

# VIRTUAL PHOTON STRUCTURE AT HERA.

P J BUSSEY

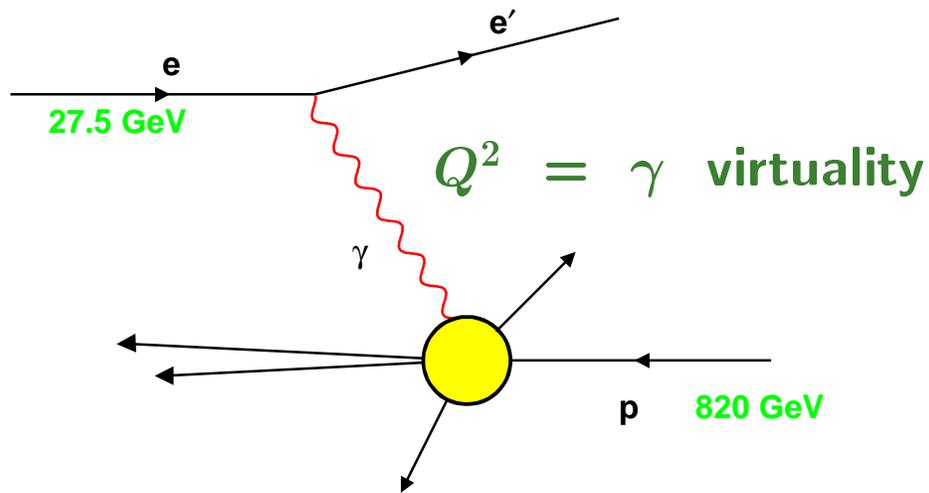
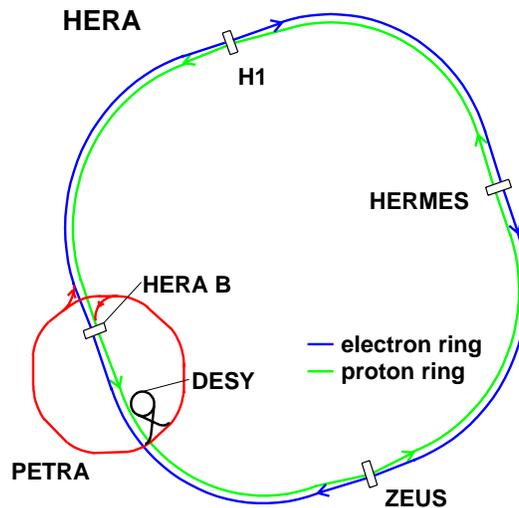


UNIVERSITY  
*of*  
GLASGOW

## CONTENTS

- (1) Introduction
- (2) Types of photon process
- (3) New results
- (4) Conclusions

# THE HERA COLLIDER



$Q^2 \approx 0$  — photoproduction  $\gamma p$

$Q^2 = O(1 \text{ GeV}^2)$  — 'Low  $Q^2$ ' transition region

$Q^2 \gg 1 \text{ GeV}^2$  — Deep Inelastic Scattering (DIS)

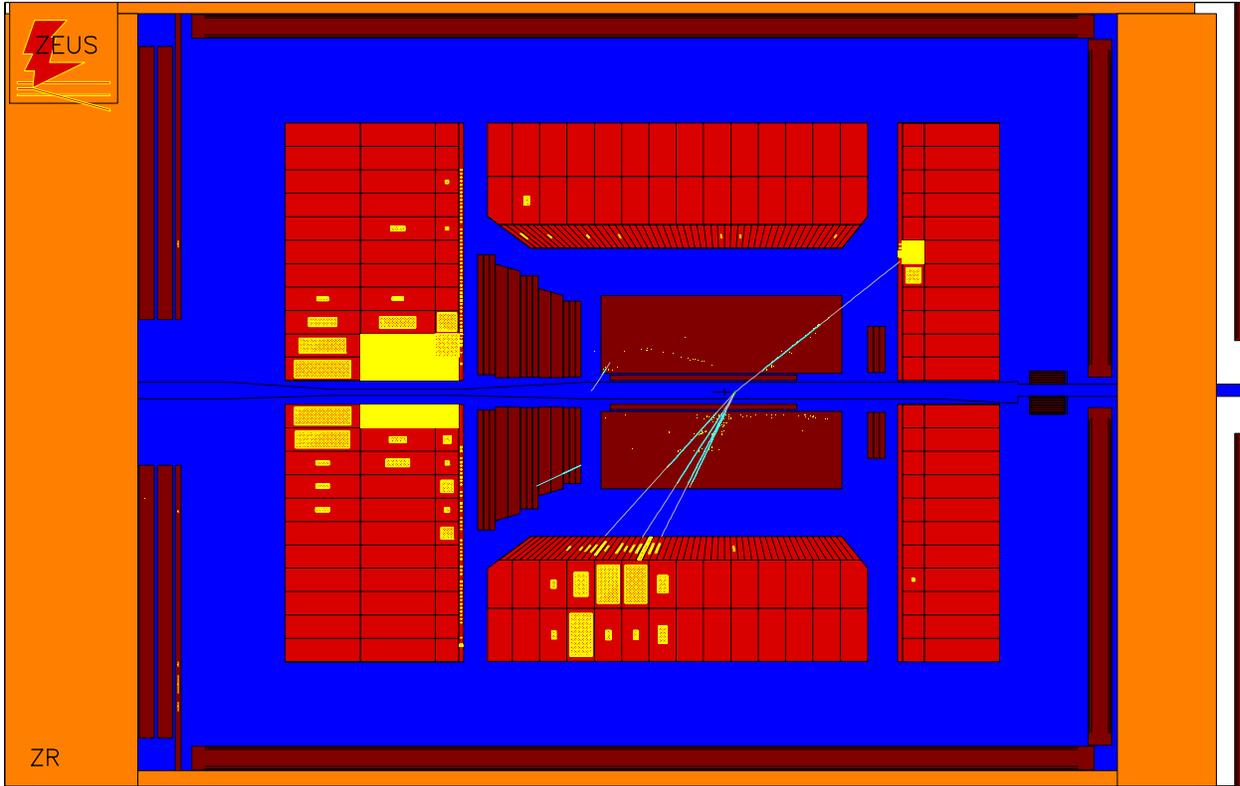
**Wish to study photon structure over these regions with special regard to transition region and DIS.**

---

# DIS AT HERA

---

## A typical DIS event



At lowest order (quark parton model) the virtual photon ejects a single quark from the proton:  $\rightarrow$  **one jet**

**No equivalent in photoproduction ( $Q^2 \approx 0$ )**

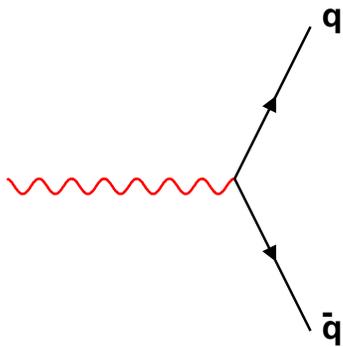
Can tag with  $e'$  for  $Q^2 \gtrsim 0.2 \text{ GeV}^2$  using low-angle 'beam-pipe' + rear calorimeters  $\rightarrow$  measure **dijet** events through transition region from photoproduction.

**Use tag to reconstruct  $\gamma^*p$  c.m. frame.**

---

# TYPES OF PHOTON COUPLING

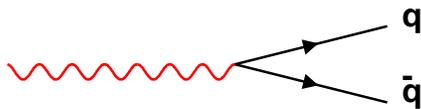
---



point-like (perturbative)



hadronic (non-perturbative)



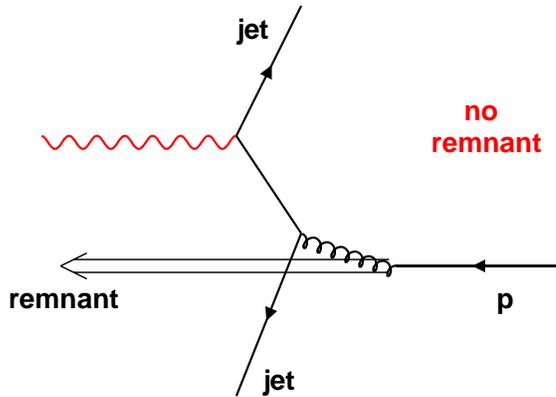
anomalous (perturbative)

These couplings can in principle be studied at all  $Q^2$  .

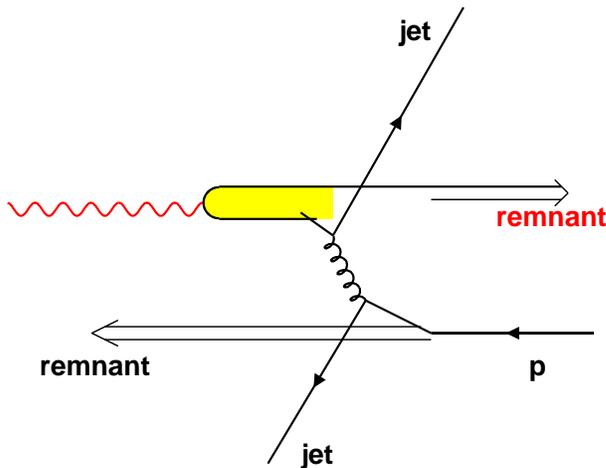
→ sensitivity to photon structure

# TYPES OF PHOTON PROCESS

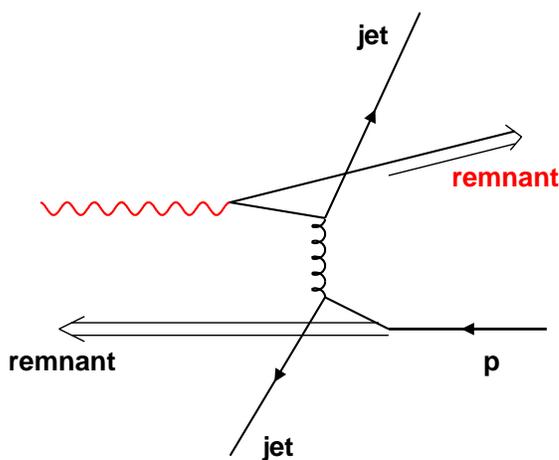
## Examples of dijet processes.



DIRECT



RESOLVED



looks resolved.

NLO since 3  
perturbative  
vertices

Direct / resolved fully distinct only at LO.

---

## THEORETICAL FRAMEWORK

---

**Aim:** to study effective structure of virtual photon as function of  $Q^2$ , with special reference to 'transition region'  $Q^2 = O(1) \text{ GeV}^2$

In perturbative QCD, assuming factorisation (OK at NLO):

$$d\sigma_{ep \rightarrow e+2\text{jets}} = \sum_{a,b} \int_0^1 dy f_{\gamma^*/e}(y, Q^2) \int_0^1 dx_{\gamma^*} f_{a/\gamma^*}(x_{\gamma^*}, Q^2, \mu_{F\gamma^*}^2) \times \int_0^1 dx_p f_{b/p}(x_p, \mu_{Fp}^2) d\sigma_{ab \rightarrow 2\text{jets}}(\mu_R)$$

where there are two contributions to the  $\gamma^*$  PDFs:

$$f_{a/\gamma^*} = f_{a/\gamma^*}^{\text{had}} + f_{a/\gamma^*}^{\text{pert}}$$

Boundary is factorisation-scale ( $\mu_F$ ) dependent.

**DISASTER++** (D. Graudenz)

used as NLO parton-level Monte Carlo.

**PYTHIA and ARIADNE**

used to convert data and NLO partons to the hadron level.

# DIJET EVENTS

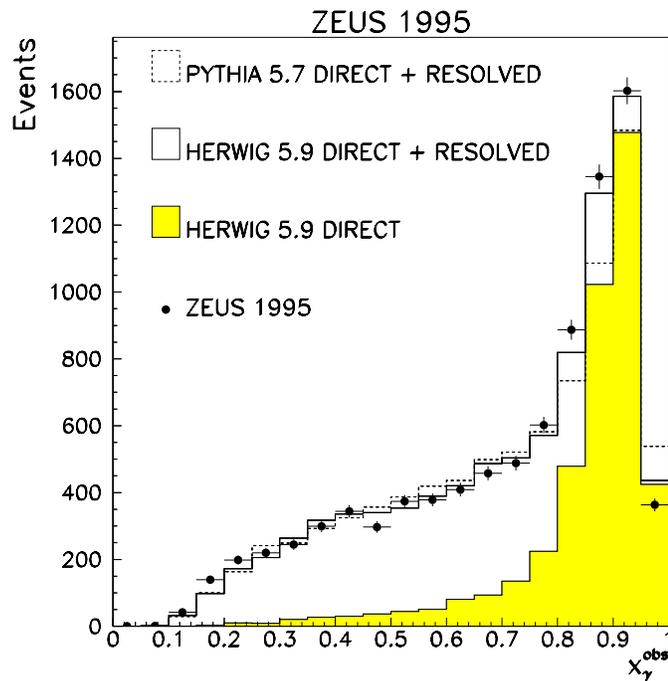
Dijet events identified in  $\gamma^*p$  c.m. frame

$$E_T^1, E_T^2 > 7.5, 6.5 \text{ GeV}, \quad -3 < \eta < 0$$

A measure of the fraction of the photon energy given to the dijet final state is:

$$x_\gamma^{\text{obs}} = (E_T^1 e^{-\eta_1} + E_T^2 e^{-\eta_2}) / 2E_\gamma$$

$x_\gamma^{\text{obs}}$  in  
photoproduction  
(typical)



$x_\gamma^{\text{obs}} < 0.75$  Resolved-dominated

$x_\gamma^{\text{obs}} > 0.75$  Direct-dominated

The ratio

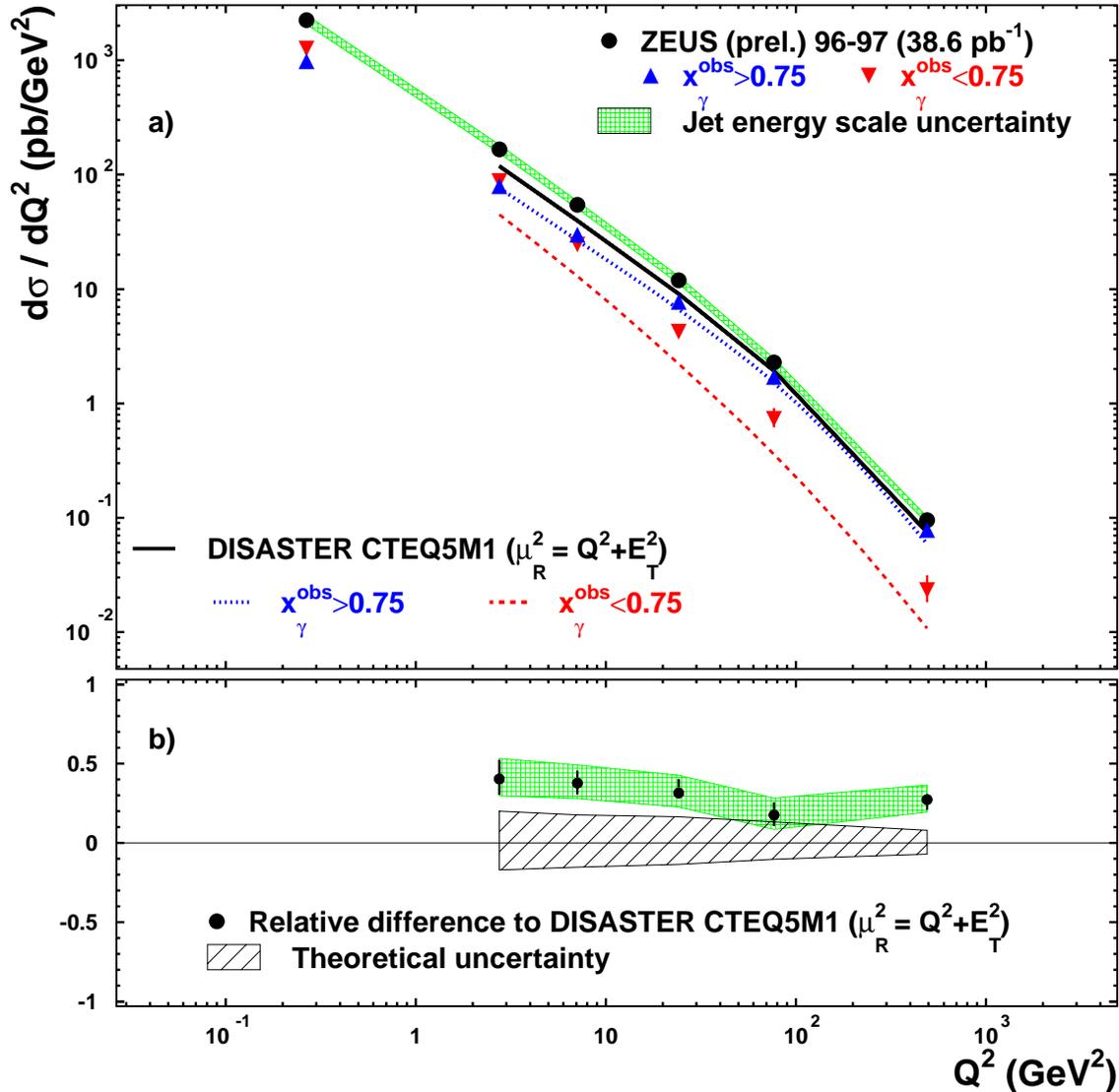
$$R = \frac{d\sigma}{dQ^2}(x_\gamma^{\text{obs}} < 0.75) / \frac{d\sigma}{dQ^2}(x_\gamma^{\text{obs}} > 0.75)$$

is a measure of the resolved fraction of the virtual photon events within the acceptance. Various experimental and theoretical uncertainties cancel in  $R$ .

# RESULTS

Results for  $d\sigma/dQ^2$

## ZEUS



Using QCD renormalisation scale  $\mu_R^2 = Q^2 + E_T^2$  :

— Data tend to lie above DISASTER NLO

— But good agreement with  $x_{\gamma}^{\text{obs}} > 0.75$  component

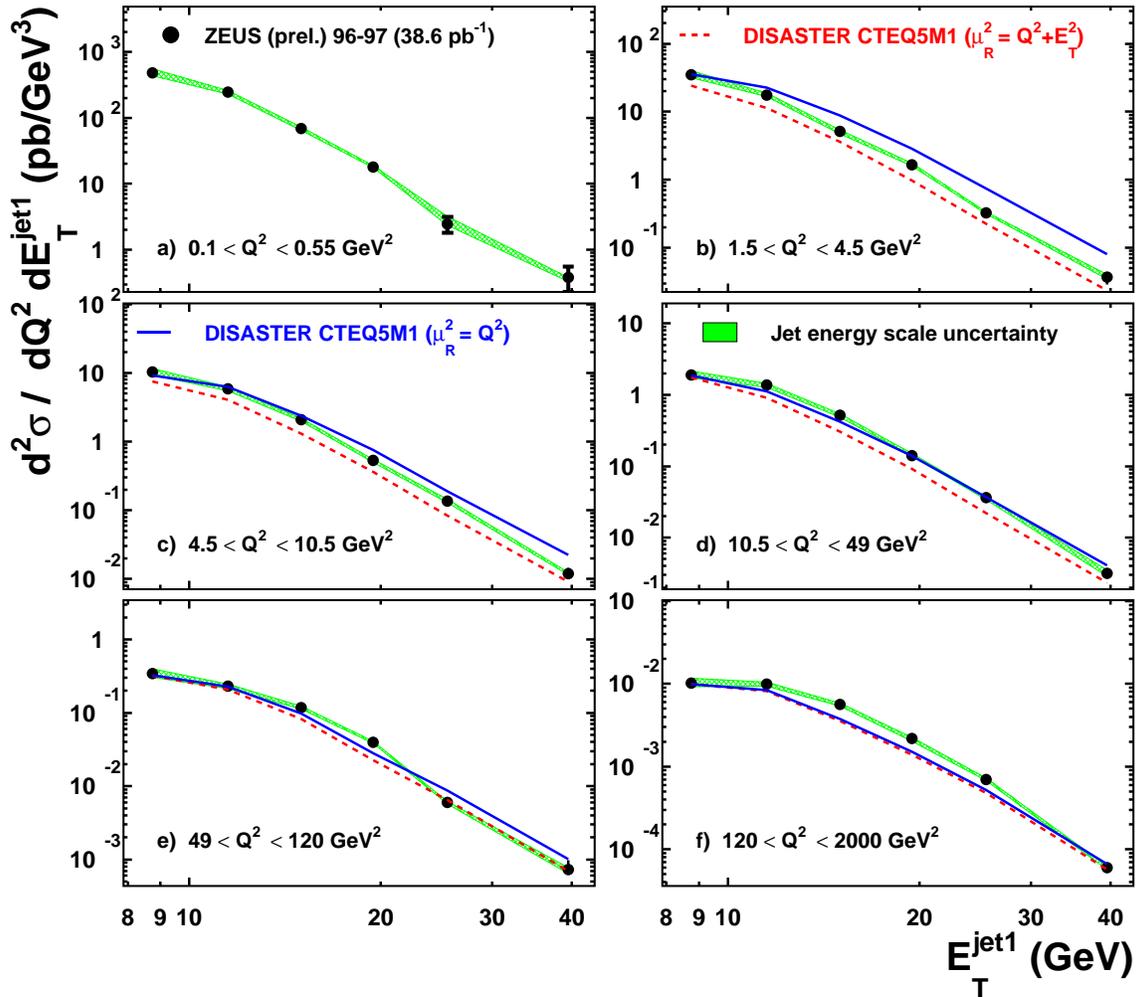
— Discrepancy mainly with the  $x_{\gamma}^{\text{obs}} < 0.75$  component

DISASTER contains only perturbative processes,  
 no hadronic component to photon.

# RESULTS

Plot against higher  $E_T^{\text{jet}}$  for varying  $Q^2$  ranges:

## ZEUS



Comparison with NLO is not strongly  $E_T^{\text{jet}}$  dependent.

Note presence of two hard QCD scales,  $Q^2$  and  $E_T^{\text{jet}}$ .

More natural to use  $\mu_R^2 = Q^2 + E_T^2$ .

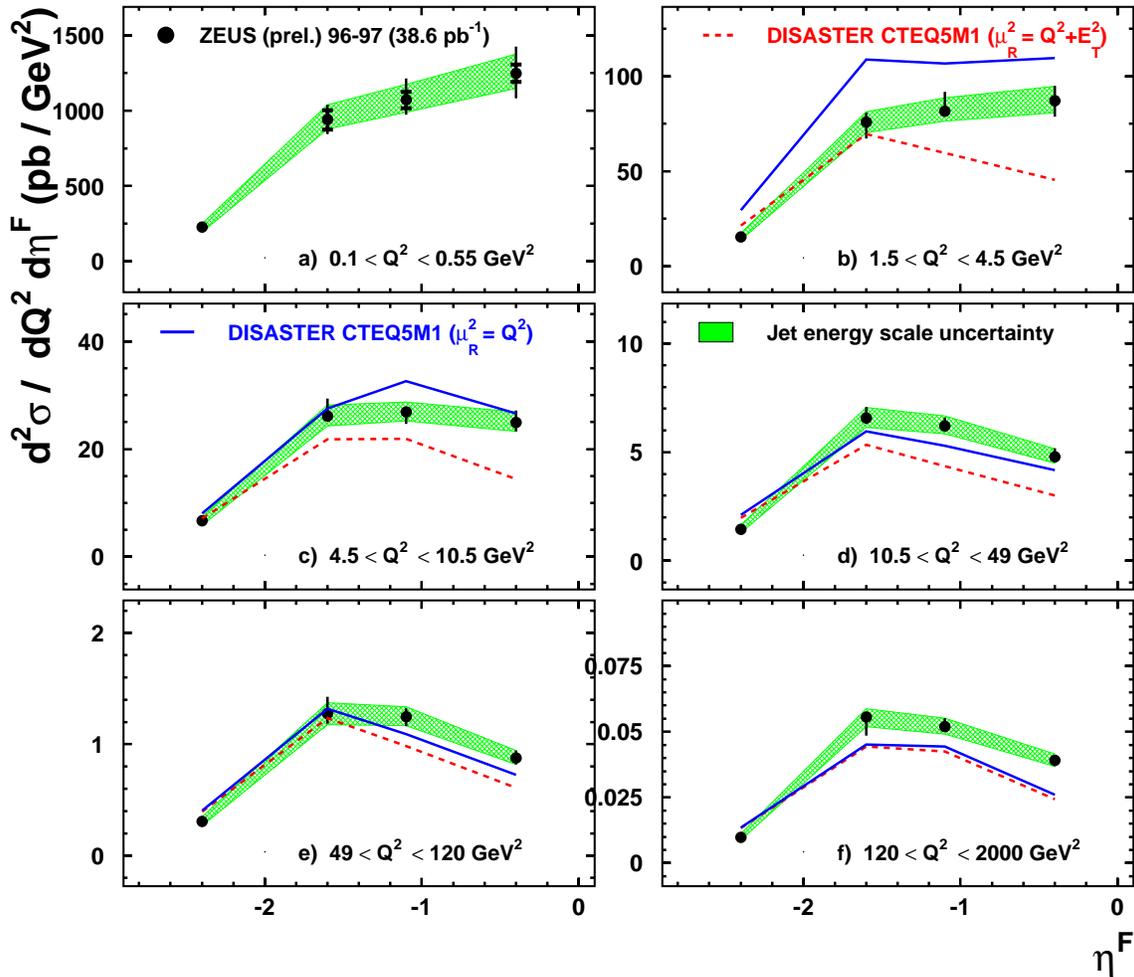
Using  $\mu_R^2 = Q^2$  raises the prediction.

Is this a better procedure?

# RESULTS

$d\sigma/dQ^2$  vs forward jet pseudorapidity  
enhances resolved contribution

## ZEUS



At low  $Q^2$ , big sensitivity to  $\mu_R^2$

At high  $Q^2$ , low sensitivity; data tend to be above theory in both cases.

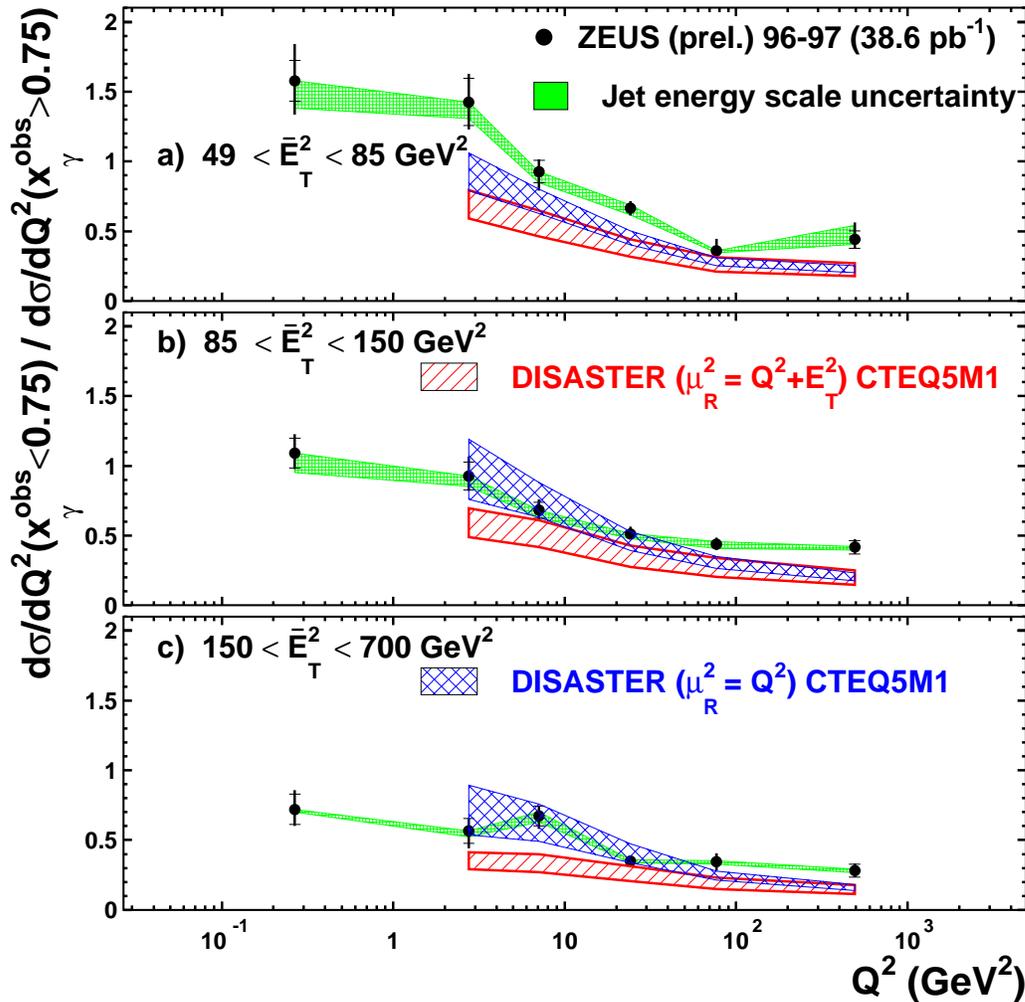
So decreasing  $\mu_R^2$  is not a successful solution.

# RESULTS

$$R = \frac{d\sigma}{dQ^2}(x_\gamma^{\text{obs}} < 0.75) / \frac{d\sigma}{dQ^2}(x_\gamma^{\text{obs}} > 0.75)$$

vs  $Q^2$  in ranges of  $E_T^{\text{jet}}$

## ZEUS



→ all resolved effects fall with  $Q^2$  and  $E_T^{\text{jet}}$

(  $R$  at  $Q^2 \approx 0$  estimated similar to  $Q^2 \approx 0.3$  )

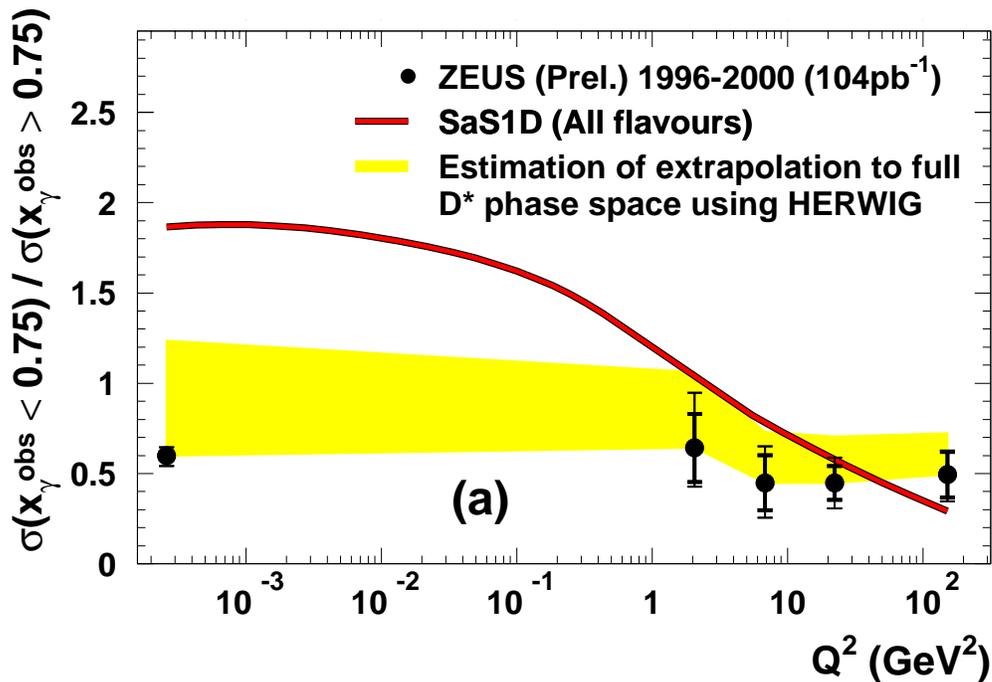
Comparison with DISASTER NLO ( $\mu_R^2 = Q^2 + E_T^2$ ) is insensitive to either hard scale. **Data lie above theory.**

→ consistent need for higher-order resolved contribn. **(hadronic!)** to photon even at high  $Q^2$

## COMPARISONS

Compare with prelim. ZEUS data with identified charm (D meson) required in event.

Jet selections similar to present analysis.



Presence of  $c$  quark forces hard scale on process.

Flat in  $Q^2$

Should all be perturbatively calculable.

---

## CONCLUSIONS

- Have measured dijet cross sections in ZEUS over a range of  $Q^2$  values crossing the transition between photoproduction and DIS.
- There is a component to the cross sections which is not modelled in NLO QCD, and is associated with resolved processes.
- Decreasing the renormalisation scale does not satisfactorily account for the discrepancy.
- Suggests presence of a hadronic part of the photon even at high  $Q^2$  values.
- The effect appears absent in dijets with heavy quarks
- Need a NLO MC which includes general photon PDFs to test these ideas further.