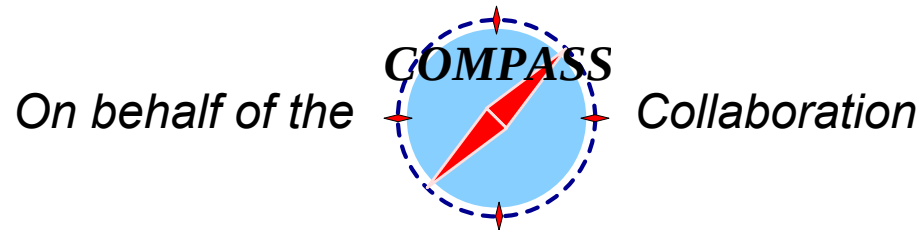


Exclusive ρ^0 production at COMPASS

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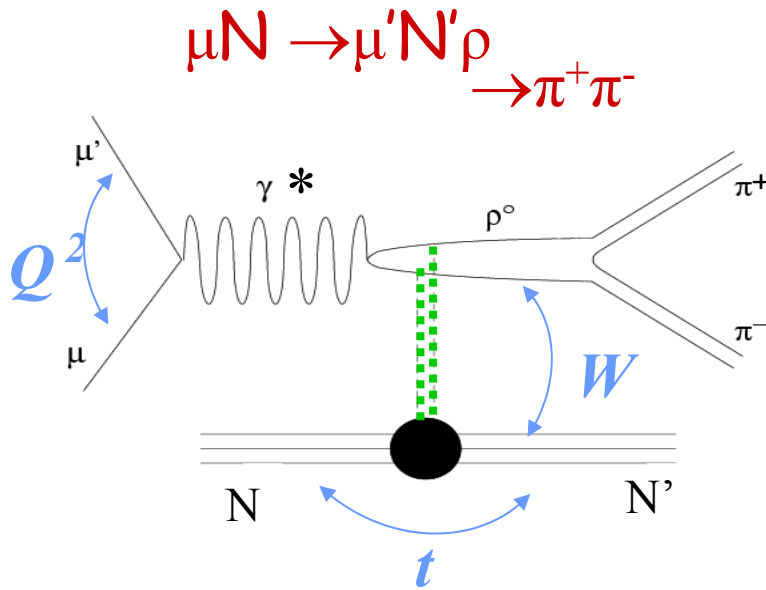
- longitudinal double-spin asymmetry A_1^ρ
- ρ^0 spin density matrix elements and $R = \sigma_L / \sigma_T$

QCD-N'06 Workshop



Villa Mondragone, Rome June 12 - 16, 2006

Physics of exclusive ρ^0 production



pQCD calculations or pQCD-inspired models with exchange of 2 quarks or 2 gluons (at large Q^2)

Diffraction

Regge theory:

at low energy ($W < 5$ GeV)

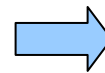
exchange of Reggeons ρ , ω , a_2 , f_2

at higher energies exchange of Pomeron

aim of the present analysis:

spin structure of cross section / helicity amplitudes for $\gamma^* N \rightarrow \rho^0 N$

- natural/unnatural parity of exchanged object
- is γ^* helicity retained by ρ^0 (SCHC) ?
- $R = \sigma_L / \sigma_T$ vs. Q^2



better understanding of Pomeron's nature

Properties of helicity amplitudes (1)

$$T_{\lambda_\rho \lambda_{N'}; \lambda_\gamma \lambda_N} = \langle \lambda_\rho \lambda_{N'} | j_{\lambda_\gamma} | \lambda_N \rangle$$

Wick and Jacob

assume $\gamma^* N \rightarrow \rho^0 N$ with t-channel exchanges of natural and unnatural parities

$$T_{\lambda_\rho \lambda_{N'}, \lambda_\gamma \lambda_N} = T_{\lambda_\rho \lambda_{N'}, \lambda_\gamma \lambda_N}^N + T_{\lambda_\rho \lambda_{N'}, \lambda_\gamma \lambda_N}^U$$

$$T_{-\lambda_\rho \lambda_{N'}, -\lambda_\gamma \lambda_N}^N = (+1) \cdot (-1)^{\lambda_\rho - \lambda_\gamma} T_{\lambda_\rho \lambda_{N'}, \lambda_\gamma \lambda_N}^N$$

natural: P, ρ, ω, a_2, f

$$T_{-\lambda_\rho \lambda_{N'}, -\lambda_\gamma \lambda_N}^U = (-1) \cdot (-1)^{\lambda_\rho - \lambda_\gamma} T_{\lambda_\rho \lambda_{N'}, \lambda_\gamma \lambda_N}^U$$

unnatural: π, a_1

How to assess contribution of unnatural exchanges ?

=> use longitudinal double spin asymmetries

$$A_1^\rho = \frac{\sigma_{1/2}^\rho - \sigma_{3/2}^\rho}{\sigma_{1/2}^\rho + \sigma_{3/2}^\rho} = \frac{\sum_{\lambda_\rho \lambda_{N'}} \{ |T_{\lambda_\rho \lambda_{N'}, ++}|^2 - |T_{\lambda_\rho \lambda_{N'}, -+}|^2 \}}{\sum_{\lambda_\rho \lambda_{N'}} \{ |T_{\lambda_\rho \lambda_{N'}, ++}|^2 + |T_{\lambda_\rho \lambda_{N'}, -+}|^2 \}}$$

$$A_1^\rho = \frac{\sum_{\lambda_\rho \lambda_{N'}} 2\text{Re} \{ T_{\lambda_\rho \lambda_{N'}, ++}^N \cdot T_{\lambda_\rho \lambda_{N'}, ++}^{*U} \}}{\sum_{\lambda_\rho \lambda_{N'}} \{ |T_{\lambda_\rho \lambda_{N'}, ++}^N|^2 + |T_{\lambda_\rho \lambda_{N'}, ++}^U|^2 \}}$$



$$\frac{2\text{Re} \{ T_{++,++}^N \cdot T_{++,++}^{*U} \}}{|T_{++,++}^N|^2 + |T_{++,++}^U|^2}$$

if SCHC

Properties of helicity amplitudes (2)

if a) nucleon unpolarized (or average over nucleon spins ≈ 0)

b) assume only natural parity exchanges contribute



only 5 independent helicity amplitudes $T_{\lambda_V \lambda_Y}$

helicity conserving T_{00}, T_{11} , single helicity-flip T_{10}, T_{01} , double helicity-flip T_{1-1}

hierarchy of helicity amplitudes

$$T_{00}, T_{11} > T_{01} > T_{1-1} > T_{10}$$

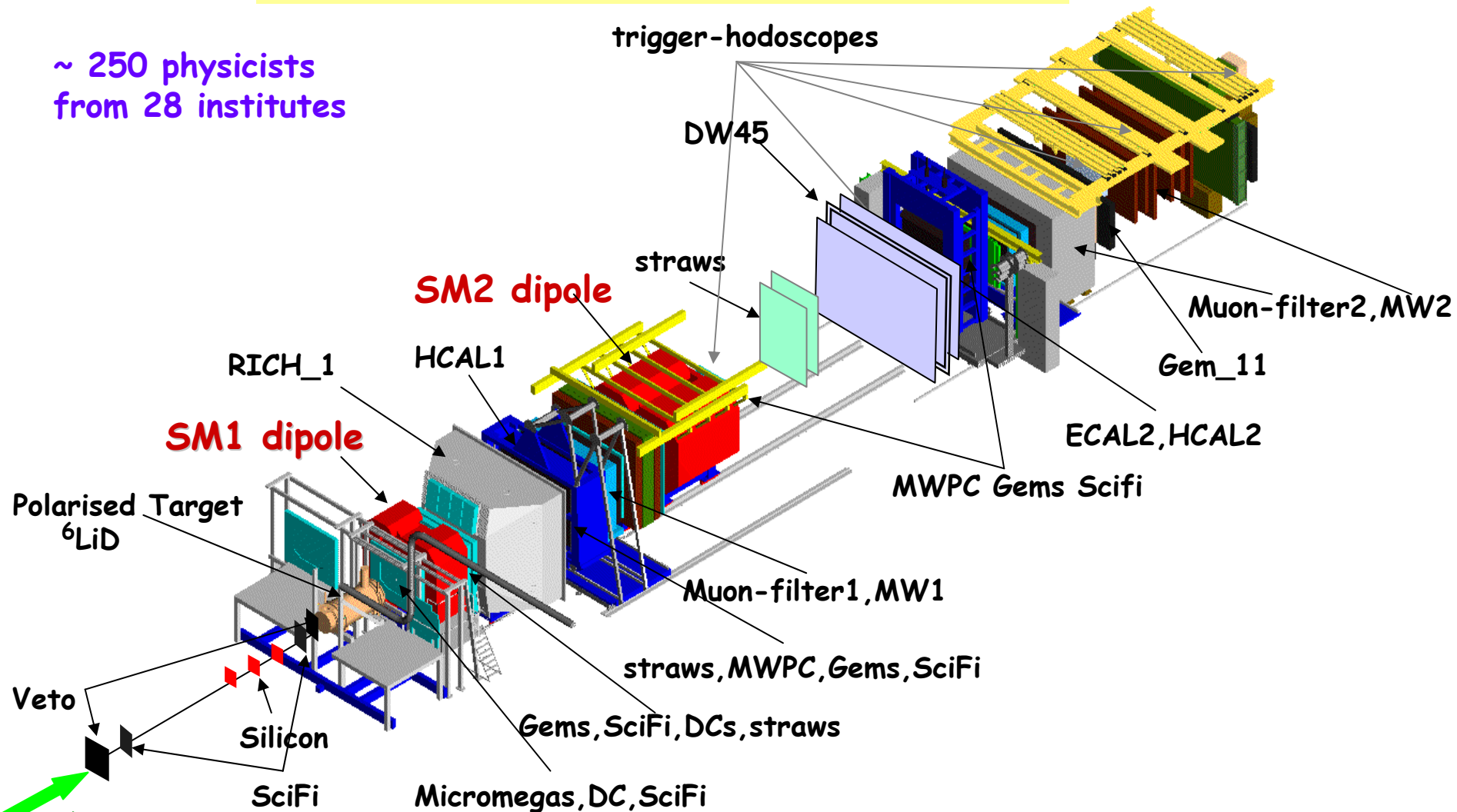
due to the angular momentum conservation
all helicity-flip amplitudes vanish at $t = 0$

st small t

$$\begin{array}{cc} T_{01}, T_{10} & T_{1-1} \\ \sim \sqrt{|t|} & \sim t \end{array}$$

COMPASS experiment

~ 250 physicists
from 28 institutes

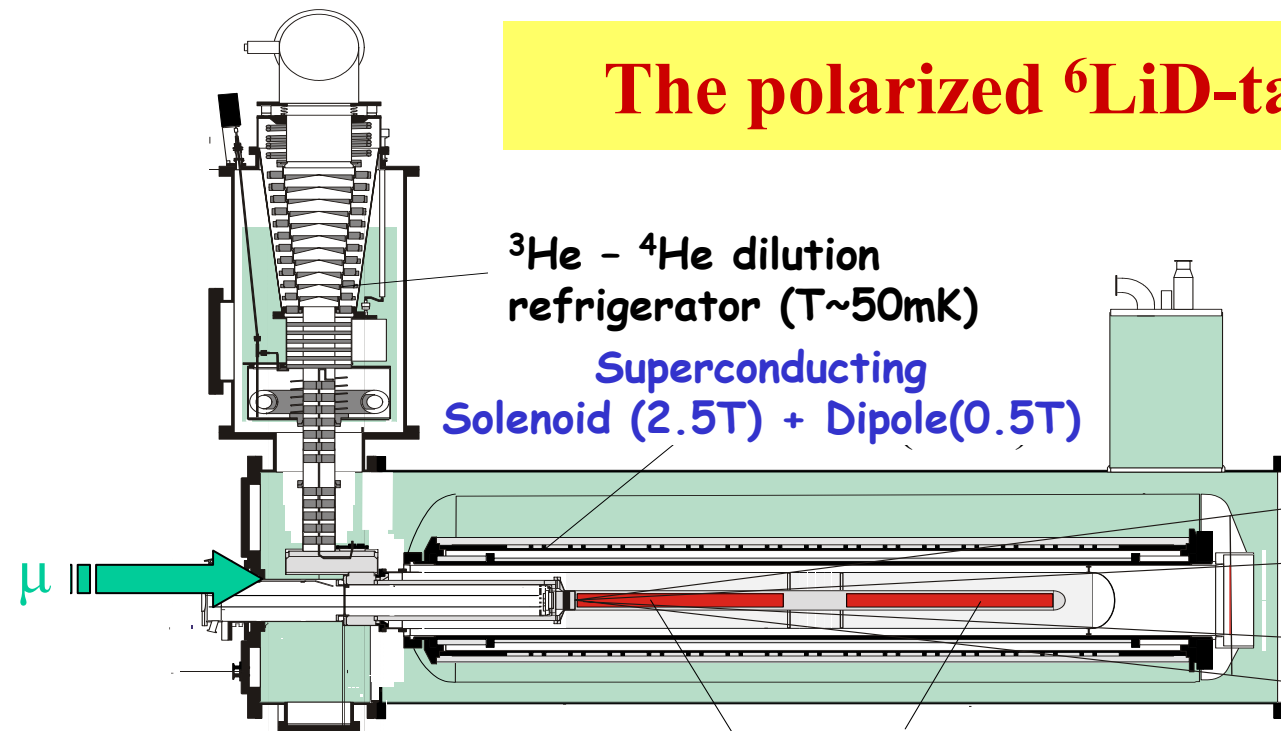


μ^- 160 GeV
from CERN SPS
Beam Polarization ~ -76%

Beam intensity
Luminosity

$2 \cdot 10^8 \mu^+/\text{spill}$ (4.8 s / 16.2 s)
 $\sim 5 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$

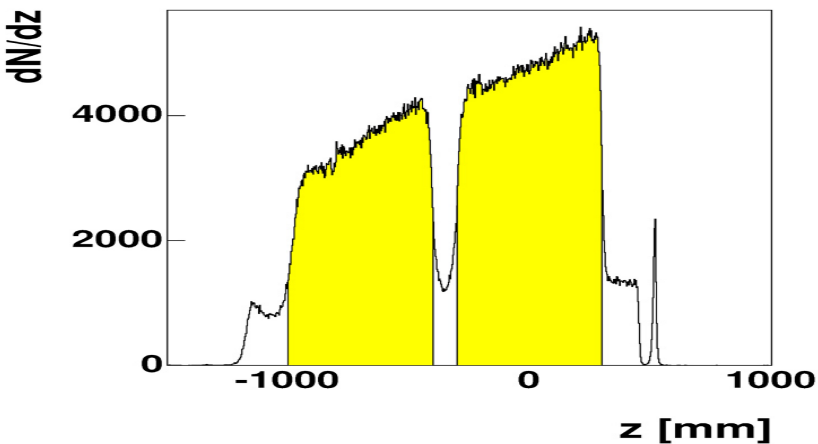
The polarized ${}^6\text{LiD}$ -target



Target Polarization
 $\sim 50\%$

Dilution factor
 $f \sim 0.36$

Two 60cm long target cells with opposite polarizations



4 possible spin combinations:

longitudinal

① $\rightarrow \leftarrow$

② $\leftarrow \rightarrow$

Reversed every 8 hours

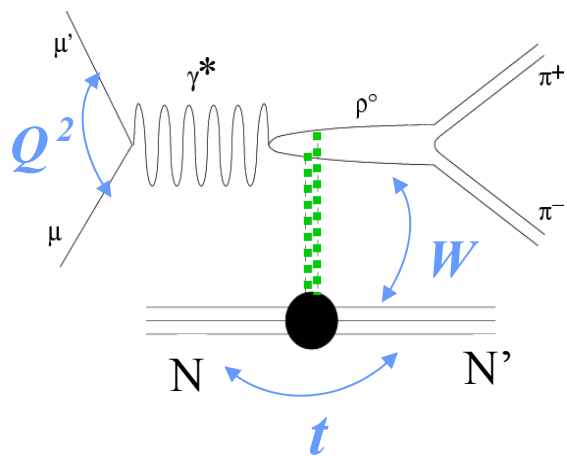
transverse

③ $\uparrow \downarrow$

④ $\downarrow \uparrow$

Reversed once a week

Incoherent exclusive ρ^0 production



^6LiD polarized target

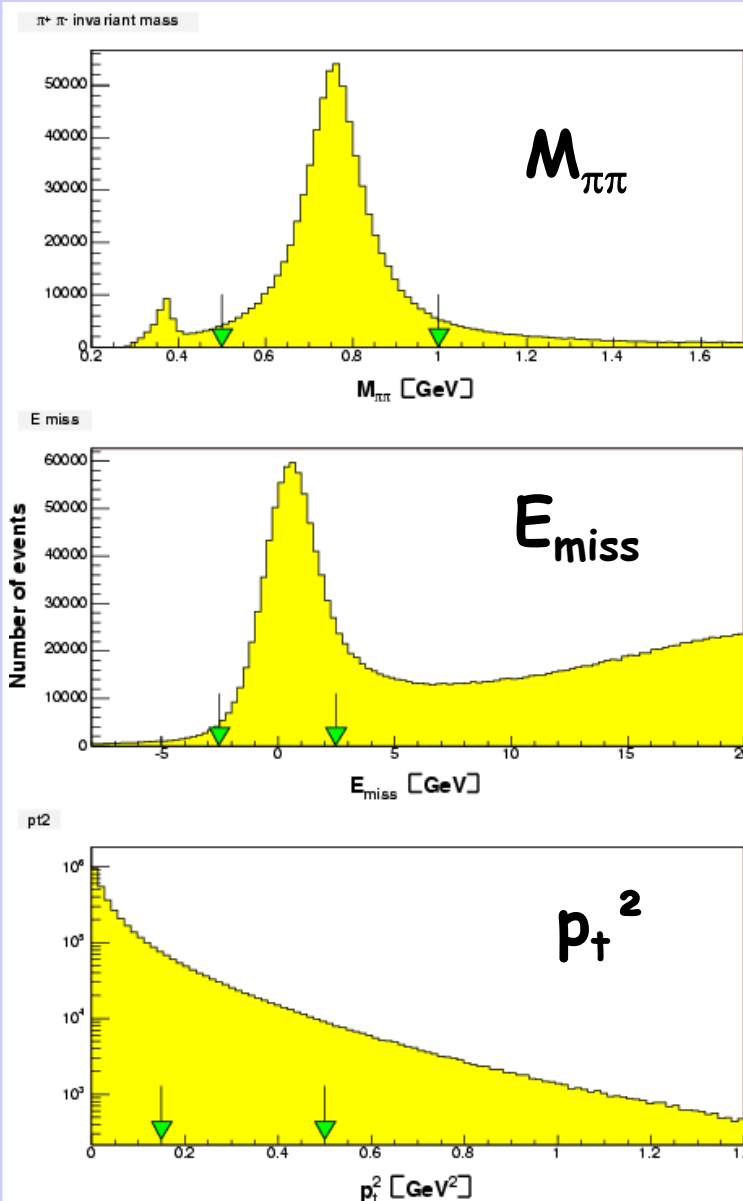
Kinematics:

$$\nu > 30 \text{ GeV}$$

$$E_{\mu'} > 20 \text{ GeV}$$

$$Q^2 > 0.01 \text{ GeV}^2$$

(Q^2 cut applied only
in SDME analysis)



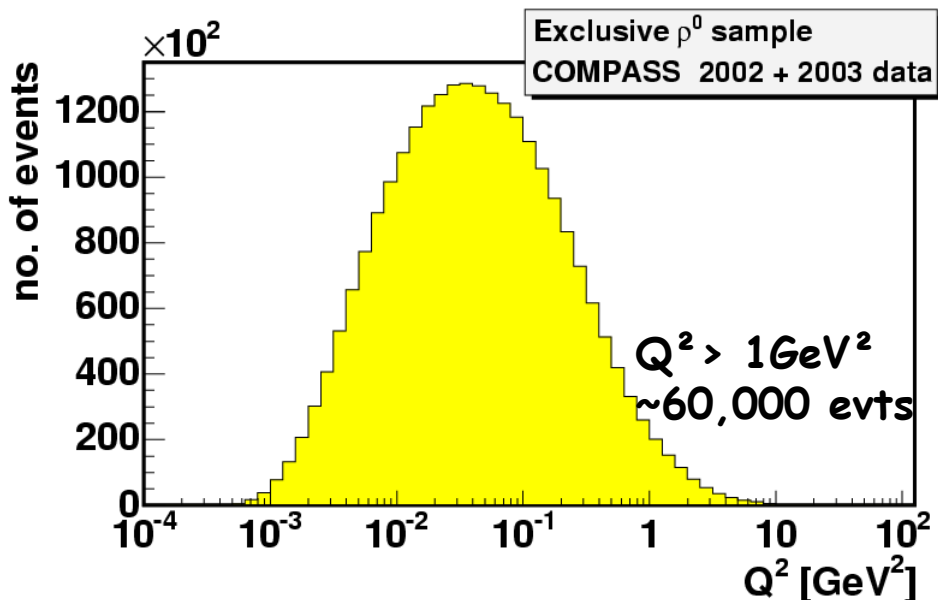
Assuming
both hadrons are π
 $0.5 < M_{\pi\pi} < 1 \text{ GeV}$

Exclusivity of
the reaction
 $E_{\text{miss}} = (M_X^2 - M_N^2) / 2M_N$
 $-2.5 < E_{\text{miss}} < 2.5 \text{ GeV}$

Incoherent production
 $0.15 < p_t^2 < 0.5 \text{ GeV}^2$
scattering off a
quasi-free nucleon

Background $\sim 12\%$

Kinematical domain of the final sample

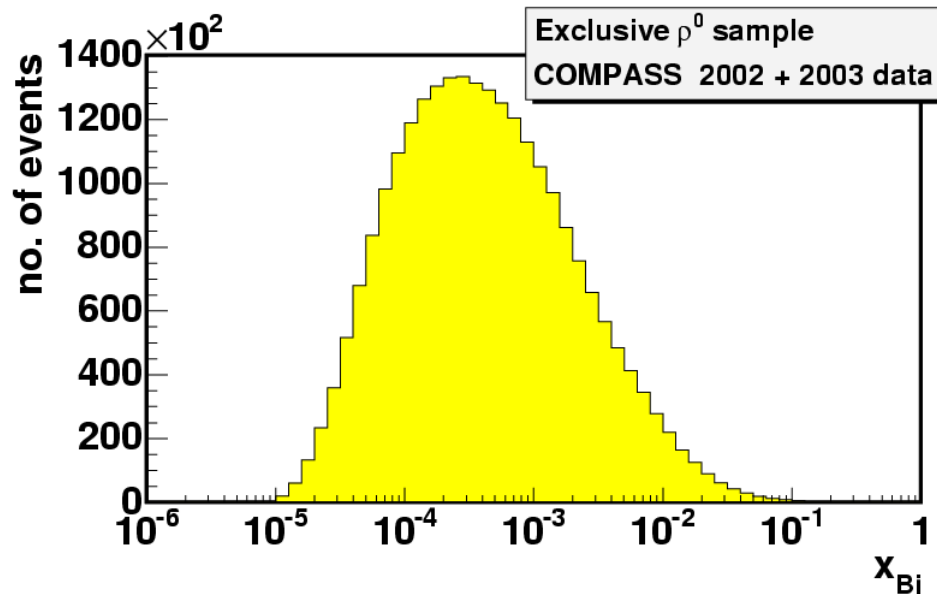
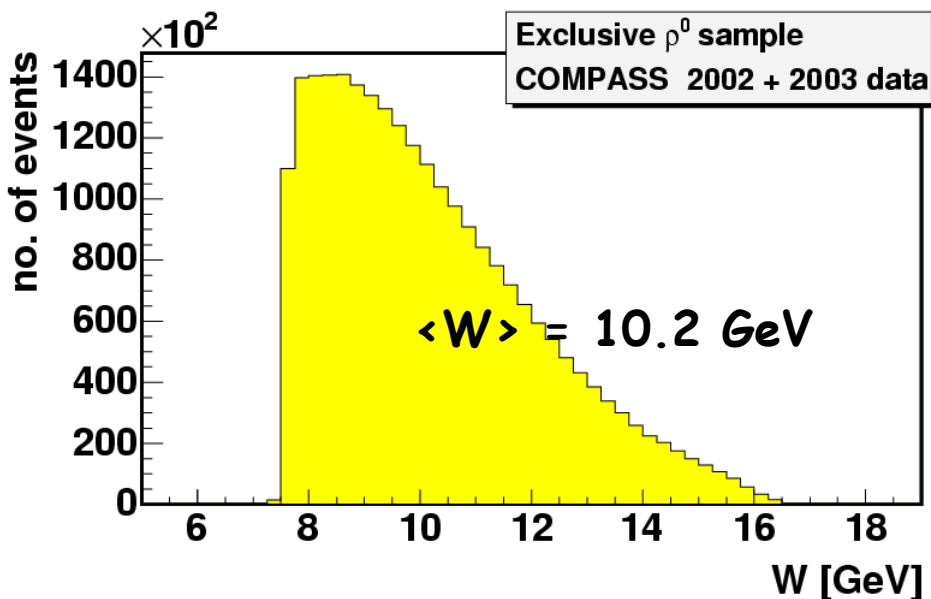


2002 : 800,000 evts

2003 : 1,600,000 evts

+ 2004 : not yet analyzed
~ will double the data sample

$$\langle p_{\text{t}}^2 \rangle = 0.27 \text{ GeV}^2$$



Longitudinal double-spin asymmetry for exclusive ρ^0 production

Both muon beam and target have to be polarized along the beam direction

in COMPASS polarized d from ${}^6\text{LiD}$

Ultimately one determines

$$A_1^\rho = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}}$$

$\sigma_{1/2}$ ($\sigma_{3/2}$) cross section for $\gamma^* N \rightarrow \rho^0 N$ for antiparallel (parallel) spin orientation of γ^* and target N

- can arise from exchange of $a_1(1260)$ trajectory in t-channel (small at 160 GeV)
- from interference of amplitudes for exchange in t-channel of Reggeons
with natural parity: $\rho, \omega, f, a_2(1320), P(\text{pomeron})$ and unnatural parity: $\pi, a_1(1260)$
(sensitivity even to small contributions of the latter ones)
- if only non-perturbative P $A_1^\rho \approx 0$
- in pQCD-inspired models possible $A_1^\rho \neq 0 \Rightarrow$ access to spin dependent GPDs
 - Ryskin
 - Goloskokov and Kroll

**Evaluation of $A_1^\rho (\gamma^* N \rightarrow \rho^0 N)$
for **incoherent exclusive ρ^0 production****

$$A_{LL}^{raw} = \frac{1}{2} \left(\frac{N_{+-}^u - N_{++}^d}{N_{+-}^u + N_{++}^d} - \frac{N_{++}^{u'} - N_{+-}^{d'}}{N_{++}^{u'} + N_{+-}^{d'}} \right)$$

$$A_{LL}(\mu N \rightarrow \mu N \rho^0) = \frac{\sigma(\mu N)_{\uparrow\downarrow} - \sigma(\mu N)_{\uparrow\uparrow}}{\sigma(\mu N)_{\uparrow\downarrow} + \sigma(\mu N)_{\uparrow\uparrow}} = \frac{1}{f} \cdot \frac{1}{P_b} \cdot \frac{1}{P_t} \cdot A_{LL}^{raw}$$

$$A_1^\rho(\gamma^* N \rightarrow \rho^0 N) \approx \frac{1}{D} A_{LL}(\mu N \rightarrow \mu N \rho^0)$$

to minimize systematic effects ‘2nd order weighted method’ used

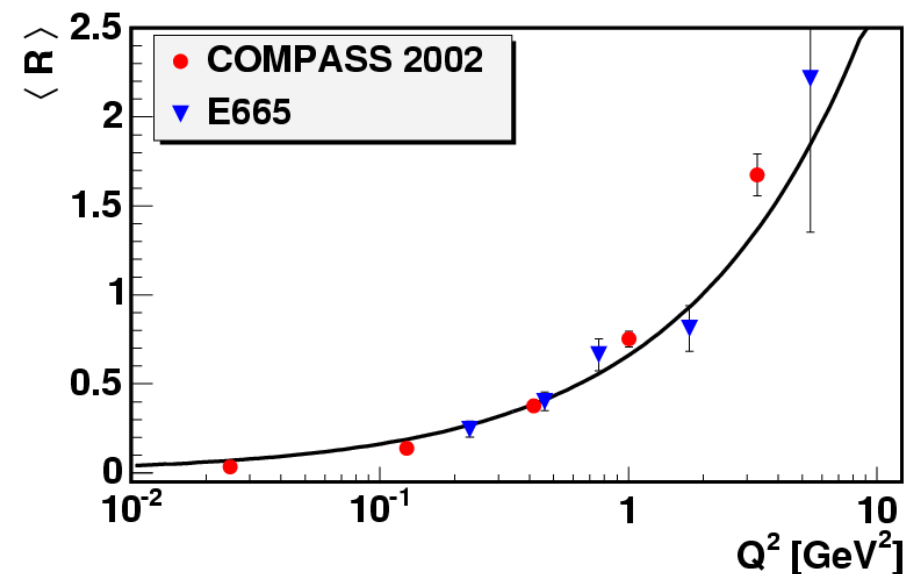
❖ each event assigned weight $w = f D P_b$ ❖ 2nd order equation for A1
with coefficients depending on weights w , P_b and acceptance ratios for u and d cells

values of **A-dependent cross sections** (for f) and **ratio R** (for D)
specific for incoherent exclusive ρ^0 production

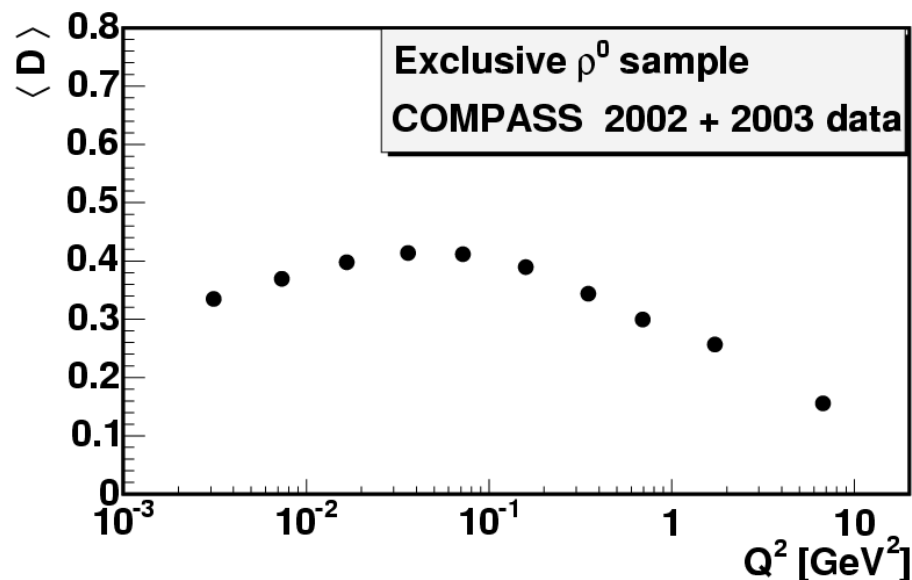
Depolarization factor **D** for incoherent exclusive ρ^0 production

$$D(y, Q^2) = \frac{y[(1 + \gamma^2 y/2)(2 - y) - 2y^2 m_\mu^2/Q^2]}{y^2(1 - 2m_\mu^2/Q^2)(1 + \gamma^2) + 2(1 + R)(1 - y - \gamma^2 y^2/4)}$$

$$\gamma^2 = Q^2 / \nu^2$$



R: E665 parametrization



Dilution factor f for incoherent exclusive ρ^0 production

$$f_0 = \frac{n_d}{n_d + \sum_A n_A (\tilde{\sigma}_A / \tilde{\sigma}_d)}$$

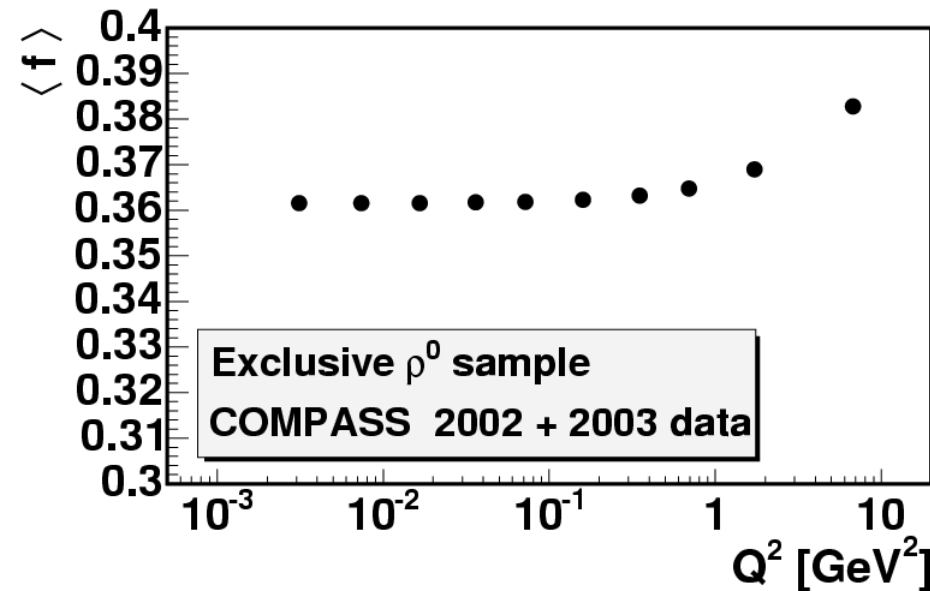
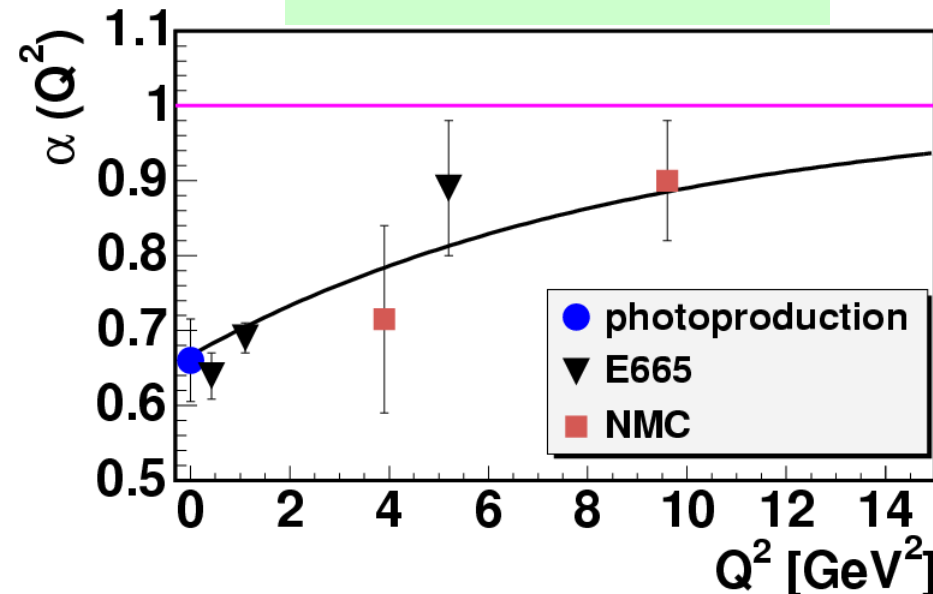
n_D, n_A : numbers of nucleons from deuteron and nucleus A in the target

$\tilde{\sigma}_A, \tilde{\sigma}_D$: cross sections **per nucleon** for **incoherent exclusive ρ^0 production** on deuteron and nucleus A

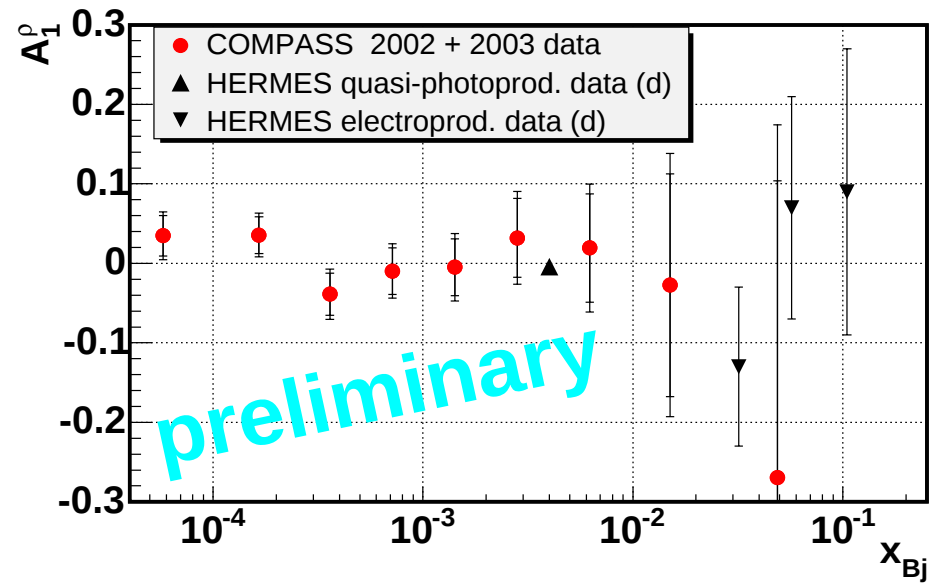
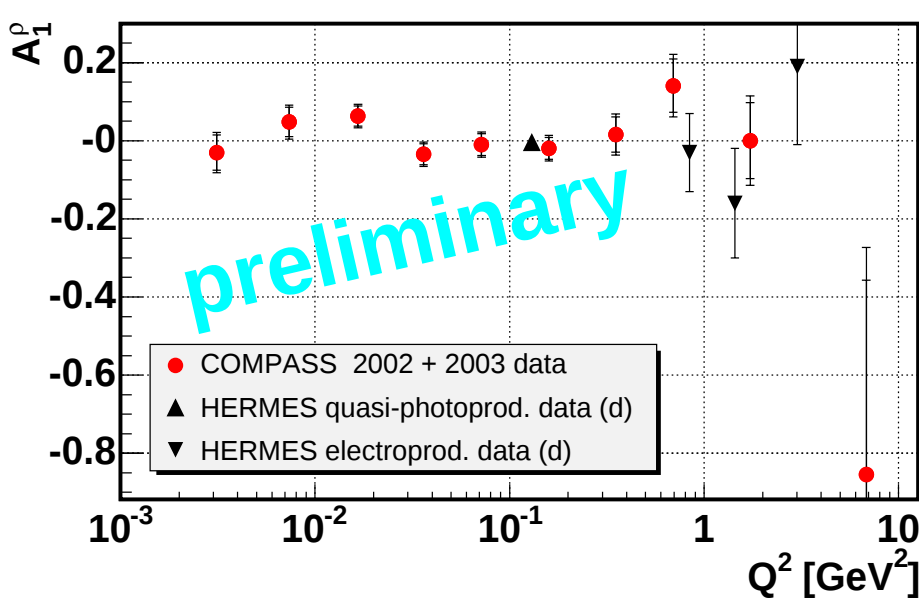
finally used $f = C_1 \cdot f_0$, where C_1 takes into account that there are 2 deuterons in ${}^6\text{LiD}$ molecule

$$\tilde{\sigma}_A = \sigma_p \cdot A^{\alpha(Q^2)-1}$$

f calculated event by event



COMPASS preliminary and HERMES results on A_1^p (d)



COMPASS results on A_1^p on polarized deuteron target consistent with 0

Extended kinematical range of COMPASS by almost 2 decades down both in Q^2 and x

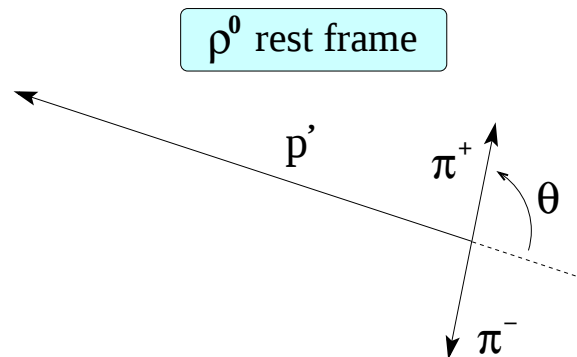
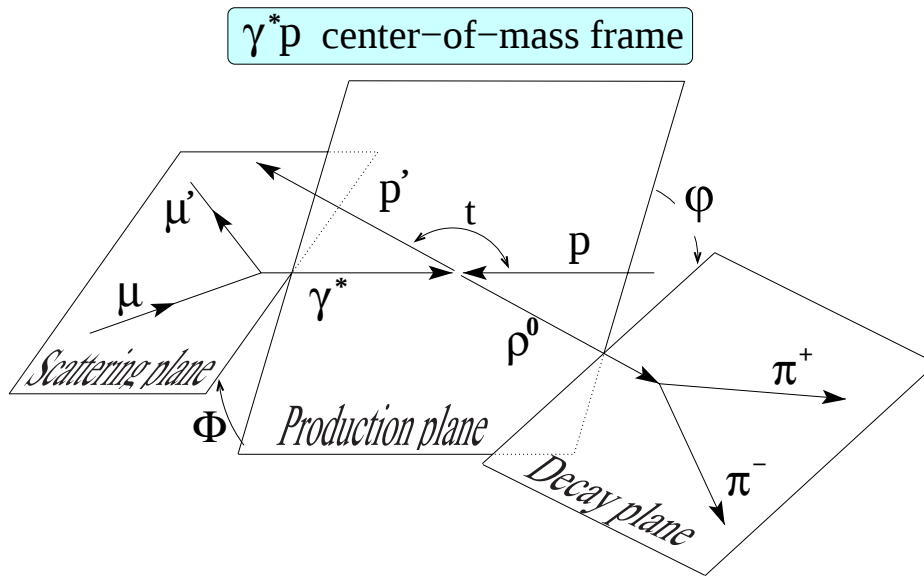
COMPASS : inner bars –stat.
outer – total errors
HERMES: total errors

➡ at small x (< 0.01) contribution of unnatural parity exchanges small
for exclusive ρ^0 production

ρ^0 angular distributions $W(\cos\theta, \varphi, \Phi)$

depend on the **spin density matrix elements (SDME)**

⇒ 23 (15) observables with polarized (unpolarized) beam



SDMEs are bilinear combinations of the helicity amplitudes

$$A(\gamma^*(\lambda_\gamma) \rightarrow \rho(\lambda_\rho)) \equiv T_{\lambda_\rho, \lambda_\gamma}$$
$$\lambda_\gamma = \pm 1, 0 \quad \lambda_\rho = \pm 1, 0$$

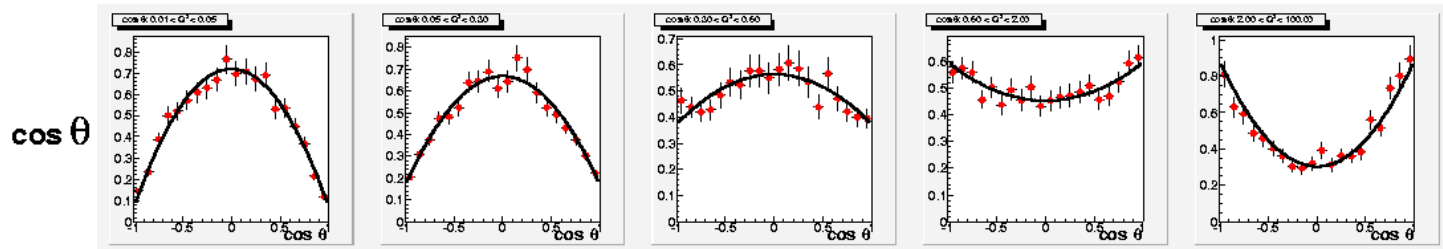
This analysis:
only one-dimensional
angular distributions

also used:

$$\psi = \varphi - \Phi$$

Measurement of r_{00}^{04}

$0.01 < Q^2 < 0.05 < Q^2 < 0.3 < Q^2 < 0.6 < Q^2 < 2.0 < Q^2 < 10 \text{ GeV}^2$



Distribution :

$$W(\cos \theta) = \frac{3}{4} [(1 - r_{00}^{04}) + (3r_{00}^{04} - 1) \cos^2 \theta]$$

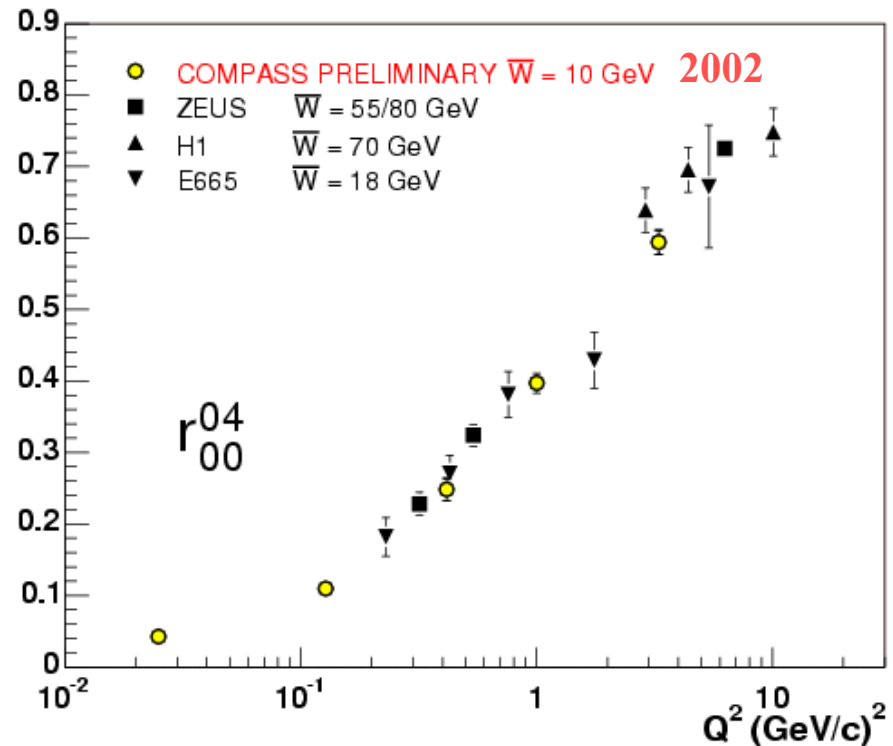
Spin density matrix element:

$$r_{00}^{04} = \frac{|T_{01}|^2 + (\varepsilon + \delta) |T_{00}|^2}{N_T (1 + (\varepsilon + \delta) R)} \xrightarrow{\text{SCHC}} \frac{\sigma_L}{\sigma_T}$$

$$R = \sigma_L / \sigma_T \quad N_T = |T_{11}|^2 + |T_{-11}|^2 + |T_{01}|^2$$

$$\Gamma_L / \Gamma_T = \varepsilon + \delta$$

$T_{\lambda p \lambda \gamma}$ helicity amplitudes
 meson photon



Determination of $R = \sigma_L/\sigma_T$

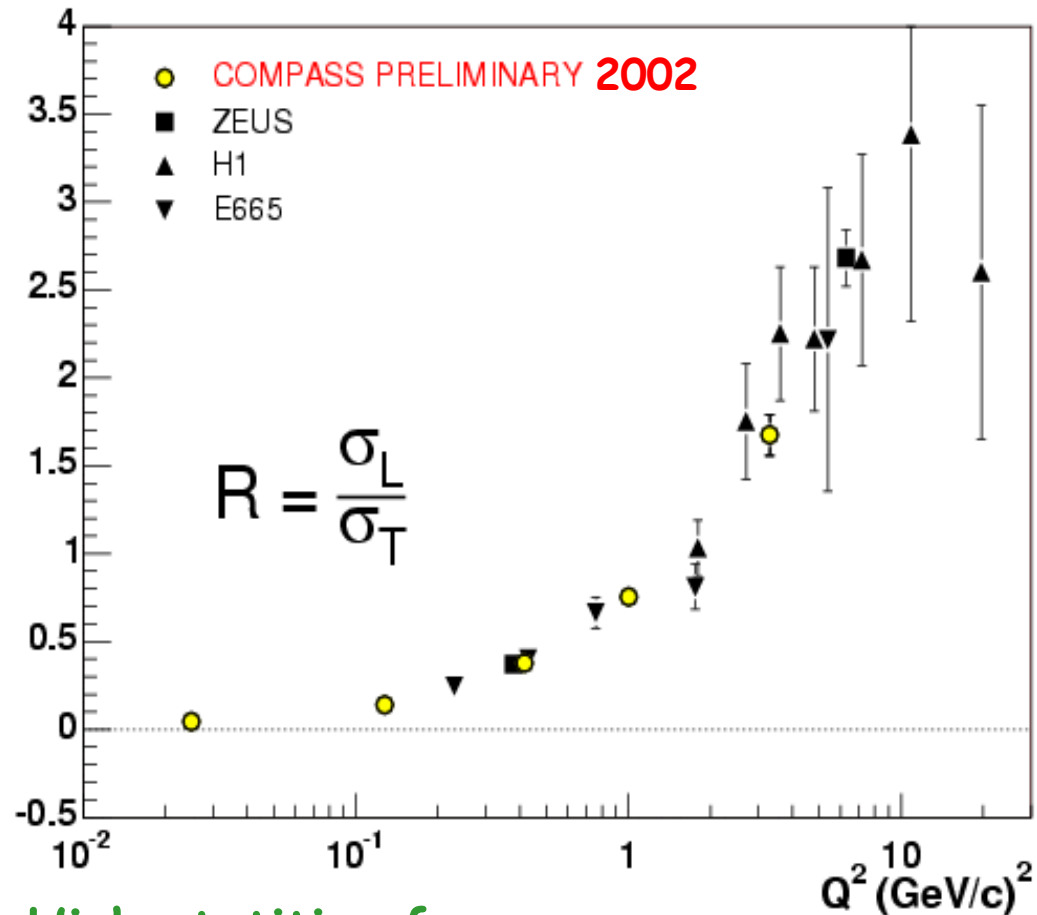
If SCHC holds :

only $T_{00} \neq 0$
 $T_{11} \neq 0$

Then :

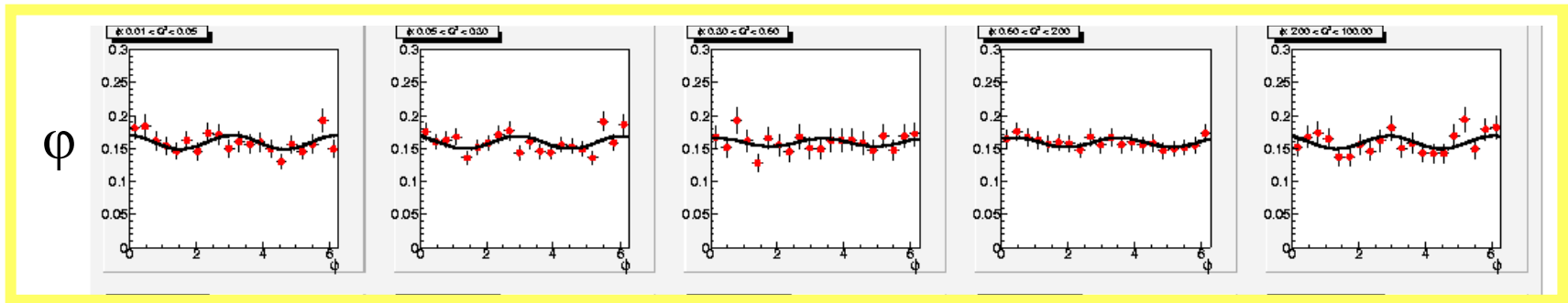
$$R = \frac{\sigma_L}{\sigma_T} = \frac{1}{(\varepsilon + \delta)} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$$

Impact on GPD study:
determination of σ_L
 σ_L is dominant at $Q^2 > 2 \text{ GeV}^2$



- High statistics from quasi-photonproduction to hard production
- Better coverage at high Q^2 with 2003 and 2004 data

Measurement of r_{1-1}^{04} and $\text{Im } r_{1-1}^3$



Distribution :

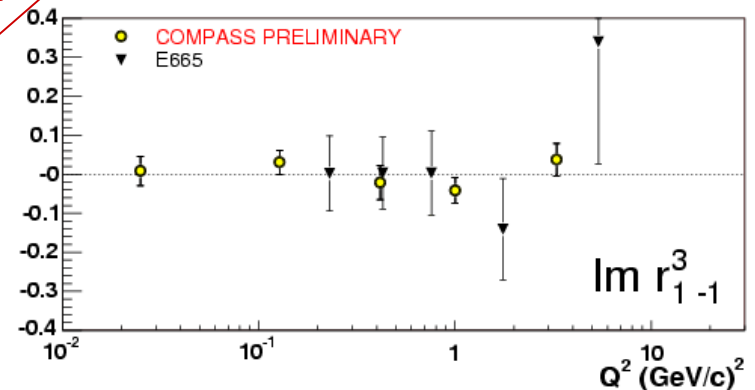
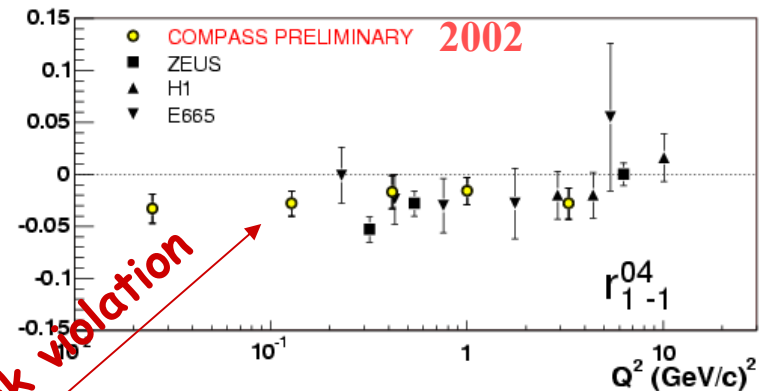
$$W(\varphi) = \frac{1}{2\pi} [1 - 2r_{1-1}^{04} \cos 2\varphi + 2\text{Im}r_{1-1}^3 P_\mu \sqrt{1-\varepsilon^2} \sin 2\varphi]$$

beam polarisation

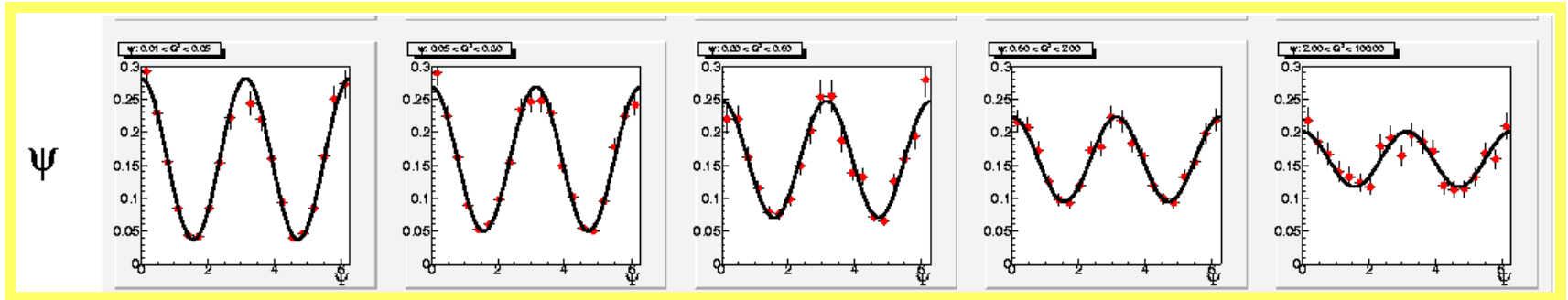
Spin density matrix elements:

$$r_{1-1}^{04} = \frac{\text{Re}(T_{11}T_{-11}^*) - (\varepsilon + \delta)|T_{10}|^2}{N_T(1 + (\varepsilon + \delta)R)} = 0$$

$$\text{Im}r_{1-1}^3 = \dots = 0 \quad \leftarrow \text{If SCHC holds}$$



Measurement of r_{1-1}^1



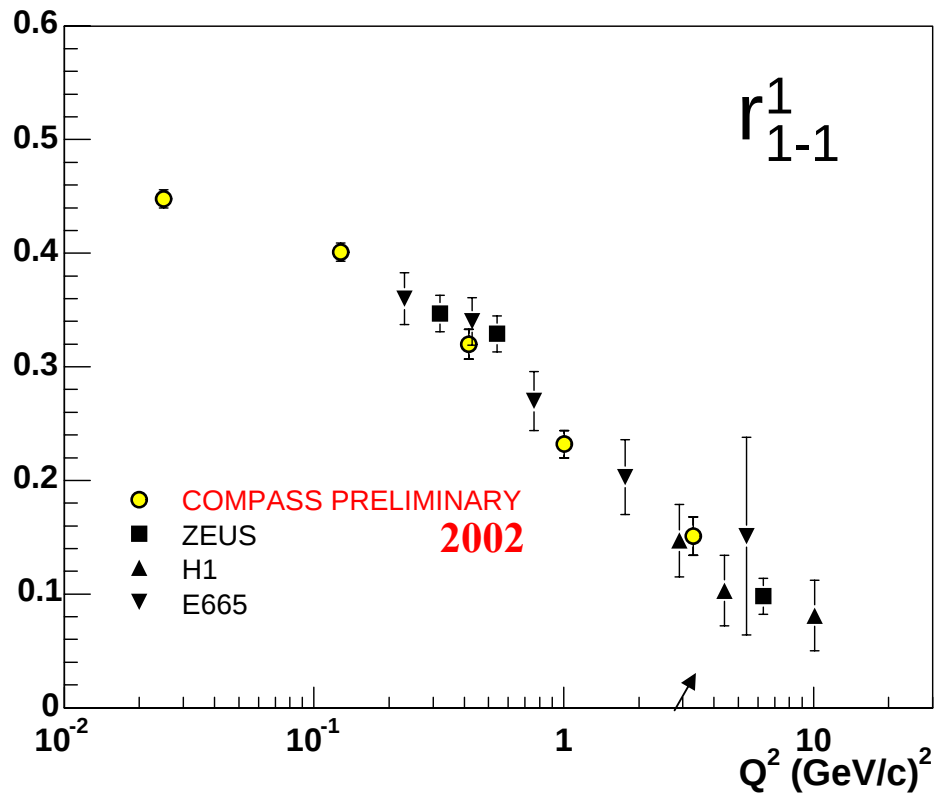
If SCHC holds

$$W(\psi) = \frac{1}{2\pi} [1 + 2r_{1-1}^1 \varepsilon \cos 2\psi]$$

general expression for r_{1-1}^1 :

$$r_{1-1}^1 = \frac{1}{2 N_T} \frac{|T_{11}|^2 + |T_{1-1}|^2}{(1 + (\varepsilon + \delta)R)}$$

in future independent of SCHC
determination of 23 SDMEs
including r_{1-1}^1 from 3D distributions



Summary and outlook

- Asymmetry A_1^p (d) consistent with zero over wide range of Q^2 and x
small contribution of unnatural parity exchanges at $x < 0.01$
first measurement at small Q^2 and small x
- High-statistics data on SDM elements and R for incoherent exclusive ρ^0 production
in a wide Q^2 range (including small Q^2 not covered previously)
weak violation of SCHC observed
- Significant improvement of accuracy expected after including 2004 (and 2003 data)
extraction of 23 SDMEs under way
- Studies of coherent exclusive ρ^0 production foreseen
- Single spin asymmetry for transversely polarized target $\rightarrow$ E/H GPDs
- Analysis of exclusive ϕ and J/ψ production in progress