Why Monte Carlo generators? Theory point of view

A. Bacchetta "TMD Monte Carlo" Frascati, 7 Nov 2011



Two types of generators

• Full event generators (Pythia, Lepto...) All final-state particles are generated • Full event generators (Pythia, Lepto...) All final-state particles are generated • Full event generators (Pythia, Lepto...) All final-state particles are generated

 Generators for single-particle (or two-particle) inclusive DIS (gmc_trans, TMDgen, ResBos...)
Only one or two final-state particles are generated



Single-particle DIS generator



Something about full event generators















Example of fragmentation model



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What are they used for?

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- Systematic studies

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- Systematic studies
- Search for new physics
- Access to quantities that are not directly measurable (i.e., W-boson mass)

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- Parameters well tuned
- Excellent coding, based on years of experience

Full description of the final state



simulation for CLAS12

Comparison with SIDIS generators



Comparison with SIDIS generators



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Example of results



Figure 6.1: Comparison between the distributions of selected DIS and SIDIS kinematic variables obtained from real events and from events generated by PYTHIA.

Thesis of L. Pappalardo

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- Based on semi-classical picture (difficult to include quantum interference)
- Difficult to modify
- Semi-inclusive DIS is not their main focus
- Computationally intensive

Possible developments

 Inclusion of spin into the microscopic fragmentation mechanism Artru, arXiv:1001.1061; Bianconi, arXiv:1109.0688, Kotzinian, hep-ph/0510359 Inclusion of spin into the microscopic fragmentation mechanism Artru, arXiv:1001.1061; Bianconi, arXiv:1109.0688, Kotzinian, hep-ph/0510359 Inclusion of spin into the microscopic fragmentation mechanism Artru, arXiv:1001.1061; Bianconi, arXiv:1109.0688, Kotzinian, hep-ph/0510359

 Artificial modulation of the final cross section based on polarized cross-section (reweighting). Often used by experimental collaborations. No publication?

Something about SIDIS generators











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- Very close to theoretical formulas and theoretical parametrizations
- Can be in principle extended to higher orders

Inclusive DIS

 $\ell(l) + N(P) \to \ell(l') + X$



Structure functions

$$\begin{aligned} \frac{d\sigma}{dx_B \, dy \, d\phi_S} &= \frac{2\alpha^2}{x_B y Q^2} \left\{ \left(1 - y + \frac{y^2}{2}\right) F_{UU,T} + (1 - y) F_{UU,L} + S_L \lambda_e \, y \left(1 - \frac{y}{2}\right) F_{LL} \right. \\ &+ \left| \boldsymbol{S}_T \right| \lambda_e \, y \sqrt{1 - y} \, \cos \phi_S \, F_{LT}^{\cos \phi_s} \right\} \end{aligned}$$

see, e.g., A.B., Diehl, Goeke, Metz, Mulders, Schlegel, JHEP093 (07)

Structure functions

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Results for inclusive DIS

Results for inclusive DIS

$$F_{UU,T} = x_B \sum_{a} e_a^2 f_1^a(x_B)$$
$$F_{UU,L} = 0$$
$$F_{LL} = x_B \sum_{a} e_a^2 g_1^a(x_B)$$
$$F_{LT}^{\cos \phi_S} = -\gamma x_B \sum_{a} e_a^2 g_T^a(x_B)$$

Semi-inclusive DIS

 $\ell(l) + N(P) \to \ell(l') + h(P_h) + X,$



A.B., D'Alesio, Diehl, Miller, PRD70 (04)

Structure functions

$$\begin{split} \frac{d\sigma}{dx\,dy\,d\phi_S\,dz\,d\phi_h\,dP_{h\perp}^2} \\ &= \frac{\alpha^2}{x\,y\,Q^2}\,\frac{y^2}{2\,(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon\,F_{UU,L} + \sqrt{2\,\varepsilon(1+\varepsilon)}\,\cos\phi_h\,F_{UU}^{\cos\phi_h} + \varepsilon\,\cos(2\phi_h)\,F_{UU}^{\cos\,2\phi_h} \\ &+ \lambda_e\,\sqrt{2\,\varepsilon(1-\varepsilon)}\,\sin\phi_h\,F_{LU}^{\sin\phi_h} + S_L\left[\sqrt{2\,\varepsilon(1+\varepsilon)}\,\sin\phi_h\,F_{UL}^{\sin\phi_h} + \varepsilon\,\sin(2\phi_h)\,F_{UL}^{\sin\,2\phi_h}\right] \\ &+ S_L\,\lambda_e\left[\sqrt{1-\varepsilon^2}\,F_{LL} + \sqrt{2\,\varepsilon(1-\varepsilon)}\,\cos\phi_h\,F_{LL}^{\cos\phi_h}\right] \\ &+ S_T\left[\sin(\phi_h - \phi_S)\left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon\,F_{UT,L}^{\sin(\phi_h - \phi_S)}\right) + \varepsilon\,\sin(\phi_h + \phi_S)\,F_{UT}^{\sin(\phi_h + \phi_S)} \\ &+ \varepsilon\,\sin(3\phi_h - \phi_S)\,F_{UT}^{\sin(3\phi_h - \phi_S)} + \sqrt{2\,\varepsilon(1+\varepsilon)}\,\sin\phi_S\,F_{UT}^{\sin\phi_S} \\ &+ \sqrt{2\,\varepsilon(1+\varepsilon)}\,\sin(2\phi_h - \phi_S)\,F_{UT}^{\sin(2\phi_h - \phi_S)}\right] + S_T\lambda_e\left[\sqrt{1-\varepsilon^2}\,\cos(\phi_h - \phi_S)\,F_{LT}^{\cos(\phi_h - \phi_S)} \\ &+ \sqrt{2\,\varepsilon(1-\varepsilon)}\,\cos\phi_S\,F_{LT}^{\cos\phi_S} + \sqrt{2\,\varepsilon(1-\varepsilon)}\,\cos(2\phi_h - \phi_S)\,F_{LT}^{\cos(2\phi_h - \phi_S)}\right] \right\} \end{split}$$

see e.g. AB, Diehl, Goeke, Metz, Mulders, Schlegel, JHEP093 (07)

Structure functions

$$\begin{aligned} \frac{d\sigma}{dx\,dy\,d\phi_{S}\,dz\,d\phi_{h}\,dP_{h\perp}^{2}} & F_{UU,T}(x,z,P_{h\perp}^{2},Q^{2}) \\ = \frac{\alpha^{2}}{x\,y\,Q^{2}}\frac{y^{2}}{2(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)}\cos\phi_{h} F_{UU}^{\cos\phi_{h}} + \varepsilon\cos(2\phi_{h}) F_{UU}^{\cos2\phi_{h}} \\ + \lambda_{e}\,\sqrt{2\varepsilon(1-\varepsilon)}\sin\phi_{h} F_{LU}^{\sin\phi_{h}} + S_{L}\left[\sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{h} F_{UL}^{\sin\phi_{h}} + \varepsilon\sin(2\phi_{h}) F_{UL}^{\sin2\phi_{h}}\right] \\ + S_{L}\,\lambda_{e}\left[\sqrt{1-\varepsilon^{2}} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_{h} F_{LL}^{\cos\phi_{h}}\right] \\ + S_{T}\left[\sin(\phi_{h} - \phi_{S})\left(F_{UT,T}^{\sin(\phi_{h} - \phi_{S})} + \varepsilon F_{UT,L}^{\sin(\phi_{h} - \phi_{S})}\right) + \varepsilon\sin(\phi_{h} + \phi_{S}) F_{UT}^{\sin(\phi_{h} + \phi_{S})} \\ + \varepsilon\sin(3\phi_{h} - \phi_{S}) F_{UT}^{\sin(3\phi_{h} - \phi_{S})} + \sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_{S} F_{UT}^{\sin\phi_{S}} \\ + \sqrt{2\varepsilon(1+\varepsilon)}\sin(2\phi_{h} - \phi_{S}) F_{UT}^{\sin(2\phi_{h} - \phi_{S})}\right] + S_{T}\lambda_{e}\left[\sqrt{1-\varepsilon^{2}}\cos(\phi_{h} - \phi_{S}) F_{LT}^{\cos(\phi_{h} - \phi_{S})} \\ + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_{S} F_{LT}^{\cos\phi_{S}} + \sqrt{2\varepsilon(1-\varepsilon)}\cos(2\phi_{h} - \phi_{S}) F_{LT}^{\cos(2\phi_{h} - \phi_{S})}\right]\right\} \end{aligned}$$

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Unpolarized sector

$$\begin{split} F_{UU,T} &= \mathcal{C} \left[f_1 D_1 \right], \\ F_{UU,L} &= \mathcal{O} \left(\frac{M^2}{Q^2}, \frac{q_T^2}{Q^2} \right), \\ F_{UU}^{\cos \phi_h} &= \frac{2M}{Q} \mathcal{C} \left[-\frac{\hat{\boldsymbol{h}} \cdot \boldsymbol{k}_T}{M_h} \left(xh H_1^{\perp} + \frac{M_h}{M} f_1 \frac{\tilde{D}^{\perp}}{z} \right) - \frac{\hat{\boldsymbol{h}} \cdot \boldsymbol{p}_T}{M} \left(xf^{\perp} D_1 + \frac{M_h}{M} h_1^{\perp} \frac{\tilde{H}}{z} \right) \right], \\ F_{UU}^{\cos 2\phi_h} &= \mathcal{C} \left[-\frac{2 \left(\hat{\boldsymbol{h}} \cdot \boldsymbol{k}_T \right) \left(\hat{\boldsymbol{h}} \cdot \boldsymbol{p}_T \right) - \boldsymbol{k}_T \cdot \boldsymbol{p}_T}{MM_h} h_1^{\perp} H_1^{\perp} \right], \end{split}$$

$$\mathcal{C}[wfD] = \sum_{a} x e_a^2 \int d^2 \boldsymbol{p}_T \, d^2 \boldsymbol{k}_T \, \delta^{(2)} \left(\boldsymbol{p}_T - \boldsymbol{k}_T - \boldsymbol{P}_{h\perp}/z \right) w(\boldsymbol{p}_T, \boldsymbol{k}_T) \, f^a(x, p_T^2) \, D^a(z, k_T^2),$$
List of structure functions

observable twist "SIDIS FT" $F_{UU,T}$ 2"SIDIS F_L" $F_{UU,L}$ 4 $F_{UU}^{\cos\phi_h}$ "Cahn" 3 $F_{UU}^{\cos 2\phi_h}$ "Boer-Mulders" 2 $F_{LU}^{\sin\phi_h}$ 3 $F_{UL}^{\sin\phi_h}$ 3 $F_{UL}^{\sin 2\phi_h}$ "Kotzinian-Mulders" $\mathbf{2}$ "SIDIS g₁" F_{LL} $\mathbf{2}$ "Polarized Cahn" $F_{LL}^{\cos\phi_h}$ 3 $F_{UT,T}^{\sin(\phi_h - \phi_S)}$ "Sivers" $\mathbf{2}$ $F_{UT,L}^{\sin(\phi_h - \phi_S)}$ 4 $F_{UT}^{\sin(\phi_h + \phi_S)}$ "Collins" $\mathbf{2}$ $F_{UT}^{\sin(3\phi_h - \phi_S)}$ "Pretzelosity" 2 $F_{UT}^{\sin\phi_S}$ 3 $F_{UT}^{\sin(2\phi_h - \phi_S)}$ 3 $F_{LT}^{\cos(\phi_h - \phi_S)}$ "Worm gear" 2"SIDIS g₂" $F_{LT}^{\cos\phi_S}$ 3 $F_{LT}^{\cos(2\phi_h - \phi_S)}$ 3

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Unpolarized structure function

 $F_{UU,T} = \mathcal{C}[f_1 D_1]$

$$\mathcal{C}[wfD] = \sum_{a} x e_{a}^{2} \int d^{2} \boldsymbol{p}_{T} d^{2} \boldsymbol{k}_{T} \, \delta^{(2)} (\boldsymbol{p}_{T} - \boldsymbol{k}_{T} - \boldsymbol{P}_{h\perp}/z) \, w(\boldsymbol{p}_{T}, \boldsymbol{k}_{T}) \, f^{a}(x, p_{T}^{2}) \, D^{a}(z, k_{T}^{2}),$$

Gaussian ansatz

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$$\begin{split} \mathcal{C}\big[wfD\big] &= \sum_{a} x e_{a}^{2} \int d^{2} \boldsymbol{p}_{T} \, d^{2} \boldsymbol{k}_{T} \, \delta^{(2)} \left(\boldsymbol{p}_{T} - \boldsymbol{k}_{T} - \boldsymbol{P}_{h\perp} / z\right) \, w(\boldsymbol{p}_{T}, \boldsymbol{k}_{T}) \, f^{a}(x, p_{T}^{2}) \, D^{a}(z, k_{T}^{2}), \\ f_{1}^{a}(x, p_{T}^{2}) &= \frac{f_{1}^{a}(x)}{\pi \langle p_{T}^{2} \rangle} e^{-\boldsymbol{p}_{T}^{2} / \langle p_{T}^{2} \rangle}, \qquad D_{1}^{a}(z, k_{T}^{2}) = \frac{D_{1}^{a}(z)}{\pi \langle K_{T}^{2} \rangle} e^{-z^{2} \boldsymbol{k}_{T}^{2} / \langle K_{T}^{2} \rangle} \\ \mathcal{C}[f_{1}D_{1}] &= \sum_{a} x e_{a}^{2} \frac{f_{1}(x) D_{1}(z)}{\pi (z^{2} \rho_{a}^{2} + \sigma_{a}^{2})} \, e^{-\boldsymbol{P}_{h\perp}^{2} / (z^{2} \rho_{a}^{2} + \sigma_{a}^{2})} \\ f_{1}^{a}(x, p_{T}^{2}) &= \frac{f_{1}^{a}(x)}{\pi \langle p_{T}^{2}(x) \rangle^{a}} e^{-\boldsymbol{p}_{T}^{2} / \langle p_{T}^{2}(x) \rangle^{a}}, \quad D_{1}^{a}(z, k_{T}^{2}) = \frac{D_{1}^{a}(z)}{\pi \langle K_{T}^{2}(z) \rangle^{a}} e^{-z^{2} \boldsymbol{k}_{T}^{2} / \langle K_{T}^{2}(z) \rangle^{a}} \end{split}$$

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- Implements several leading-twist terms of the cross section
- First attempt at tuning the parameters of the unpolarized TMDs
- Careful implementation of positivity bounds

Comparison with data

· Cohnoll



 z^2

unpublished!

 z^2

TMDgen



Extension of gmc_trans done by Steve Gliske



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- Includes non-Gaussian distributions



- Extension of gmc_trans done by Steve Gliske
- Includes non-Gaussian distributions
- Includes two-hadron inclusive DIS



- Extension of gmc_trans done by Steve Gliske
- Includes non-Gaussian distributions
- Includes two-hadron inclusive DIS
- Written in C++



Figure 3.1: Comparison of 1D kinematic distributions from TMDGen and Pythia, in 4π , for $\pi^+\pi^0$ dihadrons. Listing the rows from top to bottom, and within each row from left to right, the panels are respectively the $x, y, z, P_{h\perp}$, and M_h distributions. TMDGen data is designated with blue circles, and Pythia data designated with red open squares.

Thesis of S. Gliske

Lacks all subleading twist

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- Lacks all subleading twist
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- No QED radiative corrections
- No TMD evolution

$$f_1(x, k_T; Q) = \frac{1}{2\pi} \int d^2 b_T e^{-ik_T \cdot b_T} \left[C \otimes f_1\left(\hat{x}; \frac{2e^{-\gamma_e}}{b_T}\right) \right] e^{-S'(b_T, Q)} e^{-S'_{\rm NP}(x, b_T, Q, \alpha_i)}$$

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collinear PDF







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Up Quark TMD PDF, x = .09



T. Rogers, M. Aybat, arXiv:1101.5057 Landry, Brock, Nadolsky, Yuan, PRD67 (03) P. Schweitzer, T. Teckentrup, A. Metz, PRD81(10)

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- http://hep.pa.msu.edu/resum/index.html#SIDIS

Fourier-transformed TMDs

 Since the TMD evolution formalism is done in bT space, it may be useful to implement the Fourier-transformed formulas in the MC generator Since the TMD evolution formalism is done in bT space, it may be useful to implement the Fourier-transformed formulas in the MC generator Since the TMD evolution formalism is done in bT space, it may be useful to implement the Fourier-transformed formulas in the MC generator

• This may also useful to study Bessel-weighted extraction methods Boer, Gamberg, Musch, Prokudin, arXiv:1107.5294

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- Every collaboration and even every analysis group uses its own different solution.
- Not enough attention is devoted to publishing the ideas and share them.
- I would personally focus first on SIDIS generators, although I think the effort of modifying full event generators is extremely interesting