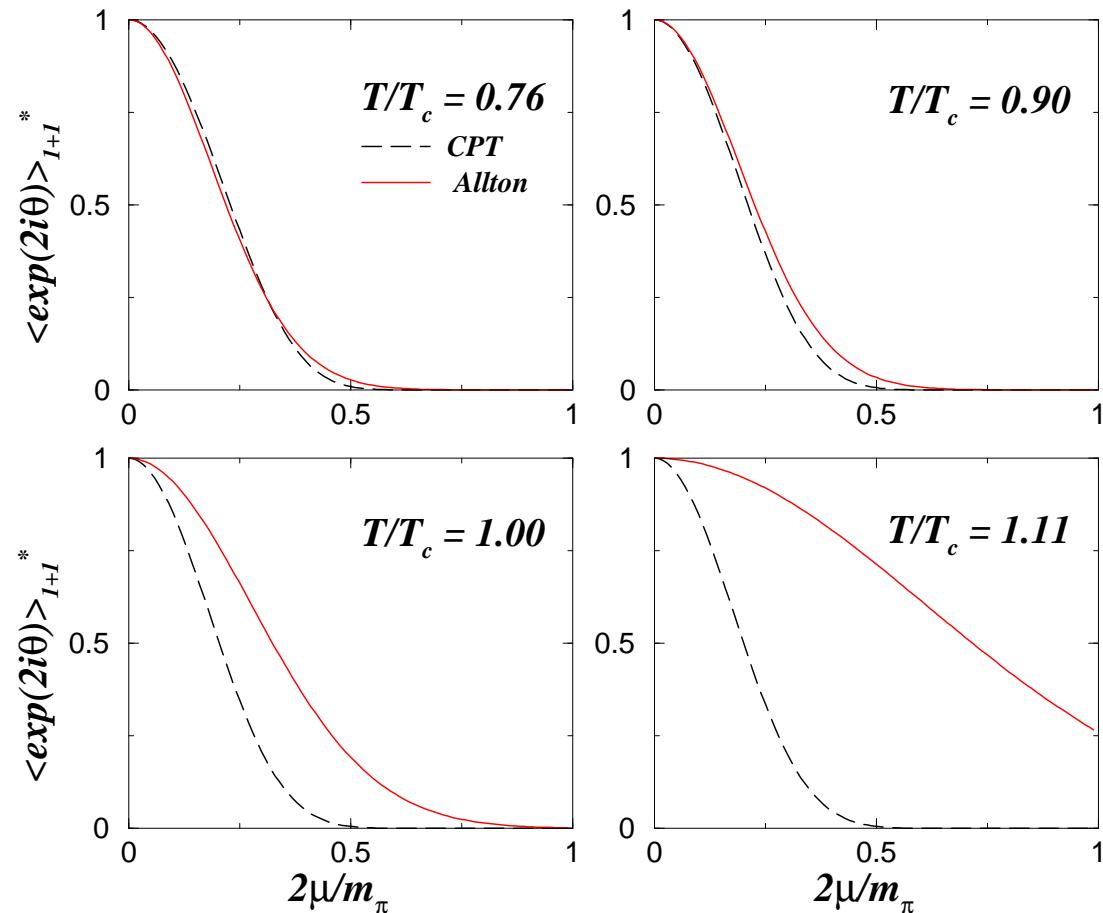
- increasing the volume for fixed temperature



# Average phase factor on the lattice



$$\langle e^{2i\theta} \rangle_{1+1^*} = e^{L_i^3 T(c_2 - c_2^I) \mu^2}$$



Allton+... Phys.Rev. D71 (2005) 054508

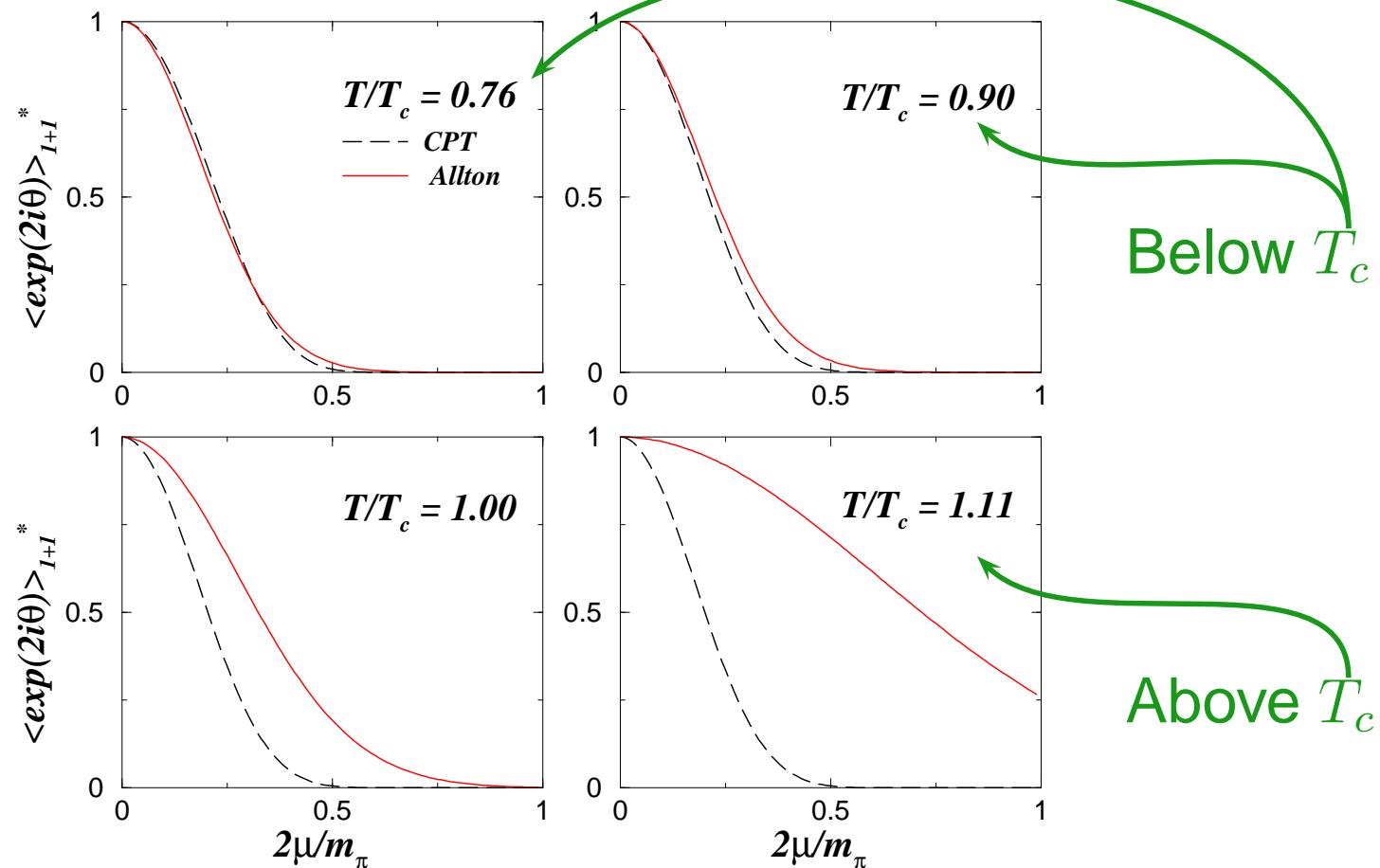
Splittorff Verbaarschot 2007



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Allton+... Phys.Rev. D71 (2005) 054508

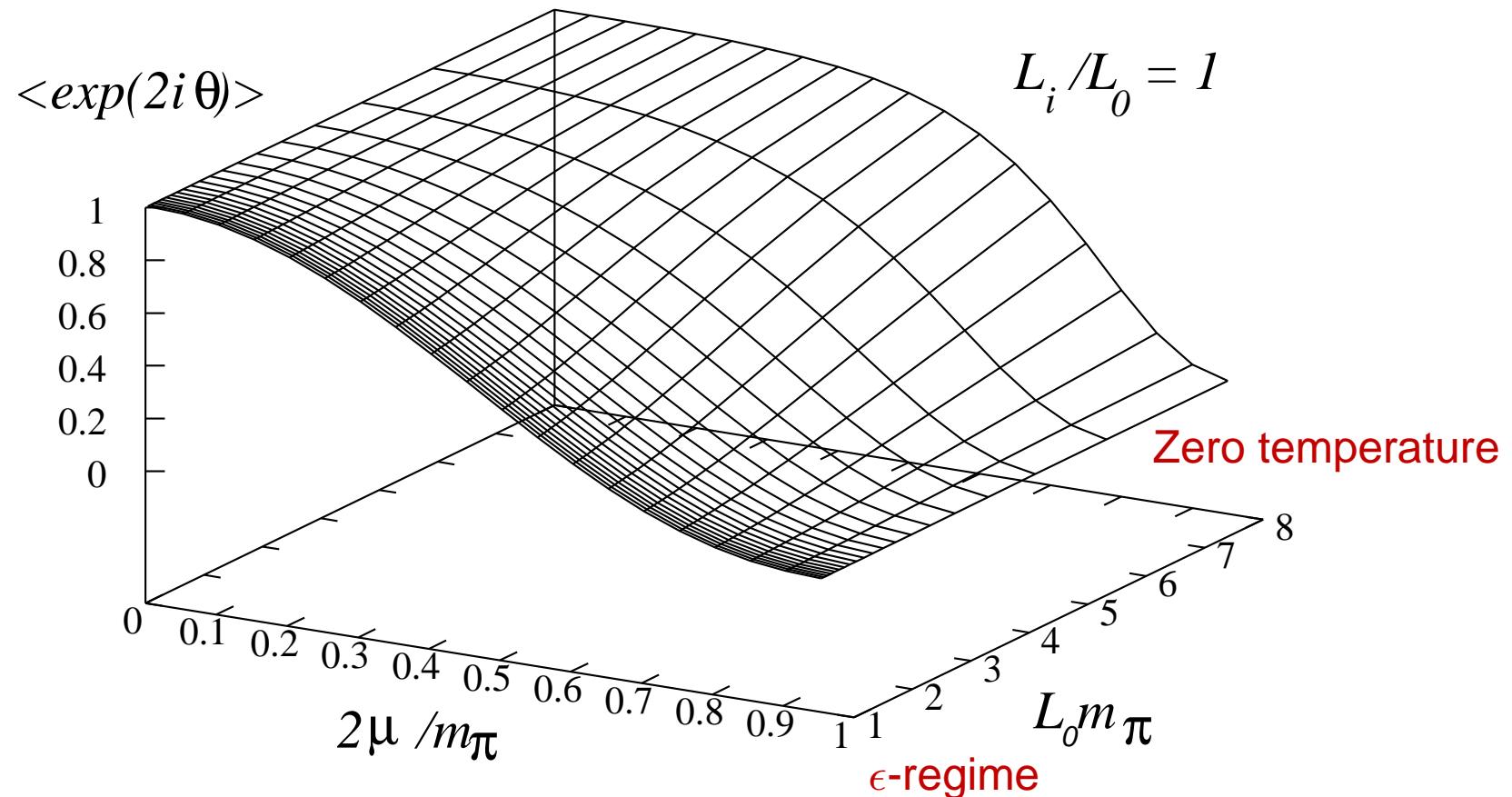
Splittorff Verbaarschot 2007



# Average phase factor $N_f = 2$ from One-Loop CPT



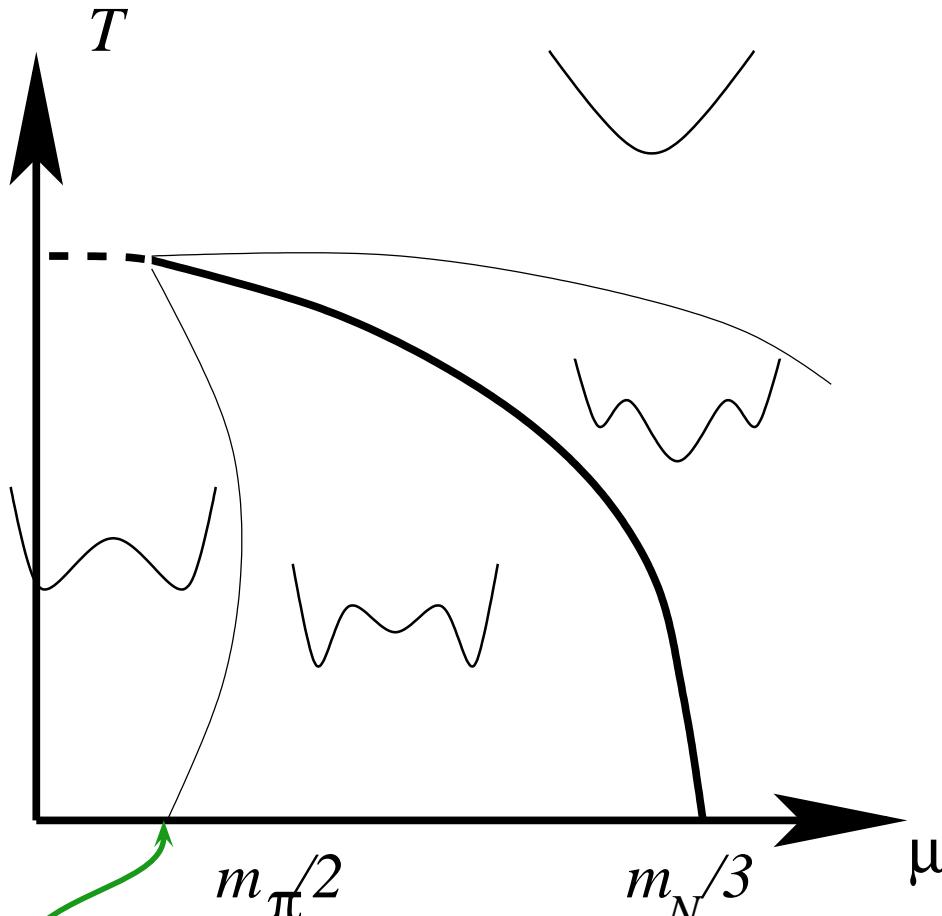
- increasing the volume for fixed  $L_i/L_0$



Ravagli Verbaarschot arXiv:0704.1111

No sign problem at zero temperature  $\mu < m_\pi/2$  ( $L_0/L_i$  fixed)

Suggests: It is possible to look for spinodal on cold lattices ( $L_0 \gg L_i$ )

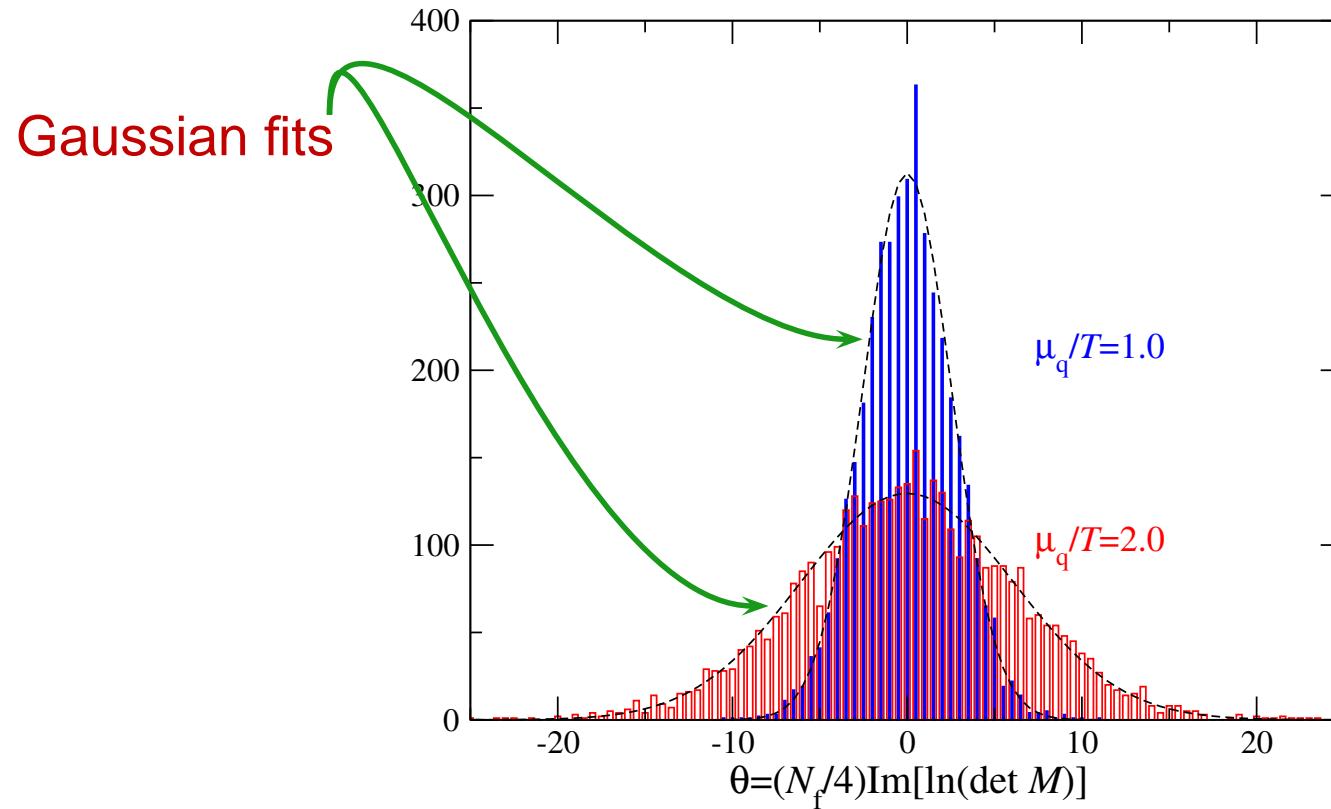


Lombardo Kogut Sinclair PRD 54 (1996) 2303

Kawamoto Miura Ohnishi Ohnuma PRD 75 (2007) 014502



# The distribution of the phase



Ejiri arXiv:0706.3549

# The distribution of the phase



$$\rho_{N_f}(\theta) \equiv \langle \delta(\theta - \theta') \rangle_{N_f}$$

from the definition

$$\rho_{N_f=2}(\theta) = e^{iN_f\theta} \rho_{1+1^*}(\theta) \frac{Z_{1+1^*}}{Z_{N_f=2}}$$



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1-loop CPT: Envelope is Gaussian                    width  $\sim \sqrt{V}$

$$\rho_{N_f}^{CPT}(\theta) = e^{iN_f\theta} \frac{1}{\sqrt{V\Delta G_0}\pi} e^{-\theta^2/(V\Delta G_0)} e^{(N_f/2)^2 V\Delta G_0}$$



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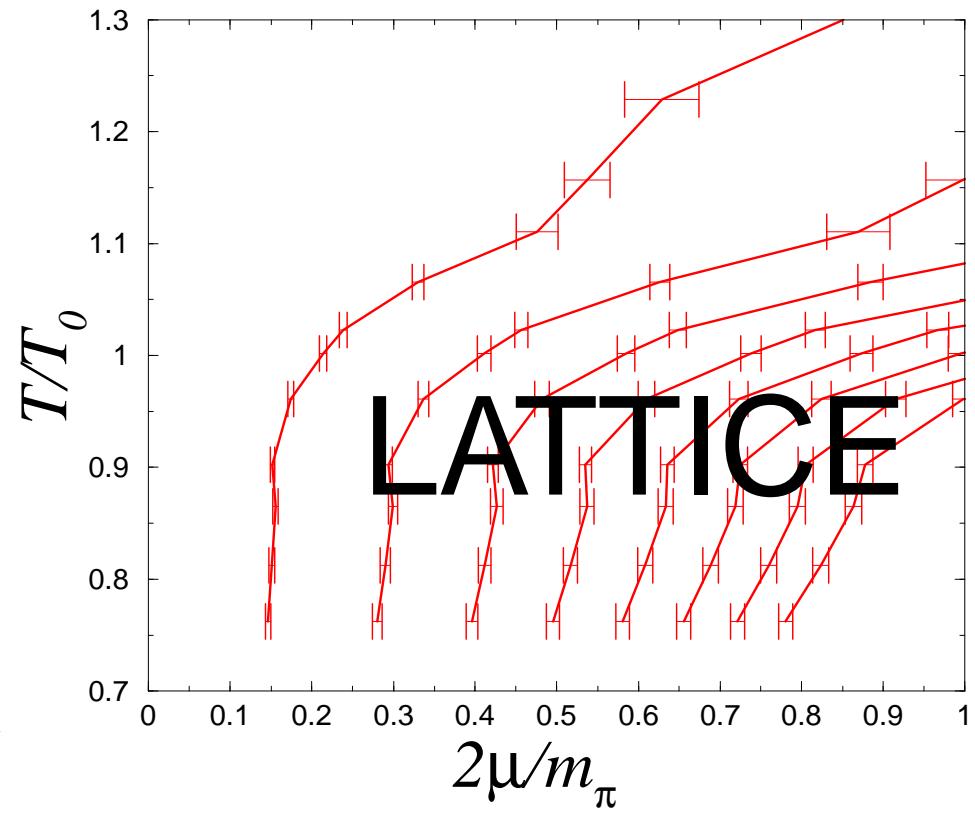
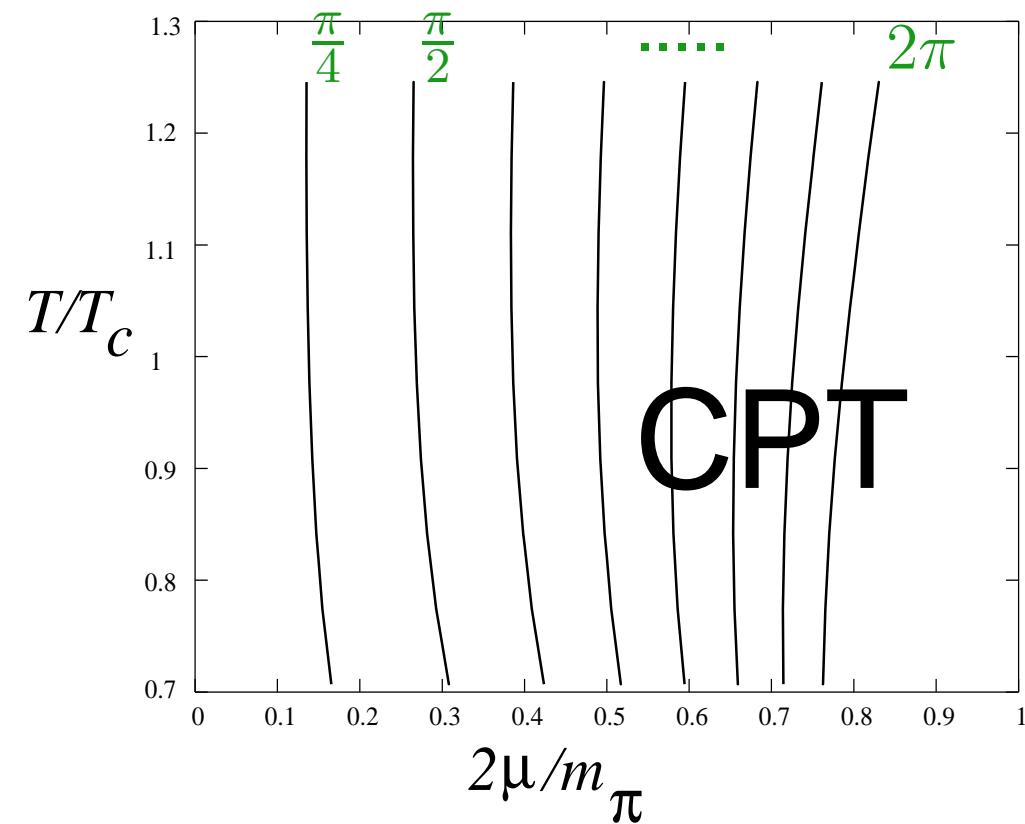
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Huge cancellations  $\rightarrow$  small non Gaussian terms big effect.  $\sim e^V$

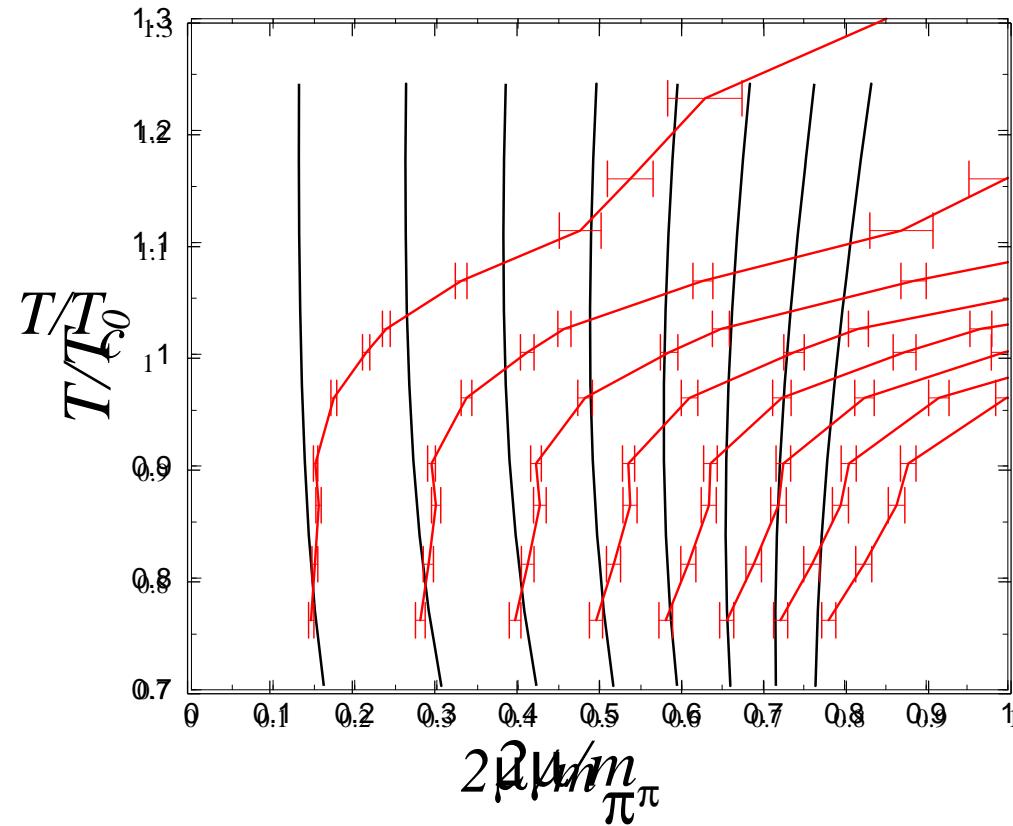


The standard deviation  $\sqrt{\langle \theta^2 \rangle - \langle \theta \rangle^2}$



Allton+... Phys.Rev. D71 (2005) 054508

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Allton+... Phys.Rev. D71 (2005) 054508

# Conclusions



Analytic understanding of the strength of the sign problem

Agreement with lattice data below  $T_c$

$\langle e^{2i\theta} \rangle \sim 1$  at low  $T$  and  $\mu < m_\pi/2$       Look for spinodal ?



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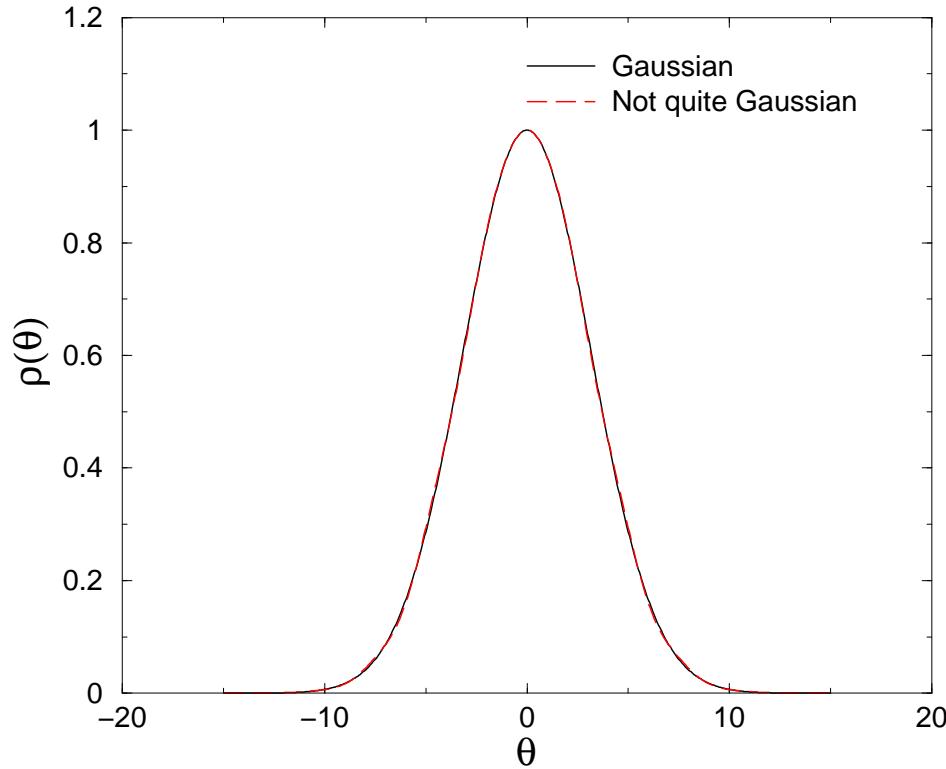
The density of angles is Gaussian at 1-loop      **WARNING**

The standard deviation of the phase    **consistent with lattice**



Huge cancellations → small terms big effect.

Example:



$$\frac{\int d\theta \rho_{\text{not-quite-gauss}}(\theta) e^{2i\theta}}{\int d\theta \rho_{\text{gauss}}(\theta) e^{2i\theta}} \sim 10^6$$