Some consideration on multidimensional likelyhood fits

Marco Mirazita

INFN – Laboratori Nazionali di Frascati

Physics environment

Extraction of the Λ polarization in SIDIS

with CLAS data

$$d\sigma = d\sigma_U \left(1 + S_{\mu} P_{\Lambda}^{\mu} \right)$$

decay proton angular distribution

$$\frac{dN}{d\cos\theta_p} \propto 1 + \alpha_w P_{\Lambda} \cos\theta_p$$

$$P_{\Lambda,z} = P_{\Lambda,z}^{I} + \lambda P_{\Lambda,z}^{T}$$

$$P_{\Lambda,z}^{I} = C_{z}^{s} \sin \varphi + C_{z}^{2s} \sin 2\varphi$$

$$P_{\Lambda,z}^{T} = D_{z}^{0} + D_{z}^{c} \cos \varphi$$

Pois
$$e^{y} \rightarrow e^{x}$$
 $p \rightarrow e^{x}$

electron scattering plane plane

The ϕ and θ_p dependencies are defined Other dependencies are encoded in the C and D parameters

hadron reaction

Unbinned likelyhood fit

$$L = \prod_{i=1}^{n} f^{norm}(X_i, \alpha) = \prod_{i=1}^{n} \frac{f(X_i, \alpha)}{N} = \frac{\prod_{i=1}^{n} f(X_i, \alpha)}{N^n}$$

Probability distribution function:

$$f^{\uparrow,\downarrow} = 1 + \alpha_{w} P_{\Lambda}^{\pm} \cos \theta_{p}$$

 N^n n=number of events X = set of kinematic variables $\alpha = set$ of parameters to be determined $f(X, \alpha) = probability$ distribution function N = normalization factor of the PDF $\uparrow \downarrow = polarization$ states

$$P_{\Lambda}^{\ \pm} = P_{\Lambda}^{I} \pm P_{\Lambda}^{T}$$

- all kinematics dependencies can be fitted at the same time (if known)
- higher statistic precision compared to binned analysis
- no needs to study binning effects

$$LogL = \sum_{i=i}^{n^{\uparrow}} Logf^{\uparrow}(X_i, \alpha) + \sum_{i=i}^{n^{\downarrow}} Logf^{\downarrow}(X_i, \alpha) - n^{\uparrow} LogN^{\uparrow} - n^{\downarrow} LogN^{\downarrow}$$

Likelyhood normalization

$$N = \int dX_i f(X_i, \alpha) \varepsilon(X_i)$$

- ensures that total probability is 1
- takes into account holes in the acceptance where the probability is zero

In some cases this term can be dropped, for example

$$\begin{cases} f(X_i, \alpha) \propto 1 + \alpha_0 \cos \varphi \\ \varepsilon(X_i) \text{ doesn't depend on } \varphi \end{cases}$$



The integral is a

Normalization using MC data

$$N \to \sum_{j=1}^{n_{MC}} f(X_i^j, \alpha)$$

- need good Monte Carlo, including polarizations
- variations of the fitted parameters through the fits should modify the angular distributions in the MC data Would the fit never converging at all?

Normalization using exp data

Prescription:

- sum over experimental events
- include ALL events in both N^{\uparrow} and N^{\downarrow}

$$N^{\uparrow,\downarrow} = \sum_{i=1}^{n^{\uparrow}+n^{\downarrow}} \left(1 + \alpha_{w} P_{\Lambda}^{\pm} \cos \theta_{p}\right)$$

Advantages:

- all the correct kinematics dependencies are there
- no MC simulation is necessary

Drawbacks:

- limited kinematics
- what are the limitations of the prescription?

Background subtraction

Good events selected within some kinematic cut contain background

$$L = \prod_{i=1}^{n} \frac{f(X_i, \alpha)}{N} = \left[\prod_{i=1}^{n_S} \frac{f(X_i, \alpha)}{N}\right] \left[\prod_{i=1}^{n_B} \frac{f(X_i, \alpha)}{N}\right]$$
SIGNAL BACKGROUND

To subtract the BKG contribution in the selected events, add n_B background events taken from outside the signal region with weight -1

$$L = \prod_{i=1}^{n} \frac{f(X_i, \alpha)}{N} = \left[\prod_{i=1}^{n_S} \frac{f(X_i, \alpha)}{N}\right] \left[\prod_{i=1}^{n_B} \frac{f(X_i, \alpha)}{N}\right] \left[\prod_{i=1}^{n_B} \frac{f(X_i, \alpha)}{N}\right]^{-1}$$

The total number of events for the fit is now n+n_B

peak events: w_i=+1

side band events: $w_i=-1$

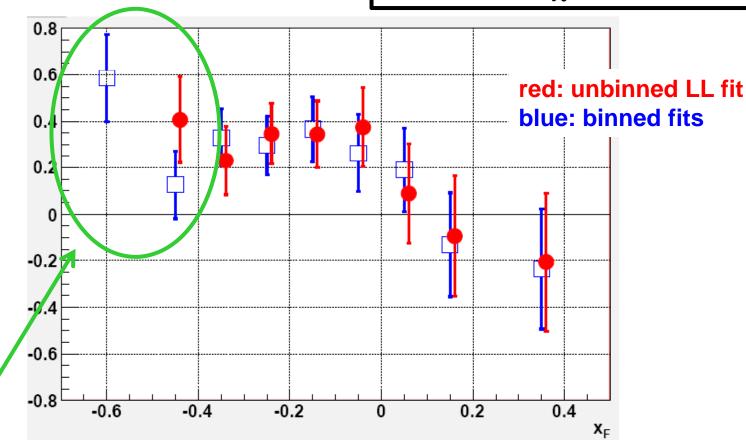


Extracting transferred polarization

Fit of the constant term $P_{\Lambda,z}^T=D_z^0$

$$P_{\Lambda,z}^T = D_z^0$$

results verified with binned analysis and χ^2 fit



where the statistics in too low, the LL fit gives better result doesn't converge

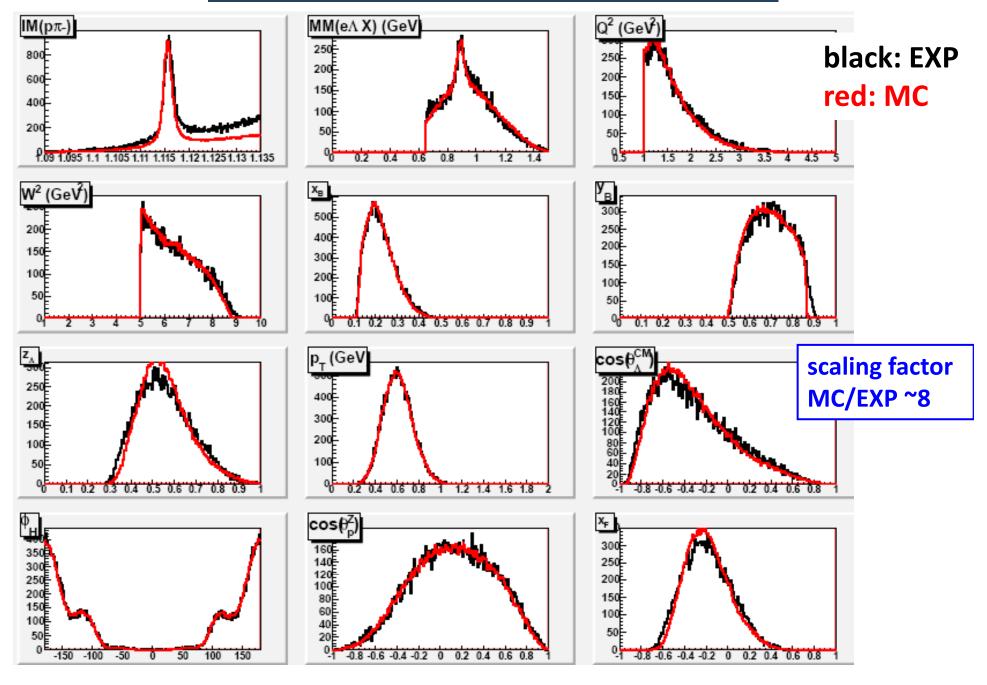
Monte Carlo simulation

clasDIS event generator based on pythia-JETSET

- adjustment needed for JETSET parameters to simulate Λ production in the target fragmentation
- standard parameters for exclusive ΛK^* events
- the two event samples are combined by fitting the experimental data

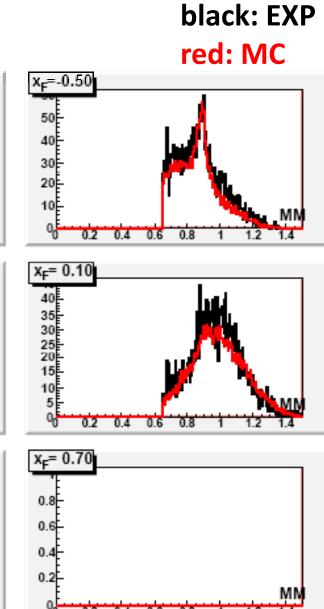
- full simulation of the CLAS detector
- analysis as for the experimental raw data

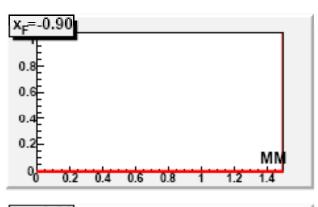
Comparison MC/EXP data

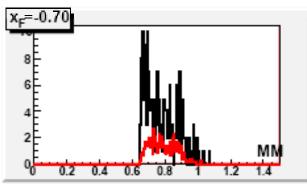


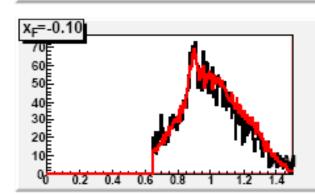
Comparison MC/EXP data

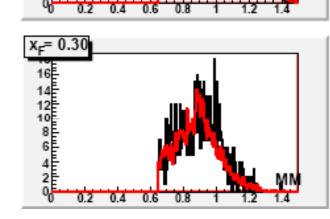
MM(e Λ X) distributions in x_F bins

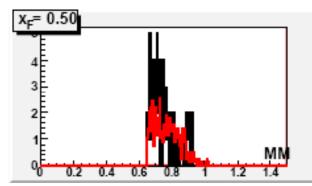










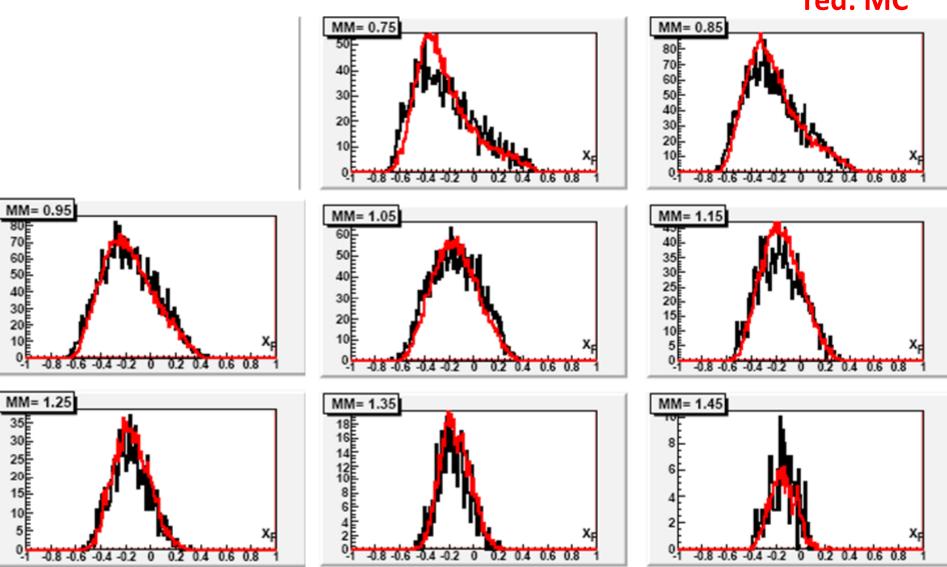


Comparison MC/EXP data

x_F distributions in MM(e Λ X) bins

black: EXP

red: MC



Polarization in the MC data

MC is unpolarized



- $\theta_{\rm p}$ generated distributions are flat (P^{l'}is zero)
 - no asymmetry between positive and negative beam helicity (P^T is zero)

For a given model of the polarization, i.e. a given (ϕ -dependent) value of P^I and P^T:

• probability of having an event at θ_n

$$\mathcal{P}(\cos \theta_p) = \frac{1 + \alpha_w P^I \cos \theta_p}{\left(1 + \alpha_w P^I\right)}$$

• probability of event produced by an electron with helicity λ

$$\mathcal{P}(\lambda) = \frac{1 + \alpha_w (P^I + \lambda P^T) \cos \theta_p}{(1 + \alpha_w P^I \cos \theta_p)}$$

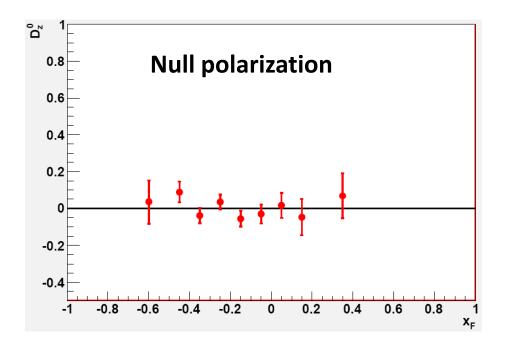
MLE of polarization in MC data

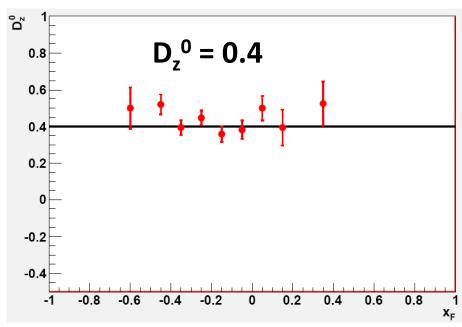
Check of the procedure using Monte Carlo data

- no false polarization from the acceptance
- put some polarization in the unpolarized MC
- verify that the likelyhood fit is able to recover the input polarization

Fit of the constant term
$$P_{\Lambda,z}^T=D_z^0$$

$$P_{\Lambda,z}^T = D_z^0$$





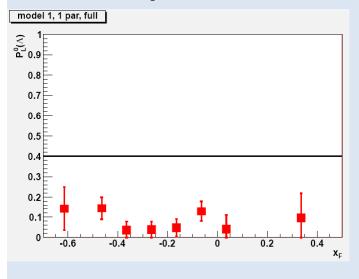
Acceptance effects

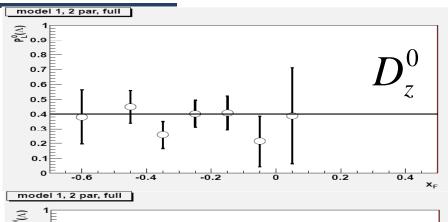
Monte Carlo data with φ dependent polarization

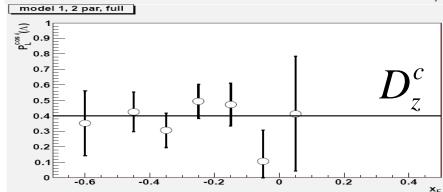
$$P_{\Lambda,z}^T = D_z^0 + D_z^c \cos \varphi$$

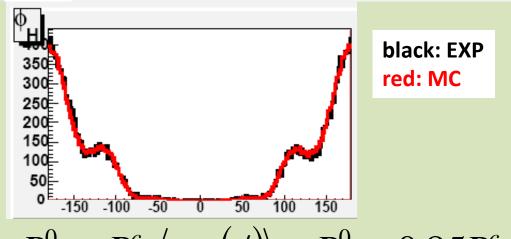
$$D_z^0 = D_z^c = 0.4$$

Likelyhood fit of the constant term only:









$$P_{\Lambda,z}^{T} = P_{\Lambda,z}^{0} + P_{\Lambda,z}^{c} \langle \cos(\phi) \rangle \approx P_{\Lambda,z}^{0} - 0.85 P_{\Lambda,z}^{c}$$

Extracting the induced polarization

consider the induced polarization only

$$P_{\Lambda,z} = P_{\Lambda,z}^{I} = C_{z}^{s} \sin \varphi + C_{z}^{2s} \sin 2\varphi$$

pdf function

$$f^{\uparrow,\downarrow} = 1 + \alpha_w P_{\Lambda,z}^{I} \cos \theta_p$$
No dependence on the beam helicity
$$f^{\uparrow} = f^{\downarrow} \qquad N^{\uparrow} = N^{\downarrow}$$

$$f^{\uparrow} = f^{\downarrow}$$
 $N^{\uparrow} = N$

likelyhood function with normalization prescription

$$LogL = \sum_{i=1}^{n} Logf(X_{i}, \alpha) - nLog\sum_{i=1}^{n} f(X_{i}, \alpha)$$

maximization of the likelyhood

$$\frac{\partial LogL}{\partial P_{\Lambda,z}^{I}} = 0$$

$$P_{\Lambda,z}^{I} = 0$$

verified with MC data with different values of the input polarization

Conclusion

- Likelyhood unbinned fits are a powerful tool for multidimensional analysis
- Monte Carlo studies are important to check the reliability of the analysis procedures
- No easy estimate of goodness-of-fit
- Normalization of the likelyhood function is a crucial point and must be studied carefully
- Can MC data be used to normalize the likelyhood function if the MC is not perfectly reproducing the experimental data?