



# Baryogenesis

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## Summary

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Is the Universe lepton  
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Ways to a Baryon  
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When the Universe  
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B and CP violations

EW phase transition  
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It is possible that far away galaxies could be made of antimatter ...

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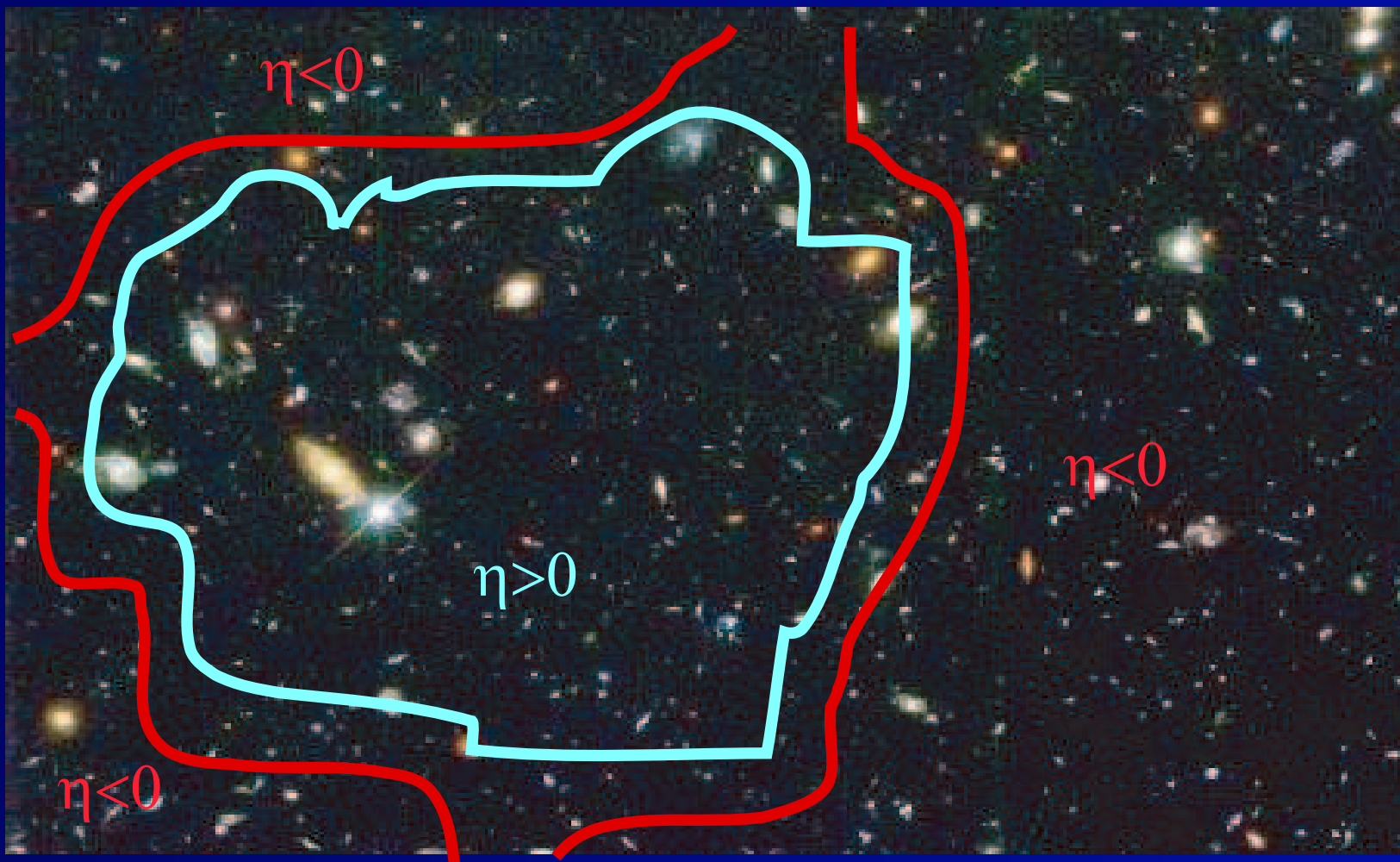
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In this case we live in a bubble where the baryon asymmetry is positive...

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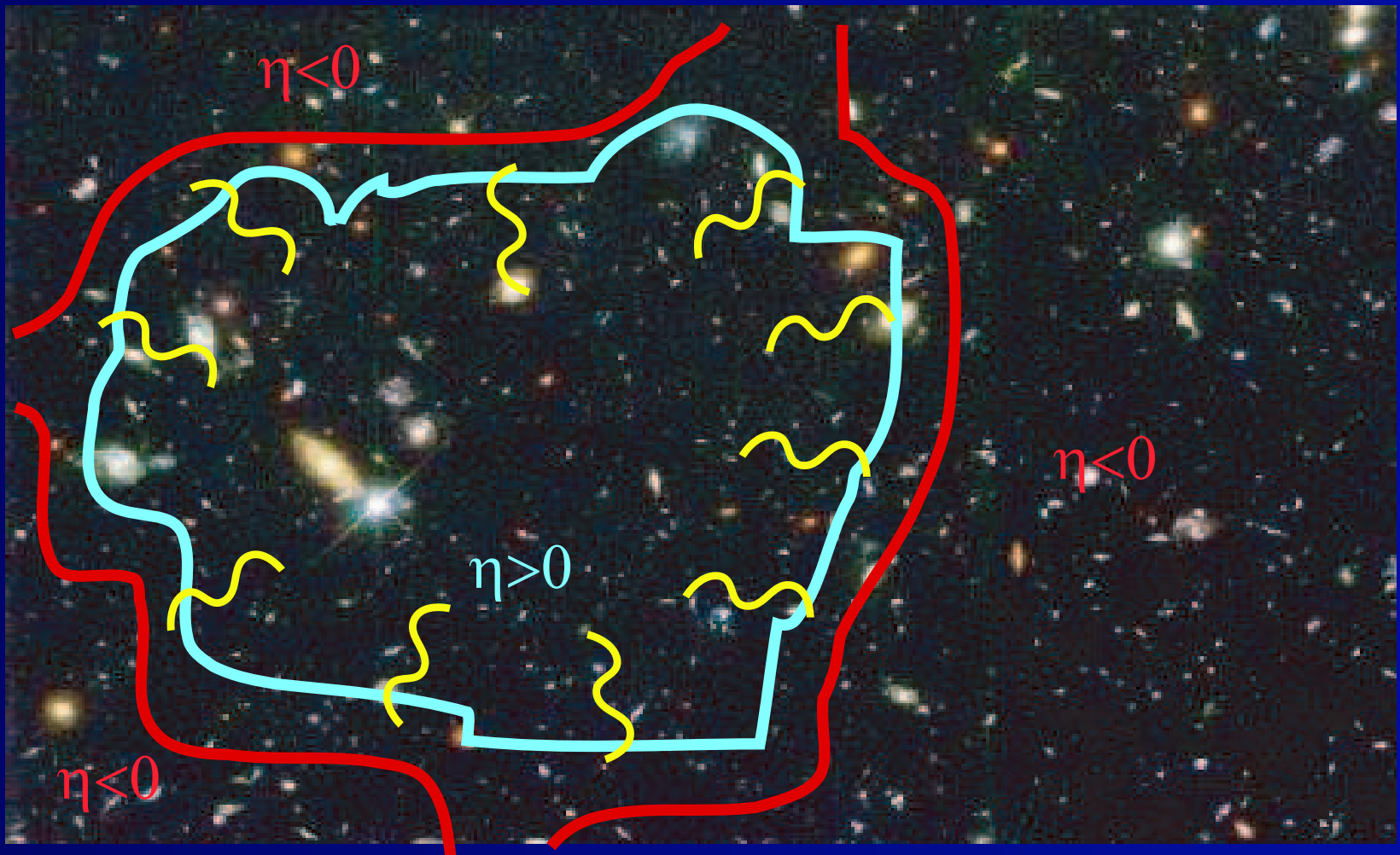
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... but annihilation  $\gamma$ -rays should be produced at the boundary of the bubble.

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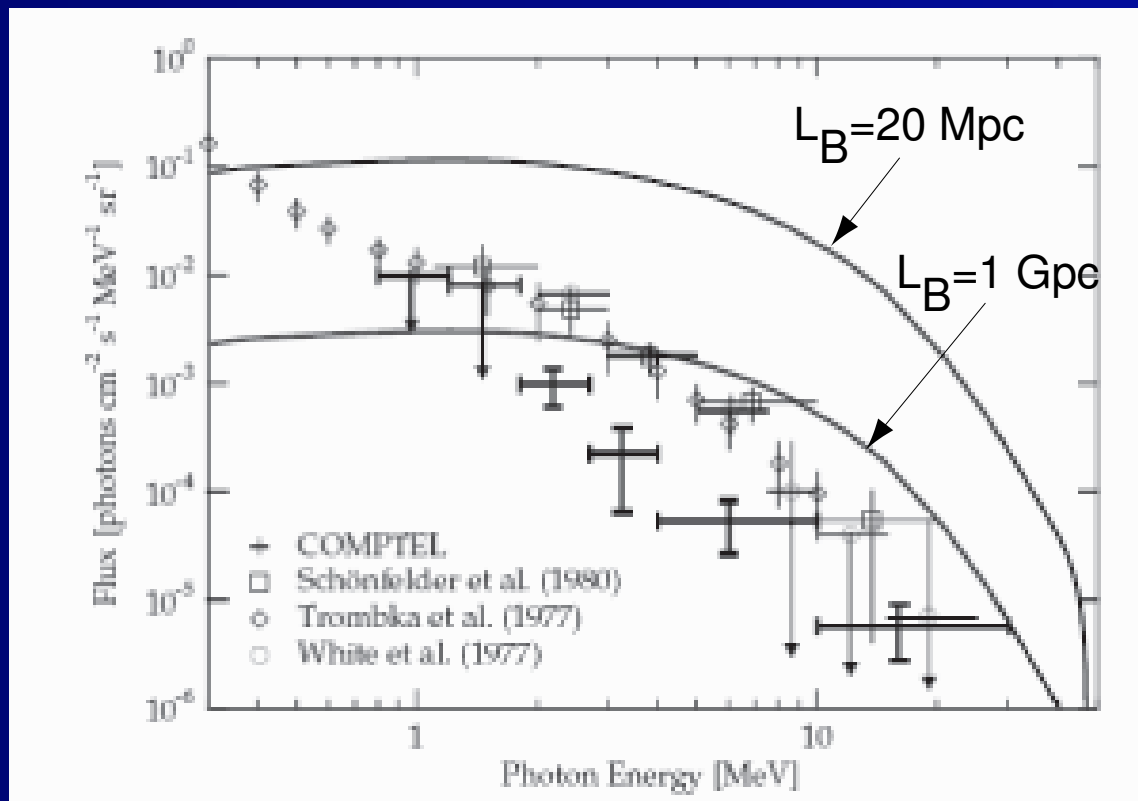
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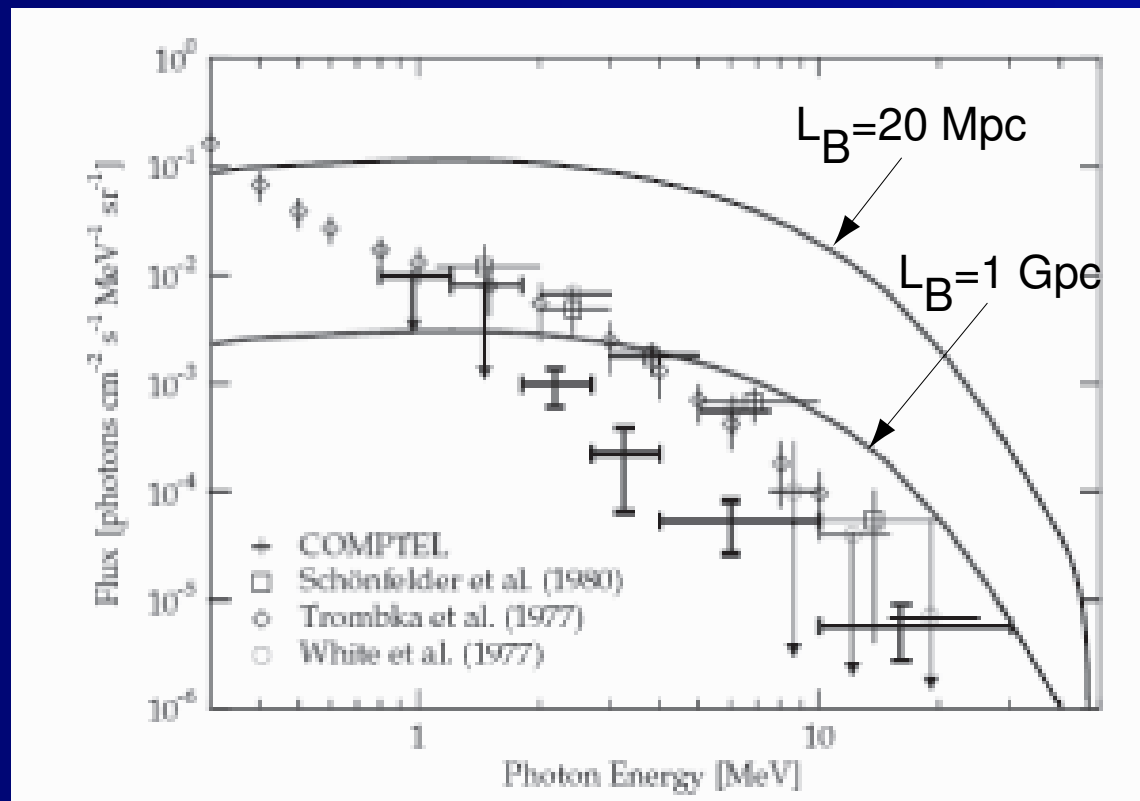
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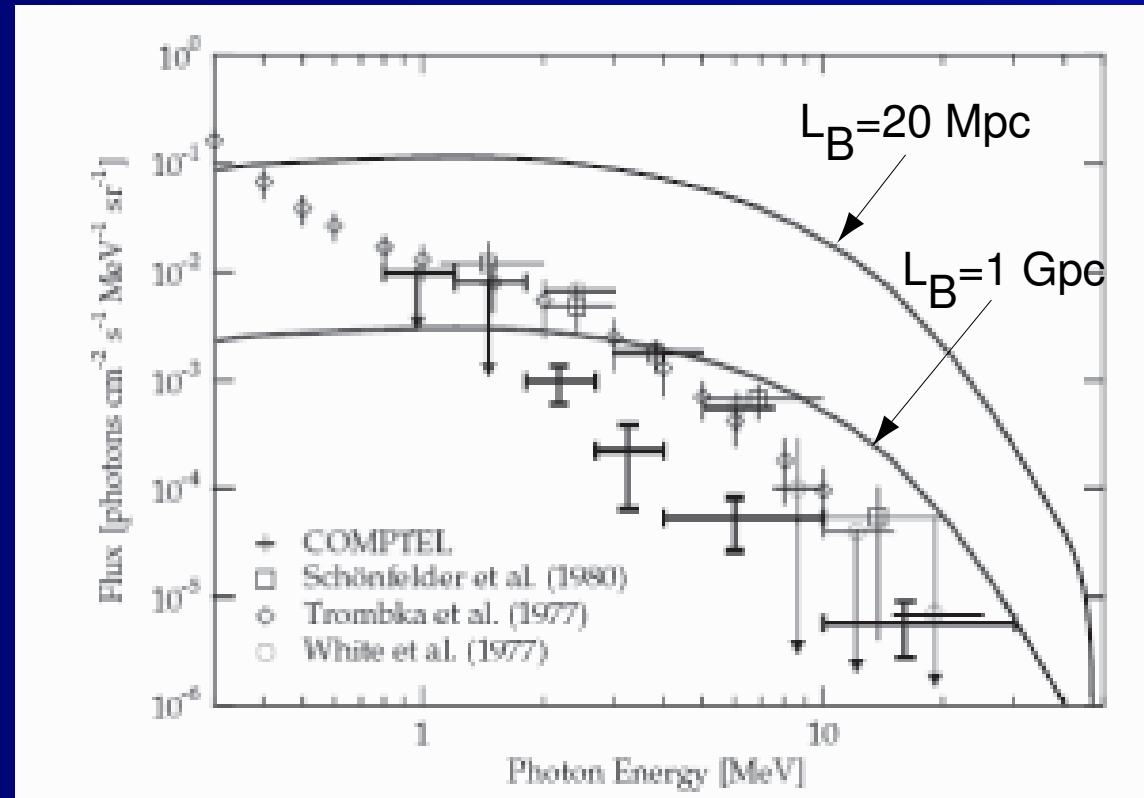
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- ➔ Diffuse  $\gamma$ -ray background from matter-antimatter annihilation at the domain wall, calculated by Cohen, Glashow & De Rujula (1998).
- ➔ Predicted flux is compatible with experimental data points if  $\ell_B \geq 1$  Gpc.



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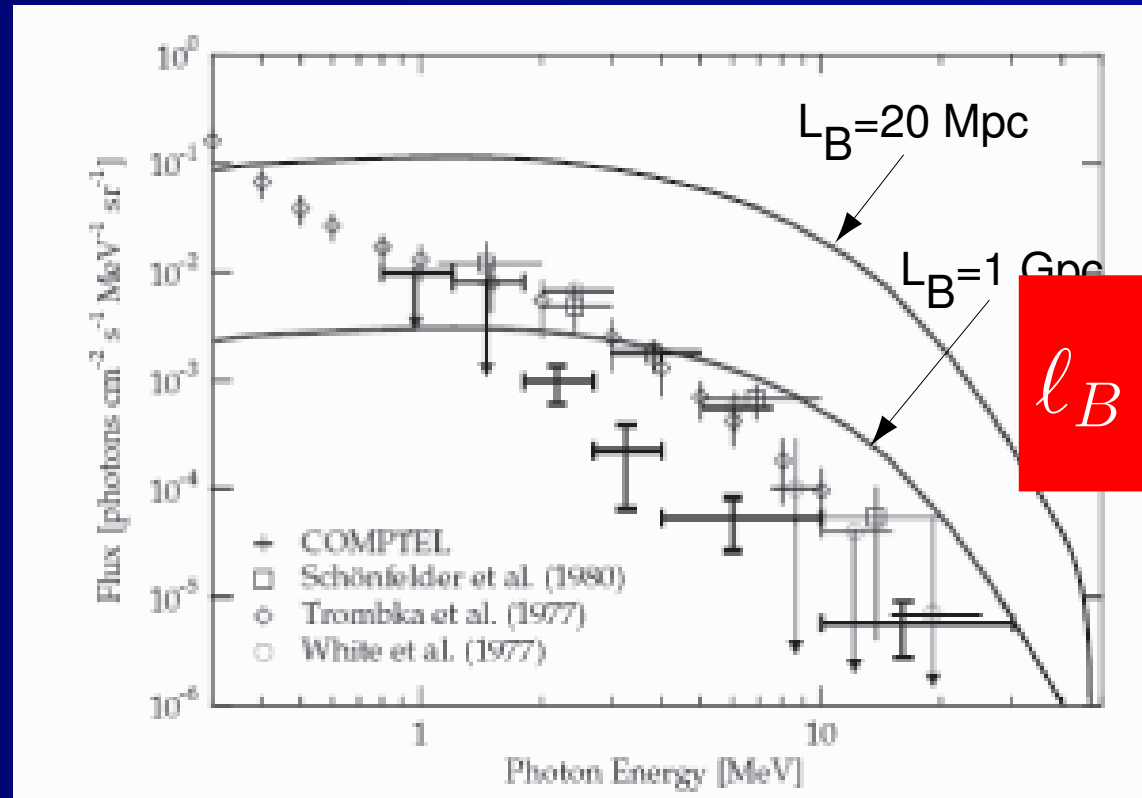
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$$\ell_B \geq \frac{1}{4} ct_0$$

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WMAP fit	$\Omega_b h^2 = 0.0223 \pm 0.0008$		Spergel <i>et al.</i> (2006)
SBBN with D/H ratio at $z > 2$	$\Omega_b h^2 = 0.0214 \pm 0.0020$		Kirkman <i>et al.</i> (2003)
Lyman- $\alpha$ forest $z > 2$	$\Omega_b h^2 \gtrsim 0.018$		Weinberg <i>et al.</i> (1997)
Visible baryons $z < 2$	E+S0 Galaxies	$1.5 \times 10^{-3}$	Persic & Salucci (2005)
	S Galaxies	$0.7 \times 10^{-3}$	
	Gas Clusters	$2.4 \times 10^{-4} h^{1.5}$	
	Intragroup gas	$1.8 \times 10^{-4} h^{1.5}$	
	Total Visible	$2.2 \times 10^{-3} + 4.2 \times 10^{-4} h^{1.5}$	
Lyman- $\alpha$ forest $z < 2$	$\Omega_b h = 0.015 \pm 0.002$		Penton <i>et al.</i> (2004)
Warm-Hot IM $z < 2$	$\Omega_b h \gtrsim 0.005 \pm 0.002$		Tripp <i>et al.</i> (2000)

➡ The number density of baryons, producing the Lyman- $\alpha$  forest, detected at redshifts  $z > 2$ , appear to be more or less consistent with CMBR fluctuations and BBN while the baryonic fraction detected at  $z < 2$  add up to just over half ( $\sim 55\%$ ) of the number seen at  $z > 2$ , meaning that about 45 percent are 'missing'.

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WHIM $z < 2$ (CHANDRA)	$\Omega_b = 0.027^{+0.038}_{-0.019}$		Nicastro <i>et al.</i> (2005)

New baryonic component detected ! (Nicastro et al. Nature 2005)

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➡ The visible baryons segregated in stars, galaxies and intracluster gas is about 10-15 % of the total

➡ Apparently most of the baryons are dispersed in filaments of the warm-hot intergalactic medium with density

$$n_B \approx (0.5 - 1) \times 10^{-6} \text{ cm}^{-3} \text{ and temperature}$$

$T \approx 10^5 - 10^7 \text{ }^\circ\text{K}$ , detectable as absorbers of soft X-rays.

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➔ The best estimate of the baryon asymmetry parameter is obtained from WMAP  $\Lambda$ CDM best fit:

$$\eta_B = \frac{n_B - n_{\bar{B}}}{s} = (4.7 \pm 0.16) \times 10^{-11}$$

consistent with BBN.

➔ Exception made for periods of the three phase transition that we have described before, a Friedmann-Lemâitre Universe evolves at constant entropy:  $S \propto g^*(T) T^3 R^3 = \text{const}$

➔ The number density of stable particles is  $n \propto R^3$  therefore the baryon asymmetry  $\eta_B$  is proportional to  $1/g^*(T)$ .

➔ Chemical potential of baryons was established by baryogenesis mechanism to be  $\frac{\mu_B}{T} \simeq 10^{-12}$  well before  $t_{BBN}$

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➔ Overall charge neutrality of the Universe imposes that the electron asymmetry should be

$$\eta_{L_e} = \eta_B$$

➔ On the contrary the neutrino asymmetry could be

$$\eta_{L_\nu} \neq \eta_B$$

➔ Lepton asymmetry implies neutrino degeneracy, because at  $T \gtrsim 1 \text{ MeV}$  is

$$\eta_L = \frac{60}{43\pi^2} \sum_i g_i \left[ \frac{\mu_i}{T} + \frac{1}{\pi^2} \left( \frac{\mu_i}{T} \right)^3 \right]$$

➔ Neutrino degeneracy has an impact on BBN, which allows to set limits to the lepton asymmetry

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- ➔ A neutrino chemical potential modifies the outcome of primordial nucleosynthesis in two different ways:
- ➔ The first effect appears only in the  $\nu_e$  sector because equilibrium is maintained by the beta process

$$\left\{ \begin{array}{l} p + e^- \rightleftharpoons n + \nu_e \\ n + e^+ \rightleftharpoons p + \bar{\nu}_e \end{array} \right\} \Rightarrow n/p \propto e^{-\frac{\mu_{\nu_e}}{T}}$$

Therefore  $\mu_{\nu_e} > 0$  decreases  $Y_p$ , the primordial  ${}^4He$  mass fraction, while  $\mu_{\nu_e} < 0$  increases it.

- ➔ A second effect is an increase of the neutrino energy density for any non-zero  $\mu_{\nu}$  which is

$$\rho_{\nu+\bar{\nu}} = \frac{7\pi^2}{120} T^4 \sum_{i=e,\mu,\tau} \left[ 1 + \frac{30}{7\pi^2} \left( \frac{\mu_{\nu_i}}{T} \right)^2 + \frac{15}{7\pi^2} \left( \frac{\mu_{\nu_i}}{T} \right)^4 \right]$$

That increases the expansion rate of the universe, thus enhancing  $Y_p$ .

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➡ However it has been shown (Dolgov *et al.* 2002) that neutrino oscillations mix all the flavors, therefore we may assume that all the chemical potentials are equal independent of flavors.

➡ The BBN limits on the electron neutrino chemical potentials are  $\mu_{\nu_e} = 0.03 \pm 0.04$ , which converts into a limit to the lepton asymmetry

$$-0.004 \leq \eta_L \leq 0.030 \quad (95\% \text{ C.L.})$$

Neutrino degeneracy could be much larger than baryon asymmetry

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➔ However it has been shown (Dolgov *et al.* 2002) that neutrino oscillations mix all the flavors, therefore we may assume that all the chemical potentials are equal independent of flavors.

➔ The BBN limits on the electron neutrino chemical potentials are  $\mu_{\nu_e} = 0.03 \pm 0.04$ , which converts into a limit to the lepton asymmetry

$$-0.004 \leq \eta_L \leq 0.030 \quad (95\% \text{ C.L.})$$

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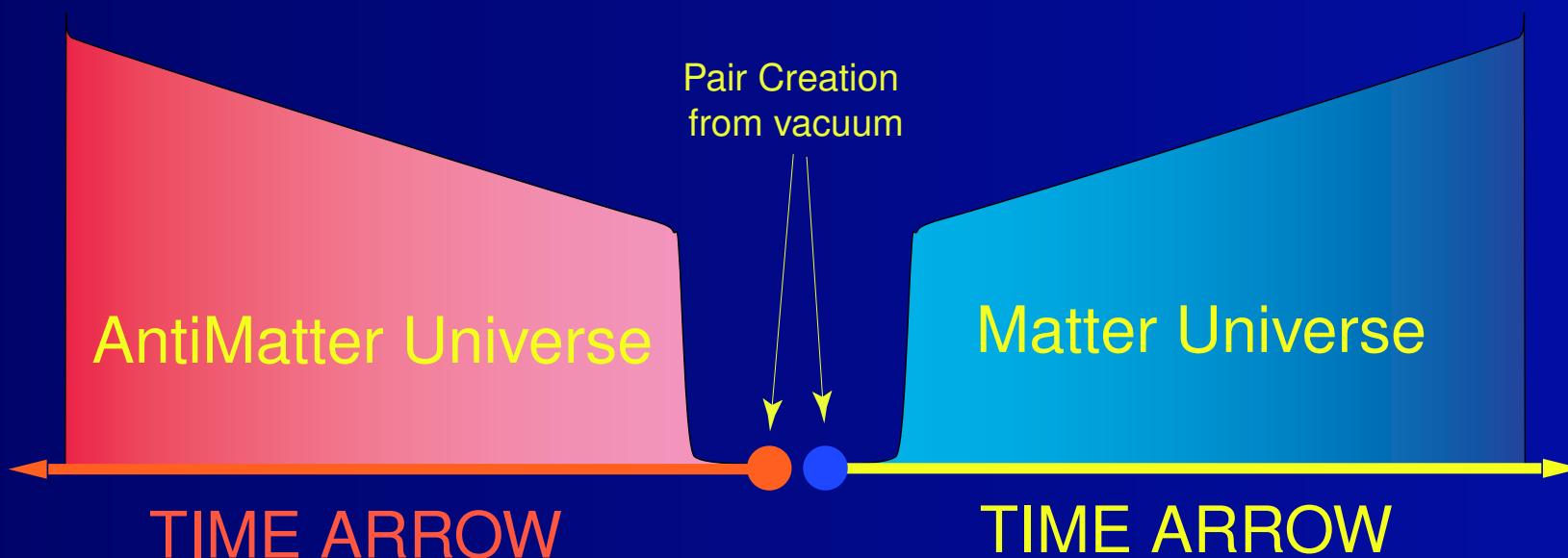
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- ➡ M. Goldhaber in a paper entitled “*Speculation on cosmogony*” published in *Science* in 1933, soon after the discovery of the anti-proton in cosmic rays, proposed a modification of the Lemâitre hypothesis of “Primeval Atom”, in which instead of one single Universe, a pair of Universe and anti-Universe were created from vacuum.
- ➡ Forgotten for a long time this type of speculations had a recent revival in the quite different form of “Parallel Universes”
- ➡ In the Goldhaber theory the two universe can coexist spatially, because they are separated in time.

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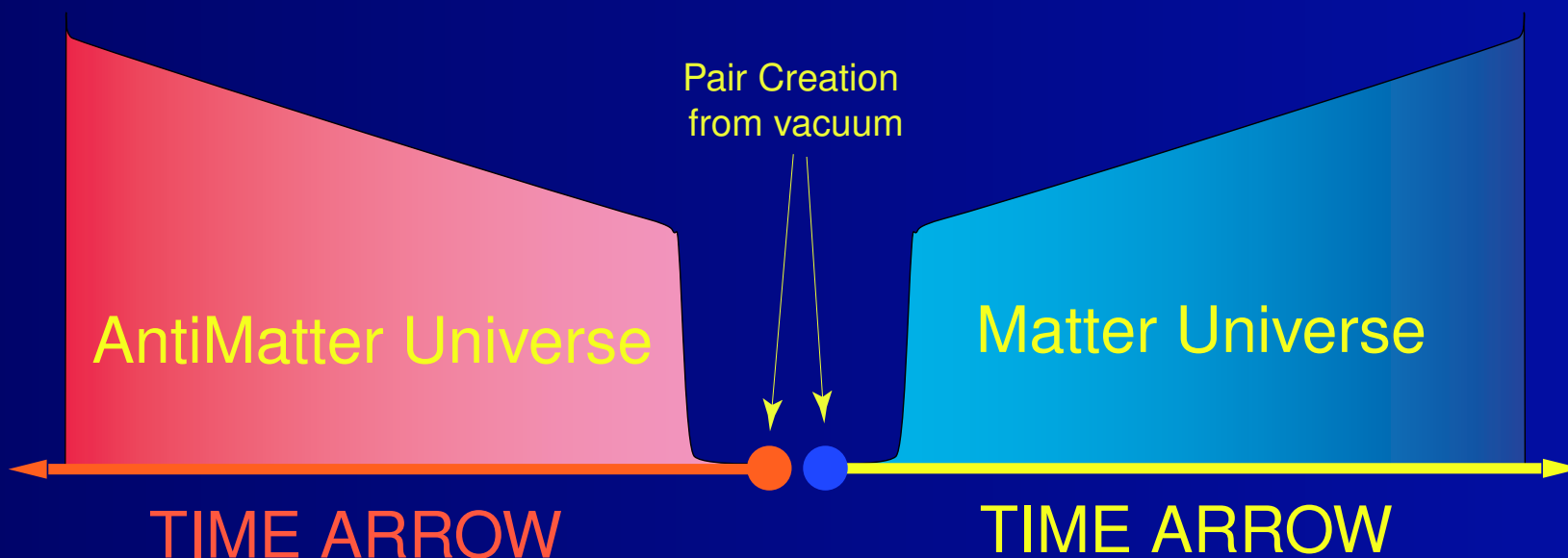
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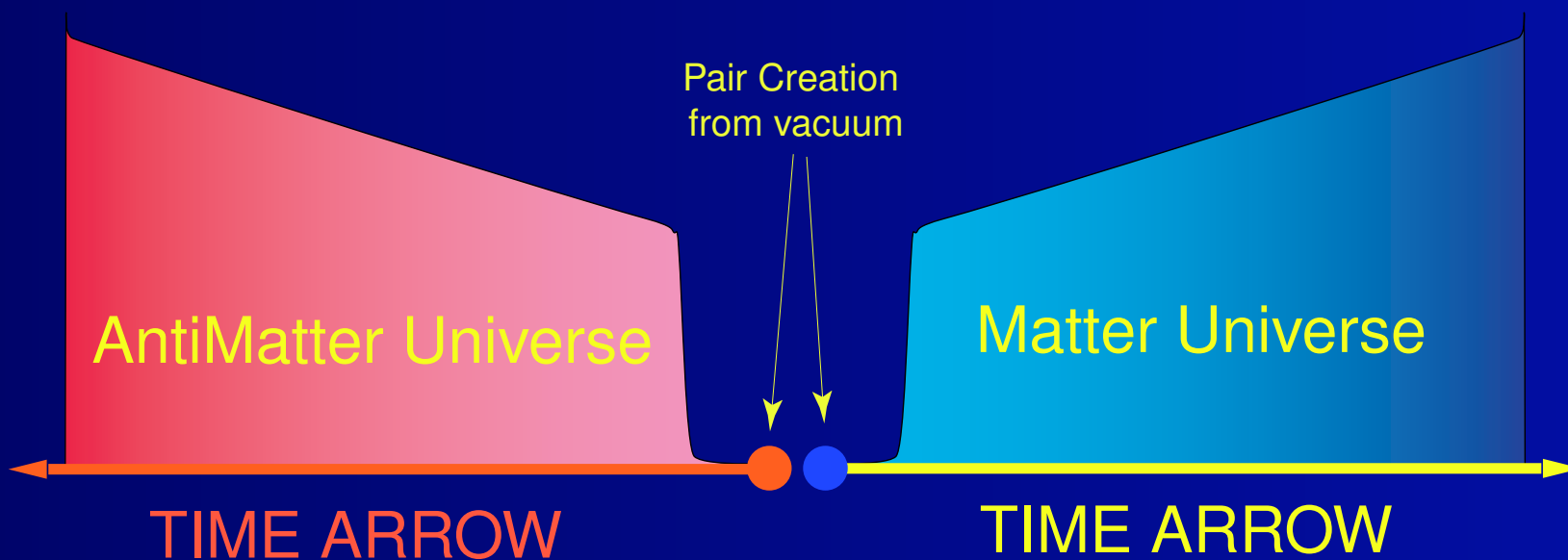
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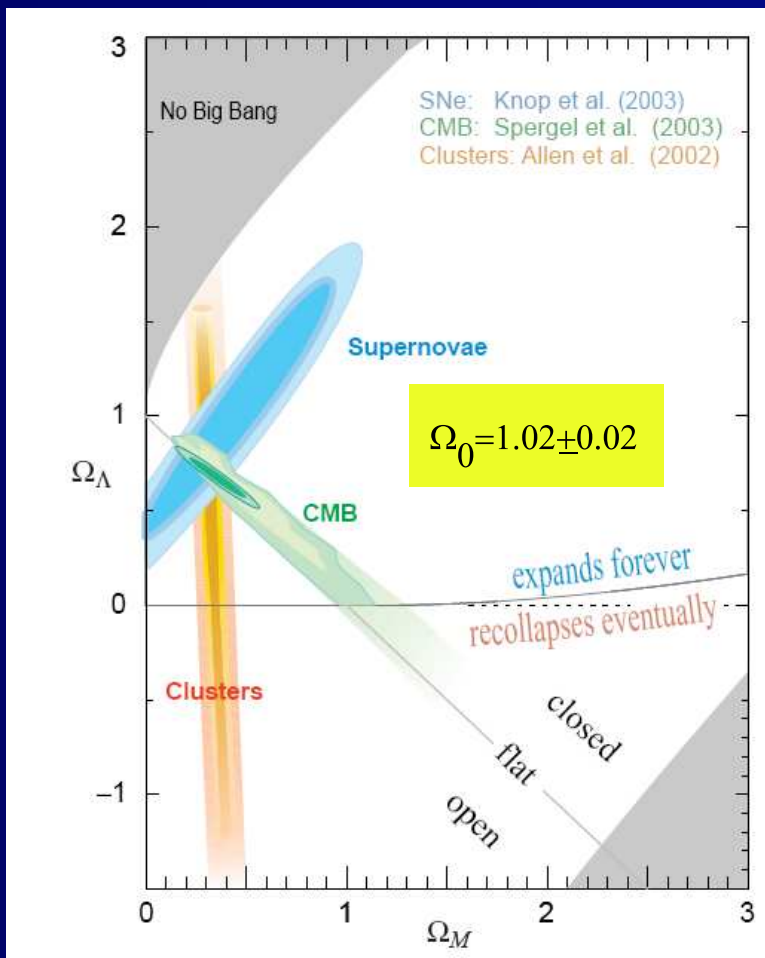
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➔ The initial condition  $B \neq 0$  is confronted with a formidable difficulty, because after WMAP is impossible to avoid inflation.

➔ In the inflationary phase any initial baryon asymmetry  $\eta_B$  would be diluted by a factor  $\propto e^{-3N} \approx 10^{-1.23N}$ , where  $N$  is the number of e-fold expansion of the scale factor.

➔ If we take the baryon asymmetry parameter  $\eta_B \approx 6 \times 10^{-11}$  we expect in simple model  $N < 14$ . A value which is firmly excluded from the requirement to form early objects weighing a million solar masses (Liddle & Leach 2003).



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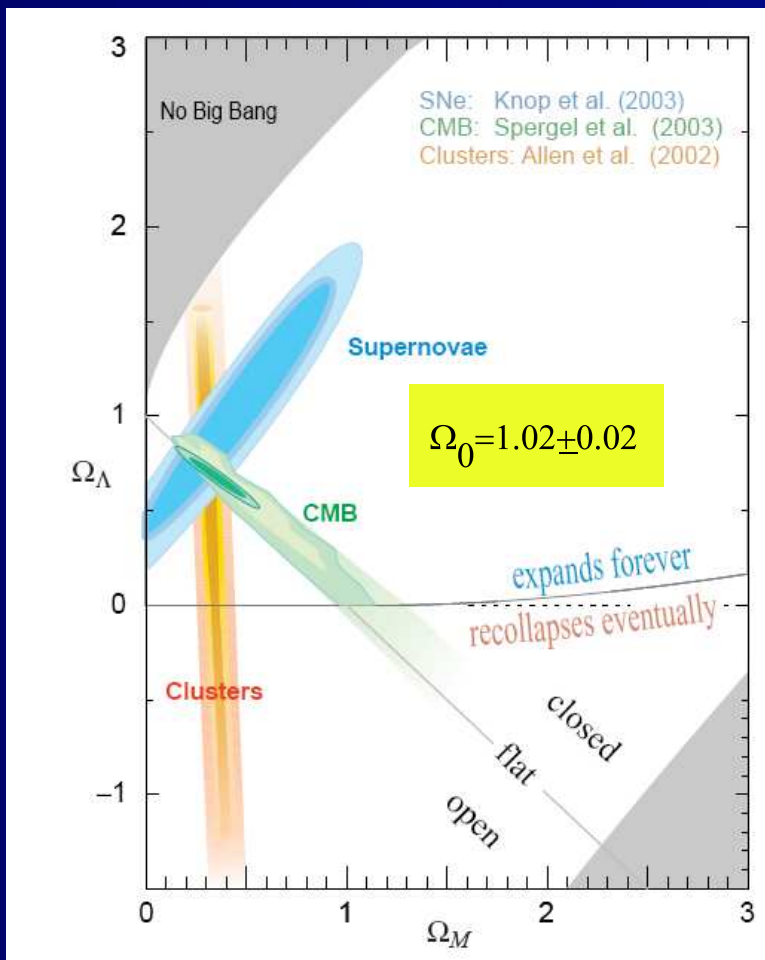
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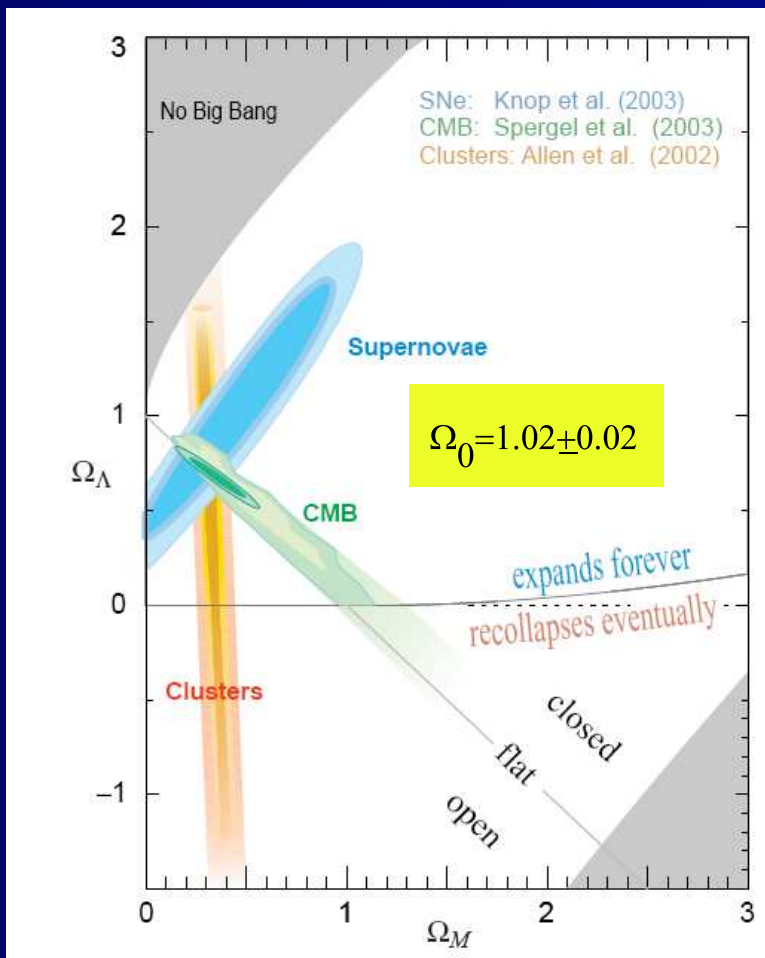
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Assuming that the universe started with a complete matter-antimatter symmetry in a standard big bang picture, one can obtain matter-antimatter asymmetry in the universe provided that the following three conditions are satisfied (Sakharov 1967):

1. Underlying theory should allow reactions with  $\Delta B \neq 0$ .
2. Charge, and CP symmetry must be violated, otherwise the excess of baryons created would be washed out by inverse reactions. So even if B is violated, one can never establish baryon-antibaryon asymmetry unless C and CP are violated.
3. However, if a baryon asymmetry has been generated at a certain time, a suppression of the B-violating processes is necessary. Otherwise if all processes, including those which violate baryon number, are in thermal equilibrium, the baryon asymmetry vanishes.

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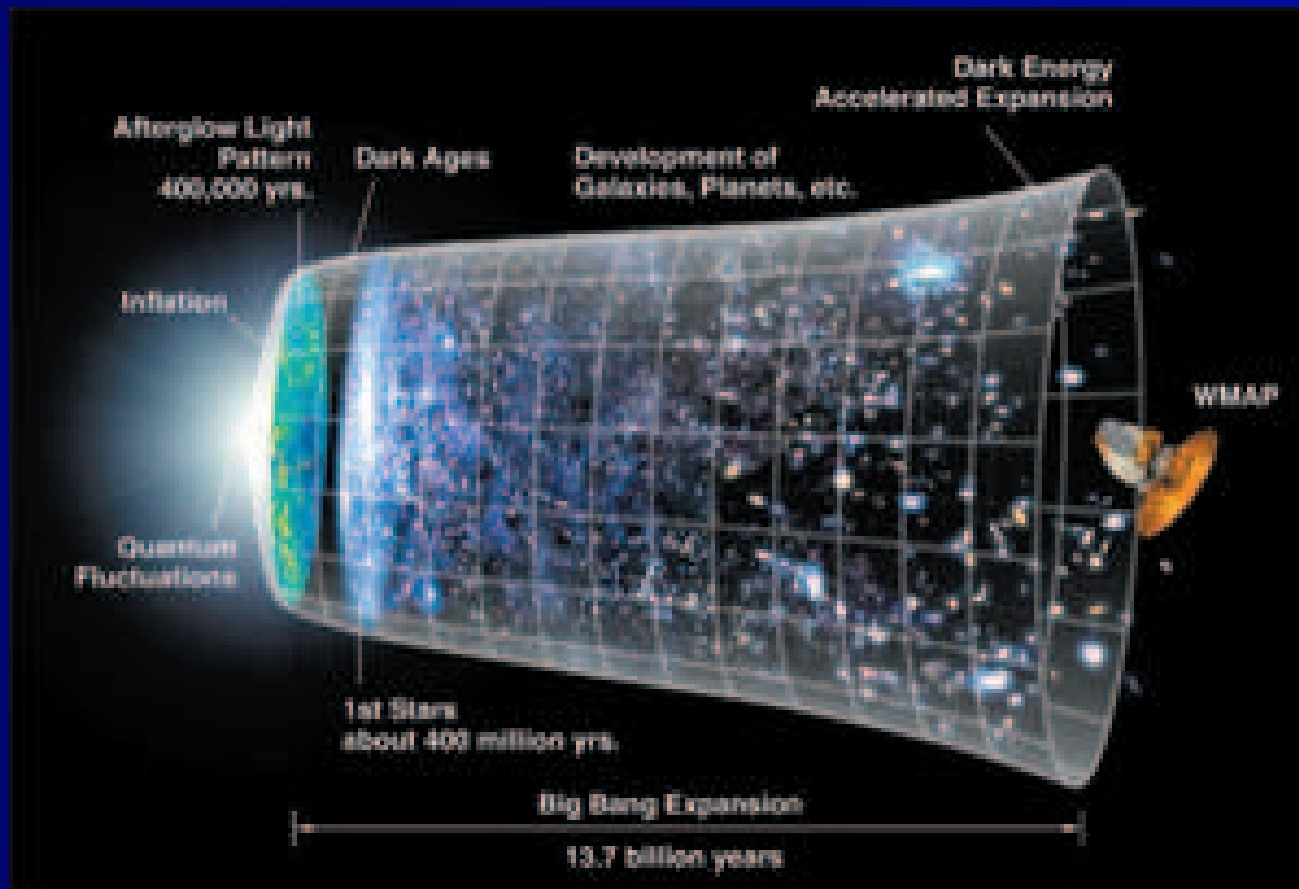
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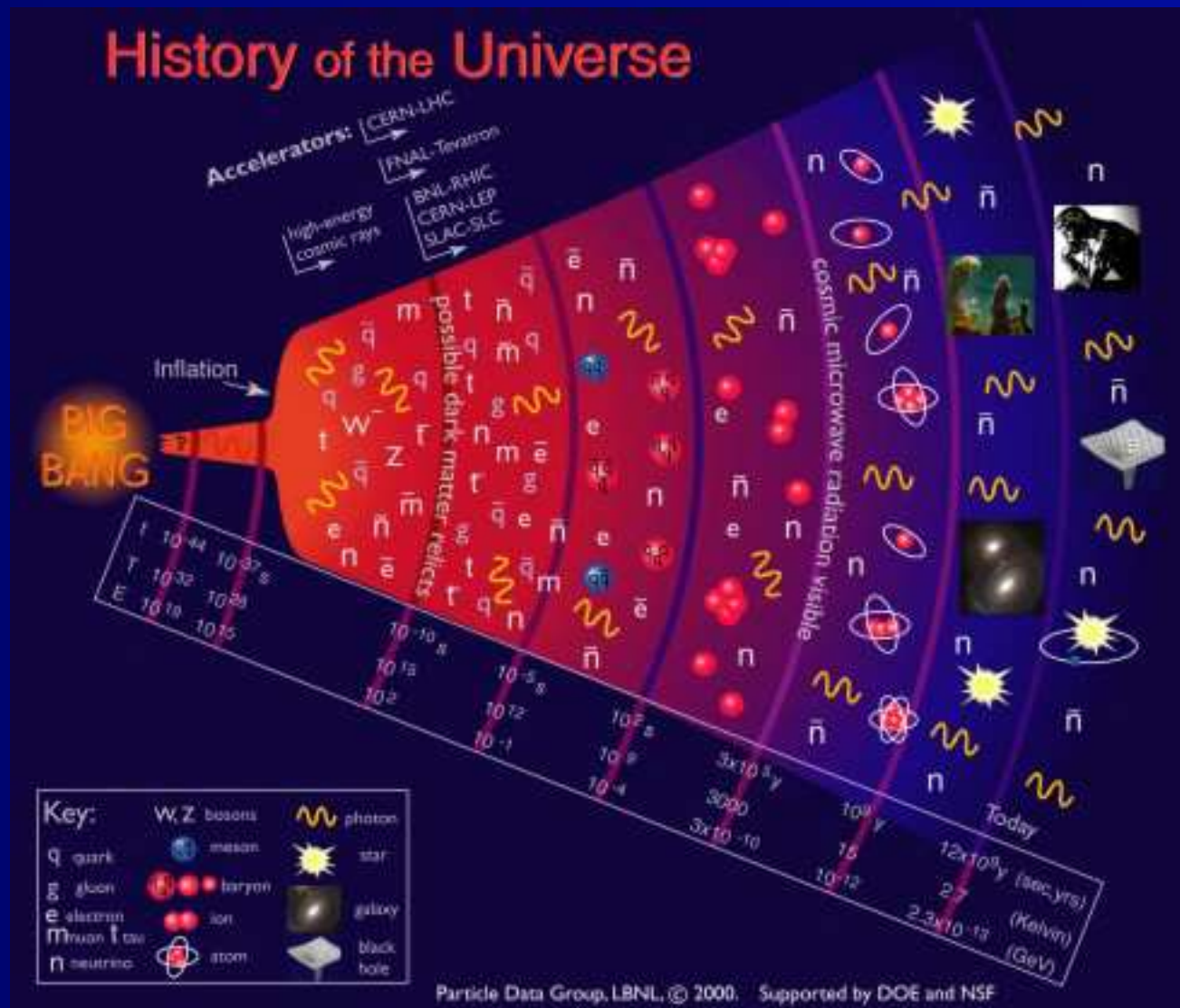
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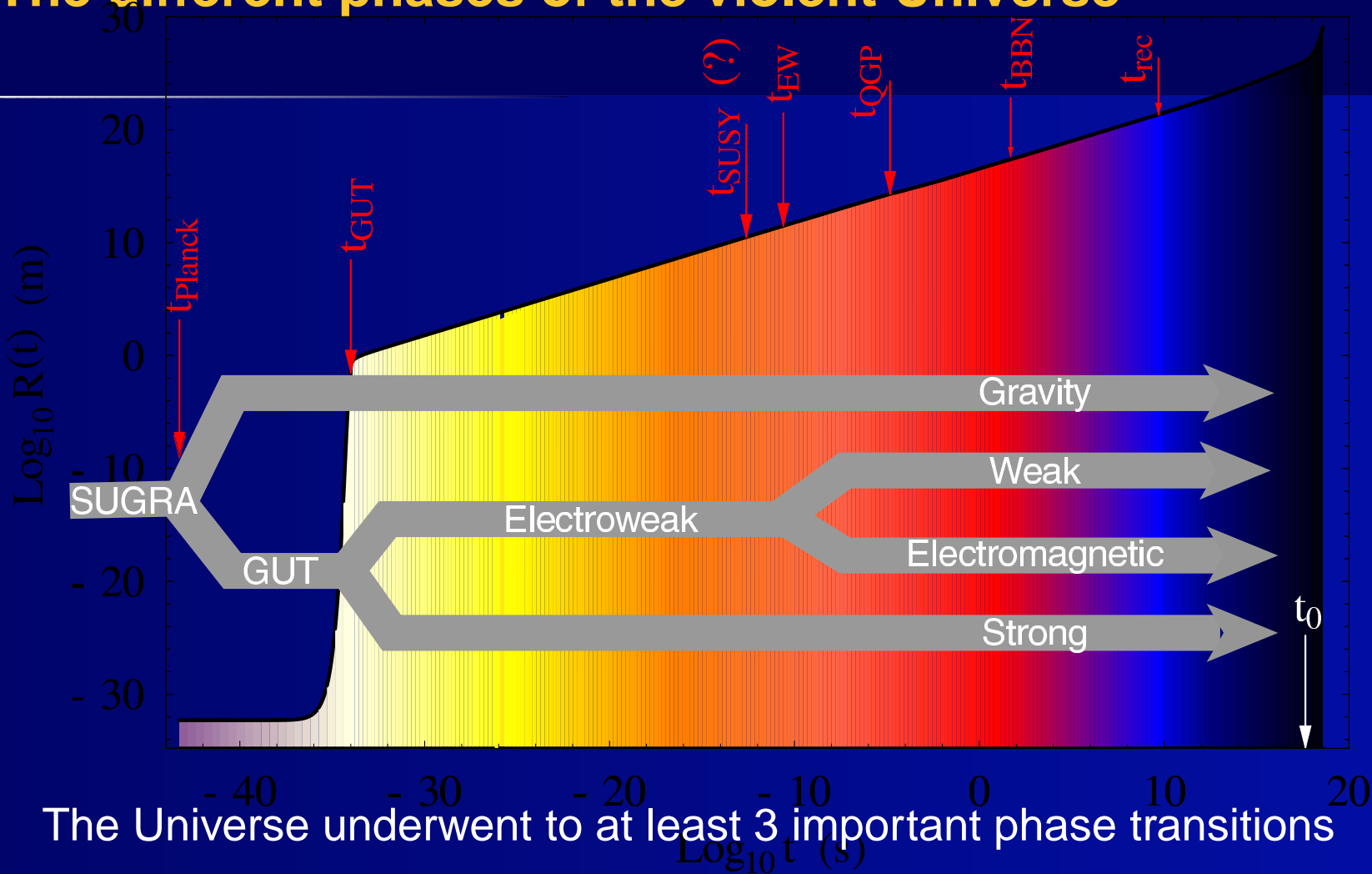
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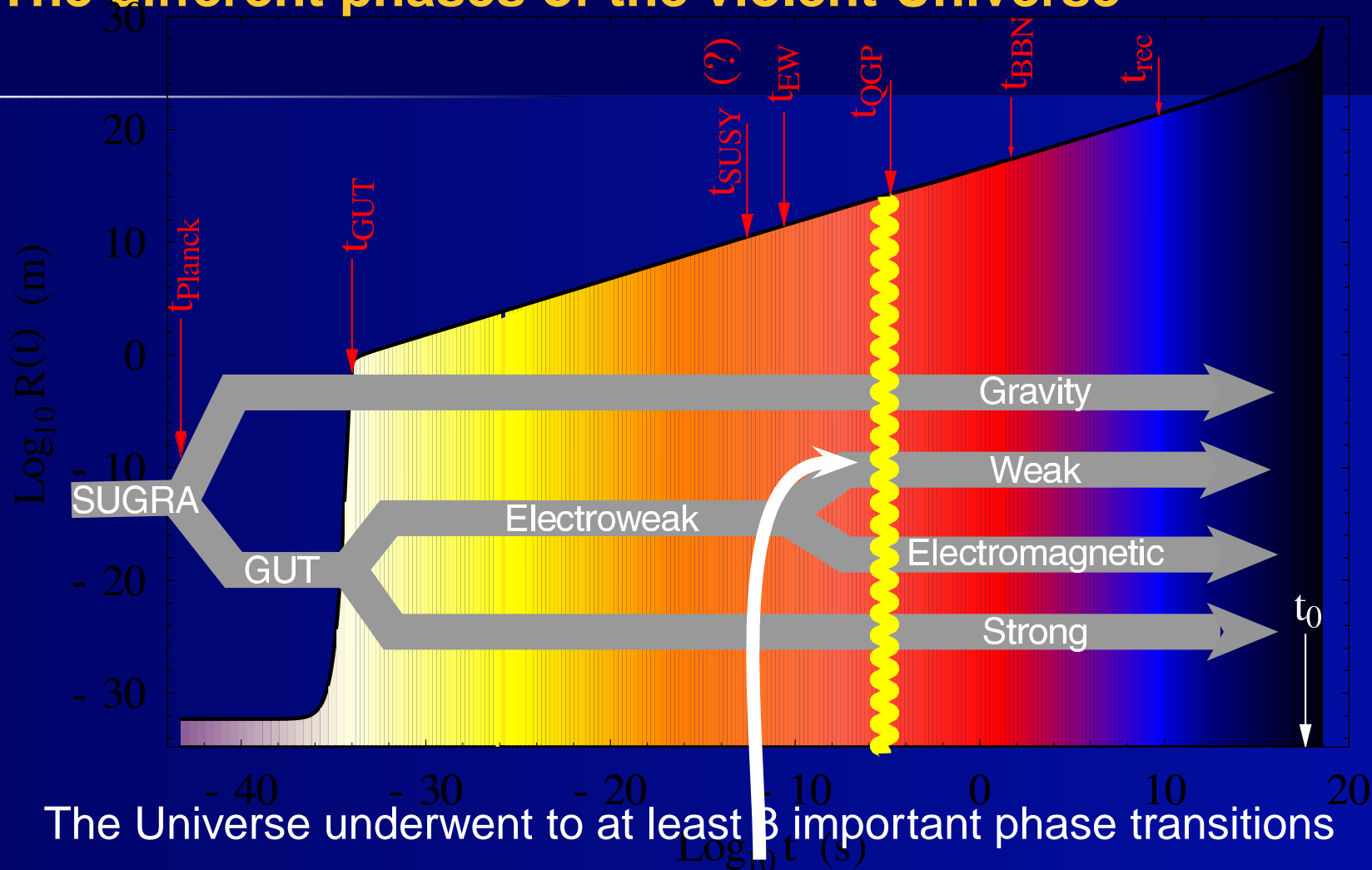
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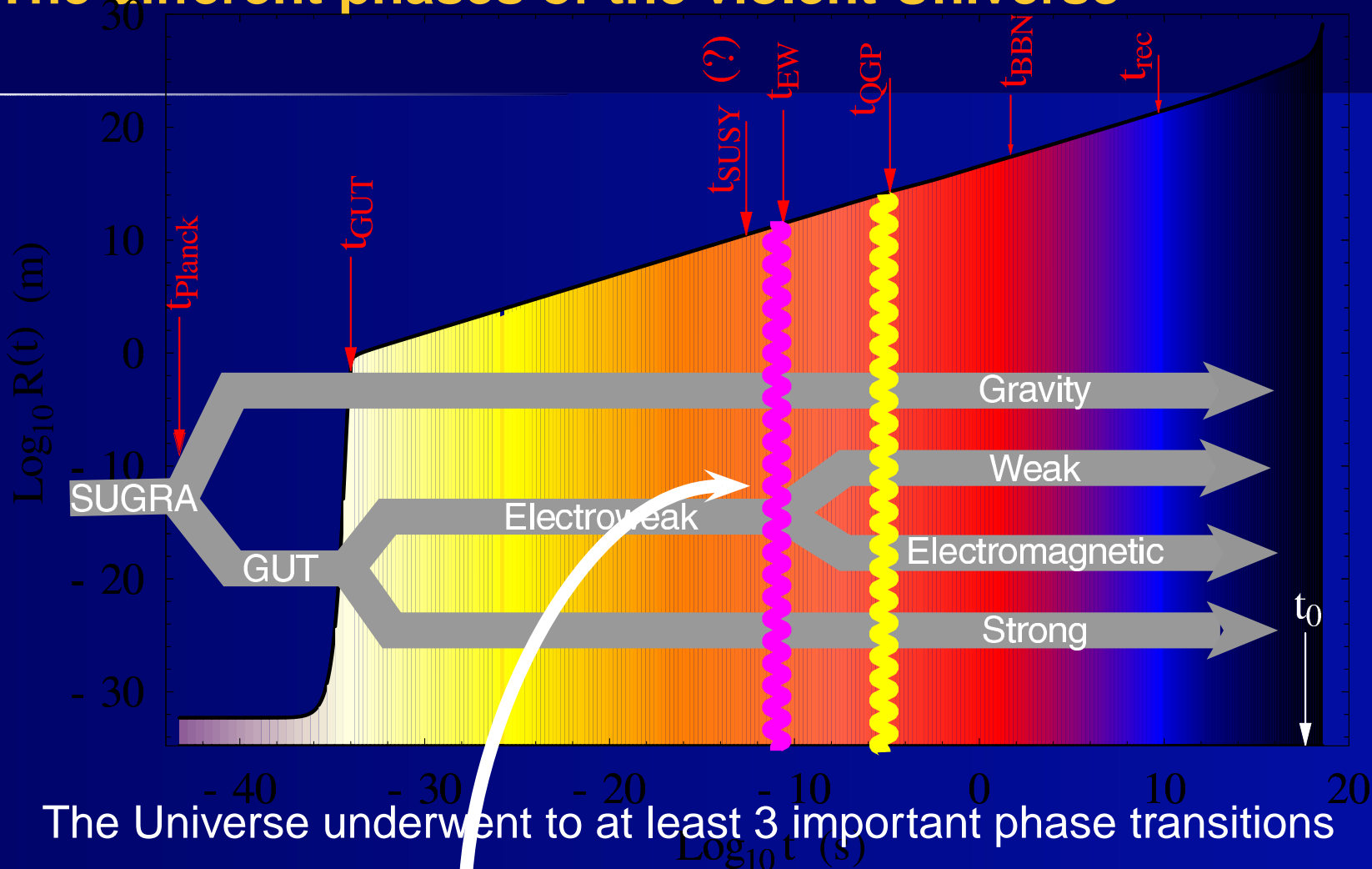
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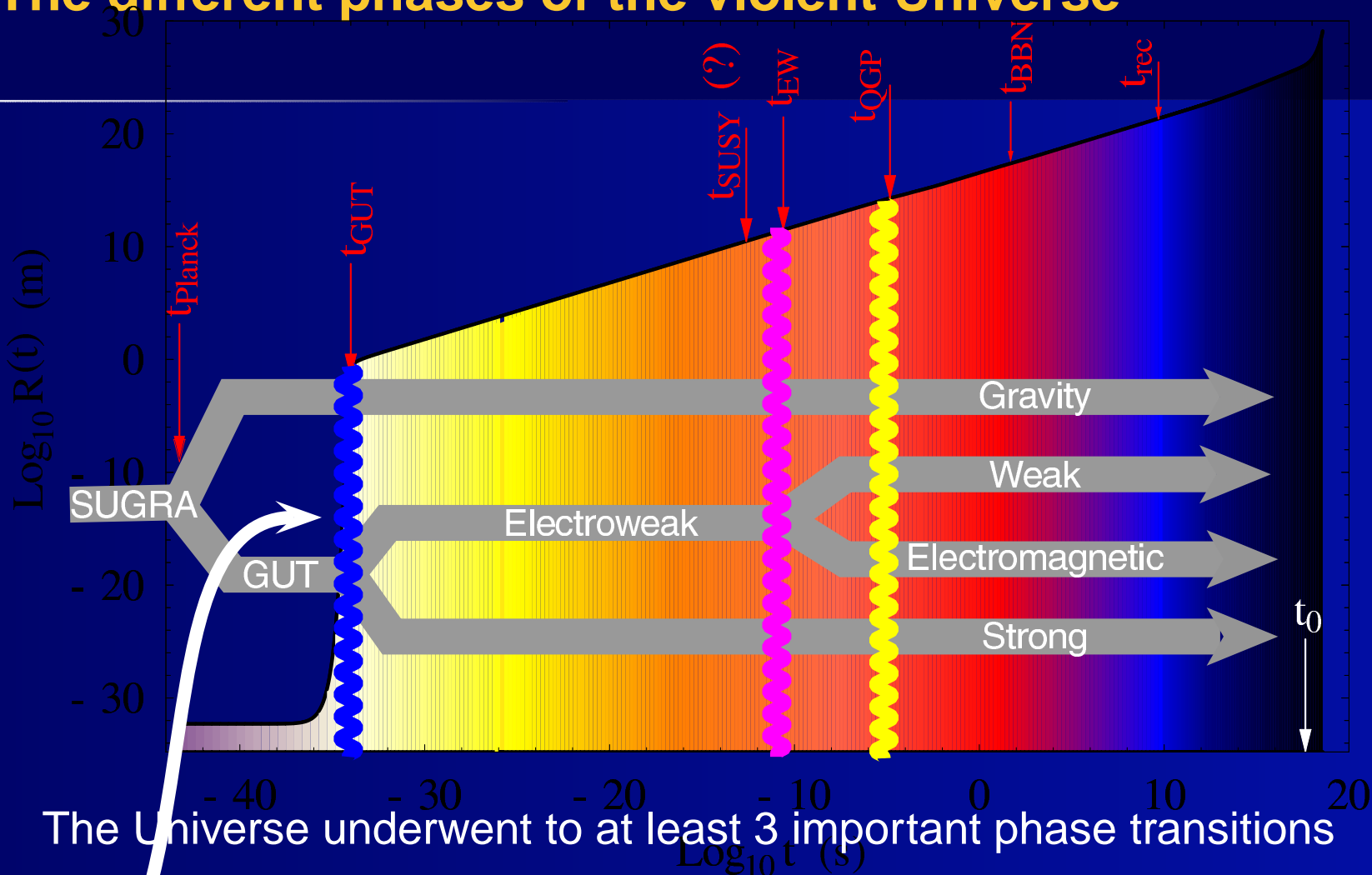
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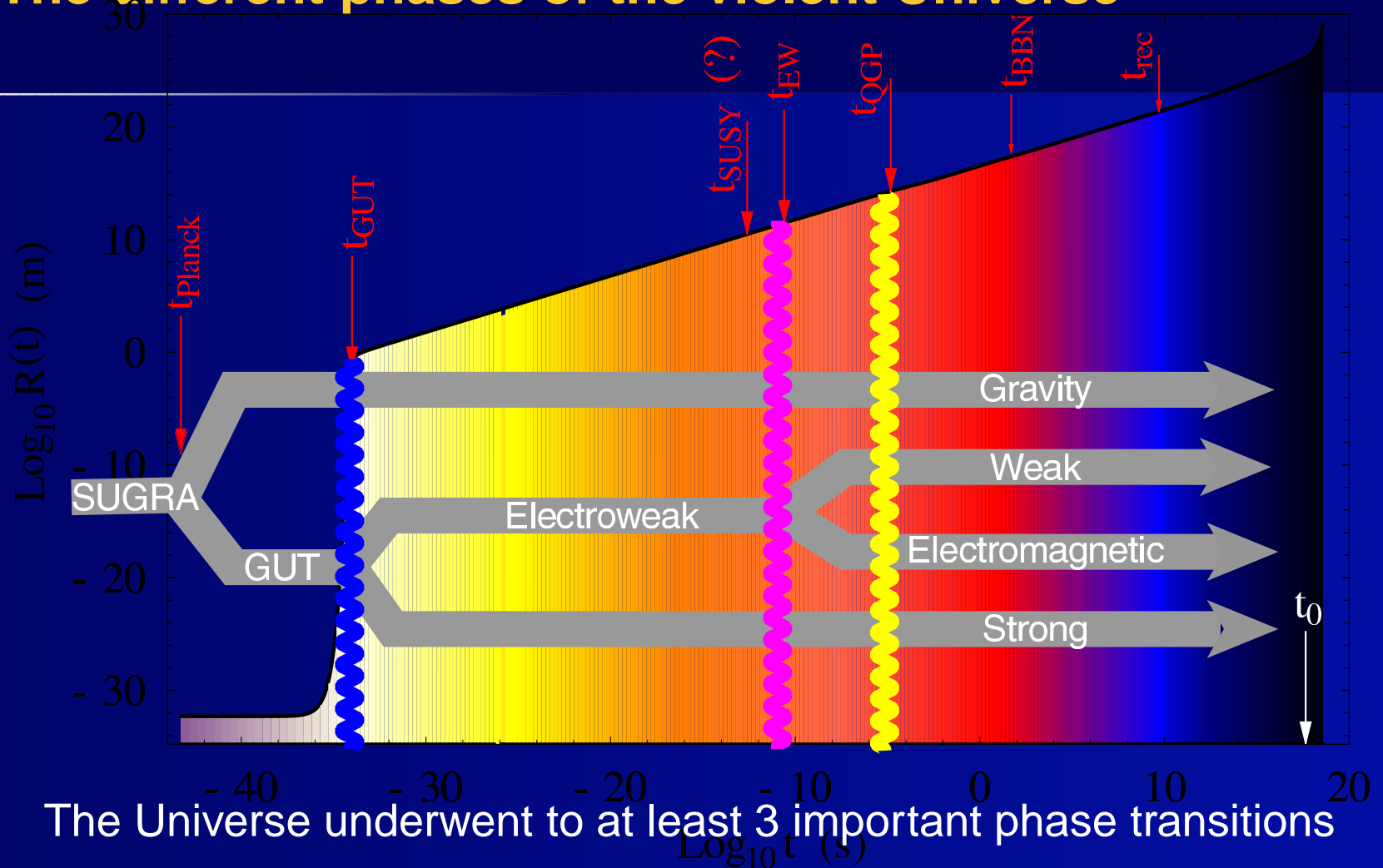
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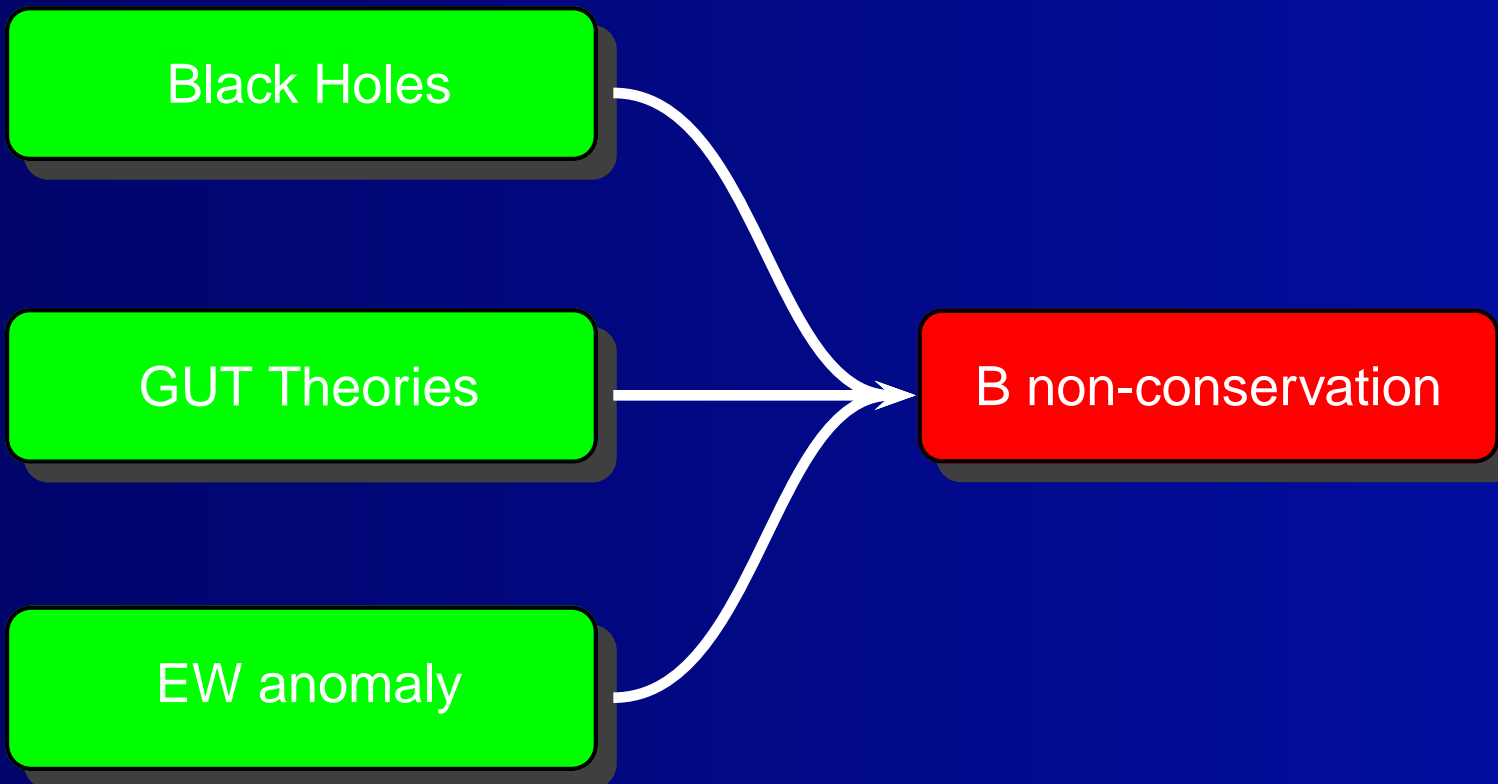
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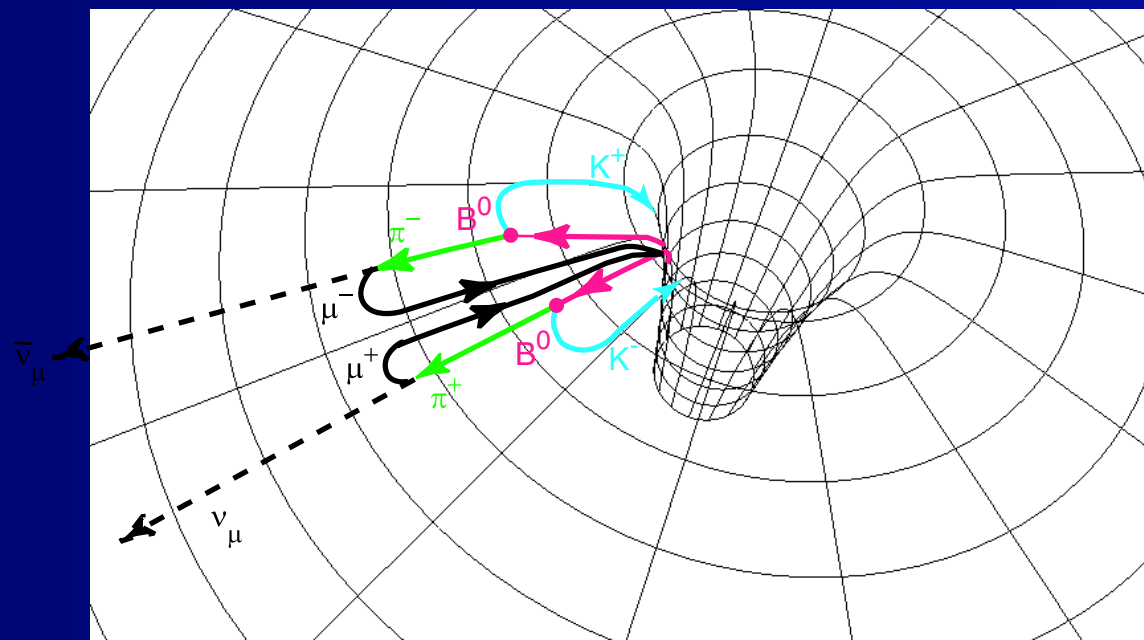
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As a “toy” model we consider the one sketched above

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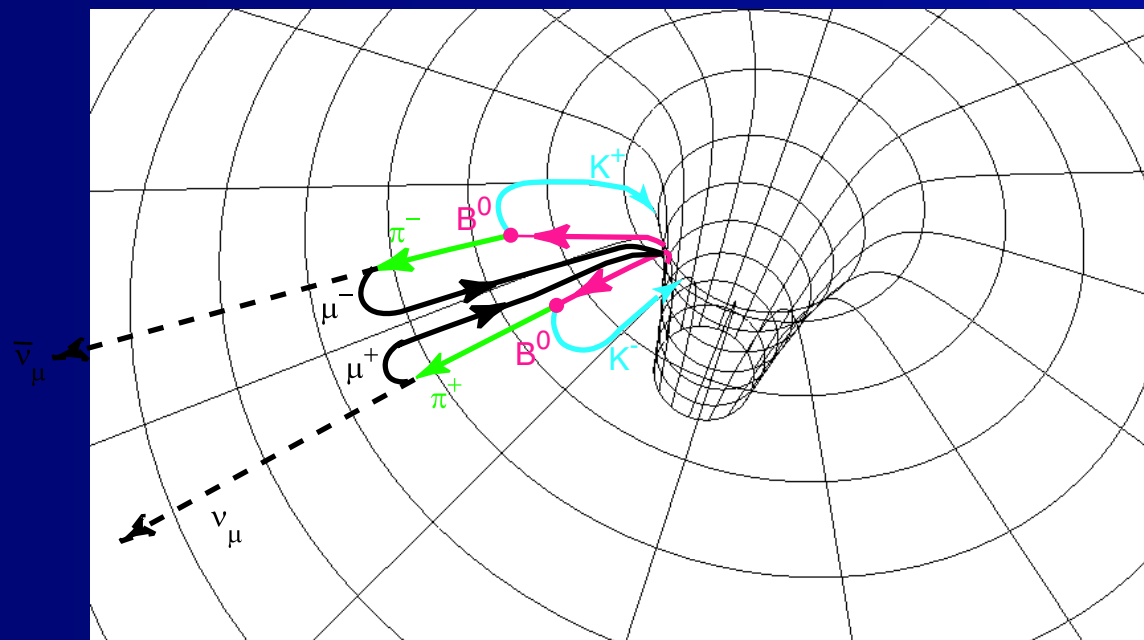
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➔ The important fact is that BH's are not effectively black, but gray due to the propagation of the produced particle in the strong gravitational field.

➔ Let us assume that a  $B^0$  is produced at the event horizon, decays respectively in the channels  $B^0 \rightarrow K^+ + \pi^-$  and  $B^0 \rightarrow K^- + \pi^+$ , the more massive kaons and muons will be preferentially recaptured, after the decay, by the BH, respect to the nearly massless neutrinos.

➔ We know from BABAR that CP-conservation is directly violated in this decay, giving

$$\frac{n_{K^+\pi^-} - n_{K^-\pi^+}}{n_{K^+\pi^-} + n_{K^-\pi^+}} = 0.133 \pm 0.030(\text{stat}) \pm 0.009(\text{syst})$$

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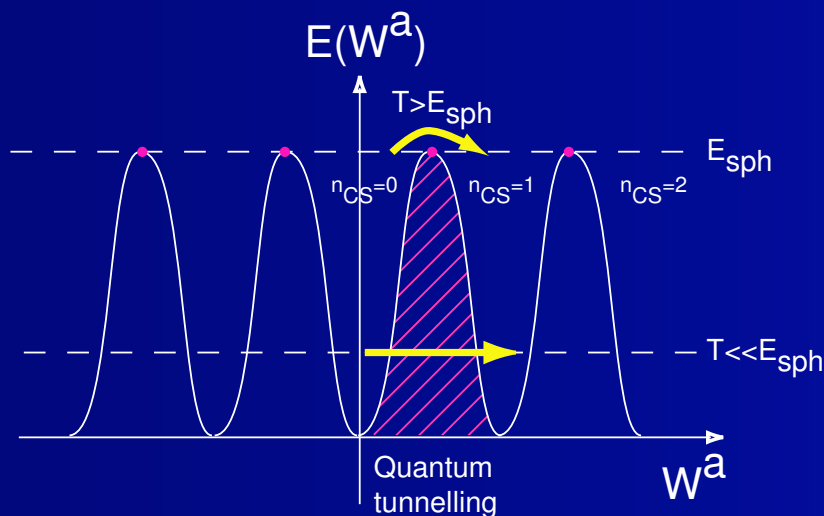
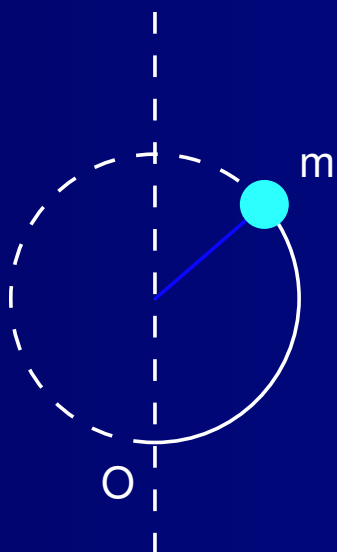
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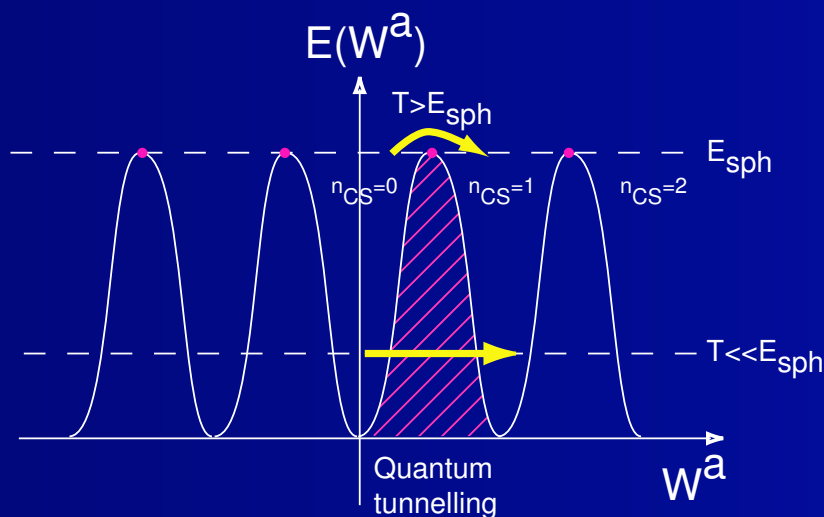
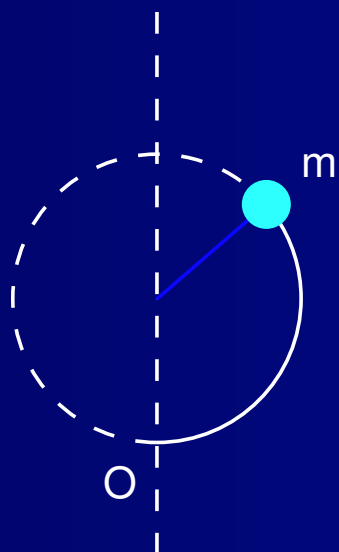
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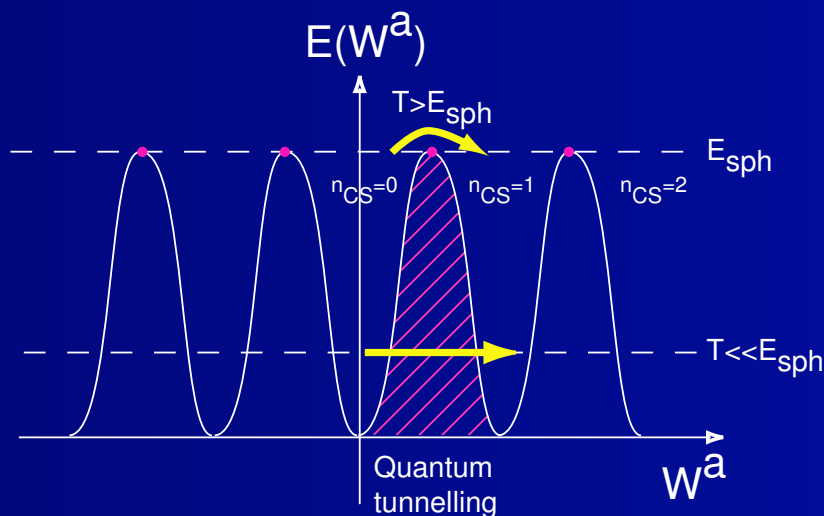
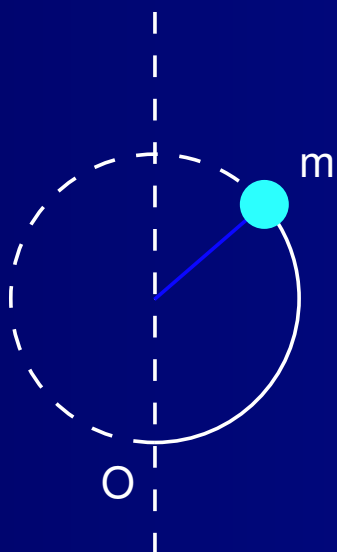
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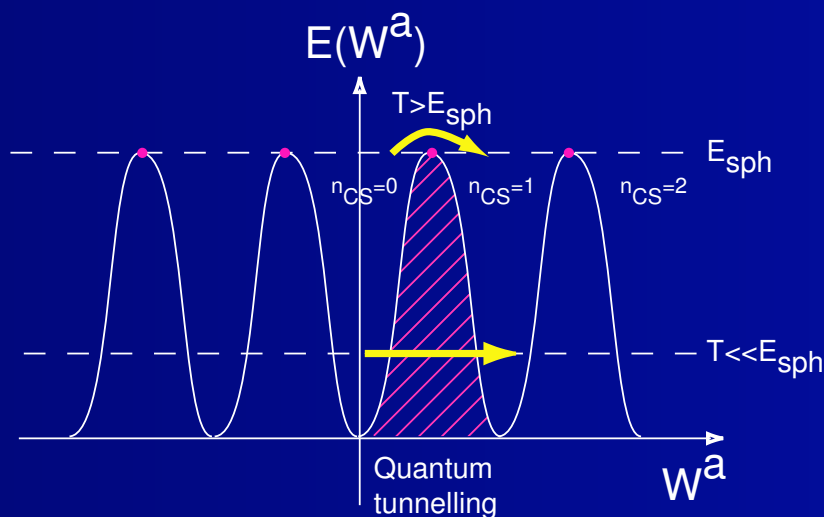
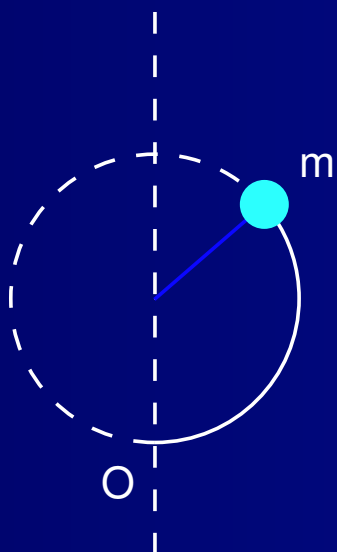
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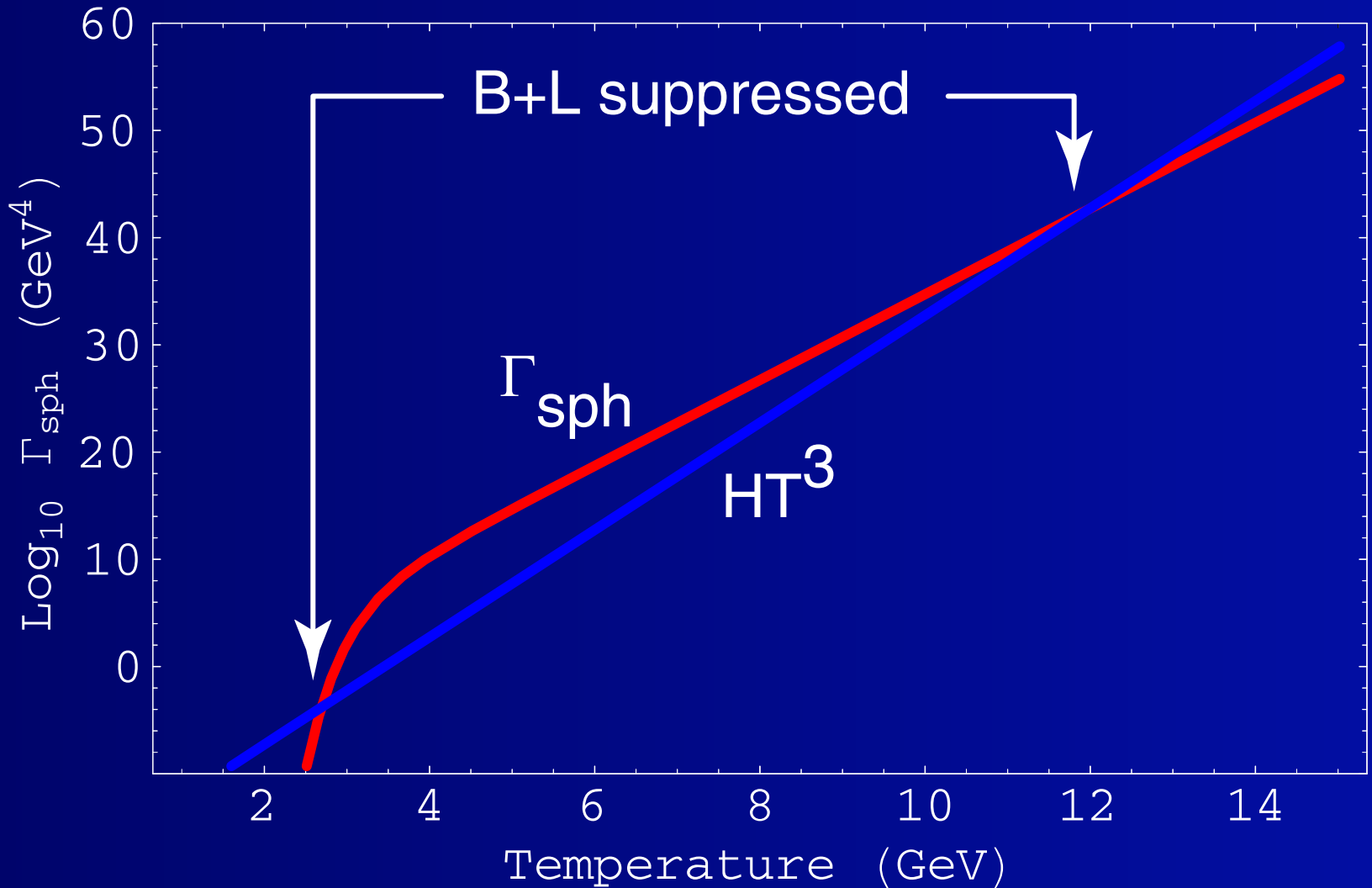
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➔ If  $(B - L)_{in} > 0$  the final baryon number will be

$$B \approx \left(\frac{1}{2}\right)(B - L)_{in} \quad \text{and} \quad L \approx -\left(\frac{1}{2}\right)(B - L)_{in}$$

Exact calculation change this coefficients to  $\frac{4}{11}$  and this to  $\frac{7}{11}$

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➔ After the discovery several interpretation has been given to the origin of the CP-Violations.

➔ There are 3 possible mechanism to violate CP:

1. In the Yukawa interaction between fermions and the Higgs boson (CKM mechanism);
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2. The “explicit” CP violation due interference of the coupling of the fermions to several Higgs bosons doublets;
3. The “spontaneous” CP violation induced by different  $\langle \phi_h \rangle_0$  of the Higgs Bosons.

➔ In the Standard Model  $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  only one Higgs doublet is allowed, therefore only CKM mechanism is possible

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➔ CP violations are a well established experimental fact since 1964 when in the Cronin-Fitch experiment (Christenson et al., 1964) observed 45 CP-violating decays  $K_L^0 \rightarrow 2\pi$  over about 20,000  $K_L^0$  produced.

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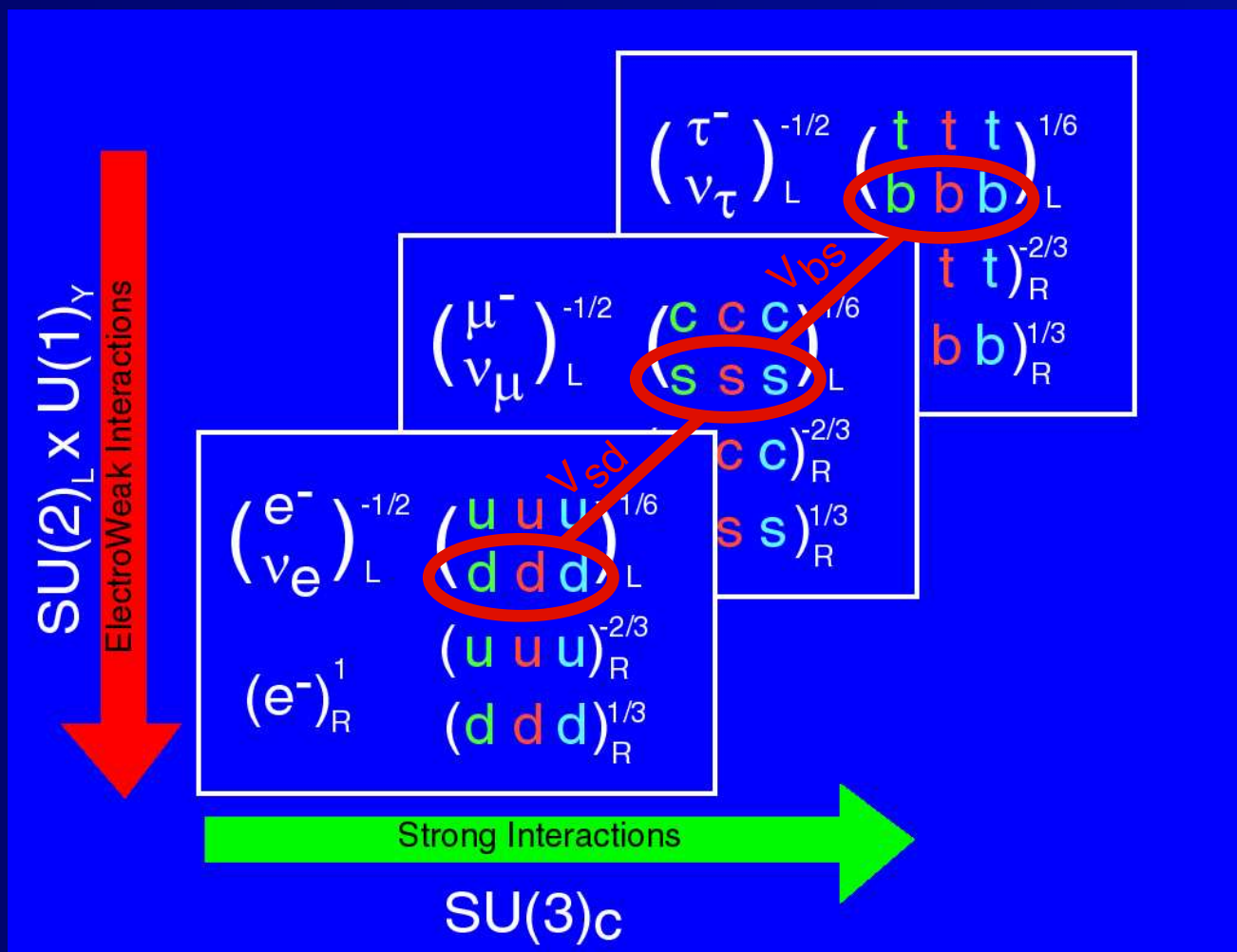
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Mass eigenstates  $\neq$  Weak Interactions (i.e. flavor's) eigenstates

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- ➔ All the observed CP violations up to now are consistent with the CKM paradigm, proposed by Kobayashy & Maskawa in 1973, originated from complex Yukawa coupling of the quarks with the Higgs field. For three generation of quarks there is only one phase  $\delta_{CP} = 70^\circ \pm 6^\circ$
- ➔ Masses of quarks are generated by the Yukawa coupling of the Higgs boson, respect to which  $D_j$  type quarks are also mixed. Therefore if any 2 of the  $D_j$  or of the  $U_i$  have the same mass the matrix  $V_{jk}$  must be real and there are no CP-violations.
- ➔ Suppression factor should be (Jarlskog 1985):

$$\epsilon_{CP} \propto (m_t^2 - m_u^2)(m_t^2 - m_c^2)(m_c^2 - m_u^2)(m_b^2 - m_s^2)(m_b^2 - m_d^2)(m_s^2 - m_d^2) J_{CP}$$

with  $J_{CP} \simeq 10^{-5}$  and  $T$  is the temperature.

- ➔ Common wisdom (with few exception see e.g. Berkooz, Nir & Volansky 2004) is that  $\epsilon_{CP} \rightarrow 0$  when  $T \rightarrow T_{EW}$ , because at  $T \geq T_{EW}$  symmetry is restored and all quark's masses  $\rightarrow 0$ .

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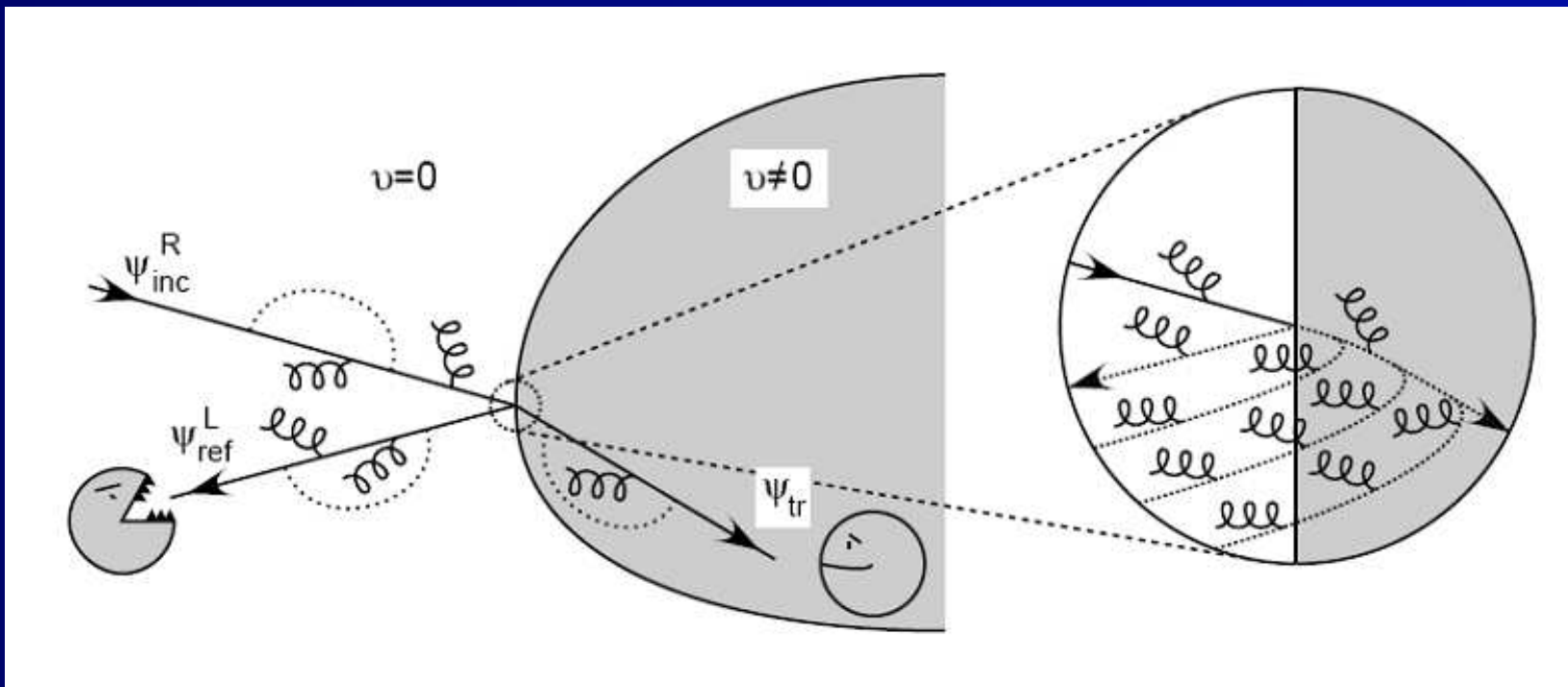
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- ➔ A first order phase transition proceeds via the formation of bubbles of the broken phase (like in supercooled gas phase).
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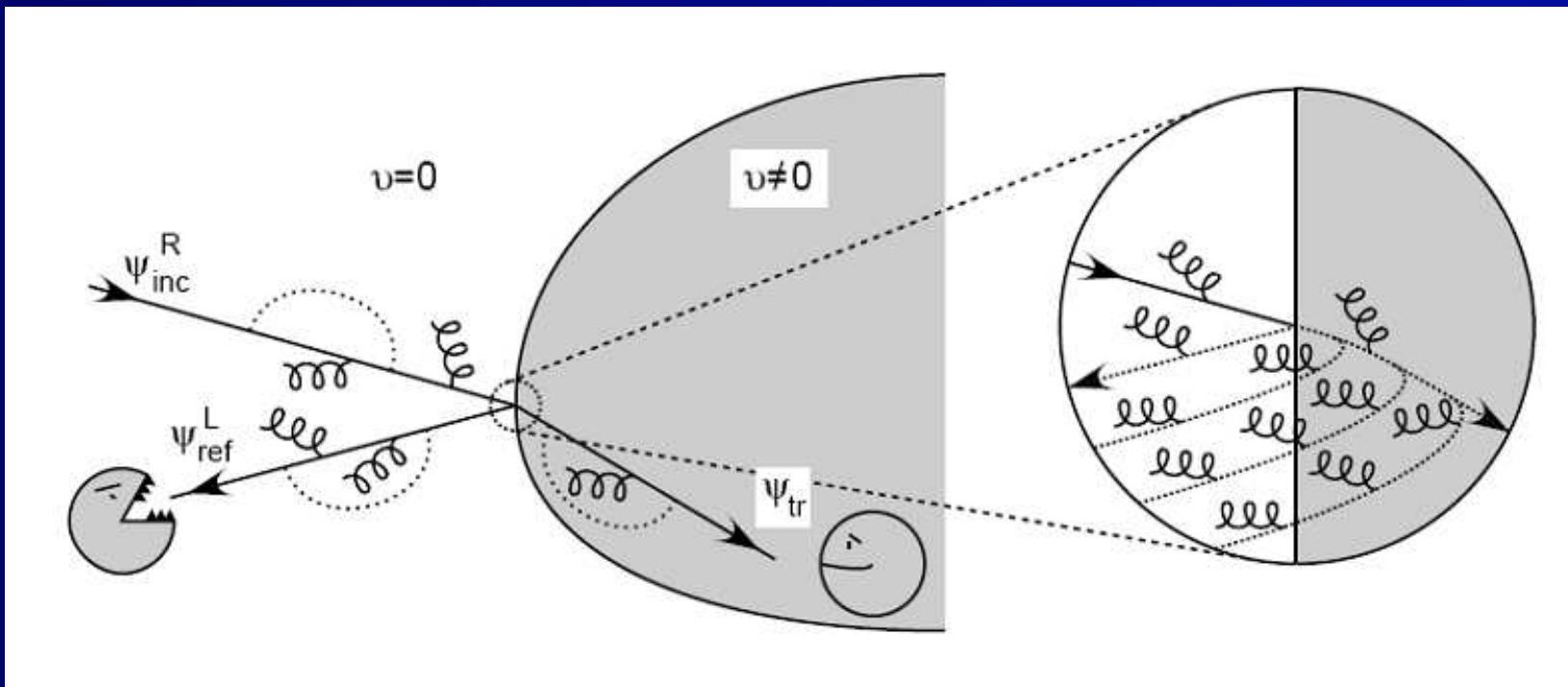
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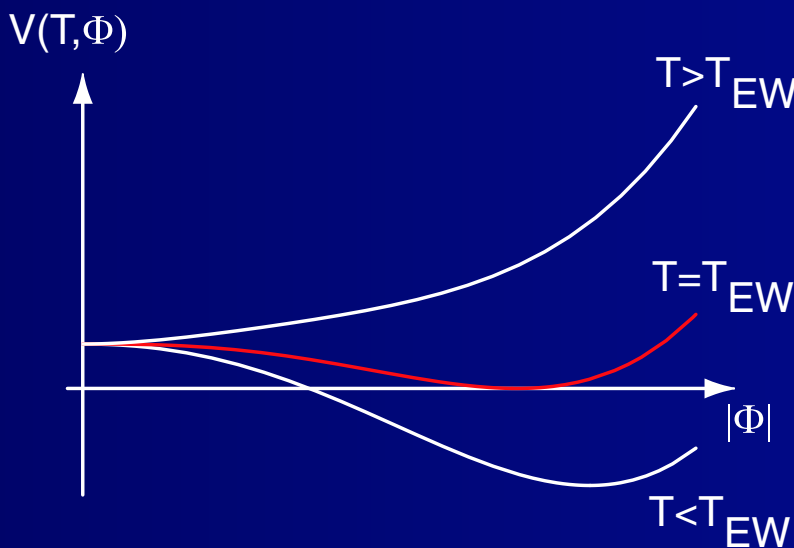
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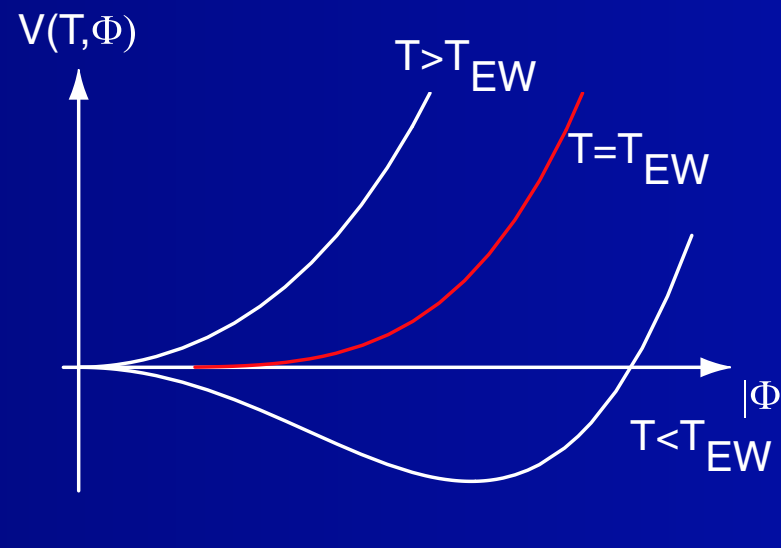
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$m_{\text{Higgs}} = 44 \text{ GeV}$



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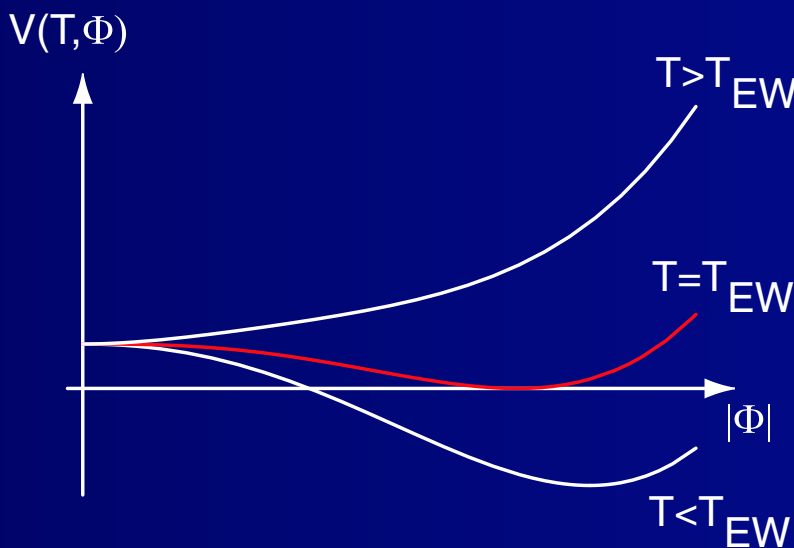
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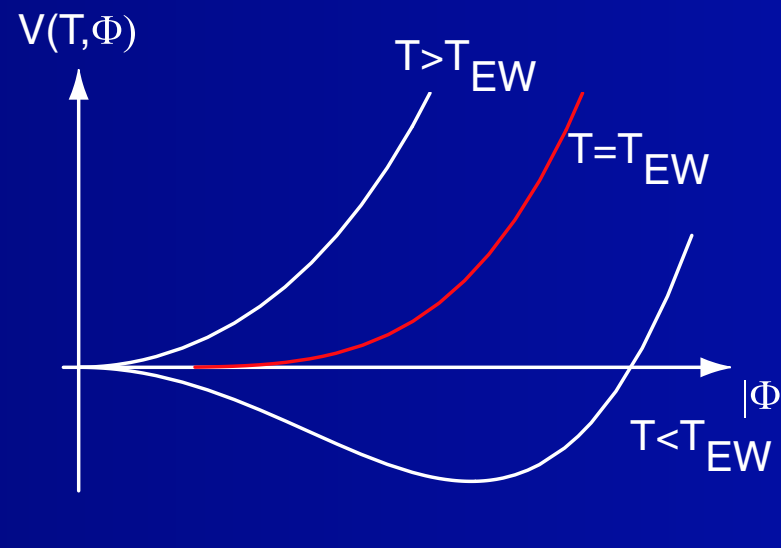
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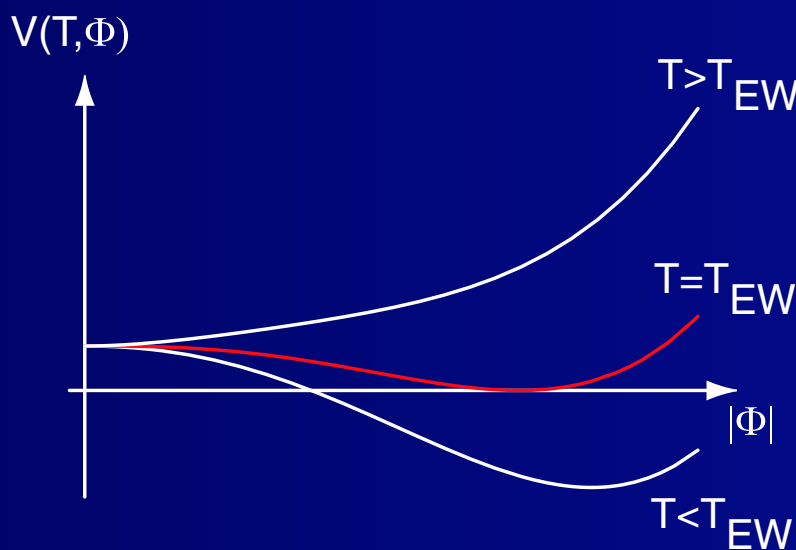
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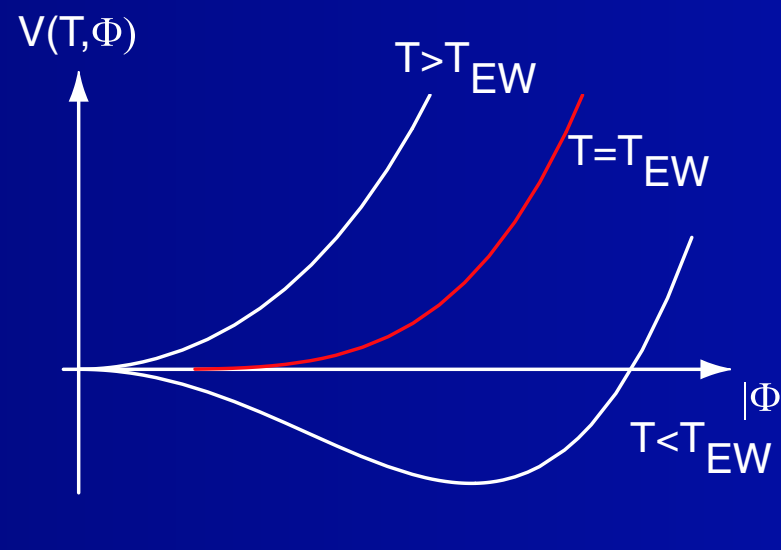
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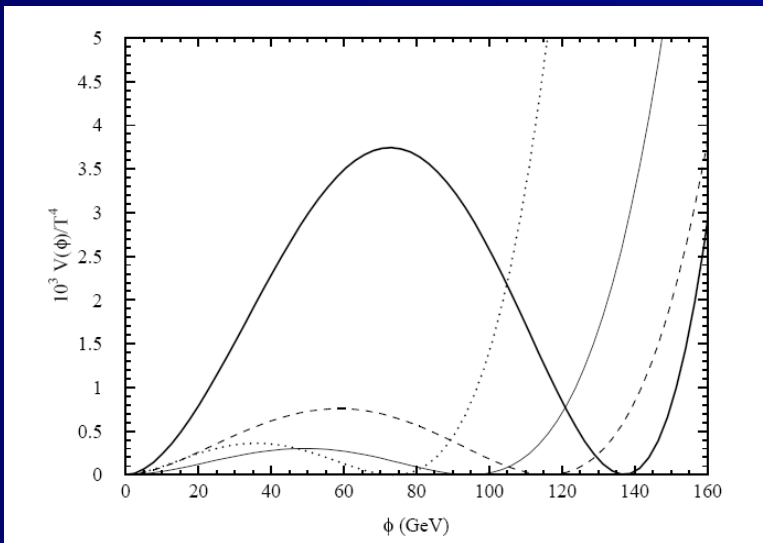
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Effective potential in the MSSM under different assumptions (Espinosa 1966)

- ➔ The expectations changes drastically in the Minimal Supersymmetric Standard Model (MSSM)
- ➔ MSSM contains multiple CP-violating complex phases. This is in marked contrast to the Standard Model which has only one phase in the CKM matrix.
- ➔ It has been shown that two-loop corrections to the finite temperature effective potential in the MSSM can have a dramatic effect on the strength of the electroweak phase transition, making it more strongly first order.



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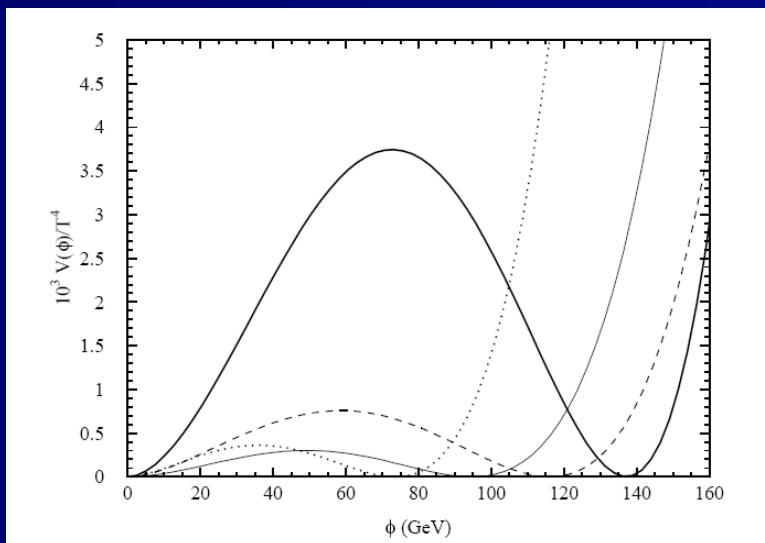
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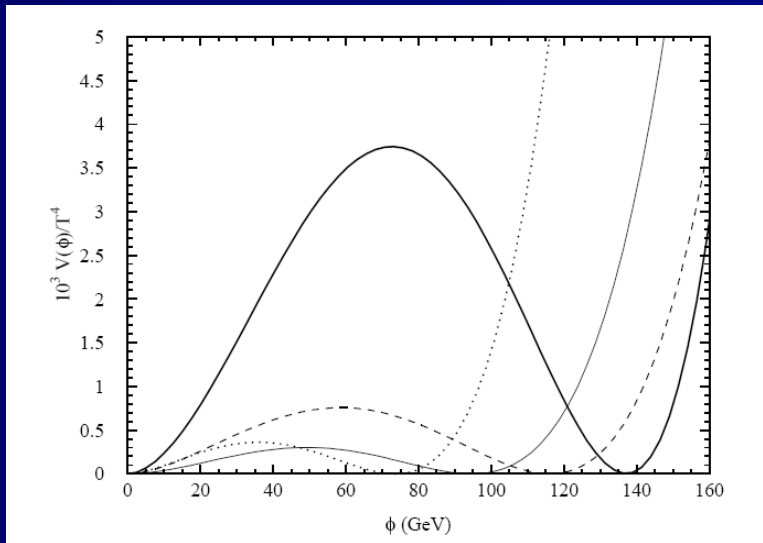
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➔ The asymmetry between matter and anti-matter, evident in our neighborhood Universe, has stimulated in the last three quarter of a century a very exciting interplay at the frontier between Astronomy, Cosmology and High Energy Particle physics.

➔ The study of CMBR, Gamma-ray Astronomy and the success of Big Bang Nucleosynthesis have given increasing evidence that presumably a positive baryon charge is a general feature of the entire visible Universe.

➔ Standard Model of Particle Physics, well supported by experiments up to now, that includes CP & B violations, suggests that the dynamical build up of the observed matter anti-matter asymmetry, could be the viable mechanism.

➔ However the minimal Standard Model alone fails to reproduce the observed asymmetry for two reasons:

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- vanishing CP phase at the EW transition;
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antimatter asymmetry

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EW phase transition in  
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