

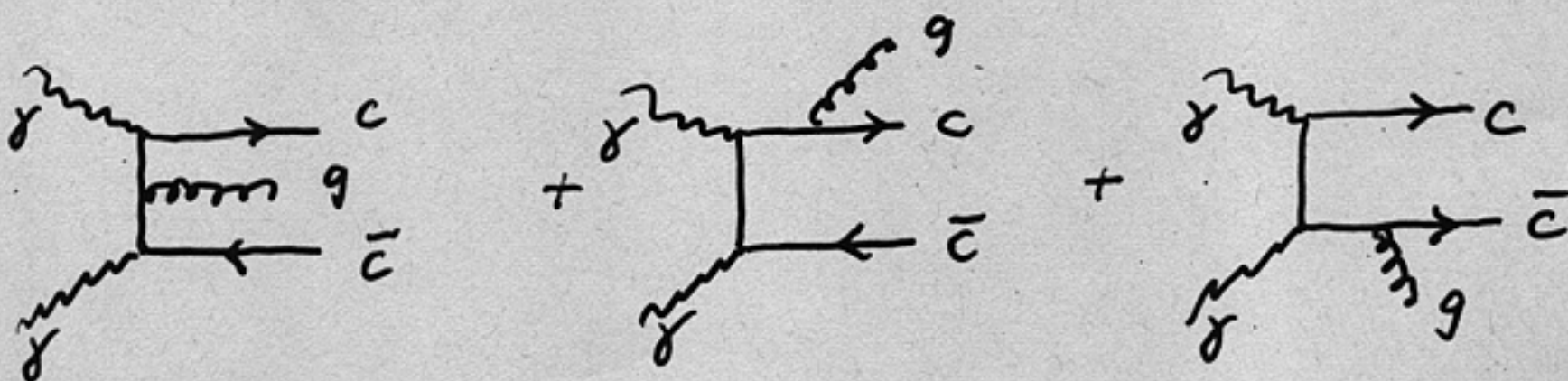
Inclusive D^* Production in Photon-Photon Collisions with Massive Quarks in NLO

Photon 2003, Frascati

Three approaches $\gamma + \gamma \rightarrow D^{*\pm} + X$

NLO: $\gamma + \gamma \rightarrow c + \bar{c} + g$

direct contribution



• (i) massless approach $m_c = 0$, $N_f = 4$

$$d\sigma(\gamma\gamma \rightarrow cX) = \frac{\alpha}{2\pi} P_{c\gamma} d\sigma(\gamma c \rightarrow cg) \ln \frac{S}{M^2}$$

M factorization scale $O(\sqrt{S})$

originates from collinear divergent piece

$$d\sigma(\gamma\gamma \rightarrow cX) = \frac{\alpha}{2\pi} \left(-\frac{1}{\epsilon}\right) \left(\frac{4\pi\mu^2}{S}\right)^\epsilon P_{c\gamma} d\sigma(\gamma c \rightarrow cg)$$

$$\rightarrow \left(\frac{4\pi\mu^2}{M^2}\right)^\epsilon \left(\frac{M^2}{S}\right)^\epsilon$$

$\rightarrow$ into NLO corr

$\rightarrow$ absorbed in f_c^γ

- (ii) massive approach $m_c \neq 0$, $N_f = 3$

$$d\sigma(\gamma\gamma \rightarrow cX) = \frac{\alpha}{2\pi} P_{c\gamma} d\sigma(\gamma c \rightarrow cg) \ln \frac{S}{m_c^2}$$

factorization scale $M = m_c$ $\rightarrow$ large

charm only in final state: $f_c^\gamma = 0$

- (iii) massive approach $m_c \neq 0$, $N_f = 4$

$$\ln \frac{S}{m_c^2} = \ln \frac{S}{M^2} + \ln \frac{M^2}{m_c^2}$$

as in massless

absorbed in f_c^γ

charm also in initial state:

single and double resolved with charm

same procedure for second and third diagram

factorization scale M_f

need fragmentation for $c \rightarrow D^*$: $D_c^{D^*}(x)$

- (i) $m_c = 0$ from start (ZM-VFN)

m_c^2/p_T^2 terms neglected

good only for large p_T^2

$\ln \frac{M^2}{\mu^2}$ terms summed in evolved f_c^γ and $D_c^{D^*}$

- (ii) $m_c \neq 0$ (FFN) $n_f = 3$

no c quarks in initial state

$f_c^\gamma = 0$, $D_c^{D^*}$ not evolved, fragmentation pert.

good for $p_T^2 \simeq m_c^2$, m_c^2/p_T^2 terms included

- (iii) $m_c \neq 0$ (Massive VFN) $n_f = 4$

c quarks in initial state

$f_c^\gamma \neq 0$, f_c^γ and $D_c^{D^*}$ evolved

$\ln \frac{M^2}{m_c^2}$ terms summed

m_c^2/p_T^2 terms included

good for $p_T^2 \gtrsim m_c^2$, where $f_c^\gamma \neq 0$, $D_c^{D^*} \neq 0$

$p_T^2 \lesssim m_c^2$ transition to FFN

massless limit of (iii) neglects m_c^2/p_T^2 terms

- (i) and $m_c \rightarrow 0$ limit of (iii) identical?

answer: no!

- (i) and (iii) differ by finite terms
do not vanish for $m_c \rightarrow 0$

- finite terms subtracted to match (i)

calculated for

direct process

single-resolved process

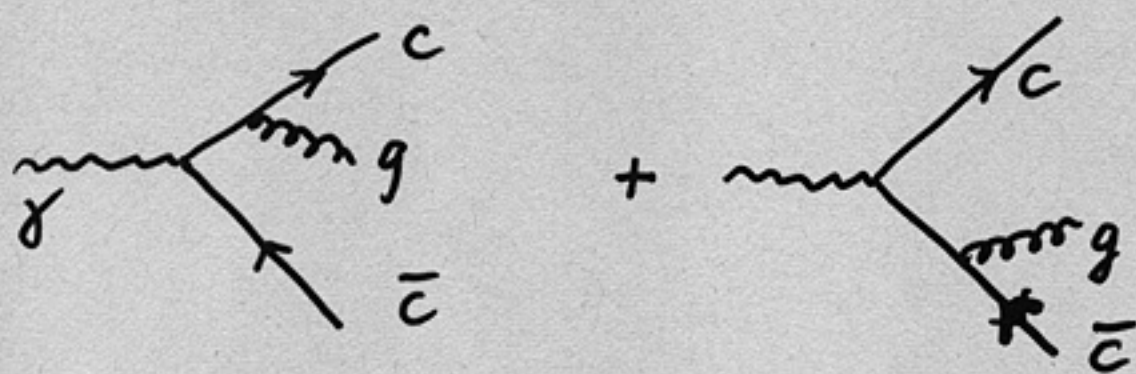
hep-ph/0109167
G.K. + H. Spiesberger
EPYC 22, 289 (2001)

hep-ph/0302081

- direct process

finite terms = LO pert. fragmentation functions

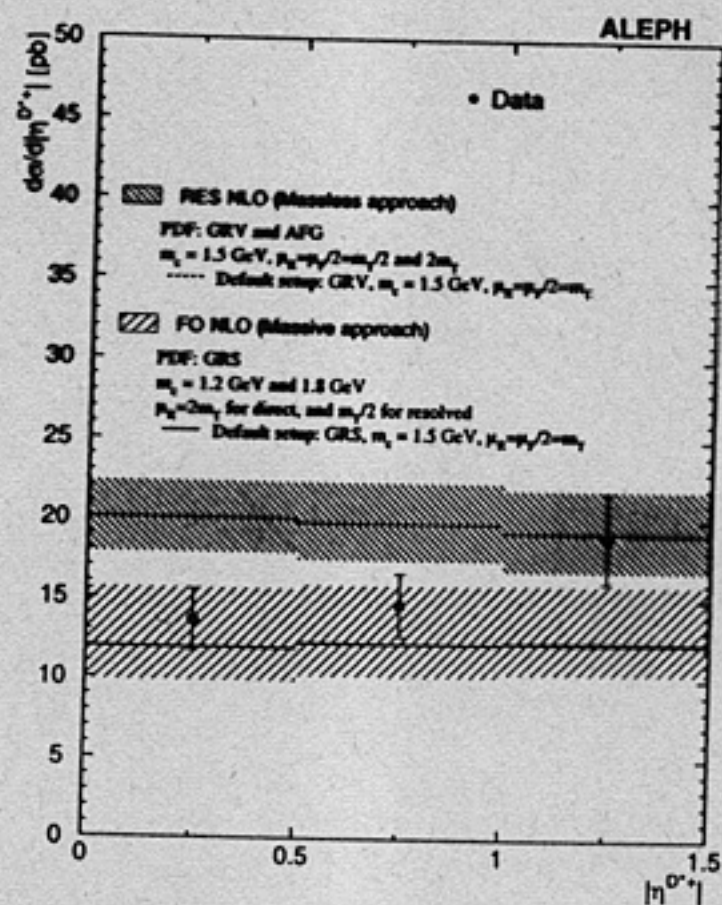
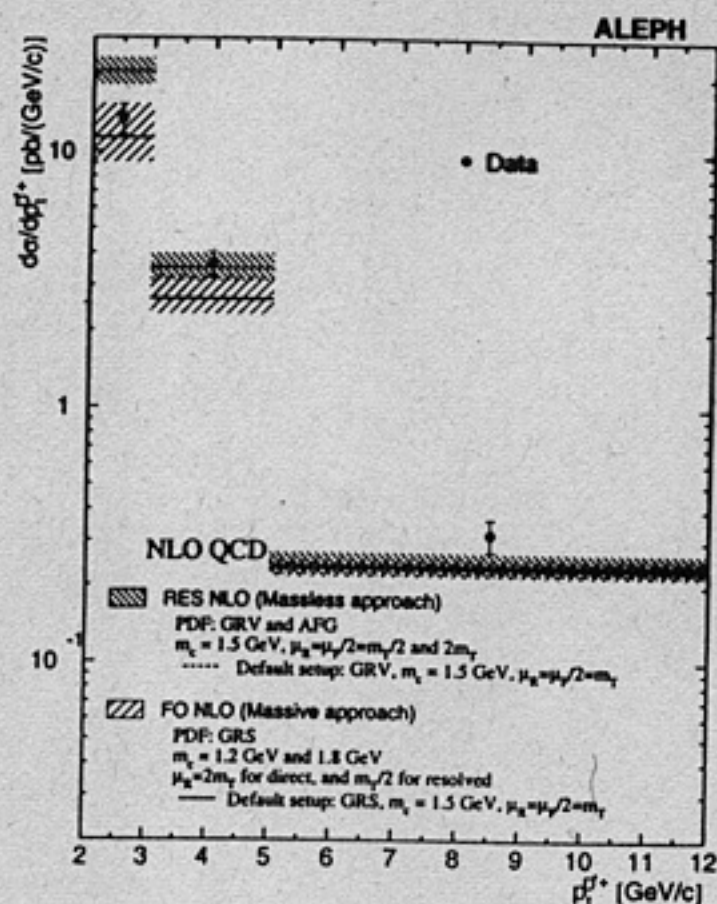
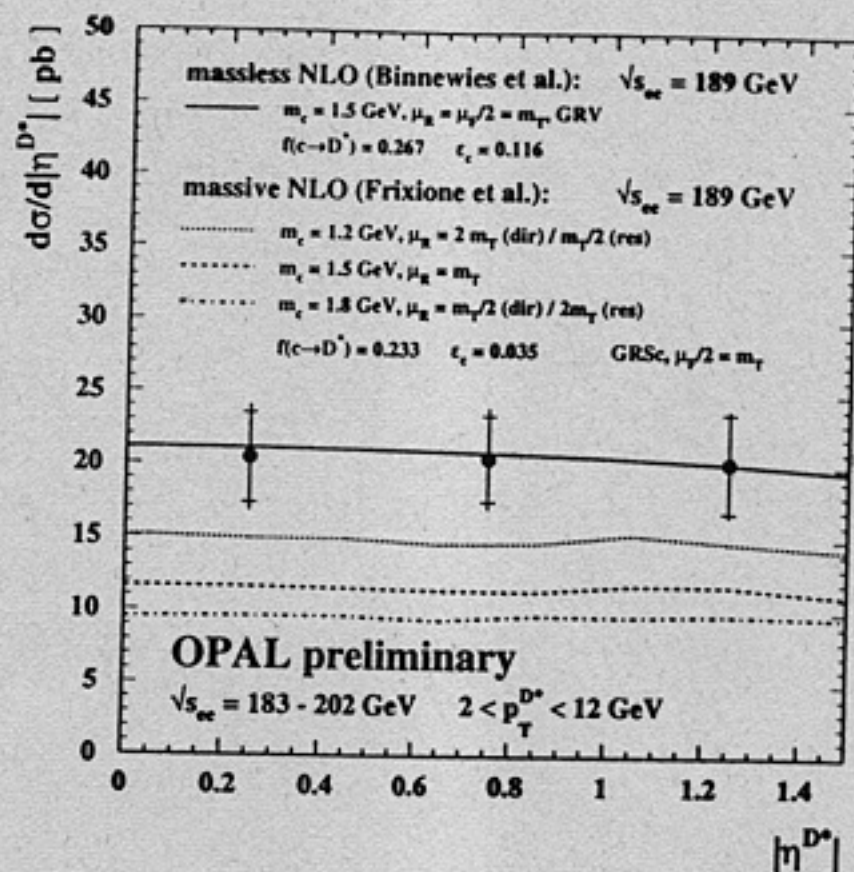
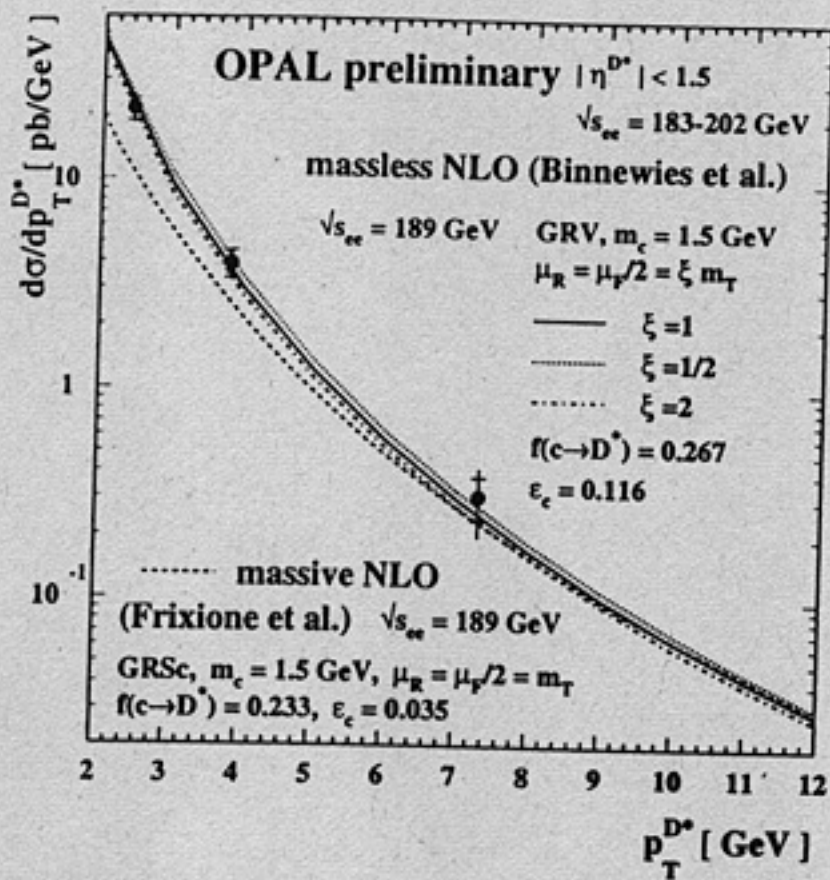
B. Mele & P. Nason, Nucl. Phys. B 361, 173 (1991)



- finite terms in (iii) subtracted to match (i)

(i) is for $\overline{MS}$ factorization scheme

scheme for PDF and FF $c \rightarrow D^*$ is $\overline{MS}$



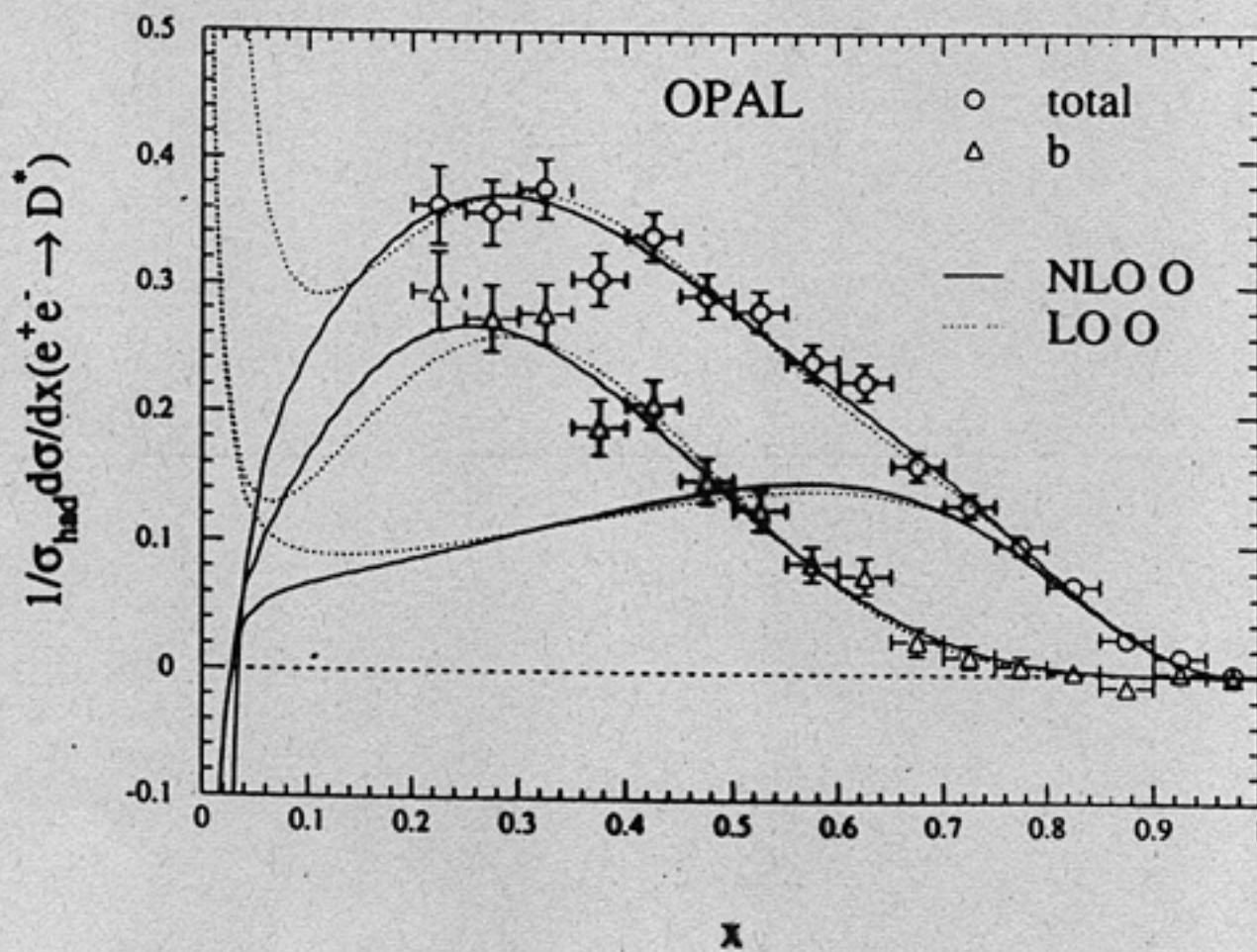
$$e^+ + e^- \rightarrow D^* + X$$

two components : $e^+ + e^- \rightarrow c + X$
 $\searrow D^*$

$$e^+ + e^- \rightarrow b + X$$

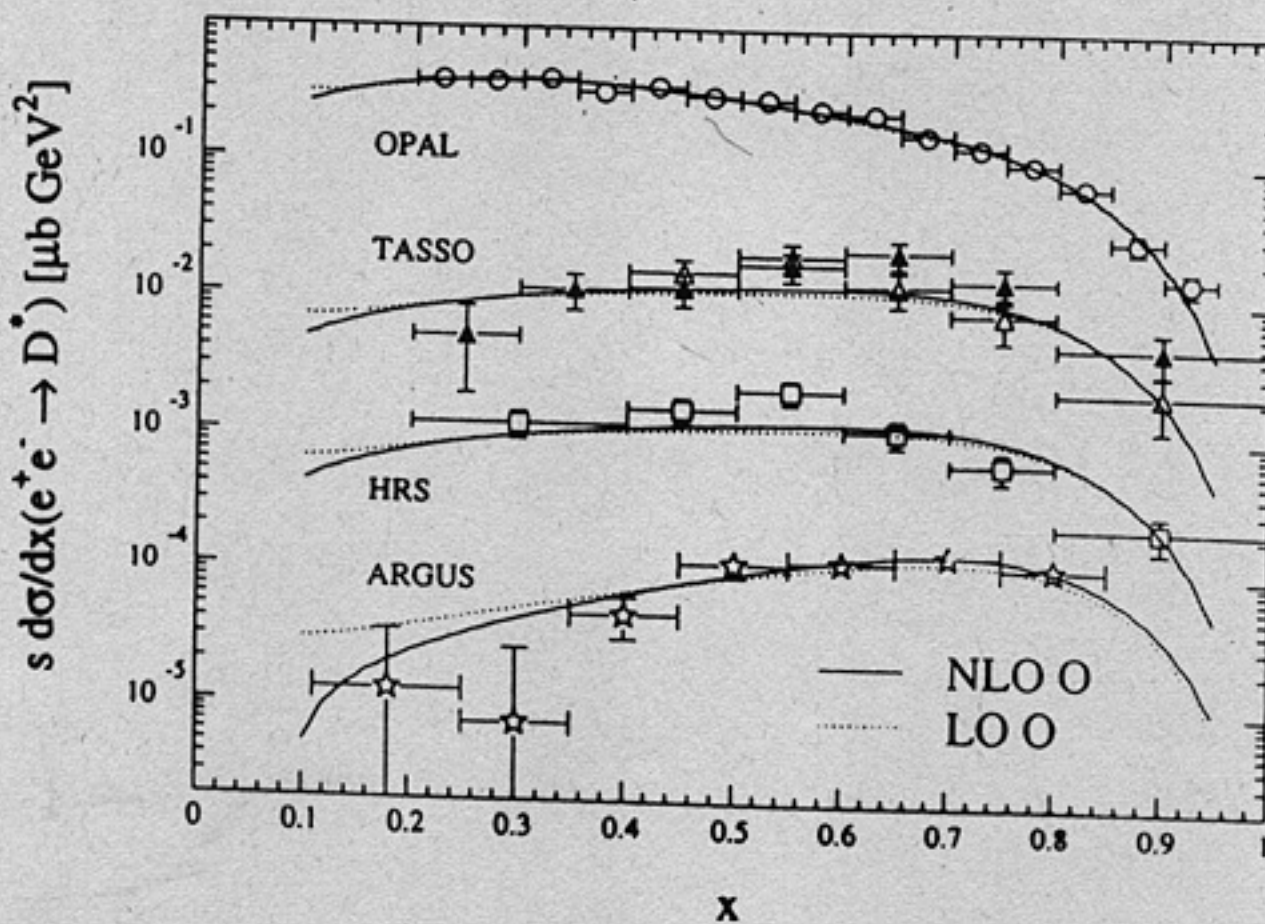
$\hookrightarrow D^*$

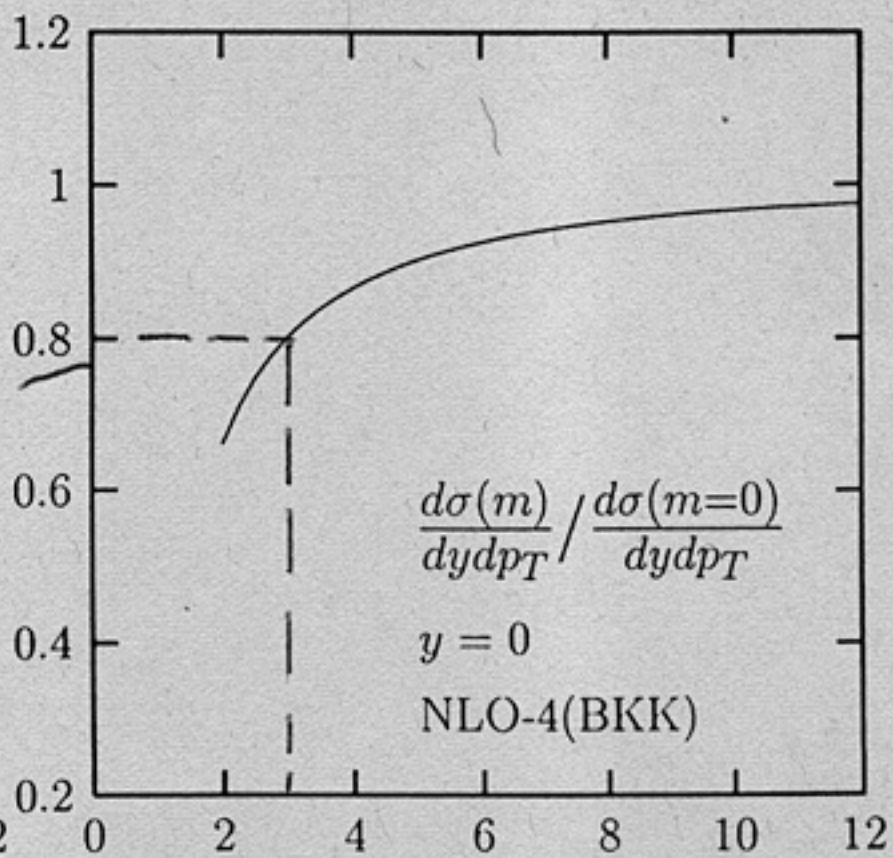
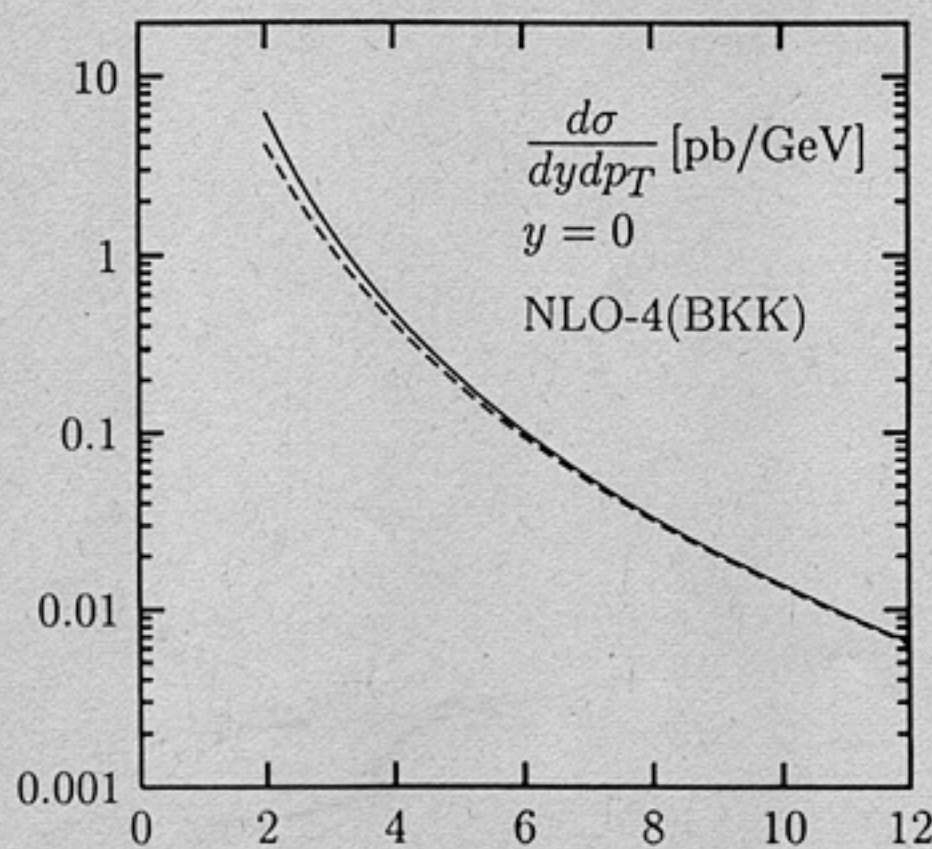
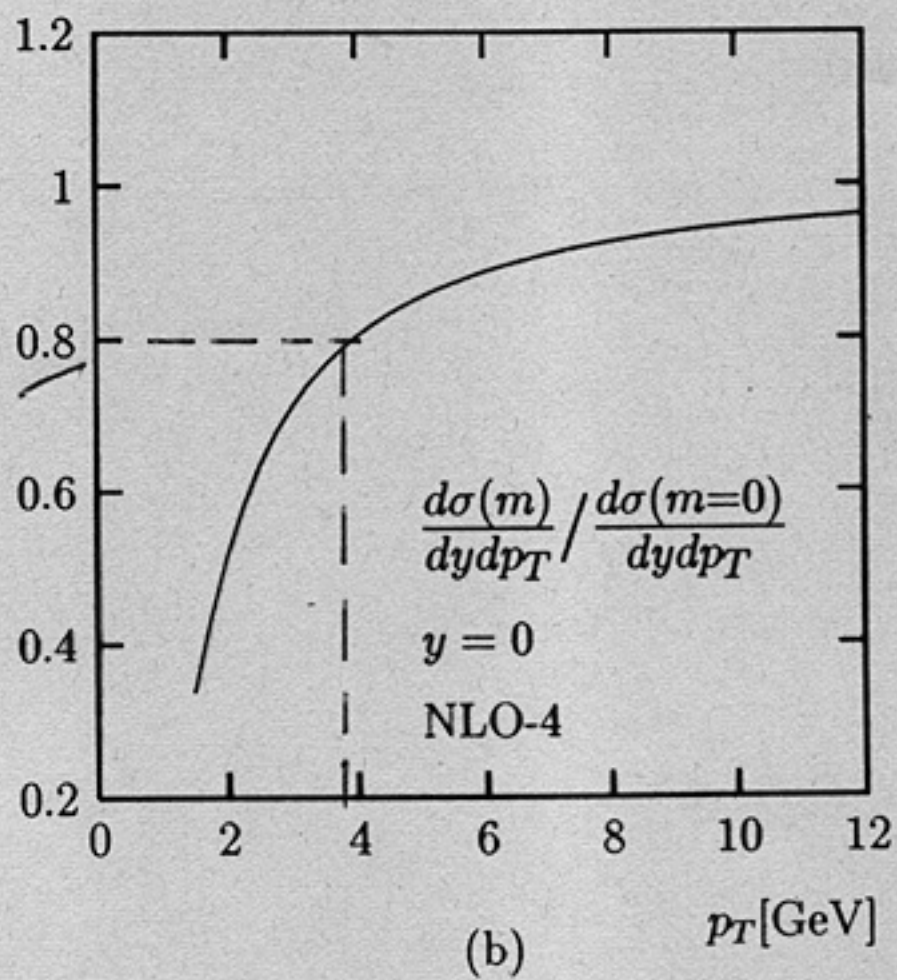
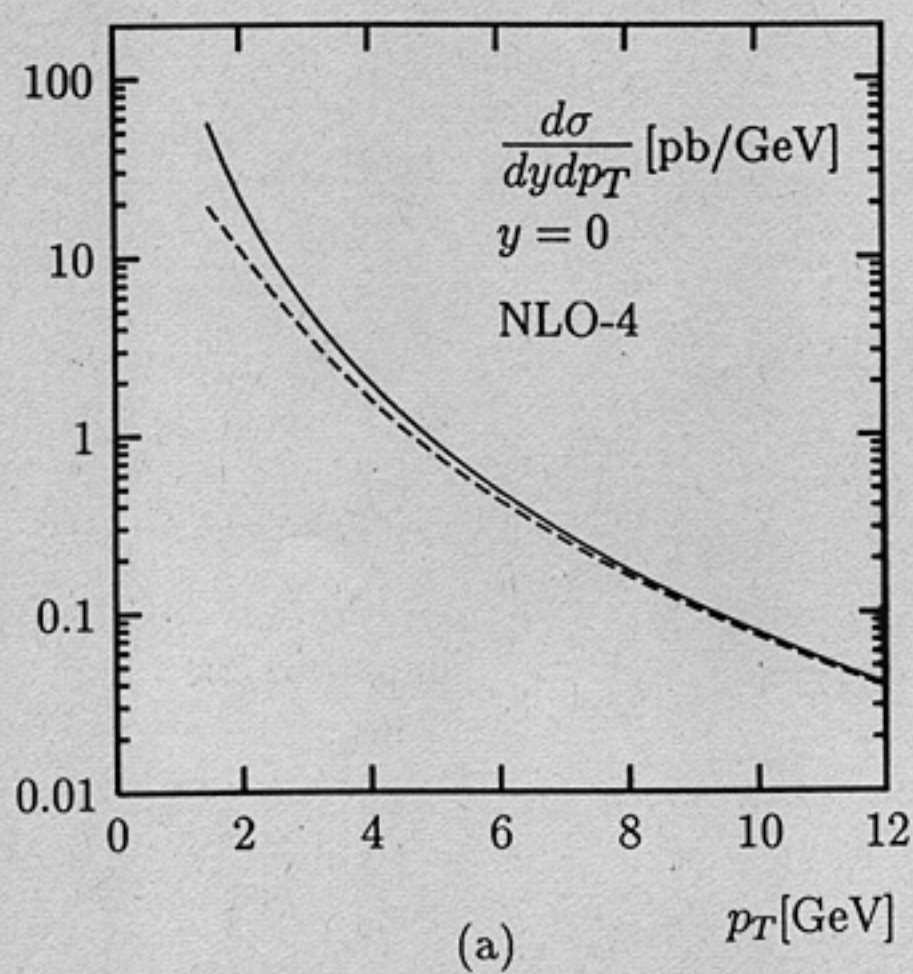
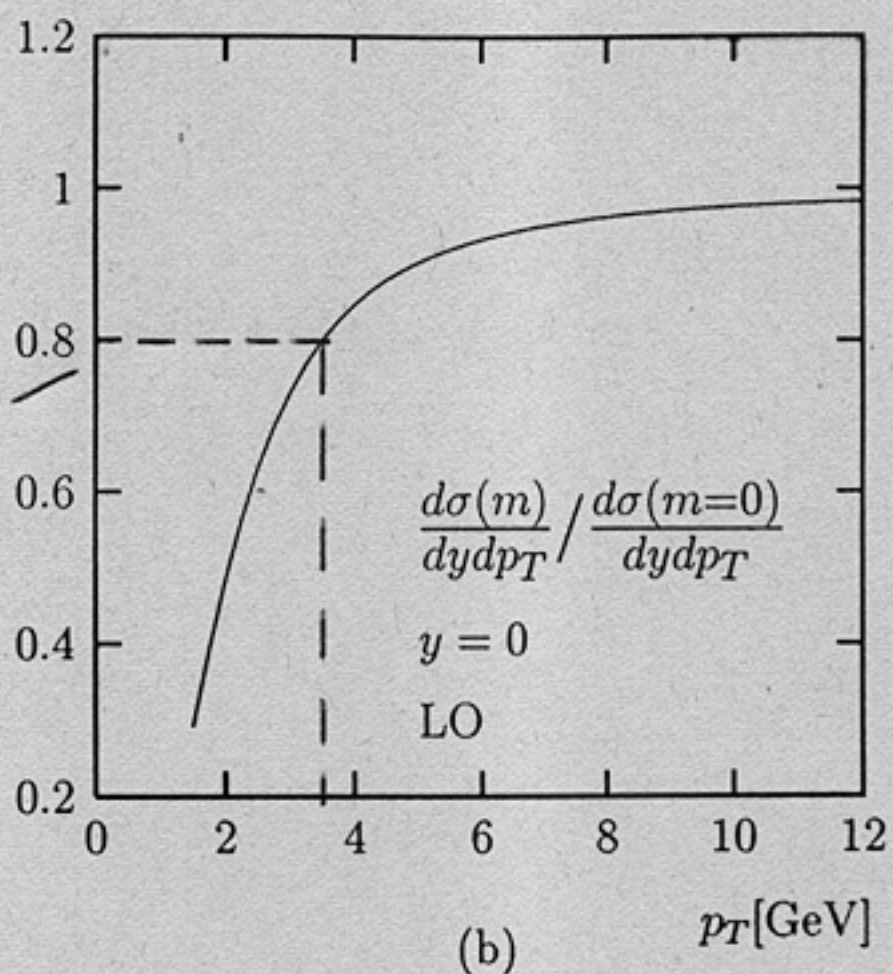
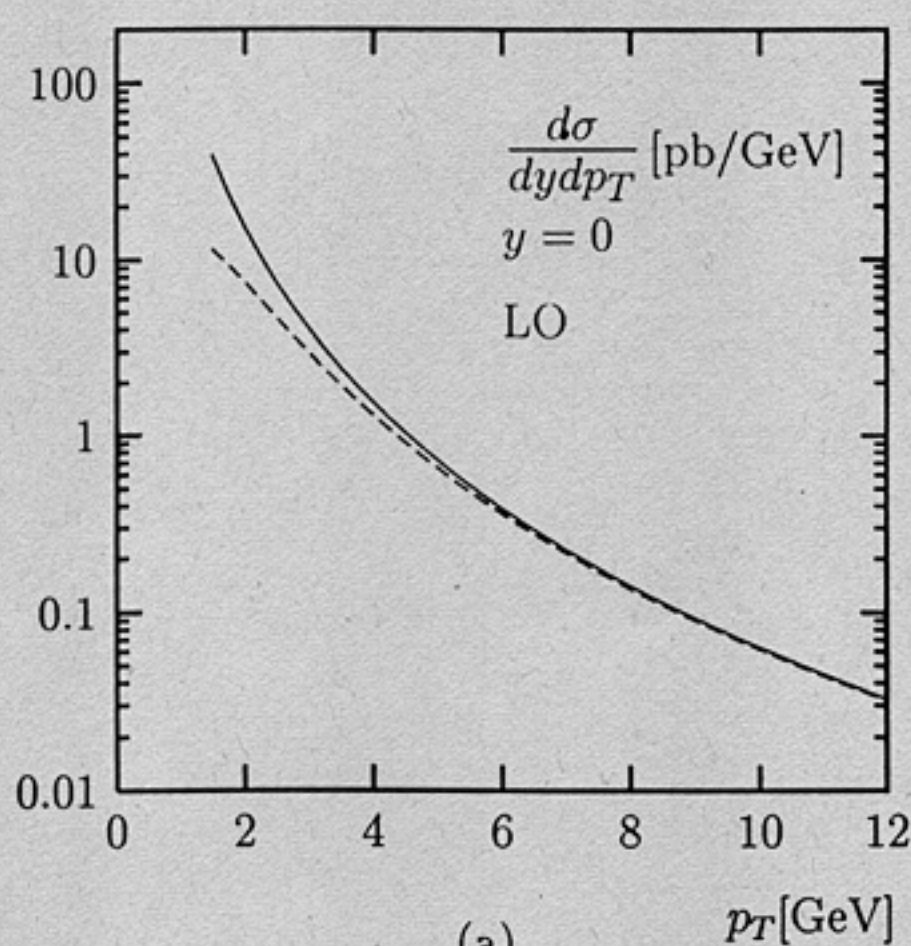
J. Binnewies, B.A. Knichl & G. K., Phys. Rev. D 58, 014014 (1998)



NLO

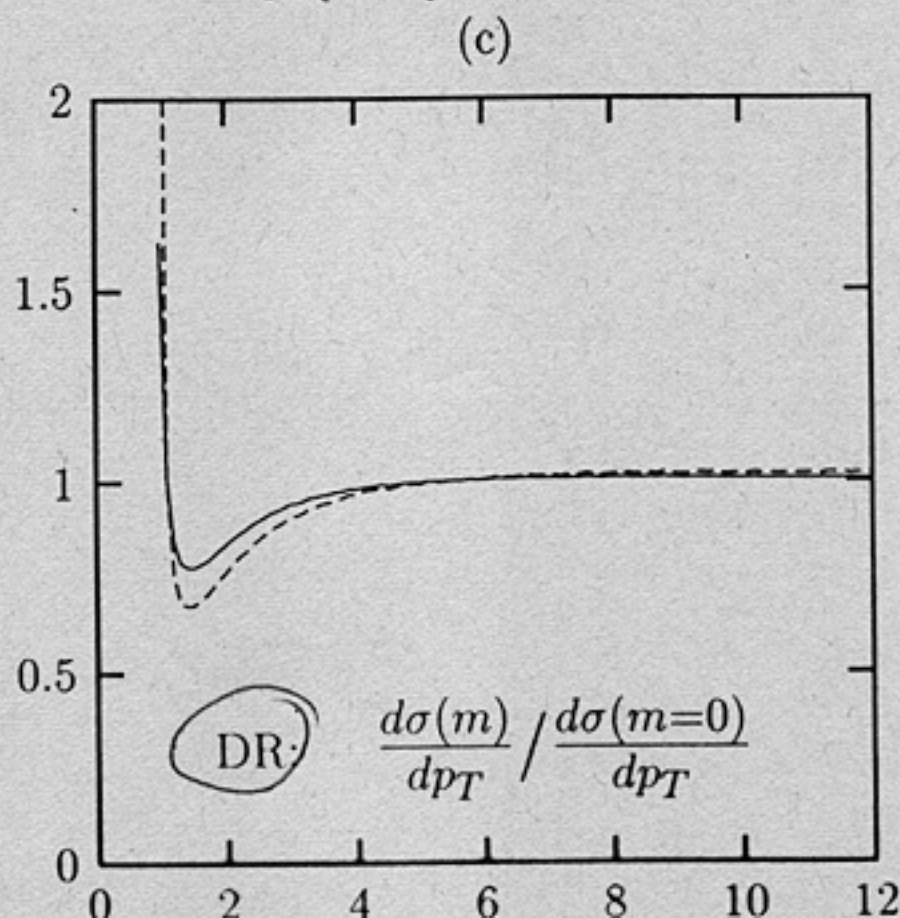
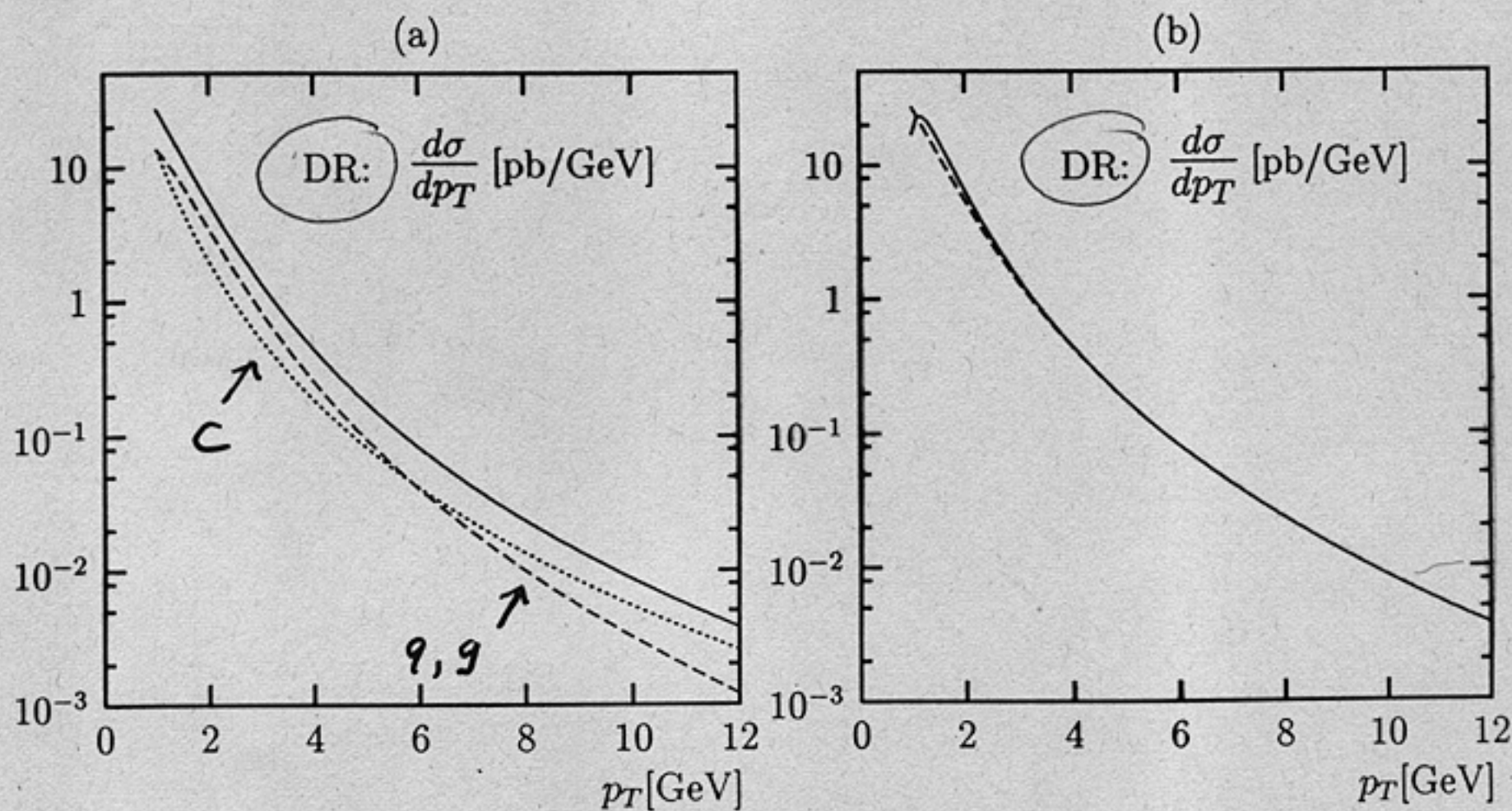
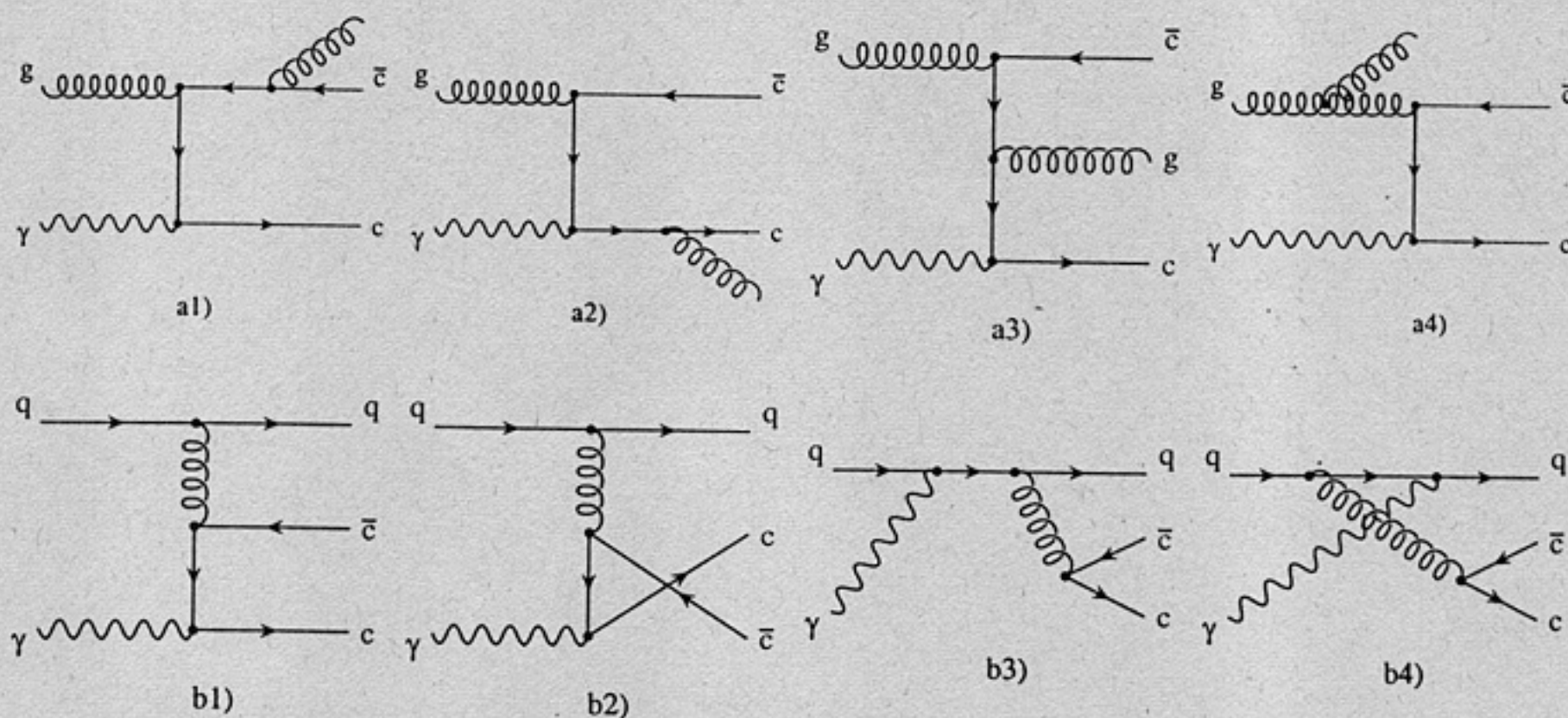
$\overline{MS}$ factorization





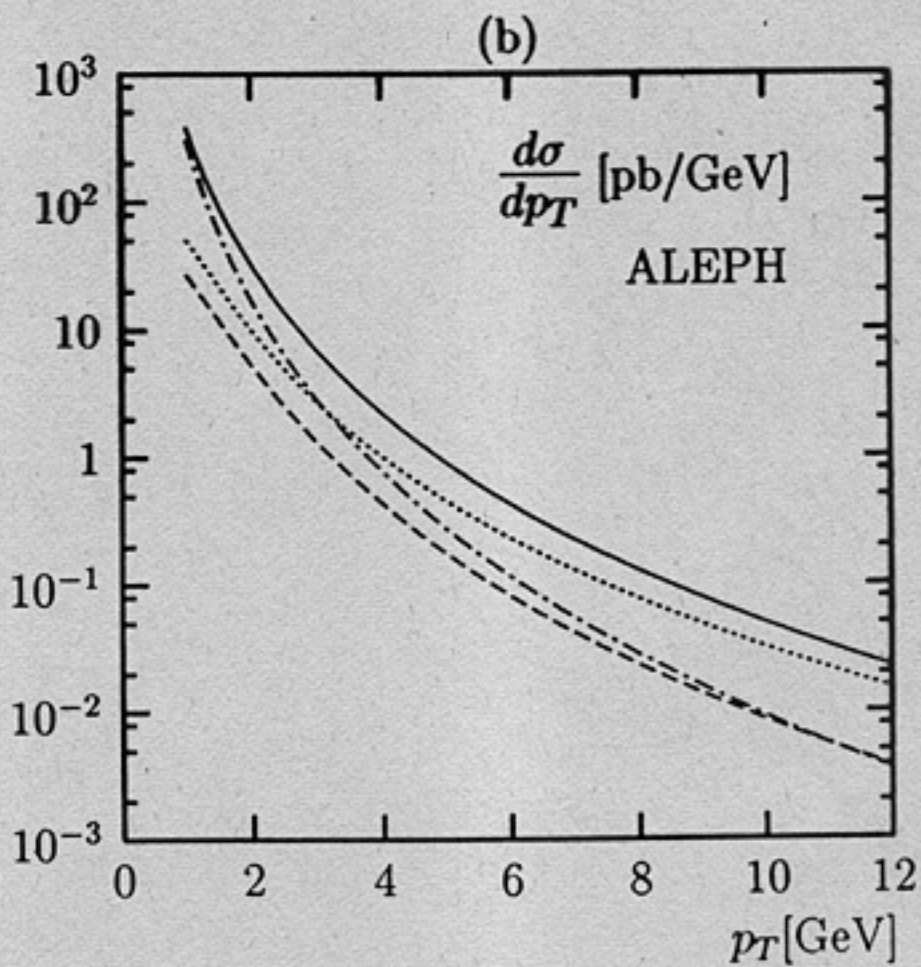
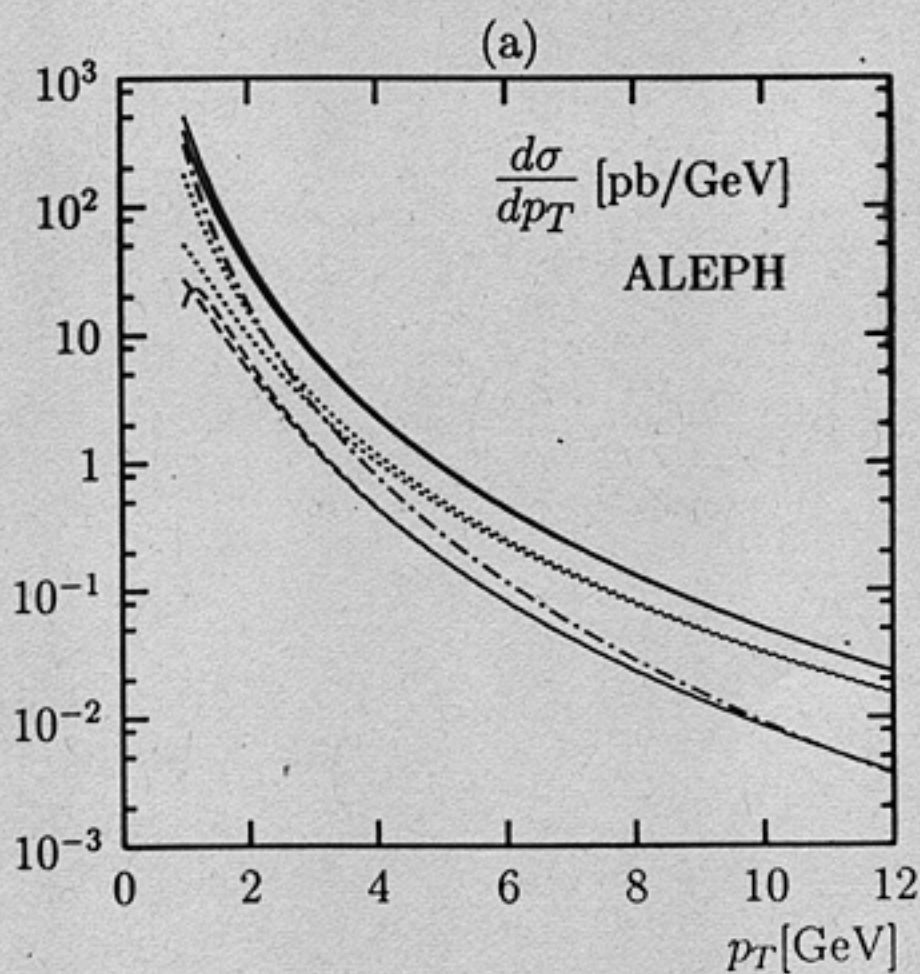
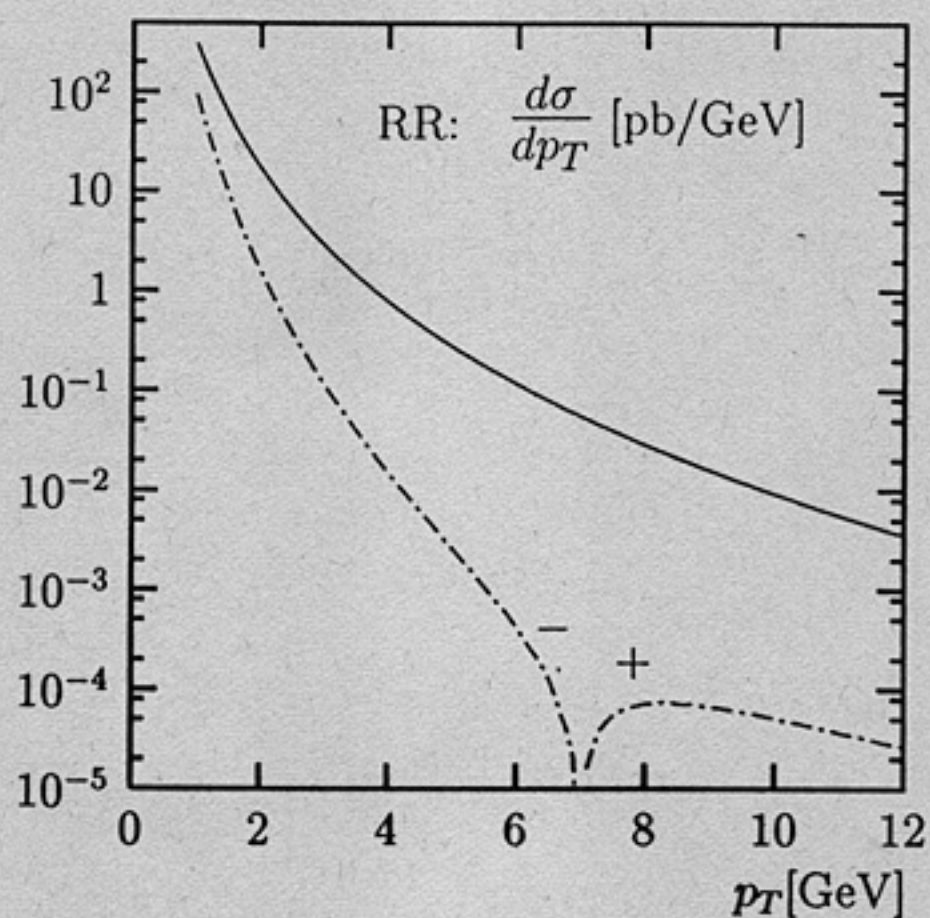
single-resolved process

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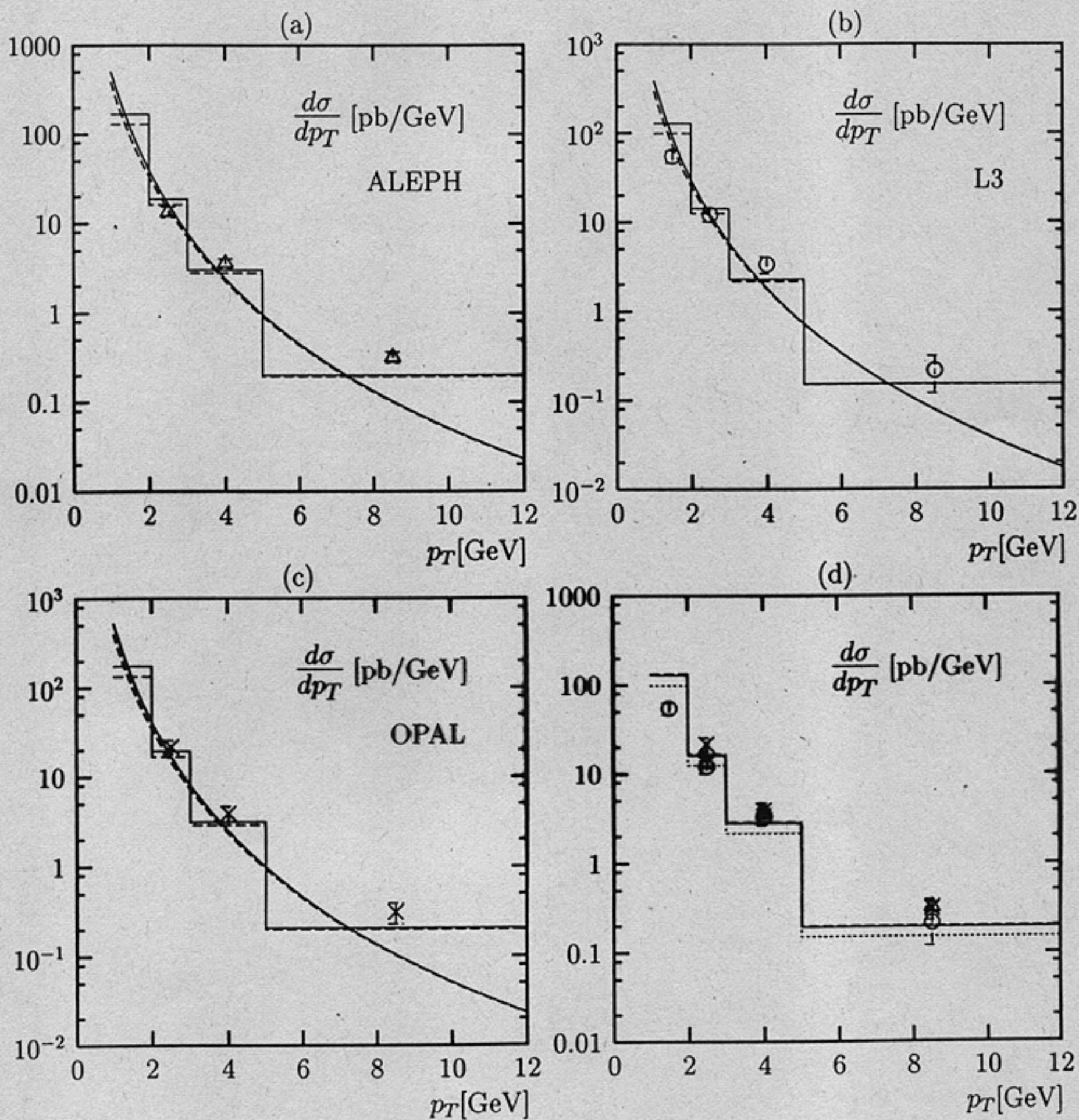
— massless
 --- massive
 --- no charm

double-resolved contribution



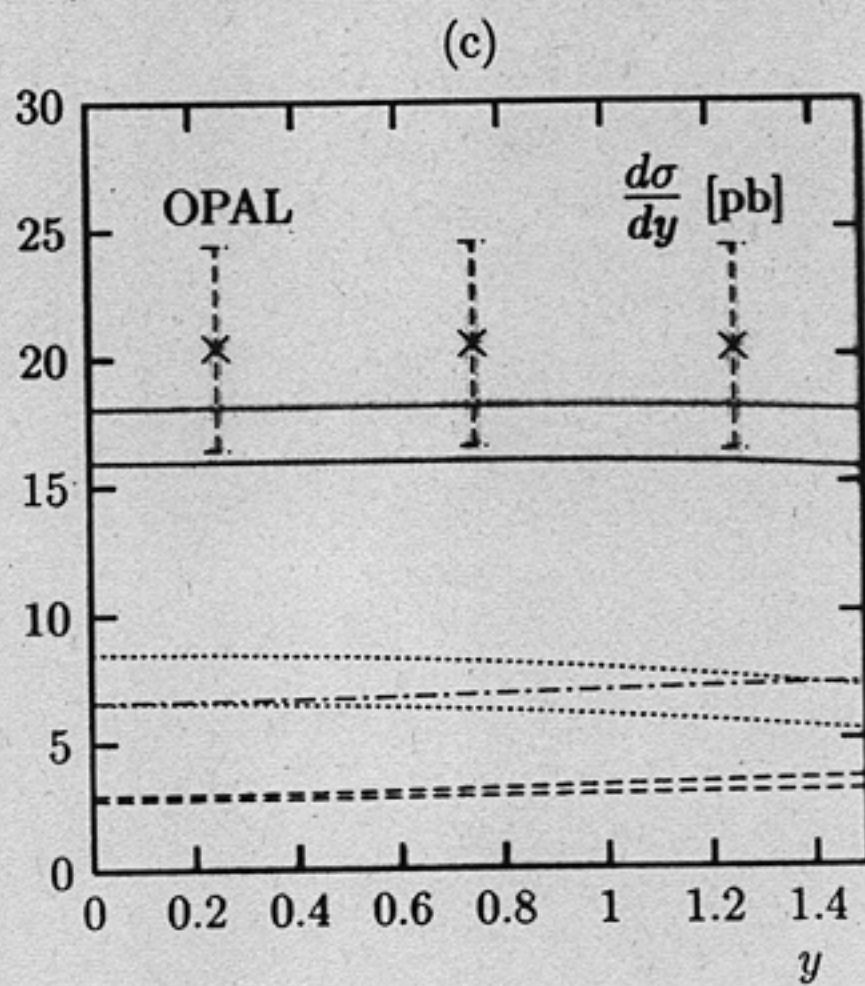
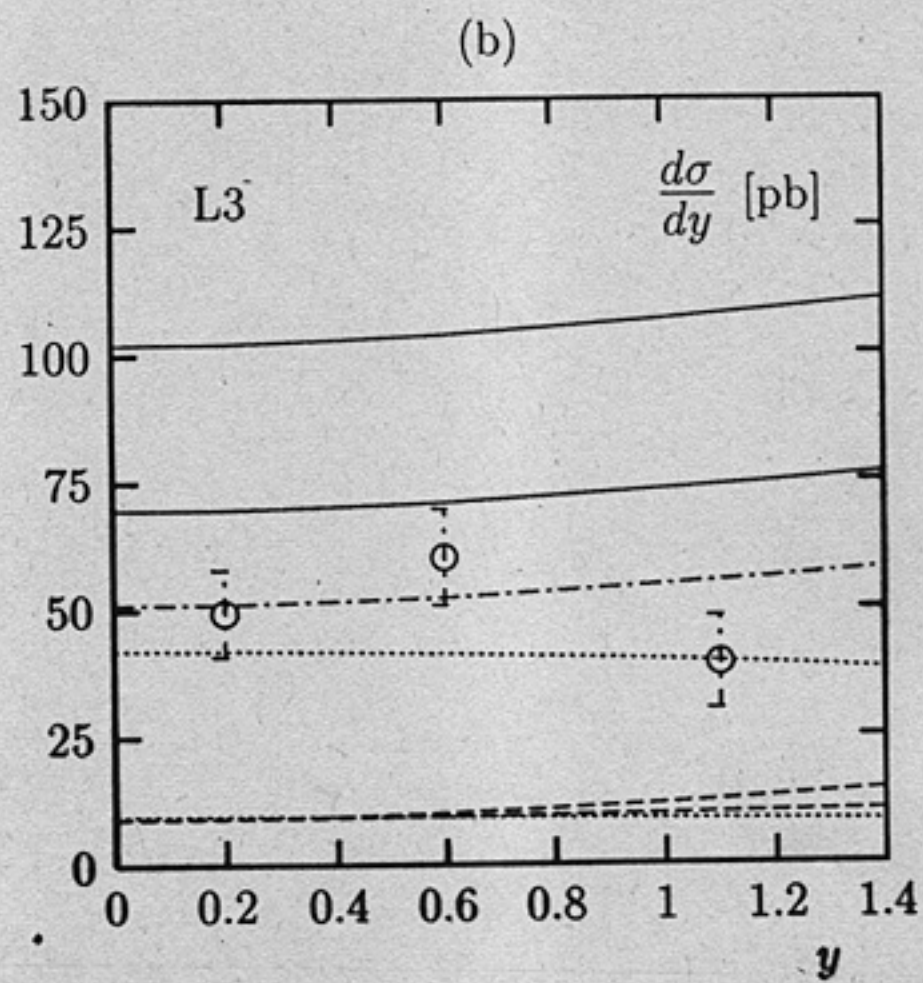
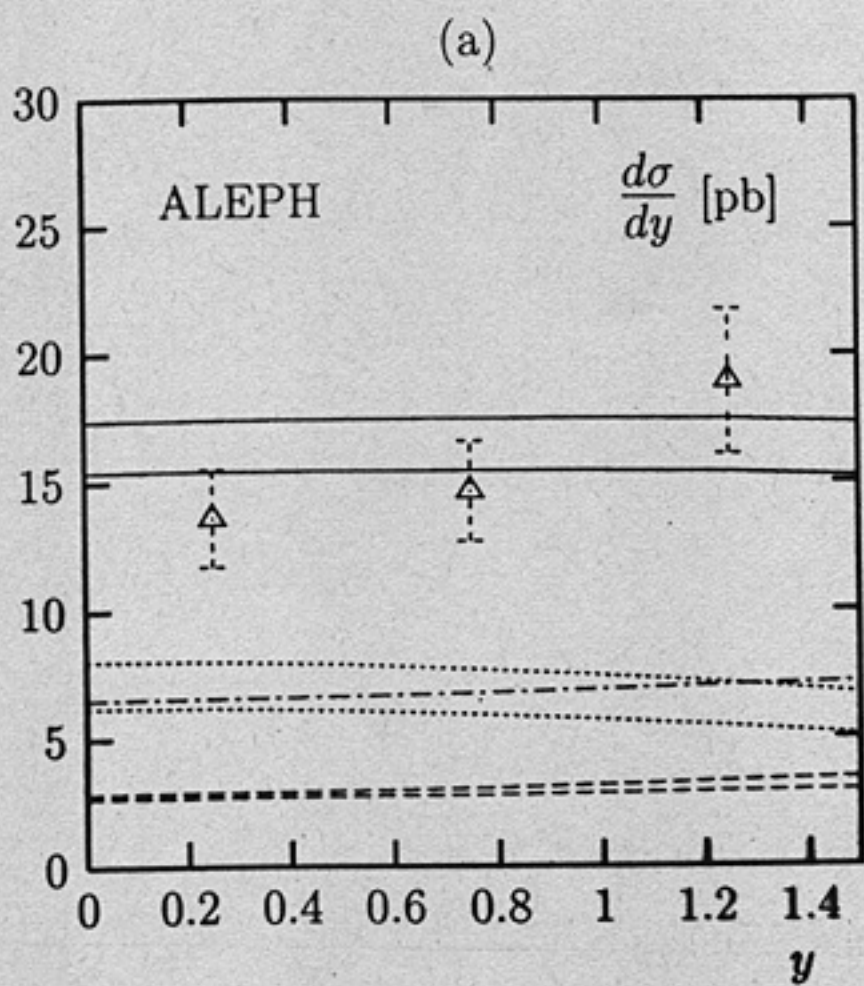
— direct + single resolved
+ double resolved

eff. of m_c^2 large small



ALEPH, L3, OPAL

different constraints on $\sqrt{s}_{av}$, tagging, η range



Conclusions

Resummed NLO calculation with $m_c \neq 0$ exists

Good agreement with ALEPH, L3 and OPA data

for $p_T > 2 \text{ GeV}$

Matching to NLO FFN, $n_f = 3$

at small p_T needs further study