

Chiral Phase Boundary of QCD with Many Flavors

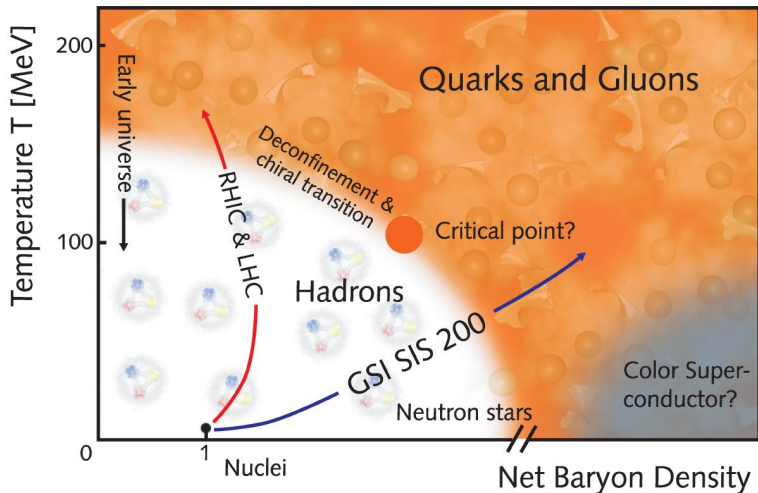
Holger Gies

Universität Heidelberg



& J. Braun, J. Jaeckel, J.M. Pawłowski, C. Wetterich

QCD Phase Diagram



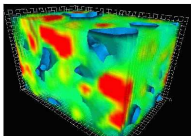
[FAIR@www.gsi.de]

“Learning by Doing”

“Learning by Deforming”

QCD

Lattice QCD



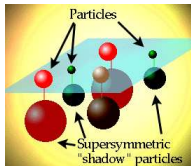
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Large N_c



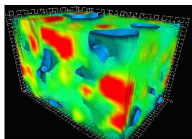
SUSY QCD



“Learning by Deforming”

QCD

Lattice QCD



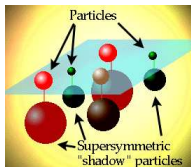
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Large N_c



SUSY QCD



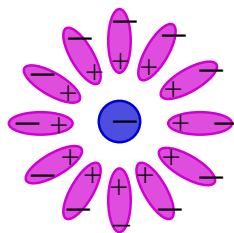
Large N_f



“many-flavor
QCD”

Many-Flavor QCD

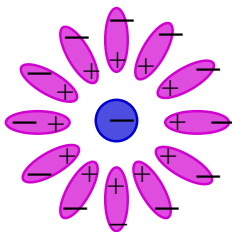
▷ charge screening:



▷ β function

$$\beta = -2 \left(\frac{11}{3} N_c - \frac{2}{3} N_f \right) \frac{g^4}{16\pi^2} - 2 \left(\frac{34N_c^3 + 3N_f - 13N_c^2 N_f}{3N_c} \right) \frac{g^6}{(16\pi^2)^2} + \dots$$

▷ charge screening:



▷ β function

$$\beta = -2 \left(\frac{11}{3} N_c - \frac{2}{3} N_f \right) \frac{g^4}{16\pi^2} - 2 \underbrace{\left(\frac{34N_c^3 + 3N_f - 13N_c^2 N_f}{3N_c} \right)}_{>0} \frac{g^6}{(16\pi^2)^2} + \dots$$

$$\text{for } N_f > \frac{34N_c^3}{13N_c^2 - 3} \stackrel{\text{SU}(3)}{\simeq} 8.05$$

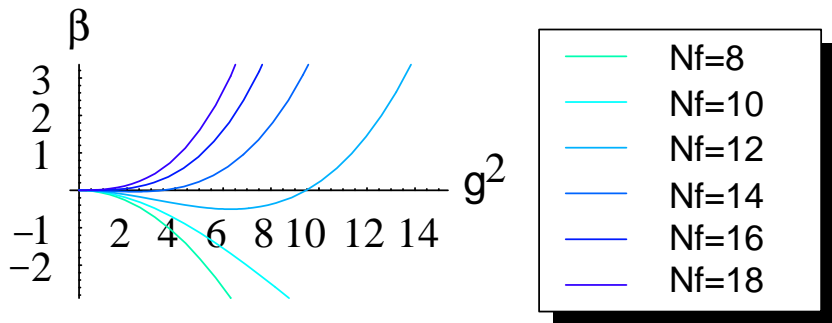
Many-flavor QCD

▷ β function

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▷ e.g., SU(3): IR fixed point α_*

(BANKS&ZAKS'82)

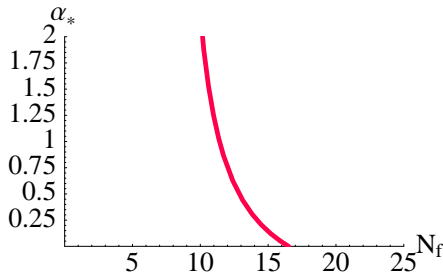
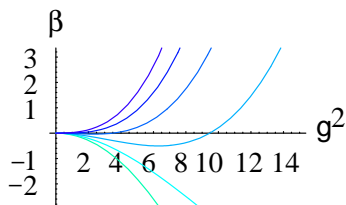


Many-flavor QCD

▷ β function

$$\beta = -2 \left(\frac{11}{3} N_c - \frac{2}{3} N_f \right) \frac{g^4}{16\pi^2} - 2 \left(\frac{34N_c^3 + 3N_f - 13N_c^2 N_f}{3N_c} \right) \frac{g^6}{(16\pi^2)^2} + \dots$$

▷ N_f dependence of α_*



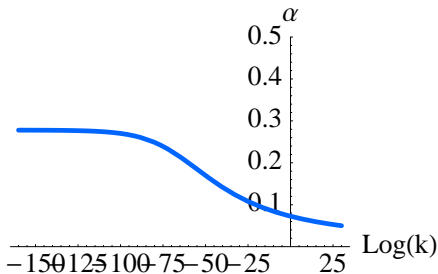
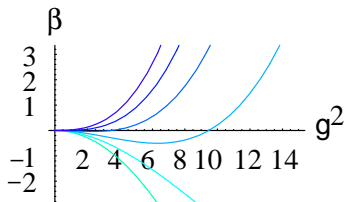
Many-flavor QCD

▷ β function

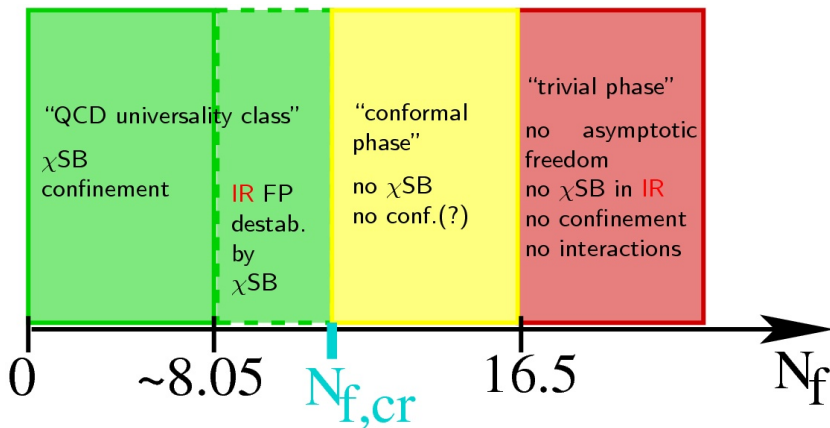
$$\beta = -2 \left(\frac{11}{3} N_c - \frac{2}{3} N_f \right) \frac{g^4}{16\pi^2} - 2 \left(\frac{34N_c^3 + 3N_f - 13N_c^2 N_f}{3N_c} \right) \frac{g^6}{(16\pi^2)^2} + \dots$$

▷ e.g. $N_f = 14$

⇒ IR fixed point



Many-flavor QCD



Many-flavor QCD

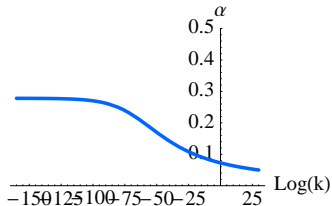
- ▷ gap equation

$$[\text{---} \overset{S}{\text{---}} \text{---}]^{-1} = [\text{---} \overset{S_0}{\text{---}} \text{---}]^{-1} + \text{---} \overset{\gamma}{\text{---}} \text{---} \overset{D}{\text{---}} \text{---} \overset{\Gamma}{\text{---}} \text{---}$$

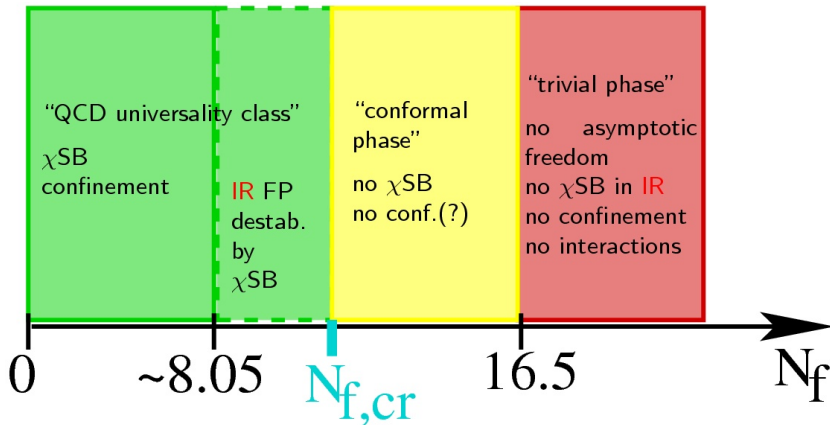
- ▷ rainbow/ladder approximation: $\Gamma \rightarrow g$

$$\Rightarrow \chi\text{SB for } \alpha > \alpha_{\text{cr}}$$

- ▷ α_* VS. α_{cr}

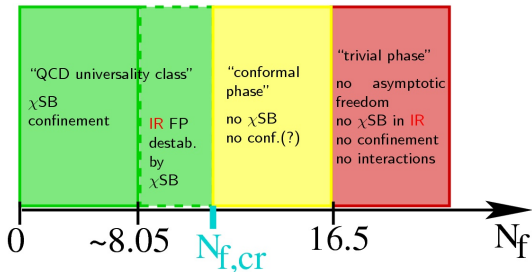


Many-flavor QCD



$$N_{f,cr} = ?$$

Many-flavor QCD



$$N_{f,cr} = \begin{cases} 5 & \text{(HARADA \& YAMAWAKI '00)} \\ 6 & \text{(IWASAKI ET AL.'03)} \\ \gtrsim 6 & \text{(VELKOVSKY \& SHURYAK '97, APPELQUIST \& SELIPSKY '97)} \\ \gtrsim 10 & \text{(SANNINO \& SCHECHTER '99)} \\ \simeq 12 & \text{(MIRANSKY \& YAMAWAKI '96, APPELQUIST ET AL.'96)} \end{cases}$$

LGT: (KOGUT \& SINCLAIR '88; BROWN ET AL.'92; IWASAKI ET AL.'96; DAMGAARD ET AL.'97)

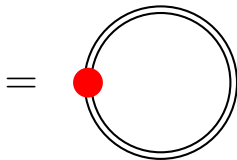
adjoint matter: (DIETRICH \& SANNINO '06)

Functional RG

Functional RG Flow Equation



$$\partial_t \Gamma_k \equiv k \partial_k \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

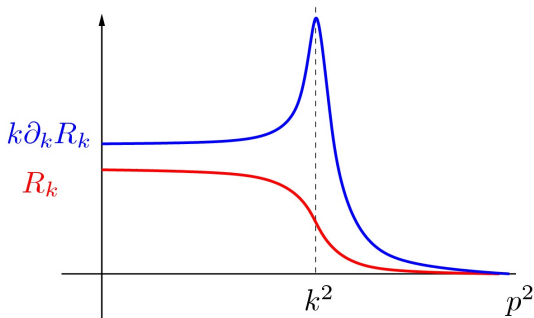


(WILSON'71; WEGNER&HOUGHTON'73; POLCHINSKI'84; WETTERICH'93)

Functional RG Flow Equation

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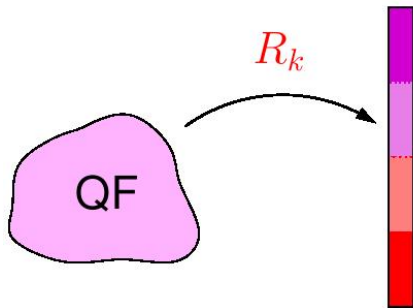
▷ regulator



Functional RG Flow Equation

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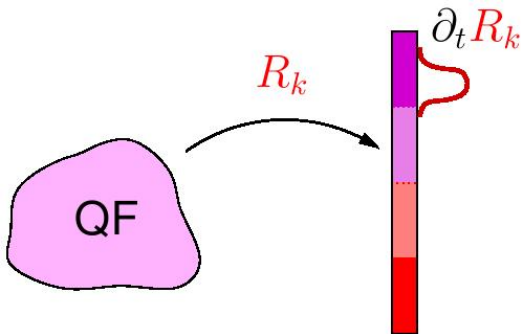
▷ quantum fluctuations:



Functional RG Flow Equation

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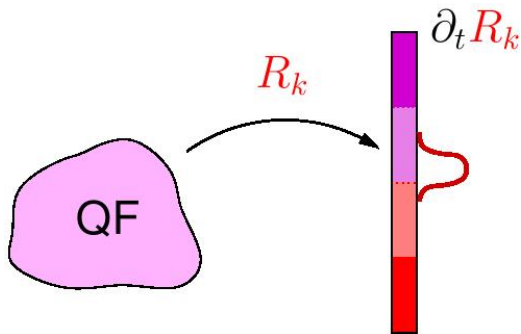
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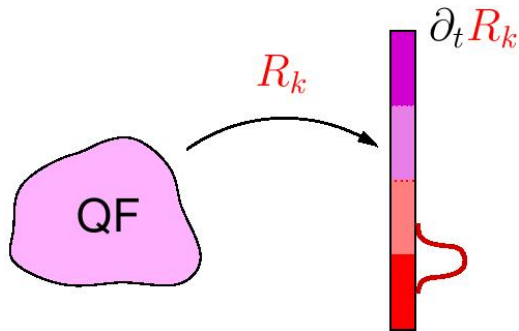
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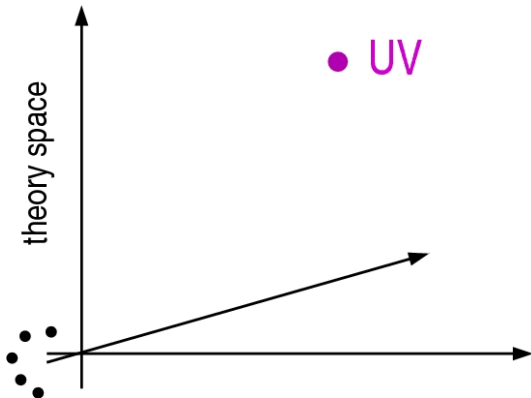
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Functional RG Flow Equation

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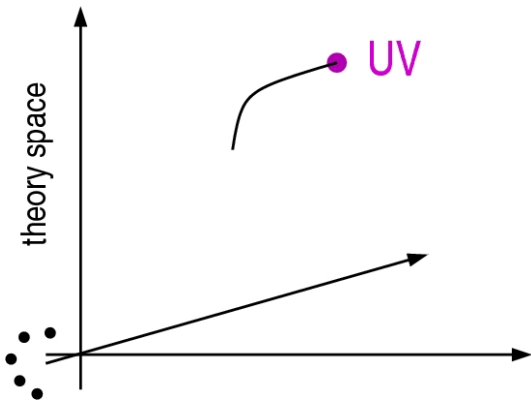
▷ RG trajectory: $\Gamma_{k=\Lambda} = S_{\text{bare}} = \int \frac{1}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \bar{\psi} (i\cancel{\partial} + gA) \psi$



Functional RG Flow Equation

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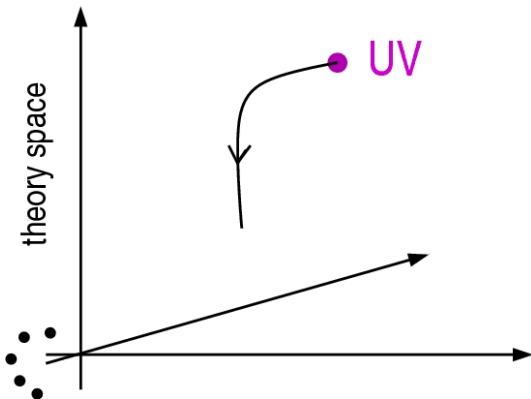
▷ RG trajectory:



Functional RG Flow Equation

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▷ RG trajectory:

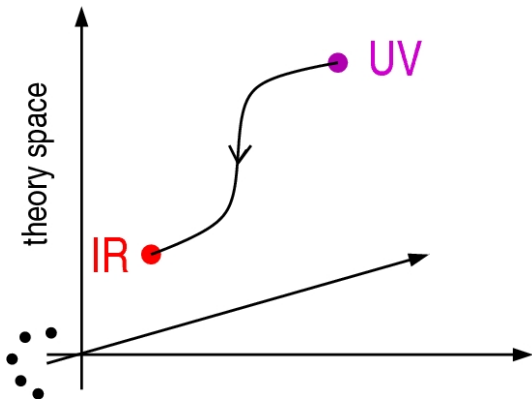


Functional RG Flow Equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

$$\Gamma_{k \rightarrow 0} = \Gamma$$

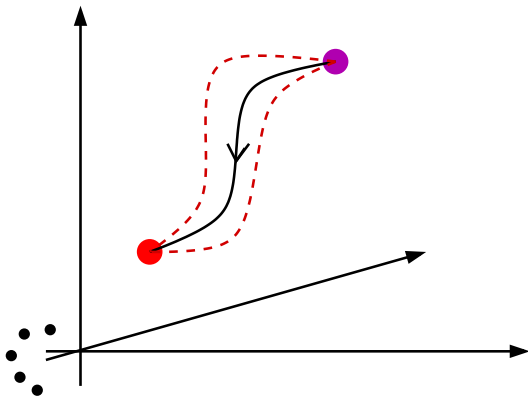


Functional RG Flow Equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

R_k scheme independence

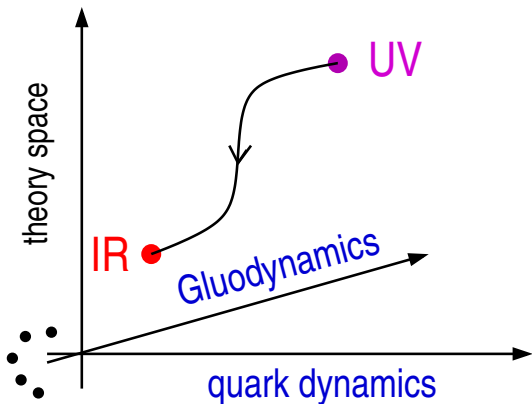


Functional RG Flow Equation

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▷ RG trajectory:

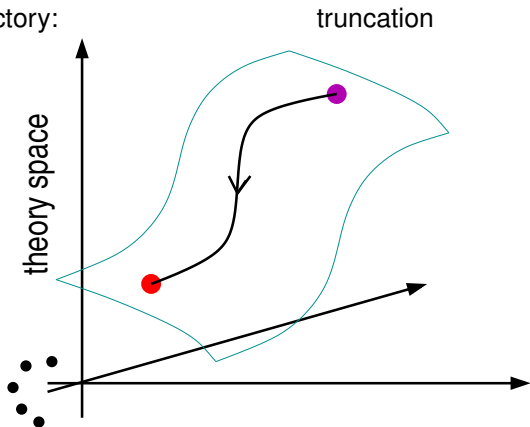
truncation



Functional RG Flow Equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

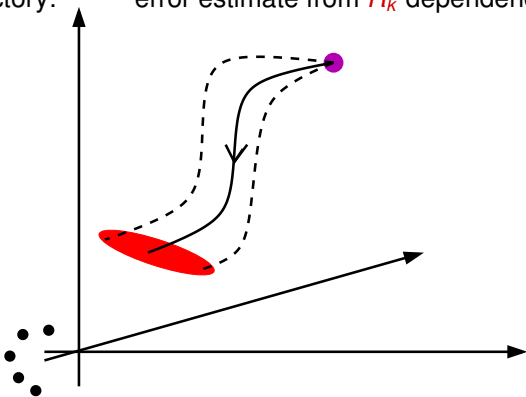


Functional RG Flow Equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

error estimate from R_k dependence



RG Flow towards the Chiral Transition

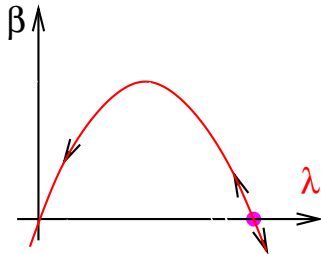
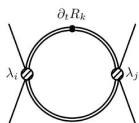
RG Flow of the Chiral Sector

▷ effective action:

$$\Gamma_k = \int \frac{1}{2} \frac{\lambda_\sigma}{k^2} [(\bar{\psi}^a \psi^b)^2 - (\bar{\psi}^a \gamma_5 \psi^b)^2]$$

▷ RG flow

$$\partial_t \lambda_\sigma = 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2$$



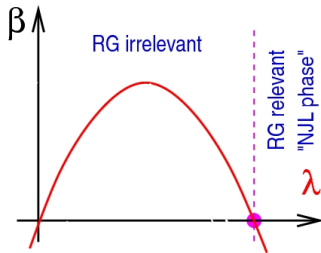
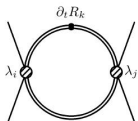
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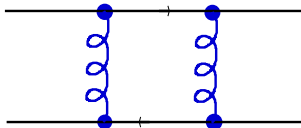
RG Flow of the Chiral Sector

▷ effective action:

$$\Gamma_k = \int \frac{1}{4} F_{\mu\nu}^Z F_{\mu\nu}^Z + \dots + \bar{\psi} (i\cancel{\partial} + \bar{g}A) \psi + \frac{1}{2} \frac{\lambda_\sigma}{k^2} [(\bar{\psi}^a \psi^b)^2 - (\bar{\psi}^a \gamma_5 \psi^b)^2]$$

▷ RG flow

$$\begin{aligned} \partial_t \lambda_\sigma &= 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2 \\ &\quad - \frac{3}{8\pi^2} \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma \\ &\quad - \frac{9}{256\pi^2} \frac{3N_c^2 - 8}{N_c} g^4 \end{aligned}$$



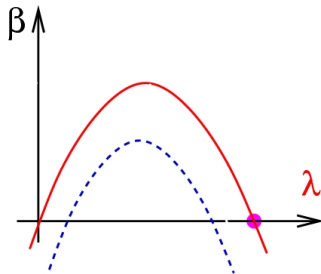
RG Flow of the Chiral Sector

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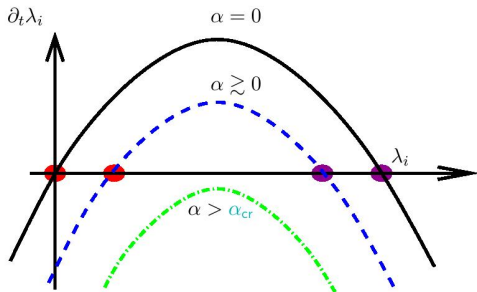
▷ RG flow

$$\partial_t \lambda_\sigma = 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2 - \frac{3}{8\pi^2} \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma - \frac{9}{256\pi^2} \frac{3N_c^2 - 8}{N_c} g^4$$



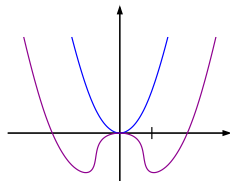
Chiral Criticality

▷ critical gauge coupling α_{cr} :



⇒ bosonization $\rightarrow \chi$ SB:

$$\text{if } \alpha > \alpha_{cr} : \quad \lambda \sim \frac{1}{m_\phi^2} \rightarrow \infty$$



RG Flow of the Chiral Sector

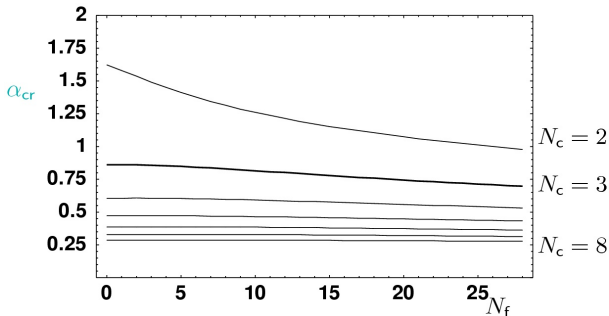
- ▷ effective action: $SU(N_c), SU(N_f)_L \times SU(N_f)_R$

$$\begin{aligned}\Gamma_k = & \int \frac{Z_F}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \dots + \bar{\psi} (iZ_\psi \not{\partial} + Z_1 \bar{g} A) \psi \\ & + \frac{1}{2} \frac{\lambda_\sigma}{k^2} (\text{S-P}) + \frac{1}{2} \frac{\lambda_{VA}}{k^2} [2(\text{V-A})^{\text{adj.}} + (1/N_c)(\text{V-A})] \\ & + \frac{1}{2} \frac{\lambda_+}{k^2} (\text{V+A}) + \frac{1}{2} \frac{\lambda_-}{k^2} (\text{V-A})\end{aligned}$$

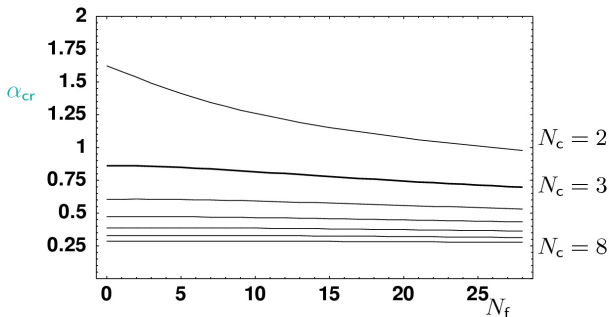
- ▷ RG flow, e.g.,

$$\begin{aligned}\partial_t \lambda_\sigma = & 2\lambda_\sigma - \frac{1}{4\pi^2} I_1^{(F)}[R_k] \left\{ 2N_c \lambda_\sigma^2 - 2\lambda_- \lambda_\sigma - 2N_f \lambda_\sigma \lambda_{VA} - 6\lambda_+ \lambda_\sigma \right\} \\ & - \frac{1}{8\pi^2} I_{1,1}^{(\text{FB})}[R_k] \left[3 \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma - 6g^2 \lambda_+ \right] \\ & - \frac{3}{128\pi^2} I_{1,2}^{(\text{FB})}[R_k] \frac{3N_c^2 - 8}{N_c} g^4\end{aligned}$$

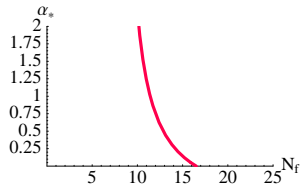
(HG, JAECKEL, WETTERICH'04)

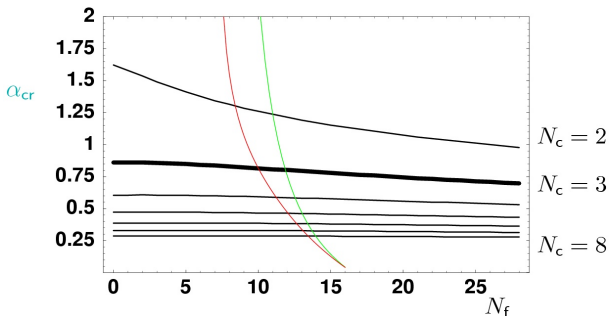


e.g., for $N_c = 3 = N_f$: $\alpha_{cr} \simeq 0.85$



▷ compare with Banks-Zaks IR fixed point:





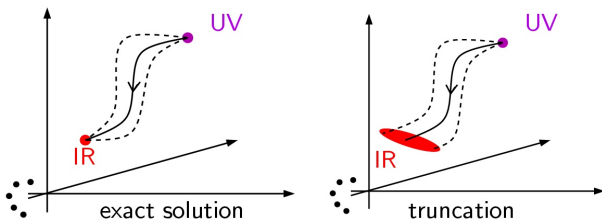
— 2-loop SU(3) β function in $\overline{\text{MS}}$ scheme

— 4-loop SU(3) β function in $\overline{\text{MS}}$ scheme $\Rightarrow N_{f,cr} \simeq 10.0$

(RITBERGEN ET AL.'97)

Error Estimate

- ▷ regulator dependence



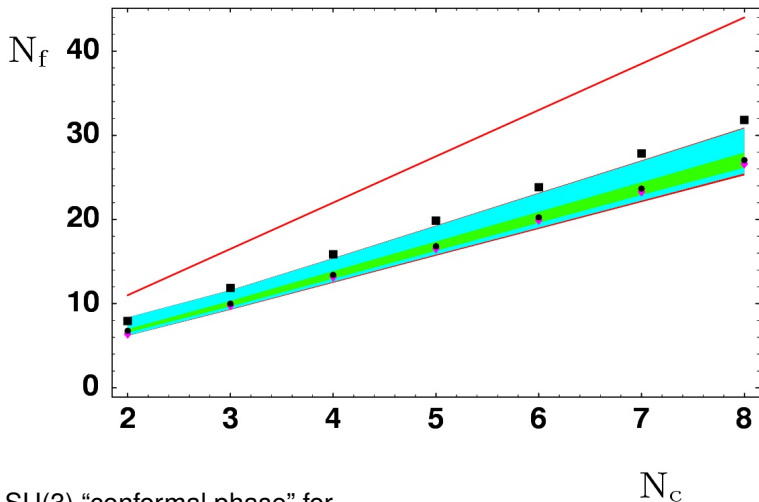
- ▷ fermion sector: “optimized” regulator vs. “sharp cutoff”

(LITIM'01)

$$l_1^{(F),4} = \frac{1}{2}, \quad l_{1,1}^{(FB),4} = 1, \quad l_{1,2}^{(FB),4} = \frac{3}{2} \quad \text{vs.} \quad l_1^{(F),4} = l_{1,1}^{(FB),4} = l_{1,2}^{(FB),4} = 1$$

- ▷ anomalous dimensions, momentum dependencies, higher-order operators $\sim \psi^8$, etc. ...
- ▷ gauge sector: 2-loop, 3-loop, 4-loop β function

$\overline{\text{MS}}$ scheme vs. RG scheme ($\sim 10, 30, 50$ % variation (?))



▷ SU(3) “conformal phase” for

$$N_{f,\text{cr}} = 10.0 \pm 0.29(\text{fermion}) \begin{matrix} +1.55 \\ -0.63 \end{matrix}(\text{gluon}) \lesssim N_f < 16.5$$

(HG, JAECKEL'05)

Lessons to be learned for “real QCD”

- fermionic screening is rather weak
- fermionic truncation (surprisingly) stable in χ symmetric phase
- phase boundary detectable with fermionic “derivative expansion”
- “real QCD” requires nonperturbative estimate of β_{g^2}

Lessons to be learned for “real QCD”

- fermionic screening is rather weak
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- phase boundary detectable with fermionic “derivative expansion”
- “real QCD” requires nonperturbative estimate of β_{g^2}

$$\Gamma_k = \int \frac{Z_F}{4} F_z^{\mu\nu} F_{\mu\nu}^z + \underbrace{\quad \cdot \quad \cdot \quad \cdot \quad}_{\uparrow}$$

RG Flow for Gluodynamics

RG Flow of Gluodynamics

▷ Operator expansion with the background-field method

(REUTER,WETTERICH'94; FREIRE,LITIM,PAWLOWSKI'00)

$$\Gamma_k[A] = \int d^d x W_k(F^2), \quad F^2 \equiv F_{\mu\nu}^a F_{\mu\nu}^a$$
$$W_k(F^2) = \frac{Z_F}{4} F^2 + \frac{W_2}{2!4^2} (F^2)^2 + \frac{W_3}{3!4^3} (F^2)^3 + \dots$$

(HG'02)

▷ running coupling:

(ABBOTT'82)

$$g^2 = Z_F^{-1} \bar{g}^2$$

▷ β function:

$$\partial_t g^2 \equiv \beta_{g^2} = -\frac{22N_c}{3} \frac{g^4}{(4\pi)^2} \dots$$

RG Flow of Gluodynamics

▷ Operator expansion with the background-field method

(REUTER, WETTERICH'94; FREIRE, LITIM, PAWLOWSKI'00)

$$\Gamma_k[A] = \int d^d x W_k(F^2), \quad F^2 \equiv F_{\mu\nu}^a F_{\mu\nu}^a$$

$$W_k(F^2) = \frac{Z_F}{4} F^2 + \frac{W_2}{2!4^2} (F^2)^2 + \frac{W_3}{3!4^3} (F^2)^3 + \dots$$

(HG'02)

▷ running coupling:

$$g^2 = Z_F^{-1} \bar{g}^2$$

(ABBOTT'82)

▷ β function:

$$\partial_t g^2 \equiv \beta_{g^2} = -\frac{22N_c}{3} \frac{g^4}{(4\pi)^2} \dots$$

1	-29.3333
2	-357.83
3	-191.32
4	15499.6
5	-1.88776 · 10 ⁶
6	1.65315 · 10 ⁷
7	2.79324 · 10 ⁹
8	-1.37622 · 10 ¹¹
9	-4.21715 · 10 ¹²
10	8.60663 · 10 ¹⁴
11	-8.05611 · 10 ¹⁶
12	5.21052 · 10 ¹⁹
13	-6.30043 · 10 ²²
14	9.35648 · 10 ²⁵
15	-1.78717 · 10 ²⁹
16	4.35314 · 10 ³²
17	-1.33397 · 10 ³⁶
18	5.08021 · 10 ³⁹
19	-2.37794 · 10 ⁴³
20	1.35433 · 10 ⁴⁷

cf. Landau gauge:

(V.SMEKAL, ALKOEFER, HAUCK'97)

(LANGFELD, REINHARDT, GATTNAR'01)

(LERCHE, V.SMEKAL'02)

(FISCHER, ALKOEFER'02)

(ZWANZIGER'02)

(PAWLOWSKI ET AL.'03)

(FISCHER, HG'04)

(OLIVEIRA, SILVA'04)

(BLOCK, CUCCHIERI, LANGFELD, MENDES'04)

(SCHLEIFENBAU, LEDERER, REINHARDT'06)

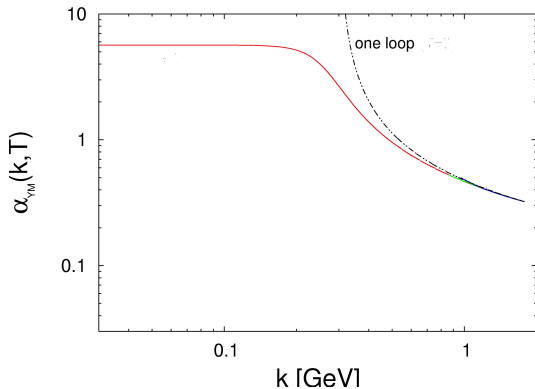
(EPPEL, REINHARDT, SCHLEIFENBAUM'07)

(MAAS'07)

(CUCCHIERI, MENDES, OLIVEIRA, SILVA'07)

(HG'02)

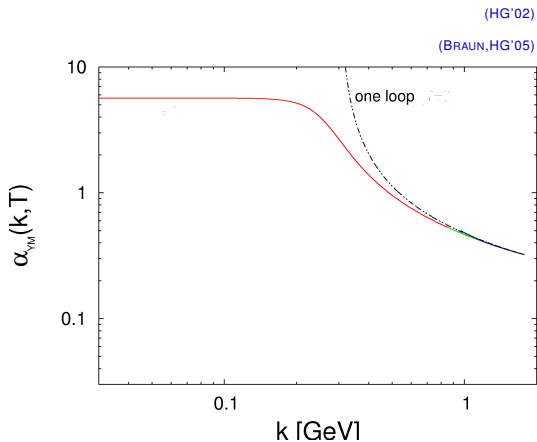
(BRAUN, HG'05)



IR fixed point: α_*

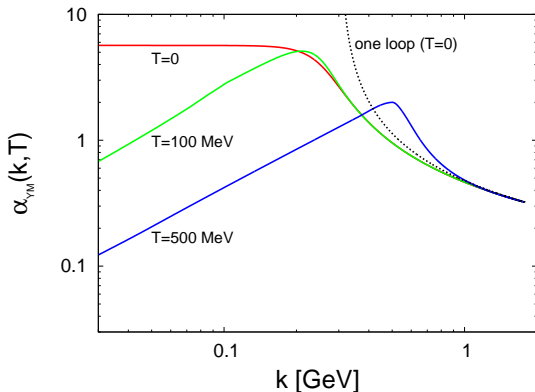
Running Coupling

- ▷ mass gap and threshold behavior:



IR fixed point α_* compatible with mass gap

Running Gauge Coupling at finite T



(HG'02)

(BRAUN, HG'05)

▷ $T/k \rightarrow \infty$: strongly interacting 3D theory

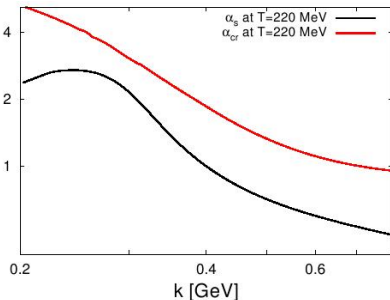
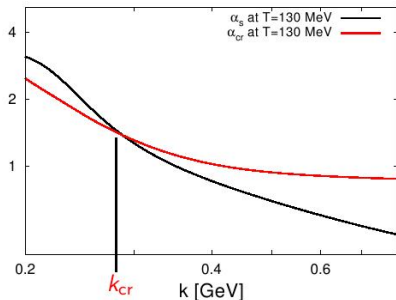
$$\alpha \rightarrow \frac{k}{T} \alpha_{3D}, \quad \alpha_{3D} \rightarrow \alpha_{3D,*} \simeq 2.7$$

cf. lattice: (CUCCHIERI, MAAS, MENDES'07)

Chiral Phase Transition



$\alpha(k, T)$ vs. $\alpha_{cr}(T/k)$



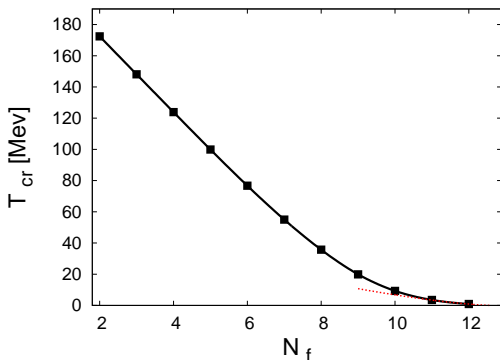
\Rightarrow χ SB triggered by α_s

single input: $\alpha_s(m_\tau) = 0.322$

T_c [MeV]	RG (BRAUN, HG'05)
$N_f=2$	172 ± 37
$N_f=3$	148 ± 32

T_c [MeV]	Lattice (BI) (CHEN ET AL.'06)	Lattice (W) (AOKI ET AL.'06)
$N_f=2+1$	$192(7)(4)$	$151(3)(3)$

Chiral Phase Boundary $T - N_f$



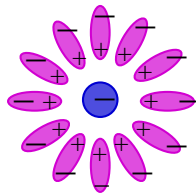
(BRAUN, HG'05,'06)

▷ small N_f : fermionic screening, $\beta_{\text{quark}} \simeq \frac{2}{3} N_f \frac{g^4}{8\pi^2}$

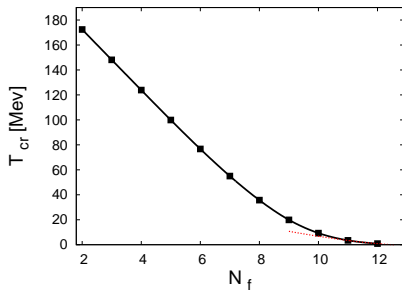
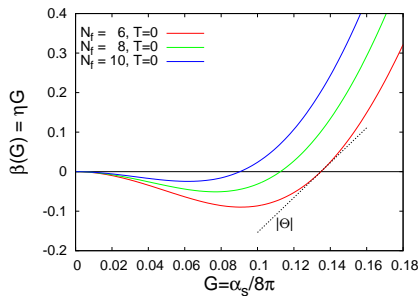
▷ critical flavor number:

$$N_f^{\text{cr}} \simeq 12$$

(CF. APPELQUIST ET AL.'96; MIRANSKI, YAMAWAKI'96; HG, JAECKEL'05)



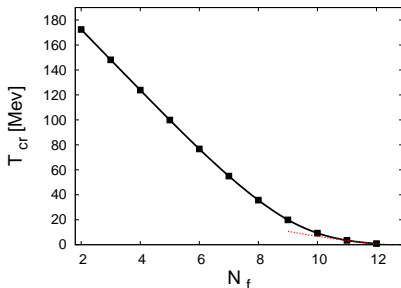
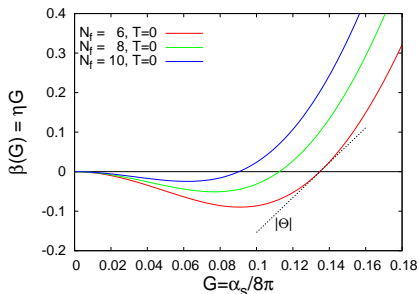
Chiral Phase Boundary $T - N_f$



▷ fixed-point regime: critical exponent Θ

$$\beta_{g^2} \simeq -\Theta (g^2 - g_*^2)$$

Chiral Phase Boundary $T - N_f$



▷ fixed-point regime: **critical exponent Θ**

$$\beta_{g^2} \simeq -\Theta (g^2 - g_*^2)$$

▷ shape of the phase boundary for $N_f \simeq N_f^{cr}$:

(BRAUN, HG'05,'06)

$$T_{cr} \sim k_0 |N_f - N_f^{cr}|^{\frac{1}{|\Theta|}}, \quad \Theta \simeq -0.71$$

- ▶ “conformal phase” in many-flavor QCD:

$$N_{f,cr} \simeq 10 - 12 < N_f < 16.5 \text{ for SU}(3)$$

... applications to walking technicolor

(DIETRICH ET AL.'06, TERAO ET AL.'07)

- ▶ relation among universal aspects:

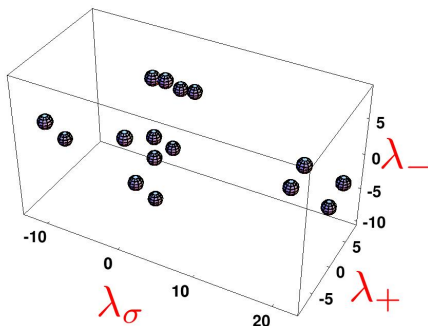
shape of the phase boundary \iff IR critical exponent

- ▶ functional RG for $\Gamma[\phi]$

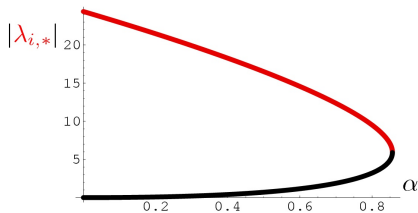
- systematic and consistent expansion schemes for QCD
- chiral symmetry ✓
- calculations “from first principles”

RG Flow of the Chiral Sector

- ▷ 2 fixed points per $\partial_t \lambda$
- ⇒ $2^4 = 16$ fixed points
- ▷ in general: 2^n FP's
for $n = \#$ of λ 's



- ▷ fixed-point annihilation
- e.g., $N_c = N_f = 3$



Chiral Criticality at Finite Temperature

▷ quark modes:

$$m_T^2 = m_f^2 + (2\pi T(n + \frac{1}{2}))^2$$

⇒ T -dependent

critical coupling:

$$\alpha_{\text{cr}}(T) \gtrsim \alpha_{\text{cr}} \simeq 0.85$$

(BRAUN, HG'05)

