

“T-Odd Correlations in Single Spin and Azimuthal Asymmetries”

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2nd Workshop on the QCD Structure of the Nucleon

- Remarks Transverse Spin and Azimuthal Asymmetries in QCD
- ★ Reaction Mechanisms: Colinear-limit ETQS-Twist Three, Beyond Co-linearity BHS-FSI Twist Two
- ★ Unintegrated PDF and FFs-ISI/FSI: “ T -odd” TMDs Distribution and Fragmentation Functions: Correlations btwn intrinsic $k_{\perp}$, transverse spin S_T
- ★ TSSA and Estimates of the Collins and Sivers Functions
- ★ Double T -odd $\cos 2\phi$ asymmetry in SIDIS & DRELL-YAN (Gary Goldstein)
- Conclusions

* G. R. Goldstein (Tufts) K.A. Oganessyan NYC, and D.S. Hwang (Sejong) Andreas Metz, Marc Schlegel (Bochum)

Transverse SSA (TSSA) and AZIMUTHAL ASYMMETRIES (AA)

★ Co-linear approximation of QCD PREDICTS
vanishingly small TSSA at large scales and leading order α_s

- Generically,

$$|\perp/\top\rangle = \frac{1}{\sqrt{2}}(|+\rangle \pm i|-\rangle) \Rightarrow A_N = \frac{d\hat{\sigma}^\perp - d\hat{\sigma}^\top}{d\hat{\sigma}^\perp + d\hat{\sigma}^\top} \sim \frac{2 \operatorname{Im} f^{*+} f^-}{|f^+|^2 + |f^-|^2}$$

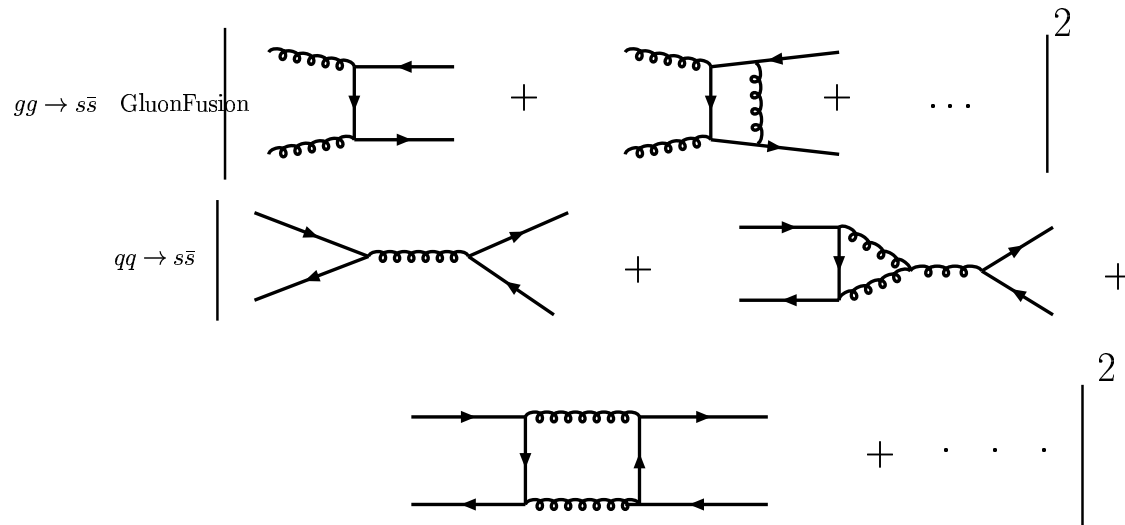
- ★ Requires *helicity flip* as well as *relative phase* btwn helicity amps
- Massless QCD conserves helicity & Born amplitudes are real!
- ★ Incorporating Interference btwn loops-tree level Kane, Repko, PRL:1978 demonstrate $A_N \sim m_q \alpha_s / \sqrt{s}$ within PQCD

Inclusive Λ Production From PQCD ($pp \rightarrow \Lambda^\uparrow X$)

PQCD contributions calculated: [Dharmartna & Goldstein PRD 1990](#)

$$P_\Lambda = \frac{d\sigma^{pp \rightarrow \Lambda^\uparrow X} - d\sigma^{pp \rightarrow \Lambda^\downarrow X}}{d\sigma^{pp \rightarrow \Lambda^\uparrow X} + d\sigma^{pp \rightarrow \Lambda^\downarrow X}}$$

- Need a strange quark to Polarize a Λ ($pp \rightarrow \Lambda^\uparrow X$)

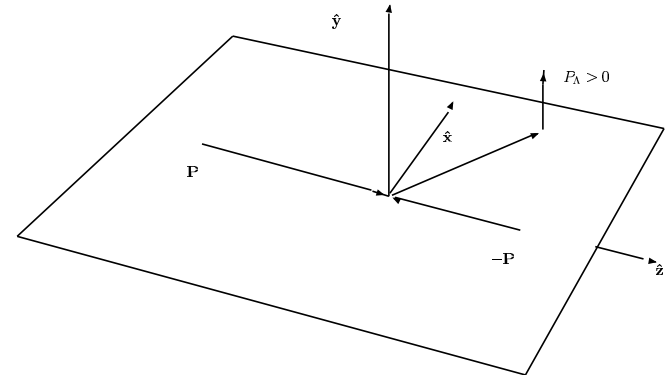
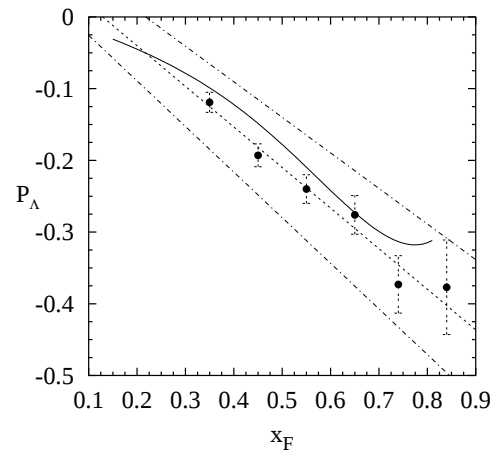


- Polarization $P_\Lambda \sim m_q \alpha_s / \sqrt{s}$ is twist 3 & small $\approx 5\%$ as predicted on general grounds, m_q is the strange quark mass

- Experiment *glaringly at odd with this result*

P_Λ in p-p scattering from Fermi Lab

$$P_\Lambda = \frac{d\sigma^{pp \rightarrow \Lambda^\uparrow X} - d\sigma^{pp \rightarrow \Lambda^\downarrow X}}{d\sigma^{pp \rightarrow \Lambda^\uparrow X} + d\sigma^{pp \rightarrow \Lambda^\downarrow X}}$$

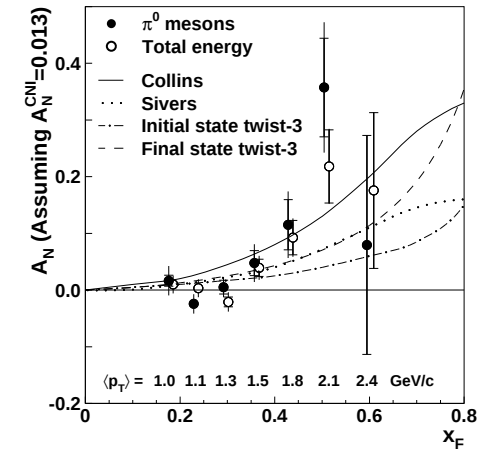
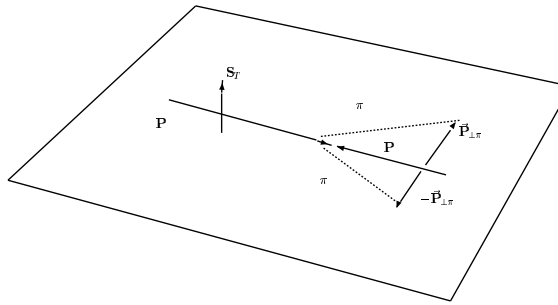


Heller,...,Bunce PRL:1983 PRL: 1983: Up-down asymmetry depicted for Λ production in p-p COM-frame.

Transverse SSA (TSSA) and AZIMUTHAL ASYMMETRIES (AA)

- LARGE TSSAS OBSERVED: E704-Fermi Lab, STAR & PHENIX

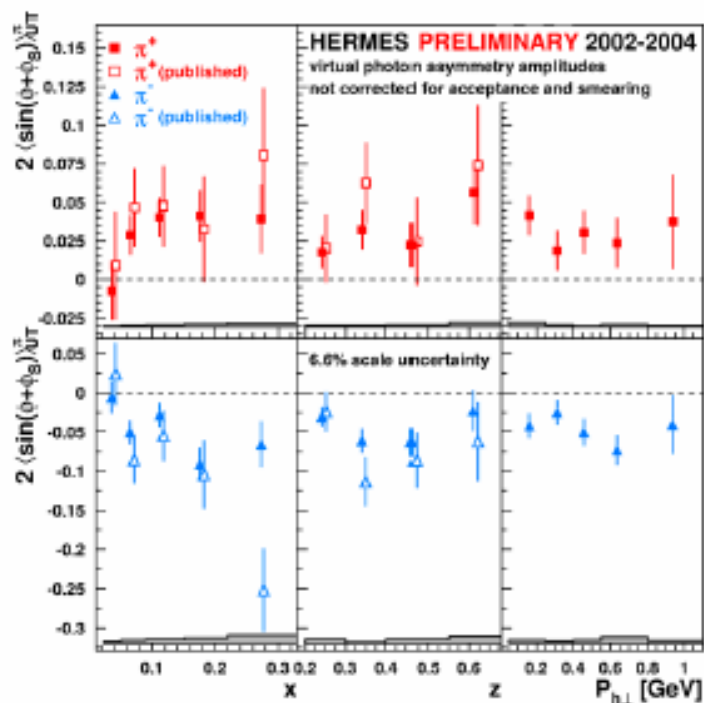
$$A_N = \frac{d\sigma^{p^\uparrow p \rightarrow \pi^- X} - d\sigma^{p^\downarrow p \rightarrow \pi^- X}}{d\sigma^{p^\uparrow p \rightarrow \pi^- X} + d\sigma^{p^\downarrow p \rightarrow \pi^- X}}$$



L-R asymmetry of π production and A_N for π_0 production at STAR : PRL 2004

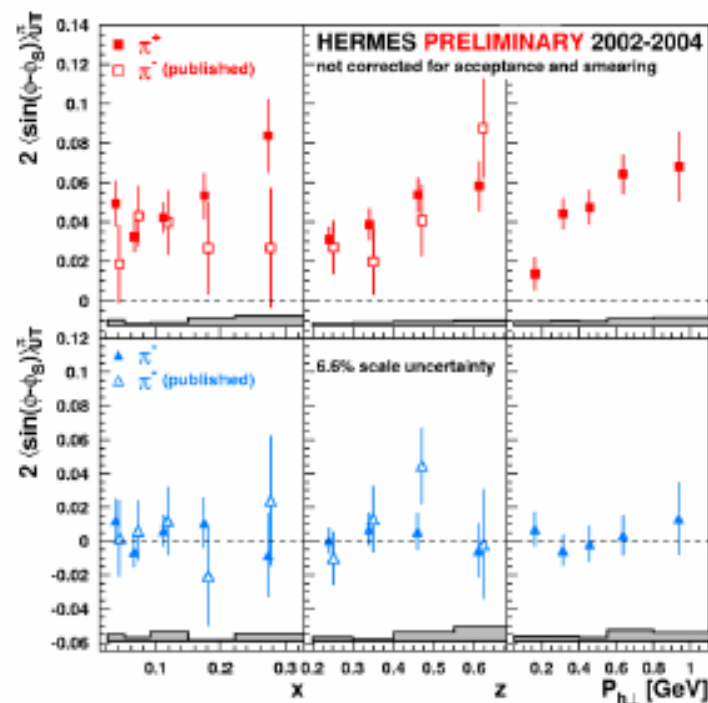
Update including Data 2004

Collins



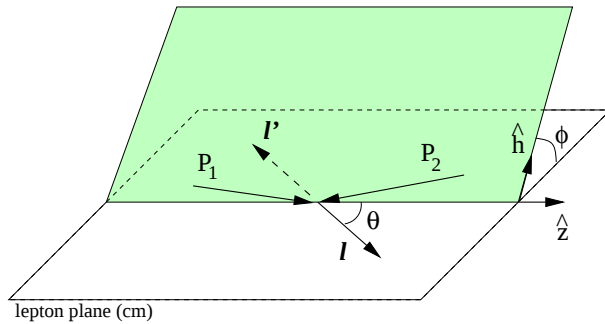
HAWAII05 T.-A. Shibata

Sivers



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Unpolarized DRELL YAN $\cos 2\phi$



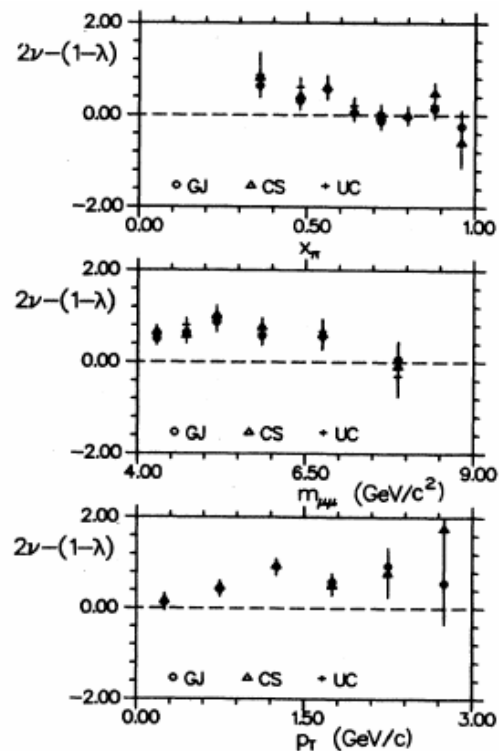
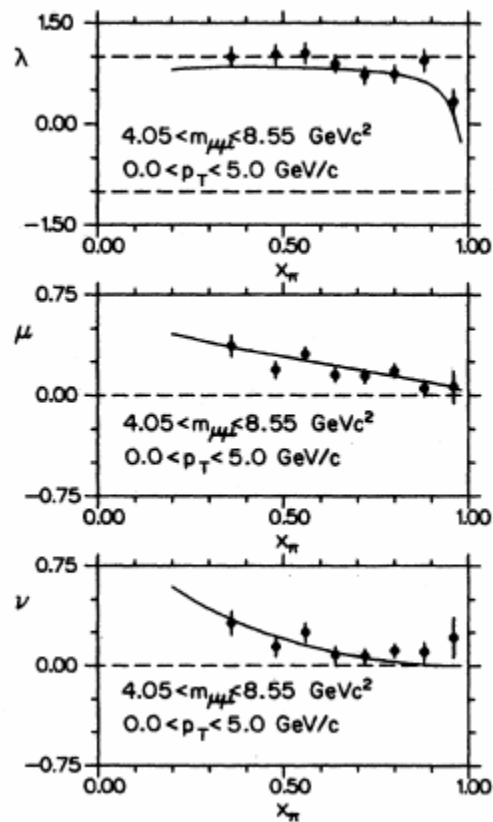
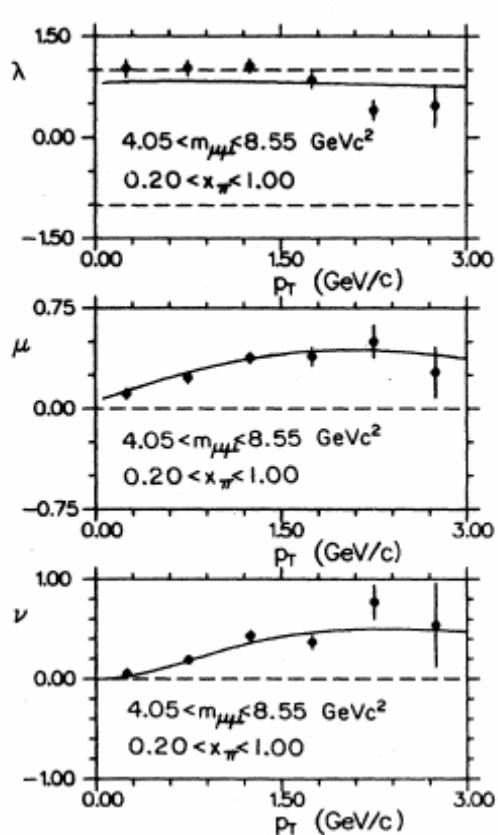
$$\bar{p} + p \rightarrow \mu^- \mu^+ + X$$

TALK of Gary Goldstein

$$\frac{dN}{d\Omega} = \left(\frac{d\sigma}{d^4q}\right)^{-1} \frac{d\sigma}{d^4q d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda + 3} \left(1 + \lambda \cos^2 \theta + \mu \sin^2 \theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi \right)$$

λ, μ, ν , depend on $s, x, m_{\mu\mu}^2, q_T$: QCD-parton model NLO and NNLO predict Lam-Tung relation $1 - \lambda - 2\nu = 0$

Measurements of $\pi^- + p \rightarrow \mu^+ + \mu^- + X$ discovered unexpectedly large values of these asymmetries [E615, NA10] compared to parton-model expectations resulting in violation of the Lam-Tung relation



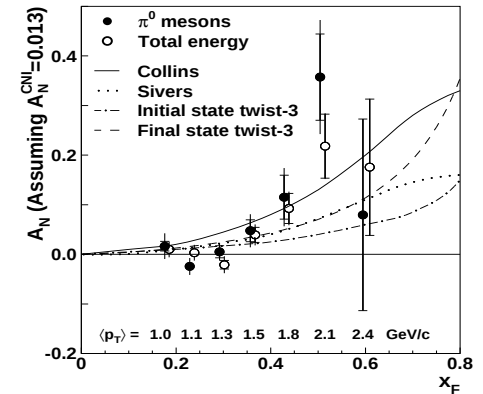
Lam-Tung Relationship Violated

ETQS-Twist Three Mechanism

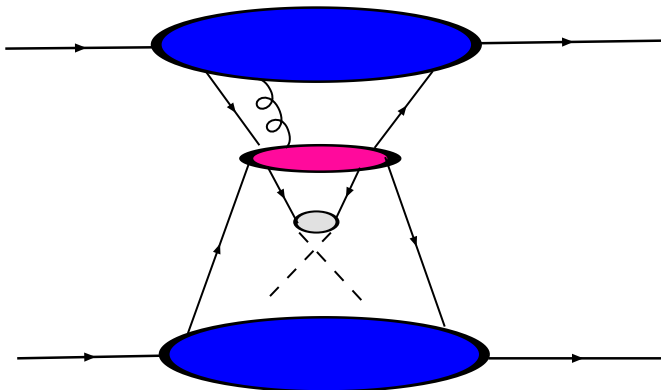
can describe TSSAs $pp^\perp \rightarrow \pi X$

@ Lg P_T , A_N twist three but phases can be generated in co-linear QCD from gluonic and fermionic poles in propagator of hard parton subprocess

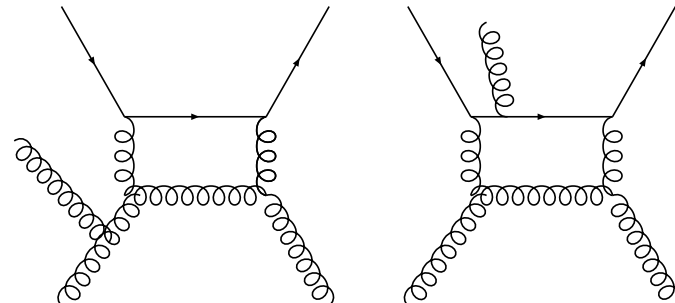
$$\frac{1}{x + s \pm i\epsilon} = P \frac{1}{x - s} \mp i\pi\delta(x - s)$$



- Qiu & Sterman :PLB 1991, 1999 & Koike & Kanazawa:PLB 2000 at Large $P_T > \Lambda_{qcd}$ get helicity flip and phases



$\Rightarrow$



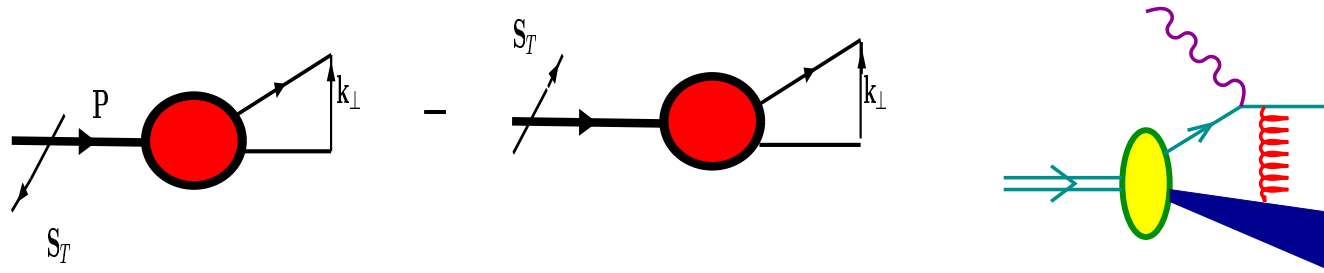
$k_{\perp} \sim \Lambda_{\text{qcd}}$ “Naive- T -Odd” Correlations thru TMDs

- Sensitivity to $k_{\perp}$ intrinsic quark momenta, associated non-perturbative transverse momentum distribution functions **TMD**

Soper, PRL:1979: $\int d\mathbf{k}_{\perp} \mathcal{P}(\mathbf{k}_{\perp}, x) = f(x)$

- TSSA indicative “ T -odd” correlations among *transverse* spin and momenta

Sivers: PRD 1990 e.g. $P P^{\perp} \rightarrow \pi X$ $i\mathbf{S}_T \cdot (\mathbf{P} \times \mathbf{k}_{\perp}) \rightarrow f_{1T}^{\perp}(x, \mathbf{k}_{\perp})$

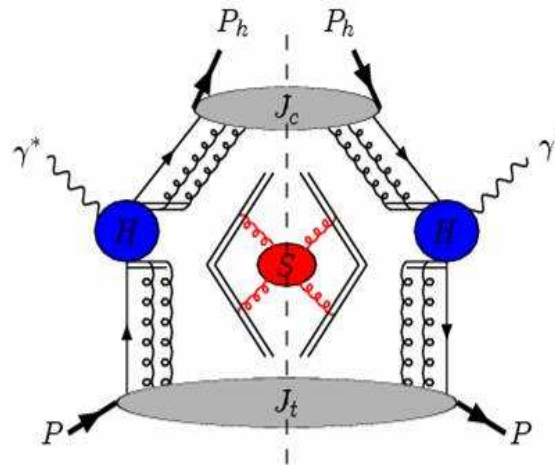


- Correlation accounts for left-right TSSA in inclusive π production
(Sivers: PRD 1990, Anselmino & Murgia PLB: 1995 ... Brodsky, Hwang, and Schmidt PLB: 2002
rescattering of a gluon could produce the necessary phase leading to nonzero SSAs at *Leading Twist*)

- Collins NPB 1993 “ T -odd” correlation of transversely polarized fragmenting quark: TSSA in lepto-production
 $\ell \vec{p} \rightarrow \ell' \pi X$ $i\mathbf{s}_T \cdot (\mathbf{p} \times \mathbf{P}_{h\perp}) \rightarrow H_1^{\perp}(x, \mathbf{p}_{\perp})$
 s_T spin of fragmenting quark, $\mathbf{p}$ quark momentum and $\mathbf{P}_{h\perp}$ transverse momentum produced pion

Factorization Demonstrated For TMD PDF and FF and Hard and Soft Parts

More recently [Ji, Ma, Yuan: PLB, PRD 2004, 2005](#) building on work of Collins-Soper NPB: 81, extended factorization theorems to 1-loop and beyond



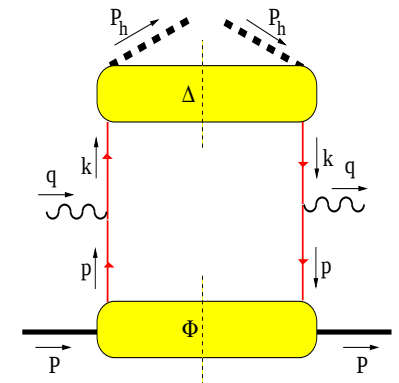
Universality & Factorization “Maximally” Correlated [Collins and Metz: PRL 2005](#)

Beyond Co-linear QCD: T -Odd Correlations

Recent Times Boer & Mulders and Co. incorporated $\mathbf{k}_\perp$ T -odd PDFs and FFs relevant to hard scattering QCD at leading twist. Adopted Factorized Description

Ellis, Furmanski, Petronzio NPB: 1982, Collins *et al.* PQCD... : 82 , J. Qui PRD: 1990, Levelt & Mulders, Mulders & Tangerman, NPB: 1994, 1996

$$\frac{d\sigma^{\ell N \rightarrow \ell' h X}}{dx dy dz d^2 P_{h\perp}} = \frac{M \pi \alpha^2 y}{2Q^4 z} L_{\mu\nu} \mathcal{W}^{\mu\nu}$$



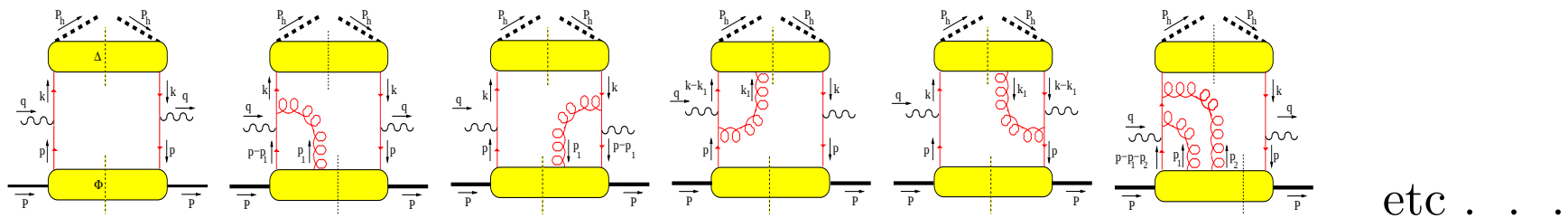
Hadronic Tensor

$$2M\mathcal{W}^{\mu\nu}(q, P, P_h) = \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^2(\mathbf{p}_T + \mathbf{q}_T - \mathbf{k}_T) \text{Tr}[\Phi(x_B, \mathbf{p}_T) \gamma^\mu \Delta(z_h, \mathbf{k}_T) \gamma^\nu] \\ + (q \leftrightarrow -q, \mu \leftrightarrow \nu)$$

T-Odd Effects to QCD Processes Naturally Built into Color Gauge Invariant Factorized QCD at “leading twist” thru-Wilson Line

- Gauge Invariant Distribution and Fragmentation Functions

Boer, Mulder: NPB 2000, Ji, Yuan & Belitsky PLB: 2002, NPB 2003, Boer, Mulder, Pijlman NPB 2003



Sub-class of loops in eikonal limit (soft gluons) sum up to yield color gauge invariant hadronic tensor factorized into the distribution Φ and fragmentation Δ operators

$$\Phi(p, P) = \int \frac{d^3\xi}{2(2\pi)^3} e^{ip \cdot \xi} \langle P | \bar{\psi}(\xi^-, \xi_\perp) \mathcal{G}_{[\xi^-, \infty]}^\dagger | X \rangle \langle X | \mathcal{G}_{[0, \infty]} \psi(0) | P \rangle |_{\xi^+ = 0}$$

$$\Delta(k, P_h) = \int \frac{d^3\xi}{4z(2\pi)^3} e^{ik \cdot \xi} \langle 0 | \mathcal{G}_{[\xi^+, -\infty]} \psi(\xi) | X; P_h \rangle \langle X; P_h | \bar{\psi}(0) \mathcal{G}_{[0, -\infty]}^\dagger | 0 \rangle |_{\xi^- = 0}$$

$$\mathcal{G}_{[\xi, \infty]} = \mathcal{G}_{[\xi_T, \infty]} \mathcal{G}_{[\xi^-, \infty]}, \quad \text{where} \quad \mathcal{G}_{[\xi^-, \infty]} = \mathcal{P} \exp(-ig \int_{\xi^-}^{\infty} d\xi^- A^+)$$

Provide source of T-Odd Contributions to TSSA and AA

- “T-odd” distribution and fragmentation functions enter transverse momentum dependent distribution and fragmentation correlators at *leading twist* Boer, Mulder: PRD 1998

$$\Delta(z, \mathbf{k}_\perp) = \frac{1}{4} \left\{ D_1(z, z\mathbf{k}_\perp) \not{n}_- + H_1^\perp(z, z\mathbf{k}_\perp) \frac{\sigma^{\alpha\beta} k_{\perp\alpha} n_{-\beta}}{M_h} + D_{1T}^\perp \frac{\epsilon_{\mu\nu\rho\sigma} \gamma^\mu n_-^\nu k_{\perp}^\rho S_{hT}^\sigma}{M_h} + \dots \right\},$$

$$\Phi(x, \mathbf{p}_\perp) = \frac{1}{2} \left\{ f_1(x, \mathbf{p}_\perp) \not{n}_+ + h_1^\perp(x, \mathbf{p}_\perp) \frac{\sigma^{\alpha\beta} p_{T\alpha} n_{+\beta}}{M} + f_{1T}^\perp(x, \mathbf{p}_\perp) \frac{\epsilon^{\mu\nu\rho\sigma} \gamma^\mu n_+^\nu p_{\perp}^\rho S_T^\sigma}{M} \dots \right\}$$

SIDIS cross section

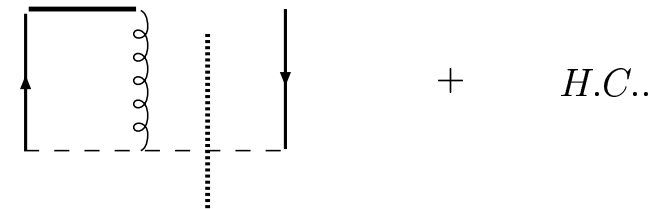
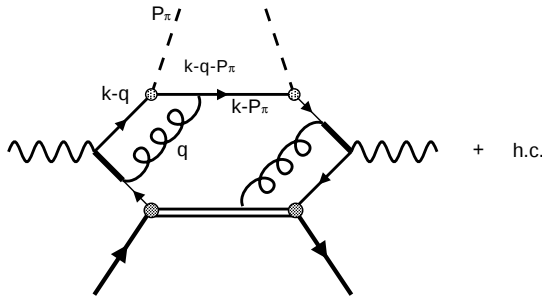
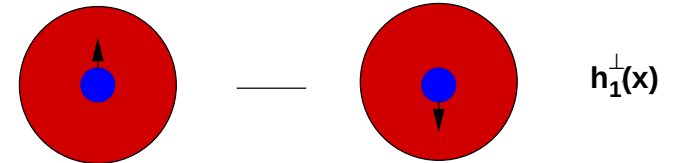
$$\begin{aligned} d\sigma_{\{\lambda, \Lambda\}}^{\ell N \rightarrow \ell \pi X} &\propto f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 + \frac{k_\perp}{Q} f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 \cdot \cos \phi \\ &+ \left[\frac{k_\perp^2}{Q^2} f_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 + h_1^\perp \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes H_1^\perp \right] \cdot \cos 2\phi \\ &+ |S_T| \cdot h_1 \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes H_1^\perp \cdot \sin(\phi + \phi_S) \quad \text{Collins} \\ &+ |S_T| \cdot f_{1T}^\perp \otimes d\hat{\sigma}^{\ell q \rightarrow \ell q} \otimes D_1 \cdot \sin(\phi - \phi_S) \quad \text{Sivers} \\ &+ \dots \end{aligned}$$

FSI Mechanism can Generate Boer-Mulders- $h_1^\perp$

Goldstein, L.G.–ICHEP-proc-hep-ph/0209085 (2002), L.G., Goldstein, Oganessyan PRD 2003

- $h_1^\perp$ “Naturally” defined from Color G.I. TMD: Convolutd with $H_1^\perp$ enters $\cos 2\phi$
- “Eikonal Feynman rules” to calculate Collins Soper: NPB: 1982

$$\Phi_{[h_1^\perp]}^{[\sigma^\perp + \gamma_5]}(x, k_\perp) = \frac{1}{2} \int dp^- \text{Tr} \left(i\sigma^{\perp+} \gamma_5 \Phi \right) = \frac{\varepsilon_{+-\perp} j^k j^\perp}{M} h_1^\perp(x, k_\perp)$$



$$\Phi^{[\Gamma]}(x, k_\perp) = \sum_X \int \frac{d\xi^- d^2\xi_\perp}{2(2\pi)^3} e^{-i\xi \cdot \vec{k}_\perp} \langle P | \bar{\psi}(\xi) \mathcal{G}_{[\xi^-, \infty]}^\dagger | X \rangle \langle X | \mathcal{G}_{[0, \infty]} \Gamma \psi(0) | P \rangle |_{\xi^+ = 0}$$

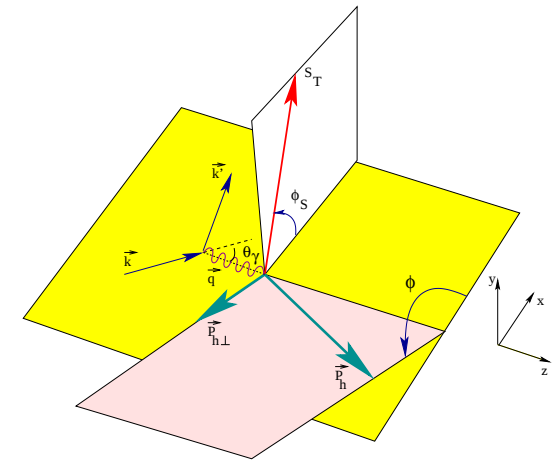
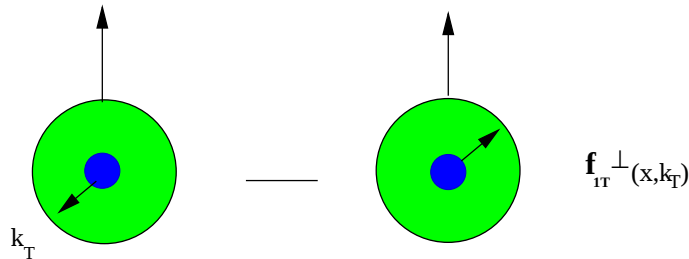
$h_1^\perp(x, k_\perp)$, represents, number density transversely polarized quarks in an unpolarized nucleons
nucleons complementary to $f_{1T}^\perp(x, k_\perp)$

SIDIS-Transversity Properties at Leading Twist

- Collins NPB:1993, Kotzinian NPB:1995, Mulders, Tangerman PLB:1995

$$\langle \frac{P_{h\perp}}{M_\pi} \sin(\phi + \phi_s) \rangle_{UT} = \frac{\int d\phi_s \int d^2 P_{h\perp} \frac{P_{h\perp}}{M_\pi} \sin(\phi + \phi_s) (d\sigma^\uparrow - d\sigma^\downarrow)}{\int d\phi_s \int d^2 P_{h\perp} (d\sigma^\uparrow + d\sigma^\downarrow)} = |S_T| \frac{2(1-y) \sum_q e_q^2 h_1(x) z H_1^{\perp(1)}(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)}$$

- Sivers PRD: 1990, Anselmino & Murgia PLB: 1995 ...

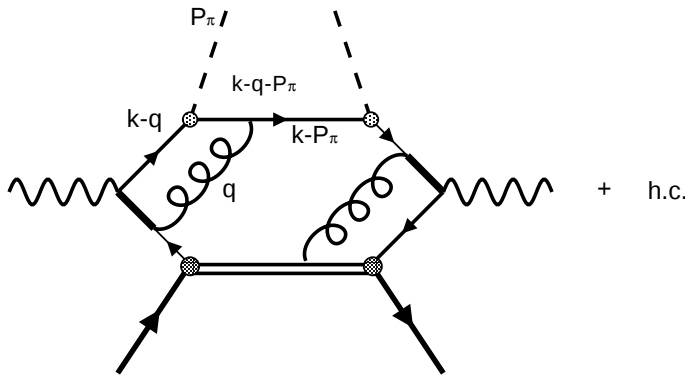


$$\langle \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) \rangle_{UT} = |S_T| \frac{(1 + (1 - y)^2) \sum_q e_q^2 f_{1T}^{\perp(1)}(x) z D_1^q(z)}{(1 + (1 - y)^2) \sum_q e_q^2 f_1(x) D_1(z)},$$

- Probes the probability for a transversely polarized target, pions are produced asymmetrically about pion production plane

$\cos 2\phi$ Asymmetry Generated by ISI & FSI thru Gauge link

Goldstein, L.G.–ICHEP-Amsterdam: 2002, hep-ph/0209085, G.G. & Oganessyan PRD:2003



$$\frac{\int d^2 P_{h\perp} \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi d\sigma}{\int d^2 P_{h\perp} d\sigma} = \langle \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi \rangle_{UU} = \frac{8(1-y) \sum_q e_q^2 h_1^{\perp(1)}(x, Q^2) z^2 H_1^{\perp(1)q}(z, Q^2)}{(1+(1-y)^2) \sum_q e_q^2 f_1^q(x, Q^2) D_1^q(z, Q^2)}$$

$$\frac{d\sigma}{dx dy dz d^2 P_\perp} \propto f_1 \otimes D_1 + \frac{k_T}{Q} f_1 \otimes D_1 \cdot \cos \phi + \left[\frac{k_T^2}{Q^2} f_1 \otimes D_1 + h_1^\perp \otimes H_1^\perp \right] \cdot \cos 2\phi$$

Leading Twist Contribution from T -Odd D. Boer, P. Mulders, PRD: 1998

Estimates of T-odd Contribution in SIDIS (HERMES, JLAB 6& 12 GeV program)

$\cos 2\phi$ Asymmetry in SIDIS:

Boer Mulders Effects Competes with Cahn Effect and Radiative corrections

- ★ The spectator model used in previous rescattering calculations assumes point-like nucleon-quark-diquark vertex, **leads to logarithmically divergent, asymmetries**

Goldstein, L.G., ICHEP 2002; hep-ph/0209085,

L.G., Goldstein, Oganessyan PRD 2003; Boer, Brodsky, Hwang, PRD: 2003(Drell-Yan)

$$\begin{aligned} h_1^{\perp(s)}(x, k_{\perp}) &= f_{1T}^{\perp(s)}(x, k_{\perp}) \\ &= \alpha_s N_s \frac{(1-x)M(m+xM)}{k_{\perp}^2 \Lambda(k_{\perp}^2)} \ln \frac{\Lambda(k_{\perp}^2)}{\Lambda(0)} \end{aligned}$$

$$\Lambda(k_{\perp}^2) = k_{\perp}^2 + x(1-x) \left(-M^2 + \frac{m^2}{x} + \frac{\mu^2}{1-x} \right)$$

- **Asymmetry involves weighted function**

$$h_1^{(1)\perp}(x) \equiv \int d^2 k_{\perp} \frac{k_{\perp}^2}{2M^2} h_1^{\perp}(x, k_{\perp}^2) \quad \text{diverges}$$

Gaussian Distribution in $k_{\perp}$

Log divergence addressed by approximating the transverse momentum dependence of the quark-nucleon-vertex by a Gaussian distribution in $k_{\perp}^2$,

L.G., Goldstein, Oganessyan, PRD 67 (2003)

$$\langle n | \psi(0) | P \rangle = \left(\frac{i}{\not{k} - m} \right) \Upsilon(k_{\perp}^2) U(P, S), \quad b \equiv \frac{1}{\langle k_{\perp}^2 \rangle}$$

where $\Upsilon(k_{\perp}^2) = \mathcal{N} e^{-b k_{\perp}^2}$.

$U(P, S)$ nucleon spinor, and quark propagator comes from untruncated quark line

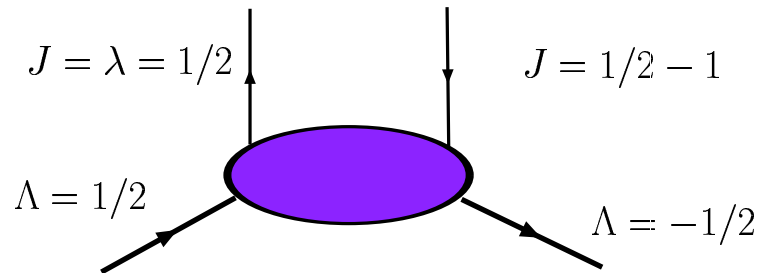
$$h_1^{\perp}(x, k_{\perp}) = \alpha_s \mathcal{N}_s \frac{M(m + xM)(1 - x)}{k_{\perp}^2 \Lambda(k_{\perp}^2)} \mathcal{R}(k_{\perp}^2, x)$$

with

$$\mathcal{R}(k_{\perp}^2, x) = \exp^{-2b(k_{\perp}^2 - \Lambda(0))} \left(\Gamma(0, 2b\Lambda(0)) - \Gamma(0, 2b\Lambda(k_{\perp}^2)) \right)$$

- $\lim < k_{\perp}^2 > \rightarrow \infty$ width goes to infinity, regain *log* result

GPDs and correlations of transverse spin and intrinsic $k_\perp$:



- Intriguing connection of Sivers effect/function $f_{1T}^{\perp(q)} \leftrightarrow -\kappa^q$ with anomalous magnetic moment of quark- q through the impact parameter space representation of the spin-flip, chirally-even GPD $\mathcal{E}(x, \mathbf{b}_\perp)$: serves to fix sign of Sivers function
- As well $k_T^q \leftrightarrow h_1^{\perp q}$ through $\tilde{H}_T(x, 0, -\Delta_\perp)$ and $E_T(x, 0, -\Delta_\perp)$ (chirally odd transversity GPDs) where κ_T governs the transverse spin-flavor dipole moment in an unpolarized target
Burkardt & Hägler, Diehl et. al
- ★ This result implies that the **up** and **down** quark Boer Mulders function are same sign. Confirms **Lg N_C** arguments of Poblitsa [hep-ph/0301236](https://arxiv.org/abs/hep-ph/0301236) Implications on $\cos 2\phi$ phenomenology in SIDIS & Drell Yan

INPUTS: Boer-Mulders $h_1^{\perp(1/2)}$ and Unpolarized Structure Function $f_1(x)$

$$f_1(x) = \frac{g^2}{(2\pi)^2} (1-x) \cdot \left\{ \frac{(m+xM)^2 - \Lambda(0)}{\Lambda(0)} - \left[2b \left((m+xM)^2 - \Lambda(0) \right) - 1 \right] e^{2b\Lambda(0)} \Gamma(0, 2b\Lambda(0)) \right\}$$

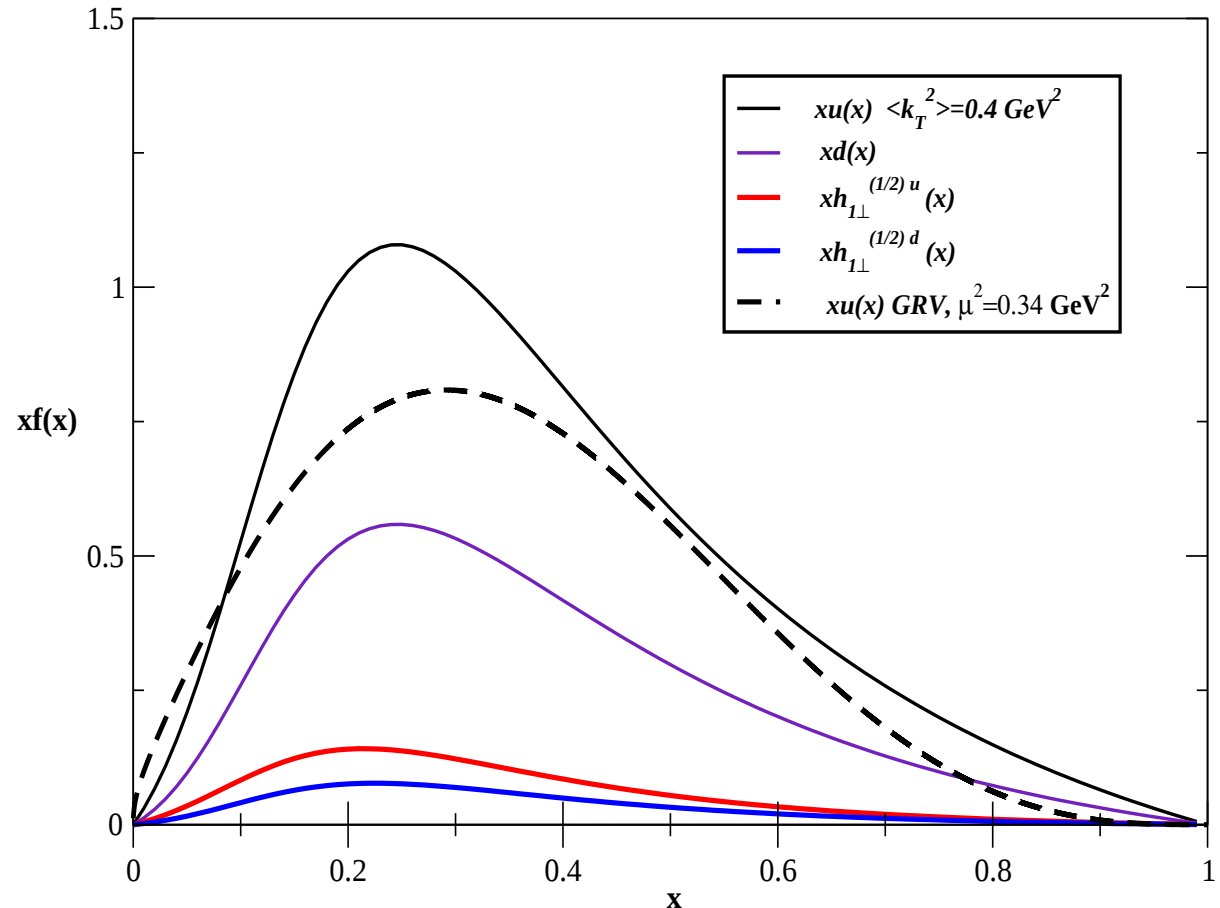
★ Valence Normalization,
 $\int_0^1 u(x) = 2, \int_0^1 d(x) = 1$

- Black curve- $xu(x)$
- Dashed curve - $xu(x)$ GRV
- Red/Blue curve $xh_1^{\perp(1/2)(u,d)}$
- axial vector diquark coupling

Jakob, Mulders, Rodrigues

NPB:1997,

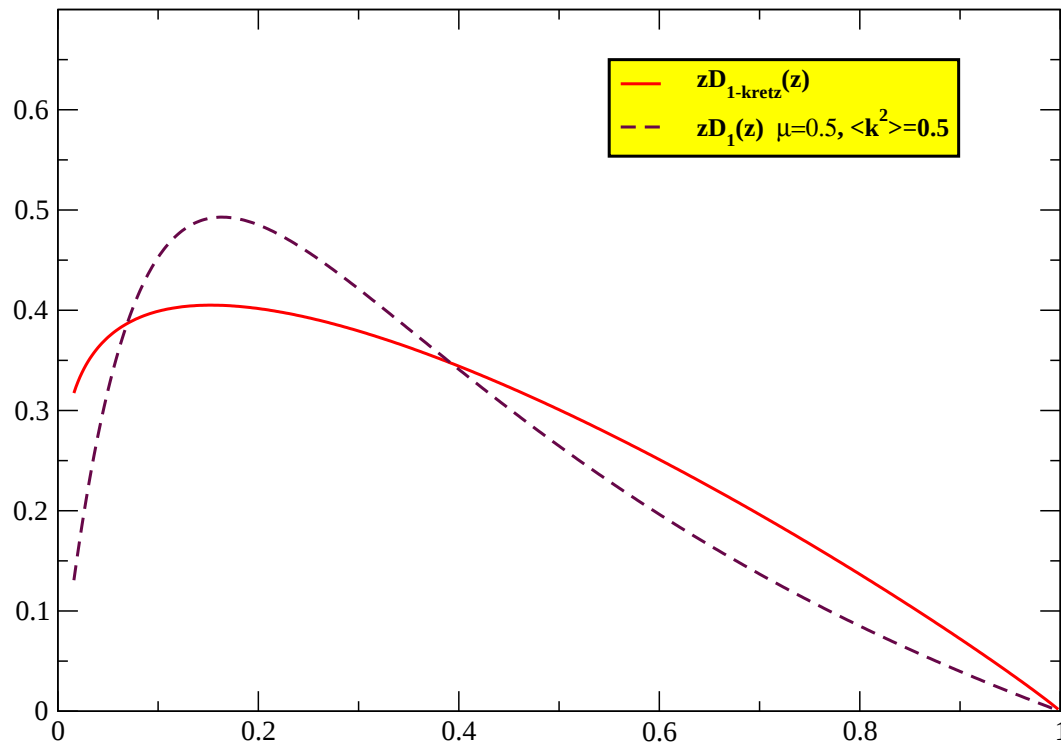
$$\gamma_5(\gamma^\mu + P^\mu/M)$$



Pion Fragmentation Function

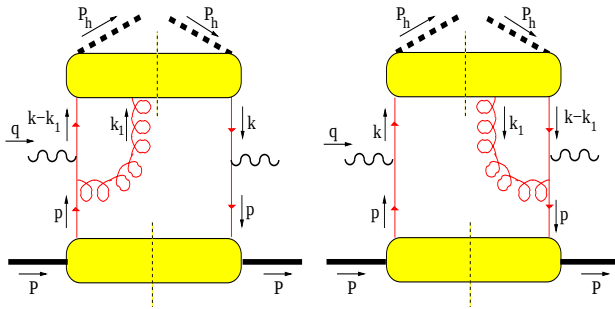
$$D_1(z) = \mathcal{N}' \frac{1}{z} \frac{(1-z)}{z} \left\{ \frac{m^2 - \Lambda'(0)}{\Lambda'(0)} - \left[2b' \left(m^2 - \Lambda'(0) \right) - 1 \right] e^{2b' \Lambda'(0)} \Gamma(0, 2b' \Lambda'(0)) \right\},$$

which, multiplied by z at $\langle k_{\perp}^2 \rangle = (0.5)^2 \text{ GeV}^2$ and $\mu = m$, estimates the distribution of [Kretzer, PRD: 2000](#)



Gauge Link-Pole Contribution to T -Odd Collins Function

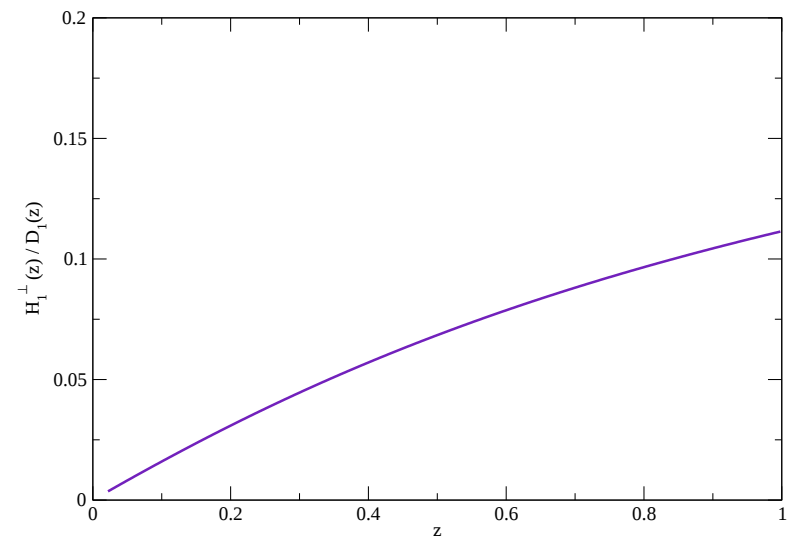
L.G., Goldstein, Oganessyan PRD68,2003 $\Delta^{[\sigma^\perp - \gamma_5]}(z, k_\perp) = \frac{1}{4z} \int dk^+ \text{Tr}(\gamma^- \gamma^\perp \gamma_5 \Delta) |_{k^- = P_\pi^- / z}$



$\frac{i}{\ell - i\epsilon}$ Eikonal Feynman Rules Collins '82' $\frac{-i}{\ell + i\epsilon}$

Motivation: color gauge .inv frag. correlator
 "pole contribution" leading twist
 T -odd pion fragmentation

$$H_1^\perp(z, k_\perp) = \mathcal{N}' \alpha_s \frac{(1-z)}{z^2} \frac{\mu - m(1-z)}{z} \frac{M_\pi}{k_\perp^2 \Lambda'(k_\perp^2)} \mathcal{R}(z, k_\perp^2)$$



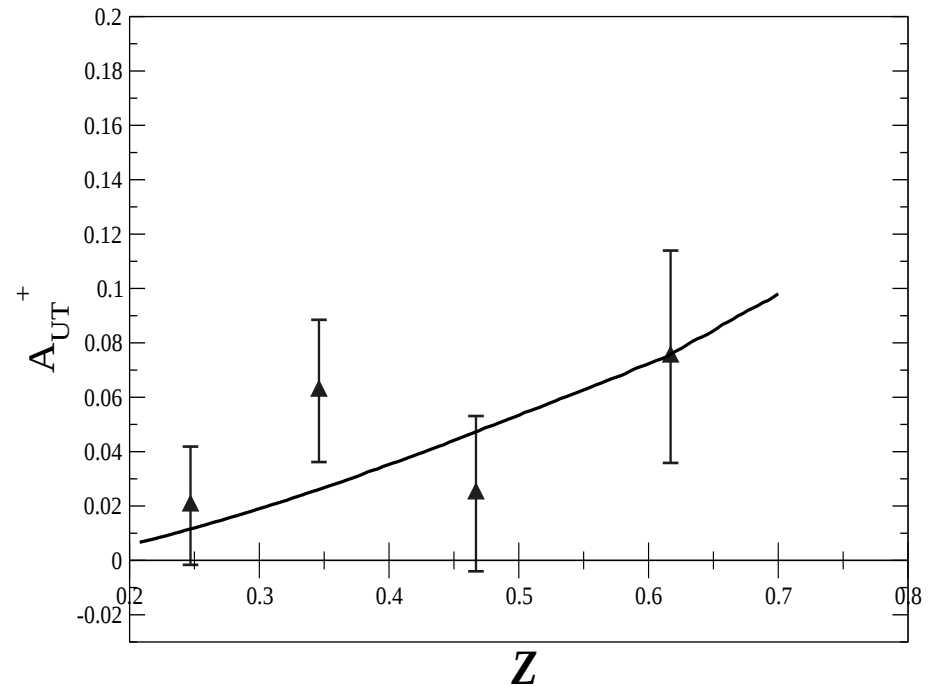
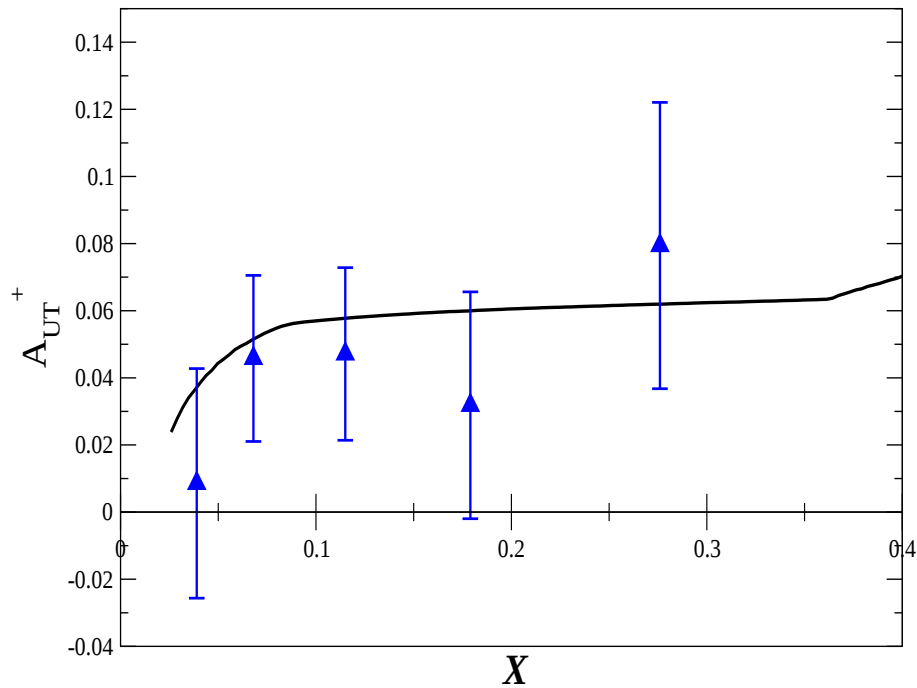
Collins Asymmetry

L.G., Goldstein, Oganessyan PRD 2003: updated For the HERMES kinematics

$1 \text{ GeV}^2 \leq Q^2 \leq 15 \text{ GeV}^2$, $4.5 \text{ GeV} \leq E_\pi \leq 13.5 \text{ GeV}$, $0.2 \leq x \leq 0.41$, $0.2 \leq z \leq 0.7$, $0.2 \leq y \leq 0.8$, $\langle P_{h\perp}^2 \rangle = 0.25 \text{ GeV}^2$

$$\left\langle \frac{P_{h\perp}}{M_\pi} \sin(\phi + \phi_s) \right\rangle_{UT} = |S_T| \frac{2(1-y) \sum_q e_q^2 h_1(x) z H_1^{\perp(1)}(z)}{(1 + (1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)}.$$

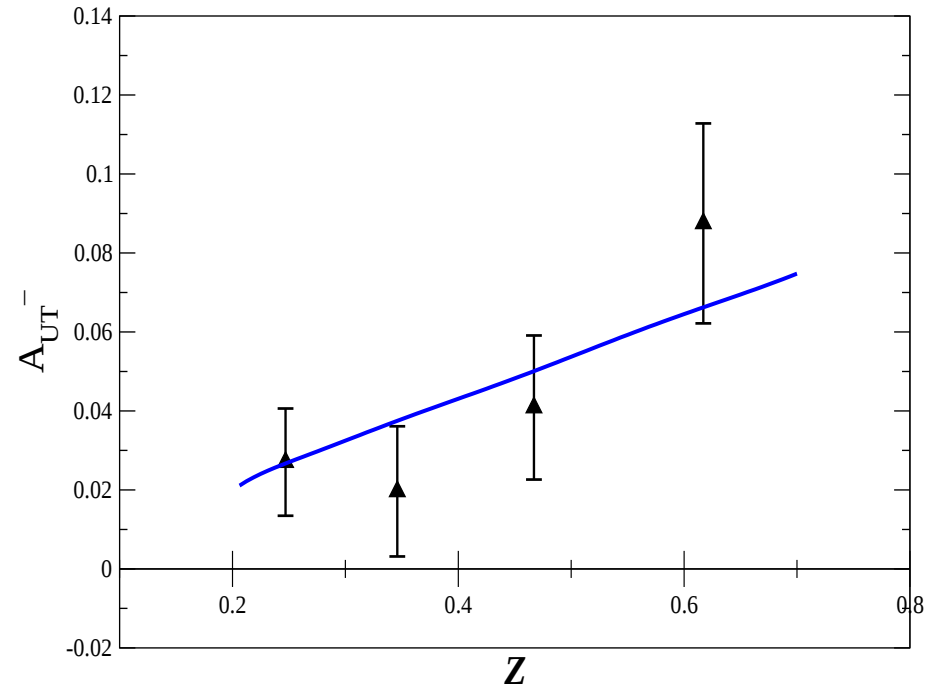
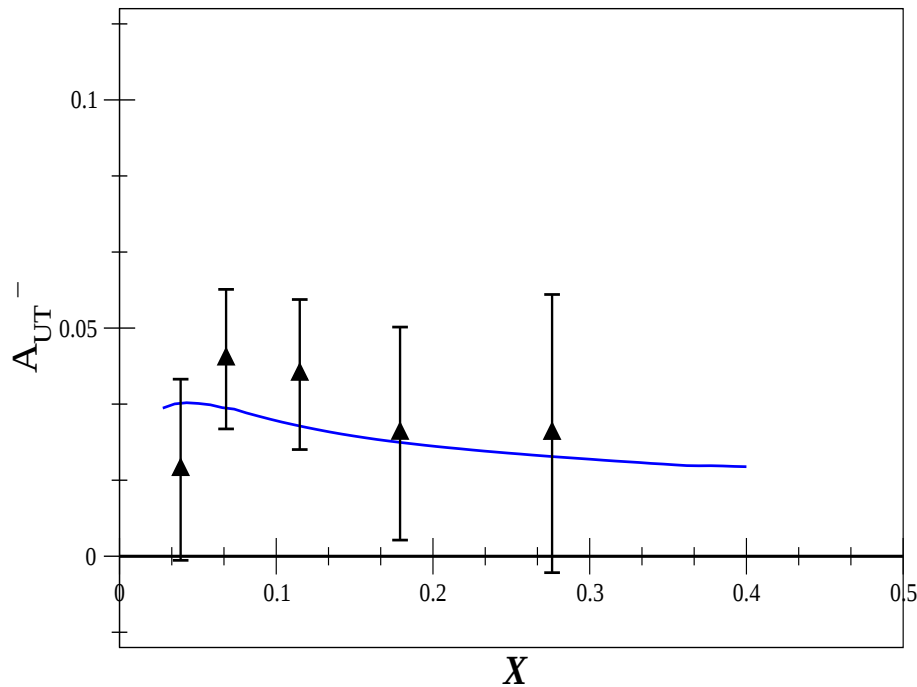
Data from A. Airapetian et al. PRL94,2005



Estimates for Sivers Asymmetry

Data from A. Airapetian et al. PRL94,2005

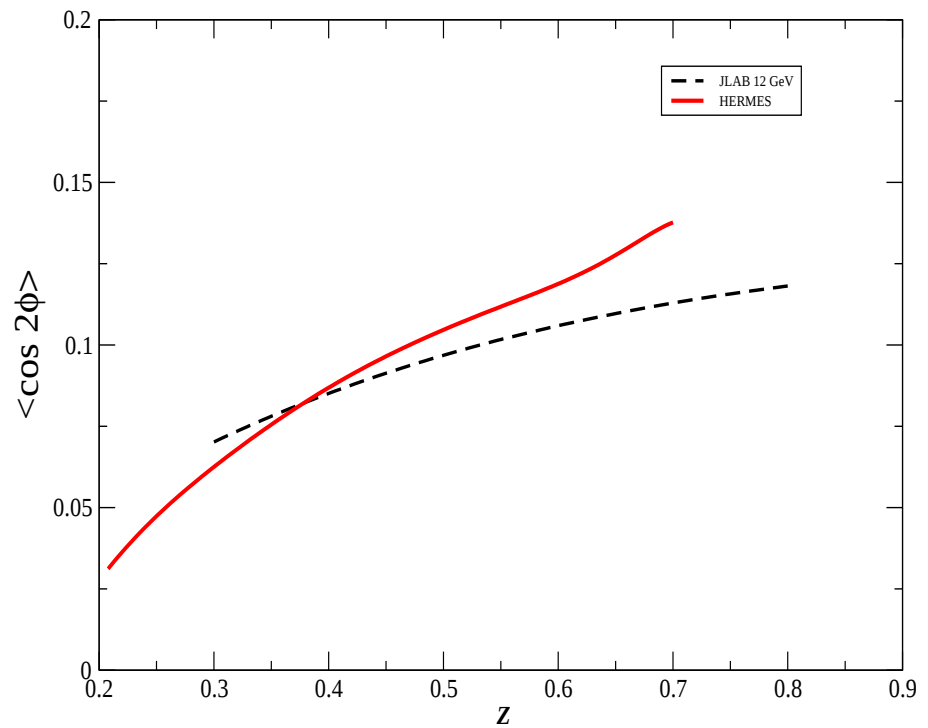
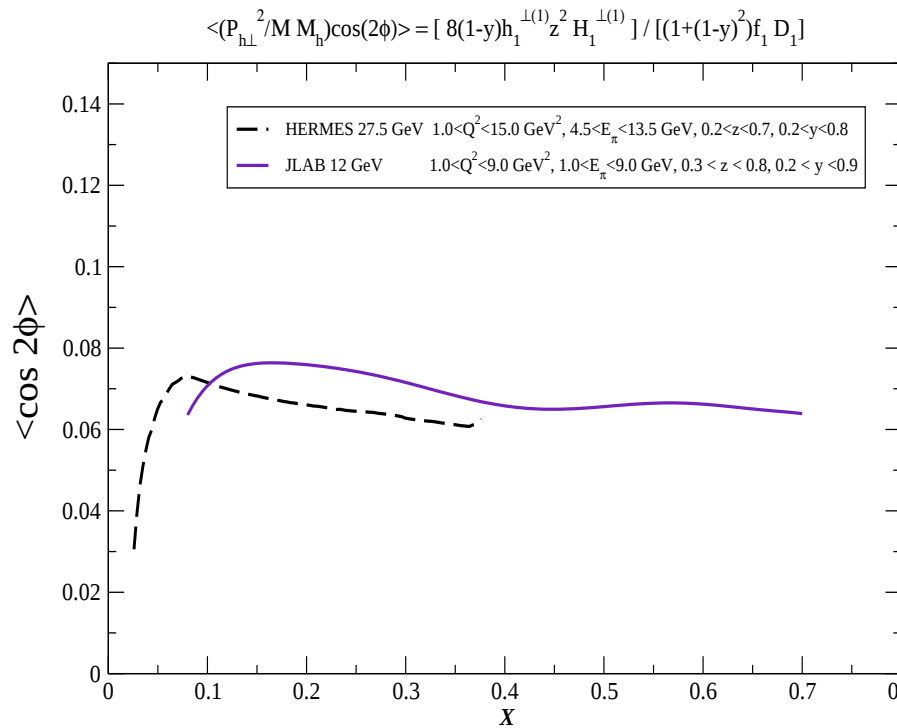
$$\left\langle \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) \right\rangle_{UT} = \frac{\int d^2 P_{h\perp} \frac{|P_{h\perp}|}{M} \sin(\phi - \phi_S) d\sigma}{\int d^2 P_{h\perp} d\sigma} = \frac{(1 + (1 - y)^2) \sum_q e_q^2 f_{1T}^{\perp(1)}(x) z D_1^q(z)}{(1 + (1 - y)^2) \sum_q e_q^2 f_1(x) D_1(z)},$$

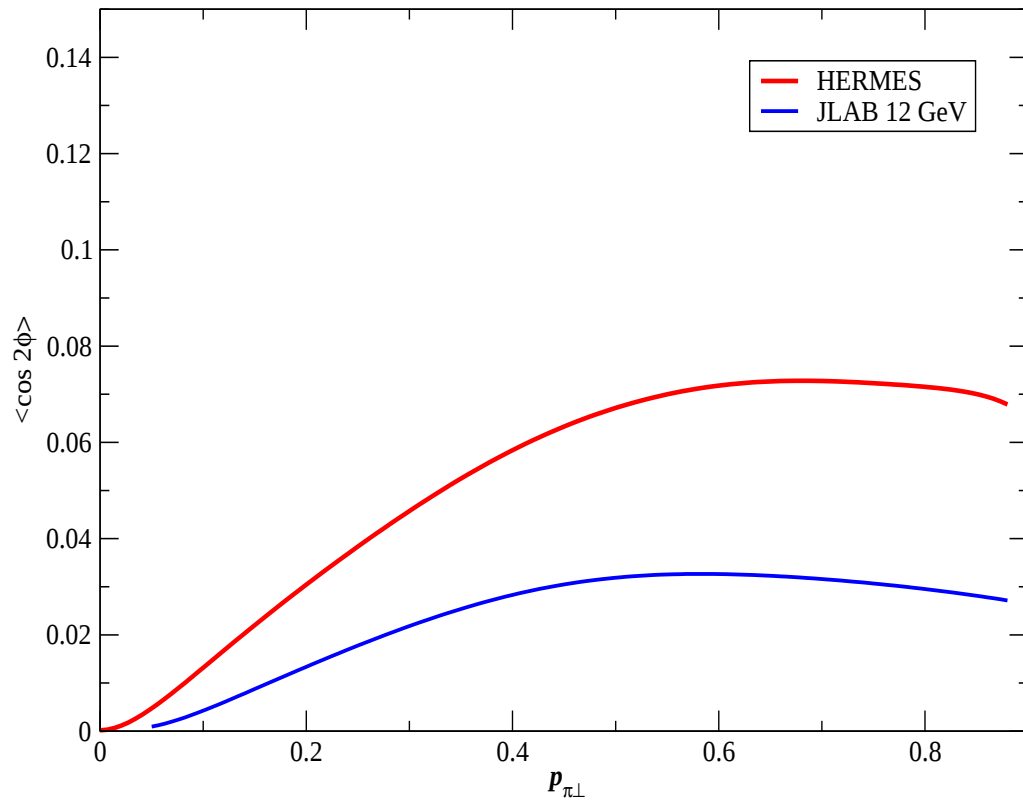


Double T-odd $\cos 2\phi$ asymmetry

Transversity of quarks inside an unpolarized hadron, and $\cos 2\phi$ asymmetries in unpolarized semi-inclusive DIS

$$\left\langle \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi \right\rangle_{UU} = \frac{\int d^2 P_{h\perp} \frac{|P_{h\perp}^2|}{MM_h} \cos 2\phi d\sigma}{\int d^2 P_{h\perp} d\sigma} = \frac{8(1-y) \sum_q e_q^2 h_1^{\perp(1)}(x) z^2 H_1^{\perp(1)}(z)}{(1+(1-y)^2) \sum_q e_q^2 f_1(x) D_1(z)}$$



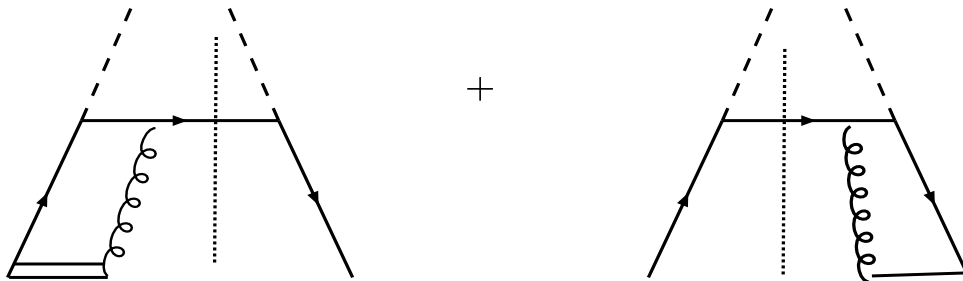
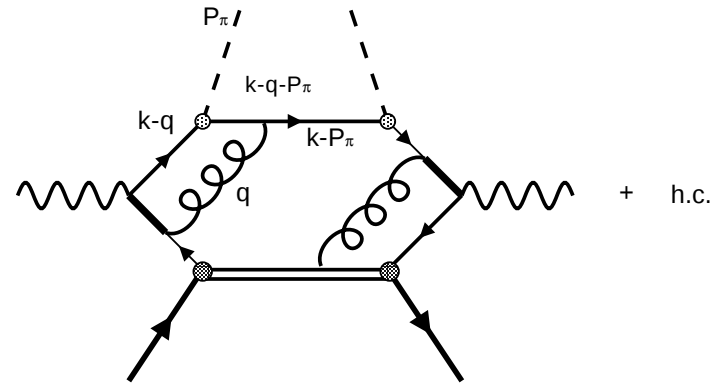
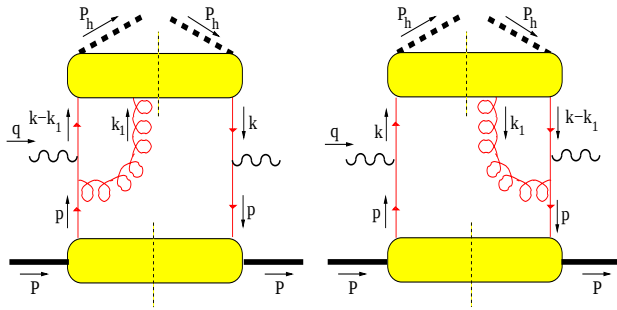


Gauge Link Contribution to Collins Function

Metz: PBL 2002, L.G., Goldstein, Oganessyan PRD: 2003: Bacchetta, Metz, Jang: PLB: 2003, Amrath, Bacchetta, Metz: PRD

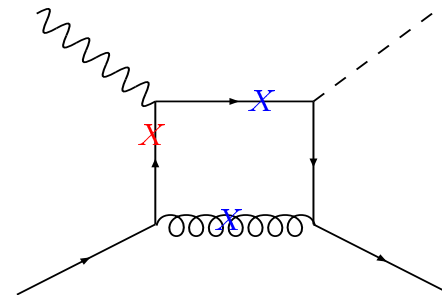
2005, L.G., Goldstein in progress

$$\Delta^{[\sigma^{\perp-}\gamma_5]}(z, k_{\perp}) = \frac{1}{4z} \int dk^+ \text{Tr}(\gamma^- \gamma^{\perp} \gamma_5 \Delta) \Big|_{k^- = P_{\pi}^- / z} \quad \text{Boer, Pijlman, Muters: NPB 2003}$$



$\frac{i}{\ell^- - i\epsilon}$ Eikonal Feynman Rules Collins '82'

$\frac{-i}{\ell^- + i\epsilon}$

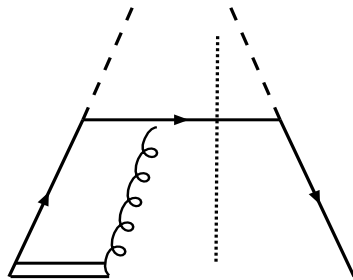


On Issues of Process Dependence: Gauge Link Contribution to Fragmentation Function

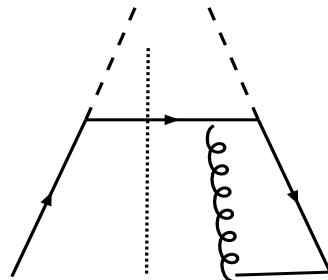
L.G., Goldstein, Oganessyan PRD: 2003: Bacchetta, Metz, Jang: PLB: 2003, Amrath, Bacchetta, Metz: PRD 2005,

L.G., G. Goldstein in progress & Como Proceedings 2006

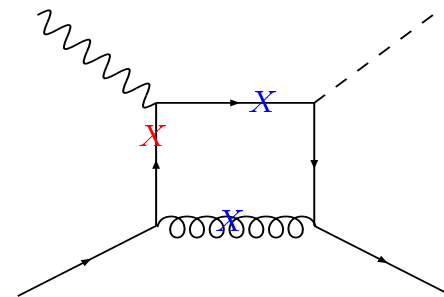
- Boer Piljman and Mulders NPB 2003: Two contributions to the Collins function.
 - ★ Gluonic Poles $" \pm B''$
 - ★ FSI $"A"$
- Does the eikonal pole contribution survive in the "T-Odd" fragmentation function Correlator?
Off shell $\gamma + q \rightarrow \pi + q'$? Are these the same "processes"?



$\frac{i}{\ell^- - i\epsilon}$ Eikonal Feynman Rules Collins '82



$\frac{-i}{\ell^- + i\epsilon}$



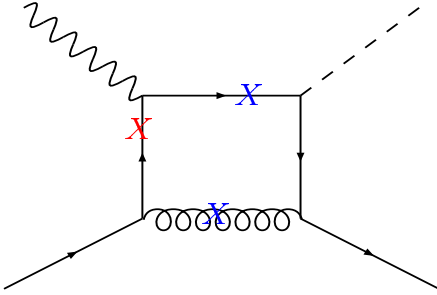
- We explored Pole Structure of correlator
 - ★ Use Cauchy's theorem to evaluate the Color Gauge invariant Correlator $\Delta^{[\sigma^{\perp} - \gamma_5]}(z, k_{\perp})$
- Analysis of pole structure in ℓ^+ indicates a *singular behavior in loop integral-looks like a "lightcone divergence"*: $\delta(\ell^-)\theta(\ell^-)$
- ★ Regulate it keep n off light cone

$$\frac{1}{n \cdot \ell \pm i\epsilon} \quad \dots$$

$n = (n^-, n^+, 0)$ (see Collins Soper NPB 1982 Ji, Yuan, Ma PLB: 2004)

- ★ Pick up poles contributions in both channels
 - On Fragmenting quark and gluon $\Rightarrow$ equivalent to cut in S -channel
 - On Eikonal and Spectator $\Rightarrow$ equivalent to cut in t -channel
- ★ This may not survive scrutiny *implying* "T-odd" Fragmentation Function universal between e^+e^- and SIDIS

S-Channel Cut-COMO Proceedings 2006



$$H_1^\perp(z, k_\perp) = \mathcal{N}'' \alpha_s \frac{M_\pi}{4z} (1-z) \frac{\mathcal{I}_1(z, P_\perp^2) + \mathcal{I}_2(z, P_\perp^2)}{\Lambda'(P_\perp^2) P_\perp^2},$$

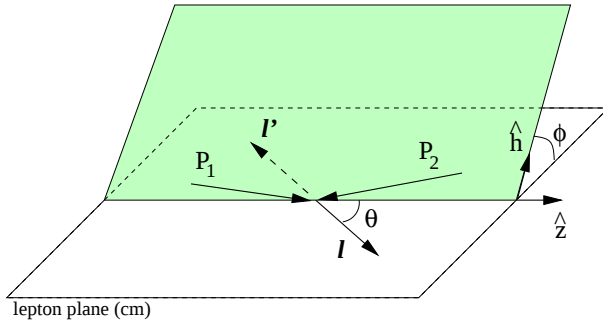
where

$$\mathcal{I}_1 = \pi(\mu - m(1-z)) \frac{E_\pi + P \cos \theta}{P + E_\pi \cos \theta} \left[\ln \frac{(P + E_\pi \cos \theta)^2}{\mu^2} - \cos \theta \ln \frac{4P^2}{\mu^2} \right]$$

$$\mathcal{I}_2 = \pi z m \frac{P \sin^2 \theta}{E_\pi - P \cos \theta} \ln \frac{4P^2}{\mu^2},$$

$P \equiv |\mathbf{P}_h|$ and $P_\perp^2 = k_\perp^2/z^2$. As in the case of the “gluonic pole” contribution, this survives the limit that incoming quark mass $m \rightarrow 0$. Both results depend the non-perturbative correlator mass μ .

Boer-Mulders Effect in Unpolarized DRELL YAN $\cos 2\phi$ (GSI & JPARC)



$$\bar{p} + p \rightarrow \mu^- \mu^+ + X$$

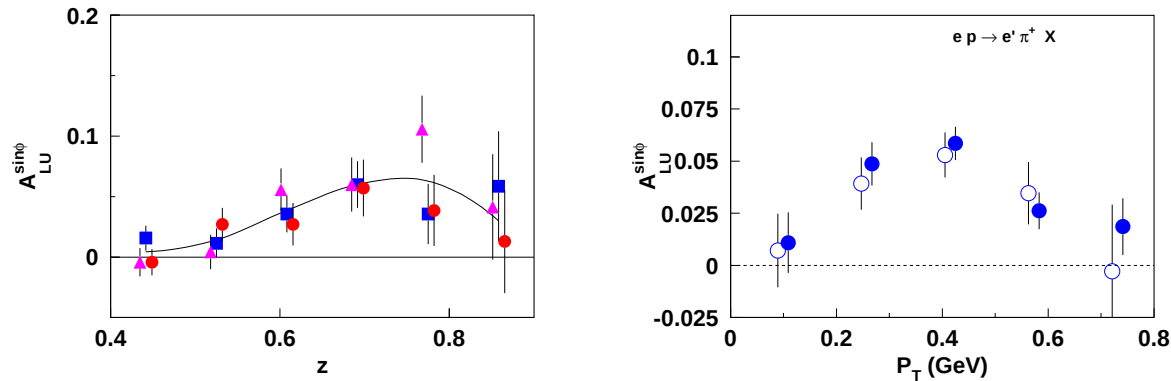
See Talk of Gary Goldstein-Thursday

SSAs& T -odd Contribution in Drell Yan (GSI & JPARC)

$$\frac{d\Delta\sigma^\uparrow}{d\Omega dx_1 dx_2 d\mathbf{q}_T} \propto \sum_a e_f^2 |\mathbf{S}_{2T}| \left\{ -B(y) \sin(\phi + \phi_{S_2}) F \left[\hat{\mathbf{h}} \cdot \mathbf{p}_{1T} \frac{\bar{h}_1^{\perp a} h_1^a}{M_1} \right] \right. \\ \left. + A(y) \sin(\phi - \phi_{S_2}) F \left[\hat{\mathbf{h}} \cdot \mathbf{p}_{2T} \frac{\bar{f}_1^a f_{1T}^{\perp a}}{M_2} \right] \dots \right\} ,$$

Beam Spin Asymmetry HERMES and CLAS $\sin \phi \rightarrow g^\perp$?!

PRD-2004 CLAS Afanasev & Carlson hep-ph/0308163, hep-ph/0603269 Bacchetta et al. PRD 2004, Metz Schlegel EJPA 2004



σ_{LU} specify the beam and target polarizations, respectively azimuthal angle ϕ is defined by a triple product:

$$\sin \phi = \frac{[\vec{k}_1 \times \vec{k}_2] \cdot \vec{P}_\perp}{|\vec{k}_1 \times \vec{k}_2| |\vec{P}_\perp|} \sim \mathbf{S}_{e^-} \cdot (\mathbf{q}_\gamma \times \mathbf{p})$$

Factorization at twist-3 is questionable L.G., Hwang, Metz, Schlegel hep-ph/0604022

$$\frac{d\sigma_{LU}}{dx_B dy dz_h d^2 P_\perp} \propto \lambda_e \sqrt{y^2 + \gamma^2} \sqrt{1 - y - \frac{1}{4} \gamma^2} \sin \phi \mathcal{H}'_{LT} \rightarrow \frac{g^\perp}{f_1} ?$$

SUMMARY

- Going beyond the collinear approximation in PQCD recent progress has been achieved characterizing transverse SSA and azimuthal asymmetries in terms of absorptive scattering.
- Central to this understanding is the role that transversity properties of quarks and hadrons assume in terms of correlations between transverse momentum and transverse spin in QCD hard scattering.
- These asymmetries provide a window to explore novel quark distribution and fragmentation functions which constitute essential information about the spin, transversity and generalized momentum structure of hadrons.
- Along with the chiral odd transversity T -even distribution function, existence of T -odd distribution and fragmentation functions can provide an explanation for the substantial asymmetries that have been observed in inclusive and semi-inclusive scattering reactions.
- We should consider the angular correlations in SDIS at 12 GeV for $\cos 2\phi$ from the standpoint of “rescattering” mechanism which generate T -odd, intrinsic transverse momentum, $k_{\perp}$, dependent *distribution and fragmentation* functions at leading twist
- Addressing issues of universality of Collins Function in spectator framework
- ★ Azimuthal asymmetries in Drell Yan and SSA measured at HERMES and COMPASS, JLAB, Belle, GSI-PAX, JPARC *may* reveal the extent to which these leading twist T -odd effects are generating the data