Beam-beam effect and degree of freedom

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- Effect of crossing angle
- Mysterious of the beam-beam limit, Non-Gaussian distribution.
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Incoherent effects

- Weak-strong model is available, if we know a self consistent distribution of beams.
- Emittance growth due to nonlinear force interacting with a fixed charged distribution,
 - 1. Gaussian, simple but not almighty,
 - 2. Non-Gaussian, determined by strong-strong simulation, discussed later.

$$\varepsilon = \varepsilon_0 (1 + D_C t)$$
 $\sigma^2 = \sigma_0^2 (1 + D_C t)$

Synchrotron radiation, excitation and damping

• Damping, τ : damping time

 $\varepsilon = \varepsilon_0 (1 - 2t/\tau) \qquad \sigma^2 = \sigma_0^2 (1 - 2t/\tau)$

- Equilibrium emittance is determined by the balance of radiation excitation and diffusion. The equilibrium emittance is ϵ_0 for D_y=2/ τ
- KEKB $D_{\gamma}=5x10^{-4}$ /turn DAFNE $D_{\gamma}=1.8x10^{-5}$ /turn
- Correlation between the diffusions due to radiation and nonlinearity. $\varepsilon(t=\infty) = (D_n + D_c)$

Non-correlation Correlation

$$\frac{\varepsilon(t=\infty)}{\varepsilon_0} = \frac{(D_{\gamma} + D_C)}{D_{\gamma}}$$
$$\frac{\varepsilon(t=\infty)}{\varepsilon_0} = \frac{D_{\gamma} + D_C + D_C}{D_{\gamma}}$$

Symplectic diffusion due to the beambeam interaction

- Calculate diffusion rate for symplectic system.
- Use Gaussian model.

$$U(x, y; \sigma_{x}(z), \sigma_{y}(z)) = -\frac{N(z)r_{e}}{\gamma} \int_{0}^{\infty} \frac{\exp\left(-\frac{x^{2}}{2\sigma_{x}^{2}+u} - \frac{y^{2}}{2\sigma_{y}^{2}+u}\right)}{\sqrt{2\sigma_{x}^{2}+u}\sqrt{2\sigma_{y}^{2}+u}} du$$

$$\mathbf{x}(s+0) = S \exp\left[-:\int_{-\Delta}^{\Delta} V_{0}^{-1}(s)U(s)V_{0}(s)ds:\right]\mathbf{x}(s-0)$$

$$V_{0}(s) = S \exp\left[-:\int_{0}^{s} \frac{p_{x}^{2}+p_{y}^{2}}{2}ds:\right] \qquad \mathbf{x} = (x, p_{x}, y, p_{y}, z, \delta)^{t}$$

- Synchro-beta mapping, by Danilov et al, and Hirata et al.
- Strong beam is sliced into some pieces with z_i. s is now function of z_i and z.

Diffusion rate in tune space

- We investigate emittance growth, diffusion, due to symplectic nonlinear force.
- These calculation is done for KEKB. Higher statistics is needed for DAFNE, because small diffusion rate should be concerned.
 KEKB D_v=5x10⁻⁴ /turn DAFNE D_v=1.8x10⁻⁵ /turn
- Increment of second order moment <y²> is evaluated. Luminosity decrement is somewhat easier than the diffusion of the second moment.

Two dimensional model

- Vertical diffusion for ξ =0.136
- $D_C << D_{\gamma}$ for wide region (painted by black).
- No emittance growth, if no interference. Actually, simulation including radiation shows no luminosity degradation nor emittance growth in the region.
- Note KEKB $D_{\gamma}=5x10^{-4}$ /turn DAFNE $D_{\gamma}=1.8x10^{-5}$ /turn



3-D simulation including bunch length ($\sigma_z \sim \beta_y$) Head-on collision





•Good region shrunk drastically.

•Synchrobeta effect near vy~0.5.

• vx~0.5 region remains safe.

- Global structure of the diffusion rate.
- Fine structure near $v_x=0.5$





Bunch length dependence ($\sigma_z \sim \beta_v/2$)













Shorter bunch is better.

Synchrotron tune

• Slow synchrotron tune is better

$$v_s = 0.004$$





Current - ξ parameter 0.070.035 ξ=0.136 0.002 0.002 0.002 0.0015 0.0015 0.0015 10 20 20 0.001 0.001 0.001 0.0005 0.0005 0.0005 15 15 20 10 nux 10 nux 20 20 nux 15 15 15 10 10 nuy nuv nuy

The diffusion rate strongly depend on the beam-beam parameter

• Note KEKB $D_{\gamma}=5x10^{-4}$ /turn DAFNE $D_{\gamma}=1.8x10^{-5}$ /turn



• Good region is only $(v_x, v_y) \sim (0.51.0.55)$.

(0.51,0.51)



5

15

10 nuy

(0.51, 0.7)





Bunch length dependence for finite crossing angle (11 mrad)





$$\sigma_z \sim \beta_y \sigma_z \sim \beta_y/4$$

$$\sigma_z \sim \beta_y/2$$

 $\sigma_z \sim \beta_y/10$





Even in very short bunch, the effect is not negligible.

Synchrotron tune dependence for crossing angle 11mrad









Current dependence for crossing
angle, 11 mradξ=0.1360.070.035



The diffusion rate strongly depend on the beam-beam parameter.

X-y coupling and dispersion at the collision point







- X-Y coupling in 3D and 2D model. Very different behavior.
- Vertical dispersion, $\eta_y = 1$ mm.
- They have strong effects.
- Note KEKB $D_{\gamma}=5x10^{-4}$ /turn DAFNE $D_{\gamma}=1.8x10^{-5}$ /turn

Fearless hypothesis in incoherent emittance growth

$$\lim_{V_x \to 0.5+} D_{C,y} = 0$$

if zero-crossing angle and no error. Discuss again later.

- Dynamic beta, and emittance $\lim_{v_x \to 0.5+} \langle x^2 \rangle < \sigma_{x,0}^2 \qquad \lim_{v_x \to 0.5+} \langle p_x^2 \rangle = \infty$
- Choice of optimum v_x

Mysterious of the beam-beam limit -Non-Gaussian distribution-

- Beam-beam limit near operating point (0.51,0.55-0.58) is very high (~0.3) in Gaussian model.
- Strong-strong simulation gives a beam-beam limit ~0.1-0.15.
- Charge distribution of the both beam distorted from Gaussian, higher Kurtosis.
- The beam-beam limit is reproduced by **2-D** model and/or weak-strong model, if the strong beam as fixed charge distribution has the self-consistent distorted distribution given by strong-strong simulation.
- Radiation excitation is essential for the emittance growth which cause the beam-beam limit.

Beam-beam limit seen in the strong-strong simulation

0

0

2000



 Beam-beam limit is caused by a change of beam charge distribution.



4000

6000

turn

8000

10000

Diffusion seen in the vertical beam size



Strong beam: distorted beam

Strong beam: Gaussian



• Weak-strong simulation using distorted and Gaussian distributions.

- Emittance growth is induced by radiation excitation for the distorted strong beam.
- In the Gaussian strong beam, no strong emittance growth.

Solver: Error function & PIC

50000

Why such big difference appears?

- Poincare plot for particles with x=0 interacting with the distorted strong beam.
- This figure shows an interesting feature, but it is not essential for the strong diffusion actually.



What is the diffusion source?

- Particles with x=0 obeys 1 dimensional dynamics. Motion may be chaotic, but KAM surface limits the diffusion in time dependent 1 dimensional system.
- Investigate an effect of the radiation excitation, $D_{\gamma,x}$, $D_{\gamma,y}$, individually.
- $D_{\gamma,x}$ is essential for the emittance growth.
- Perhaps coupling term like yⁿx^m contribute the vertical diffusion via yⁿ<Δx^m>.



>0.5

- For $v_x=0.5$, particles experience the same y potential every turns, because the potential is symmetric for x, that is, it is 1 degree of freedom system, therefore no emittance growth.
- Needless to say, <x'2> increases as the side effect.



Luminosity degradation due to noise in KEKB -Feedback noise and beam-beam effect-

- In 2005 spring operation, luminosity boosted up 1.35x10³⁴ to 1.58x10³⁴ cm⁻²s⁻¹.
- It is due to that the gain of the transverse bunch-by-bunch feedback system was optimized (weakened but kept a sufficient strength to suppress the coupled bunch instability).





External diffusion: Vertical offset noise (simulation)

- Since the beam-beam system is chaotic, such noise enhances the diffusion of the system.
- Luminosity degradation for the noise without correlation between turns.

$$\langle \Delta y(t) \Delta y(t') \rangle = \Delta y^2 \delta(t-t')$$



Orbit offset (static) (simulation)

- Static vertical offset. Tolerance is easier than the fast noise.
- For slower variation than radiation damping time, emittance can be an adiabatic invariant.

 $\langle \Delta y(t) \Delta y(t') \rangle = \Delta y^2 \exp(-t/\tau)/2\tau$



1/20 compare than that for fast noise

Estimation of feedback noise (Hiramatsu, K.O. & Tobiyama)

- Twp-tap filter and vector composition with two position monitors
- Phase space position at kicker, vector composition with two position monitor

$$\begin{pmatrix} X \\ P \end{pmatrix}_{K} = N \begin{pmatrix} X_{1} \\ X_{2} \end{pmatrix} \qquad N = \frac{1}{\sin \psi_{21}} \begin{pmatrix} \sin \psi_{2K} & \sin \psi_{K1} \\ -\cos \psi_{2K} & \cos \psi_{K1} \end{pmatrix}$$

 Offset noise due to kicker error (δE) and monitor error(δP(δX₁,δX₂)))

$$\left\langle X^{2}\right\rangle = \frac{3\left\langle \delta E^{2}\right\rangle + (3K_{1}^{2} + 3K_{2}^{2} + 2K_{1}K_{2}\cos\mu)\left\langle \delta P^{2}\right\rangle}{-6K_{1}\cos\mu - 2K_{2}\cos 2\mu}$$

Kicker noise measurement (LER)

 (7/14/05) Kicker output depending on feedback gain.

Feedback Gain set	Out Vp-p	OutVrms	kick(Vrms	δE(x10 ⁻⁷)
-13.4dB(3.77 oper.)	107mV	16.3mV	1001	0.62
-14.3dB(final value)	87.3mV	13.3mV	83V	0.51
-10.4dB(Bad lum.)	196.9mV	30mV	188V	1.16
FB switch OFF	4.2mV	0.6mV	4.1V	0.025

 $\delta E = \beta^{1/2} \delta k / E_0$ E₀=3.5 GeV

Speculated beam noise for the kicker noise



Effect on the beam-beam performance of the phase jitter of cavity and crab RF's in KEKB

- Luminosity and beam size as functions of δx .
- Correlation time of the jitter, 1 or 10 turns, is important for the degradation.
- Since Q=200,000 and H=5120, the correlation time will be larger than 10 turns.
- Tolerance is 0.05 degree.





Coherent motion

- Operating point, (0.52,0.58).
- For short bunch $\sigma_z << \beta_y$, vertical coherent motion arise. No coherent motion in horizontal in this operating point.
- The growth rate is not very large.



Tune difference to moderate the coherent motion



Summary

- Sympletic diffusion has been studied using weak-strong Gaussian model. Longitudinal degree of freedom and crossing angle affect the diffusion strongly.
- Golden operation point, v_x→0.5. For head-on collision, the system is regarded as 1 degree of freedom approximately, therefore no emittance growth.
- Diffusion seen in Non-Gaussian interaction couple to radiation excitation determines final beam-beam limit.
 v_x→0.5 is a nice charm again.
- External diffusions, feedback noise, crab cavity noise and others, limit the beam-beam performance.