

Hadronic Light by Light Contribution to Muon $g - 2$: Status and Prospects

Joaquim Prades

CAFPE and Universidad de Granada



PHIPSI08 Workshop, 7-10 April 2008, Frascati

FlaviA
net

Plan

⇒ Introduction

Plan

⇒ Introduction

⇒ “Old” Calculations: 1995-2001

Plan

- ⇒ Introduction
- ⇒ “Old” Calculations: 1995-2001
- ⇒ New Short-Distance Constraints: 2003-2004

Plan

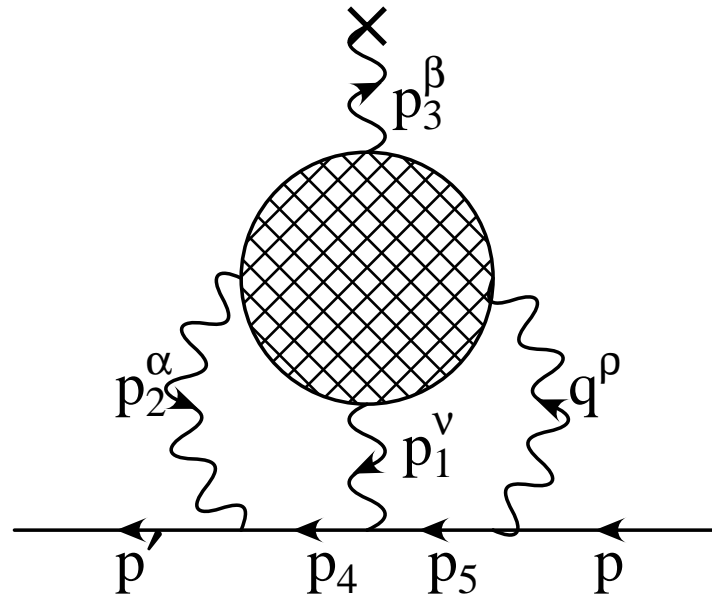
- ⇒ Introduction
- ⇒ “Old” Calculations: 1995-2001
- ⇒ New Short-Distance Constraints: 2003-2004
- ⇒ Comparison

Plan

- ⇒ Introduction
- ⇒ “Old” Calculations: 1995-2001
- ⇒ New Short-Distance Constraints: 2003-2004
- ⇒ Comparison
- ⇒ Conclusions and Prospects

Introduction

Hadronic light-by-light contribution to muon $g - 2$



$$\mathcal{M} = |e|^7 A_\beta \int \frac{d^4 p_1}{(2\pi)^4} \int \frac{d^4 p_2}{(2\pi)^4} \frac{1}{q^2 p_1^2 p_2^2 (p_4^2 - m^2) (p_5^2 - m^2)} \\ \times \underline{\Pi^{\rho\nu\alpha\beta}(p_1, p_2, p_3)} \bar{u}(p') \gamma_\alpha (\not{p}_4 + m) \gamma_\nu (\not{p}_5 + m) \gamma_\rho u(p)$$

Introduction

Need

$$\Pi^{\rho\nu\alpha\beta}(p_1, p_2, p_3) = i^3 \int d^4x \int d^4y \int d^4z \exp^{i(p_1 \cdot x + p_2 \cdot y + p_3 \cdot z)} \times \\ \times \langle 0 | T [V^\rho(0) V^\nu(x) V^\alpha(y) V^\beta(z)] | 0 \rangle$$

with $V^\mu(x) = [\bar{q} \hat{Q} \gamma^\mu q](x)$ and $\hat{Q} = \frac{1}{3} \text{diag}(2, -1, -1)$

full four-point function with $p_3 \rightarrow 0$ •

Using gauge invariance

$$\Pi^{\rho\nu\alpha\lambda}(p_1, p_2, p_3) = -p_{3\beta} \frac{\delta \Pi^{\rho\nu\alpha\beta}(p_1, p_2, p_3)}{\delta p_{3\lambda}}$$

one just needs derivatives at $p_3 = 0$ •

Introduction

★ Many scales involved: Impose low energy and several OPE limits \Rightarrow Not full first principle calculation at present ●

(Two lattice groups just starting: Not clear final uncertainty) ●

Large N_c and CHPT counting:

Organizes different degrees of freedom contributions ●

E. de Rafael

- Goldstone boson exchange: $\mathcal{O}(N_c)$ and $\mathcal{O}(p^6)$ ●
- Quark Loop and non-Goldstone boson exchange: $\mathcal{O}(N_c)$ and $\mathcal{O}(p^8)$ ●
- Goldstone bosons Loop: $\mathcal{O}(1)$ in $1/N_c$ and $\mathcal{O}(p^4)$ ●

Introduction

Based on this counting:

- Two full calculations
J. Bijnens, E. Pallante, J.P. (BPP)
M. Hayakawa, T. Kinoshita, A. Sanda (HKS)
- Dominant pseudo-scalar exchange: Extensive analytic analysis ●
M. Knecht, A. Nyffeler (KN)

Found sign mistake ✓

M. Knecht, A. Nyffeler, M. Perrottet, E. de Rafael

Introduction

★ New four-point form factor short-distance constraint:

K. Melnikov, A. Vainshtein

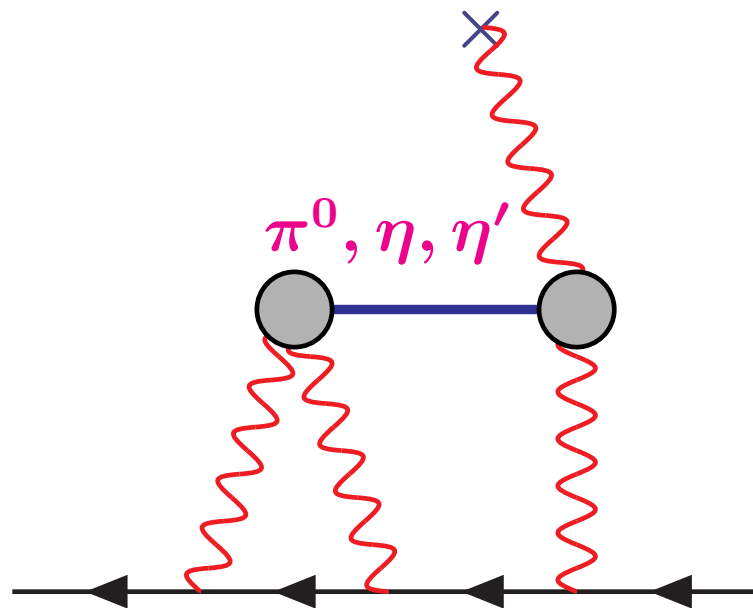
(see also M. Knecht, S. Peris, M. Perrottet, E. de Rafael)

Model:

Full light-by-light saturated by pseudo-scalar and pseudo-vector pole exchanges ●

“Old” Calculations: Pseudo-Scalar Exchange

Dominant contribution \Rightarrow pseudo-scalar exchange •



Here, I discuss work in J. Bijnens, E. Pallante, J.P. •

“Old” Calculations: Pseudo-Scalar Exchange

We used a variety of $\pi^0 \gamma^* \gamma^*$ form factors

$$\mathcal{F}^{\mu\nu}(p_1, p_2) = \frac{N_c}{6\pi} \frac{\alpha}{f_\pi} i\varepsilon^{\mu\nu\alpha\beta} p_{1\alpha} p_{2\beta} \underline{\mathcal{F}(p_1^2, p_2^2)}$$

fulfilling as many as possible QCD constraints •
(Short-distance, data, $U_A(1)$ normalization and slope at the origin). In particular,

$$\mathcal{F}(Q^2, Q^2) \rightarrow \frac{A}{Q^2}$$

$$\mathcal{F}(Q^2, 0) \rightarrow \frac{B}{Q^2}$$

for Q^2 Euclidean and very large

“Old” Calculations: Pseudo-Scalar Exchange

All form factors we used converge for $\mu \sim (2 - 4)$ GeV and the numerical difference between them is small ✓

Somewhat different $\pi^0 \gamma^* \gamma^*$ form factors used in

M. Hayakawa, T. Kinoshita, A. Sanda and M. Knecht, A. Nyffeler •

Results agree very well (after correcting a mistake in the sign of the phase space) •

	$10^{10} \times a_\mu$
Adding π^0 , η and η' contributions	BPP (8.5 ± 1.3)
	HKS (8.3 ± 0.6)
	KN (8.3 ± 1.2)

“Old” Calculations: Pseudo-Vector Exchange

Need $a_1^0 \gamma \gamma^*$ and $a_1^0 \gamma^* \gamma^*$ form factors •

⇒ related to $\pi^0 \gamma \gamma^*$ and $\pi^0 \gamma^* \gamma^*$ by anomalous Ward identities ✓

Pseudo-vector exchange

	$10^{10} \times a_\mu$
BPP	(0.25 ± 0.10)
HKS	(0.17 ± 0.10)

“Old” Calculations: Scalar Exchange

Need $S^0\gamma\gamma^*$ and $S^0\gamma^*\gamma^*$ form factors •

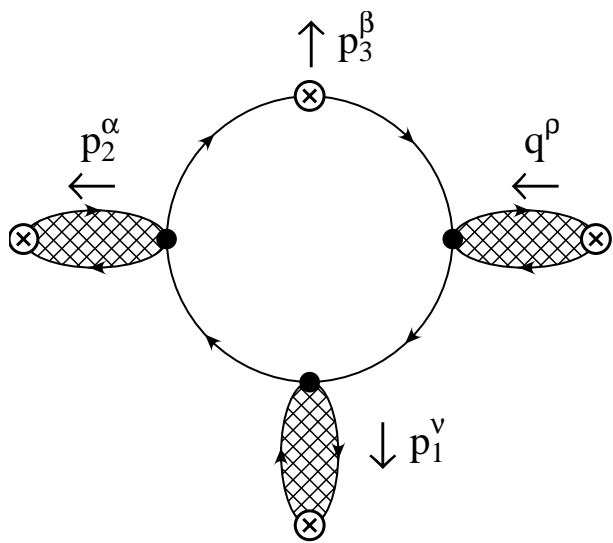
They are constrained by CHPT at $\mathcal{O}(p^4)$: L_i 's reproduced ✓

Within ENJL: Ward identities impose relations between Quark loop and Scalar exchange •

$$a_\mu(\text{Scalar}) = -(0.7 \pm 0.2) \cdot 10^{-10}$$

Not included by M. Hayakawa, T. Kinoshita and A. Sanda
nor by K. Melnikov and A. Vainshtein •

“Old” Calculations: Non-Meson Exchange “Quark-Loop”

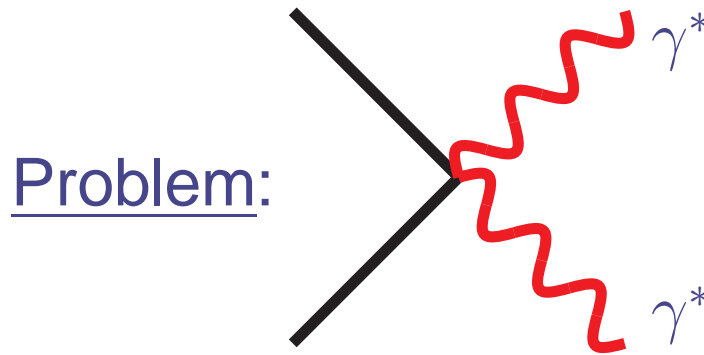
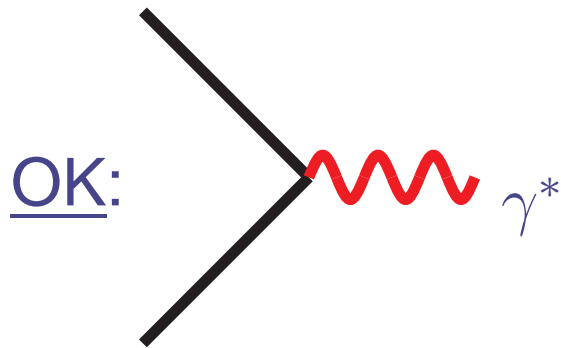


Λ [GeV]	$10^{10} \times a_\mu$
0.7	2.2
1.0	2.0
2.0	1.9
4.0	2.0

- Low Energy (0 to Λ): ENJL model ●
- High Energy (Λ to ∞): Bare heavy quark loop with $m_Q = \Lambda$ ●
- Numerical matching ✓

“Old” Calculations: Pion- and Kaon-Loop

Leading contribution in chiral counting, suppressed by $1/N_c$



No $\gamma^*\gamma^* \rightarrow \pi\pi$ data available: Models needed !

Model for $\pi\pi\gamma(\gamma)$	$10^{10} \times a_\mu$
BPP (Full VMD)	-1.8
HKS (HGS)	-0.4

Kaon loop is much smaller: -0.05×10^{-10} ●

New Short Distance Constraints

K. Melnikov and A. Vainshtein

New short-distance constraint on four-point function form factor

$$\langle 0 | T [V^\nu(p_1) V^\alpha(p_2) V^\rho(-(p_1 + p_2 + p_3))] | \gamma(p_3 \rightarrow 0) \rangle$$

using OPE with $-p_1^2 \simeq -p_2^2 \gg -(p_1 + p_2)^2$ Euclidean and large,

$$T[V^\nu(p_1) V^\alpha(p_2)] \sim \frac{1}{\hat{p}^2} \varepsilon^{\nu\alpha\mu\beta} \hat{p}_\mu [\bar{q} \hat{Q}^2 \gamma_\beta \gamma_5 q](p_1 + p_2)$$

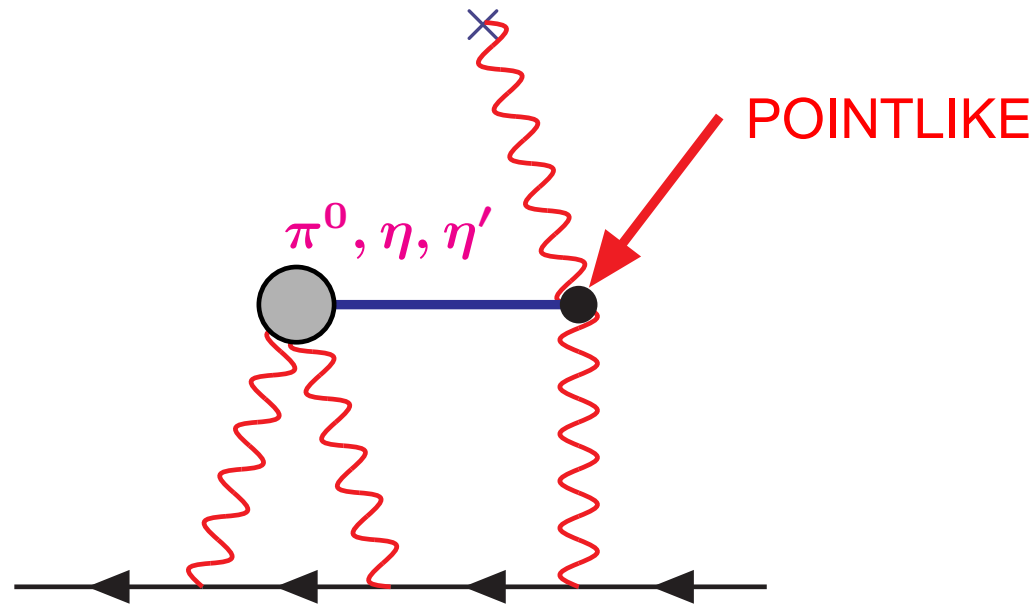
with $\hat{p} = (p_1 - p_2)/2 \simeq p_1 \simeq -p_2$

New OPE Constraint: Pseudo-scalar exchange

New OPE constraint saturated by pseudo-scalar exchange

⇒ Model uses a point-like vertex when $p_3 \rightarrow 0$ •

Not all OPE constraints satisfied: Negligible numerically •



New OPE Constraint: Axial-Vector exchange

Axial-Vector exchange depends very much on the resonance mass mixing •

K. Melnikov and A. Vainsthein:

Ideal mixing for $f_1(1285)$ and $f_1(1420)$ •

Mass mixing	$10^{10} \times a_\mu$
No New OPE (Nonet symmetry)	0.3 ± 0.1
M=1.3 GeV (Nonet symmetry)	0.7
M= M_ρ (Nonet symmetry)	2.8
Ideal mixing	2.2 ± 0.5

Comparison: Leading order in N_c

Leading order in N_c :

Quark Loop + Pseudo-Scalar + Pseudo-Vector + Scalar Exchanges •

Total at $\mathcal{O}(N_c)$	$10^{10} \times a_\mu$
BPP (Nonet symmetry)	$(10.9 \pm 1.9) + \underline{-(0.7 \pm 0.1)} = (10.2 \pm 1.9)$
HKS (Nonet symmetry)	$(9.4 \pm 1.6) + ??\text{Scalar}??$

MV: Hadronic model saturated by pole exchanges:

Cannot compare individual contributions •

Total at $\mathcal{O}(N_c)$	$10^{10} \times a_\mu$
MV (Nonet symmetry)	$(12.1 \pm 1.0) + ??\text{Scalar}??$
MV (Ideal mass mixing)	$(13.6 \pm 1.5) + ??\text{Scalar}??$

Masses produce main difference in pseudo-vector exchange •

Comparison: Momenta Regions for π^0

Study of momenta regions contribution for π^0 exchange

$$P_1^2 = -p_1^2, \quad P_2^2 = -p_2^2, \quad Q^2 = -(p_1 + p_2)^2$$

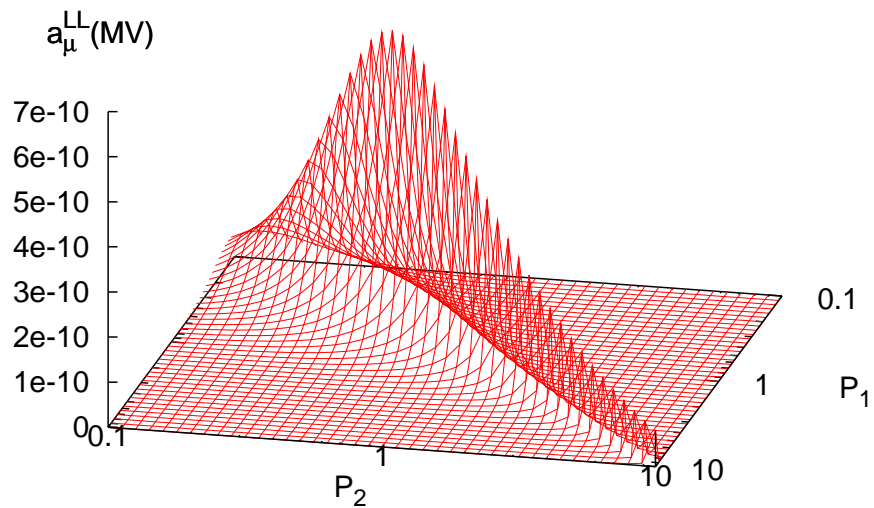
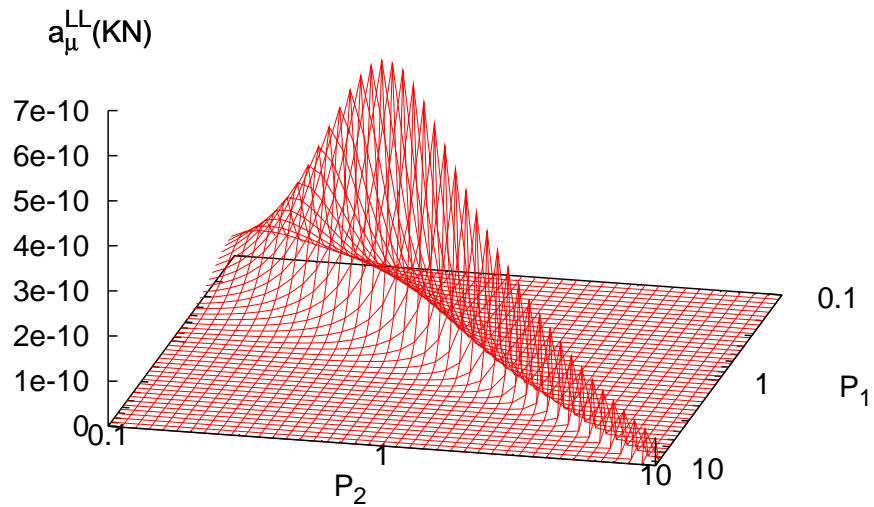
$$a_\mu^{\text{lbl}} = \int dP_1 dP_2 a_\mu^{PP}(P_1, P_2)$$

$$= \int dl_1 dl_2 a_\mu^{LL}(l_1, l_2)$$

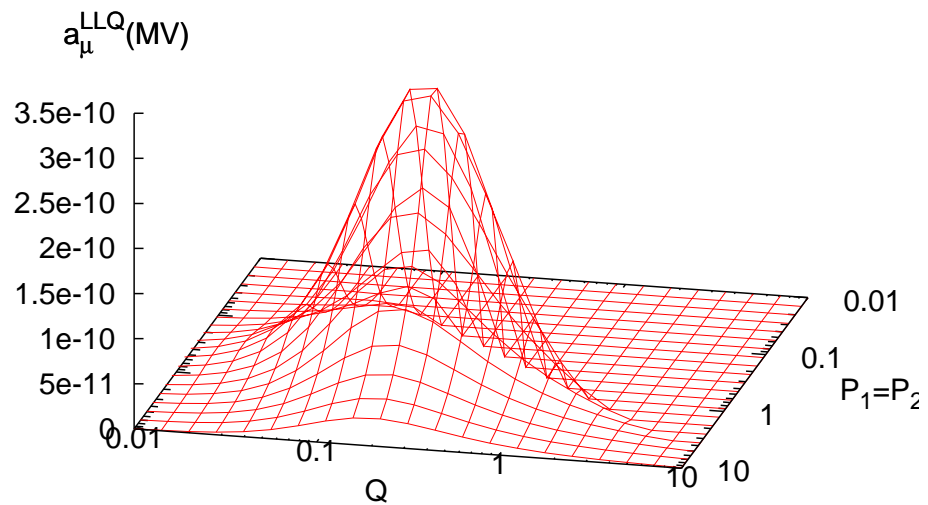
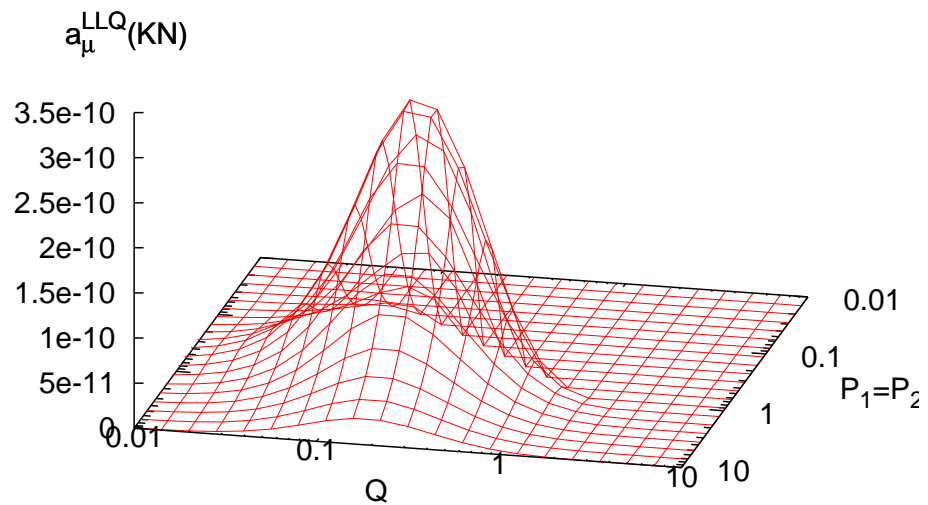
$$= \int dl_1 dl_2 dq a_\mu^{PPQ}(l_1, l_2, q)$$

with $l_1 \equiv \ln(P_1/\text{GeV})$, $l_2 \equiv \ln(P_2/\text{GeV})$ and $q \equiv \ln(Q/\text{GeV})$

Comparison: Momenta Regions for π^0



Comparison: Momenta Regions for π^0



Comparison: Momenta Regions for π^0

Conclusions of this comparison:

Cut-off in Q \Rightarrow 58 % of numerical difference come from Melnikov-Vainshtein OPE constraint violating regions •

Fixing $P_1 = P_2$ \Rightarrow Numerical difference comes from low values of Q and moderate values of $P_1 = P_2$ •

\Rightarrow Important to control energy regions below 2 GeV •

Main ENJL quark-loop contribution is from that region •

Comparison: NLO in $1/N_c$

Next to leading order in $1/N_c$ contributions:

Charged Pion and Kaon Loop •

Model for $\pi\pi\gamma(\gamma)$	$10^{10} \times a_\mu$
BPP (Full VMD)	-1.9 ± 0.5
HKS (HGS)	-0.45 ± 0.8

K. Melnikov and A. Vainshtein:

Full NLO in $1/N_c$ estimate

$$a_\mu = (0 \pm 1) \cdot 10^{-10}$$

Comparison

BPP vs HKS:

Full	$10^{10} \times a_\mu$
BPP	8.3 ± 3.2
HKS	8.9 ± 1.7

No scalar exchange, different quark loop and different pion and kaon loops almost compensate ●

Comparison

BPP vs MV:

Full	$10^{10} \times a_\mu$
BPP	8.3 ± 3.2
MV	13.6 ± 2.5

Several order $1.5 \cdot 10^{-10}$ differences,
in addition to new OPE effects ●

$-1.5 \cdot 10^{-10}$ (Different pseudo-vector mass mixing)

$-0.7 \cdot 10^{-10}$ (No scalar exchange)

$-1.9 \cdot 10^{-10}$ (No pion+kaon loop)

= $-4.1 \cdot 10^{-10}$ ●

Final [BPP-MV] difference: $-5.3 \cdot 10^{-10}$ ●

Conclusions and Prospects

Unsatisfactory situation: Needs new evaluation(s) of the full hadronic light-by-light contribution •

★ At $\mathcal{O}(N_c)$: Study relevant reduced full four-point function with large N_c techniques •

➡ First step: Analytic expressions • Granada-Marseille

➡ Second step: Implement as many short-distance and low energy constraints as possible •

(possible problems J. Bijnens, E. Gámiz, E. Lipartia, J.P.)
Granada-Marseille-Lund-València

Conclusions and Prospects

★ At NLO in $1/N_c$:

⇒ Goldstone bosons at one loop: Need $\gamma^* \gamma^* \rightarrow \pi^+ \pi^-$ vertex ●
Experiment can help here ●

⇒ Non-Goldstone bosons at one loop: Little is known ●
(recent related work by A. Pich, I. Rosell, J.J. Sanz-Cillero) ●

More theoretical work is clearly needed here ●

Conclusions and Prospects

At present, large N_c results agree within 1 σ ✓

$$a_{\mu}^{\text{lbl}} = (11.0 \pm 4.0) \times 10^{-10}$$

More work needed to have the hadronic light-by-light contribution to muon $g - 2$ with reduced uncertainty •

Goal: To have under control model dependences •