

# Large Non-perturbative Effects of Small $\Delta m_{21}^2 / \Delta m_{31}^2$ and $\sin \theta_{13}$

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## Main Results

1. Non-perturbative effects of small  $\Delta m_{21}^2 / \Delta m_{31}^2$  and  $\sin \theta_{13}$  become **more than 50%** at the MSW resonance point.
2. We **propose a method to include these effects** in approximate formulas for oscillation probabilities.

## Two Small Parameters

- Atmospheric Neutrino

$$\theta_{23} \simeq \frac{\pi}{4}, \quad |\Delta m_{31}^2| \simeq 2.0 \times 10^{-3} \text{eV}^2$$

- Solar Neutrino

$$\theta_{12} \simeq \frac{\pi}{6}, \quad \Delta m_{21}^2 \simeq 8.3 \times 10^{-5} \text{eV}^2$$

- CHOOZ

$$\theta_{13} \leq 0.23$$

- ★ Two small parameters

1.  $\alpha = \Delta m_{21}^2 / \Delta m_{31}^2 \simeq 0.03$

2.  $s_{13} = \sin \theta_{13} \leq 0.23$

- ★ Double suppression of CP violating effects

$$P(\nu_e \rightarrow \nu_\mu) = A \cos \delta + B \sin \delta + C$$

$$\lim_{\alpha \rightarrow 0} A = 0 \quad \lim_{s_{13} \rightarrow 0} A = 0$$

$$\lim_{\alpha \rightarrow 0} B = 0 \quad \lim_{s_{13} \rightarrow 0} B = 0$$

$$A = O(\alpha s_{13}) \quad B = O(\alpha s_{13})$$

Remark

It is **important** to **measure** CP violation we take into account **corrections due to small parameters**.

## Second Order Perturbation

★ Second order formula for  $P(\nu_e \rightarrow \nu_\mu)$

$$A \simeq 2\alpha s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \Delta \frac{\sin \bar{a} \sin(\bar{a} - 1) \Delta}{\bar{a} (\bar{a} - 1)}$$

$$B \simeq 2\alpha s_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \frac{\sin \bar{a} \sin(\bar{a} - 1) \Delta}{\bar{a} (\bar{a} - 1)}$$

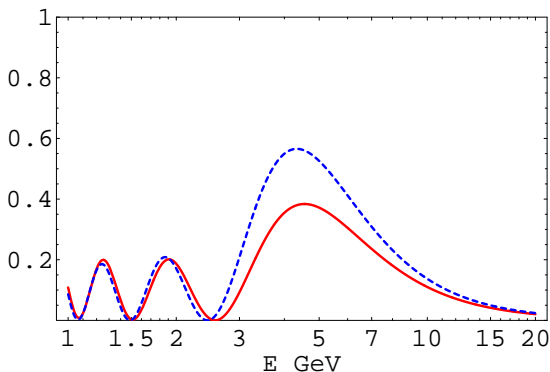
$$C \simeq \alpha^2 \sin^2 2\theta_{12} s_{23}^2 \frac{\sin^2 \bar{a} \Delta}{\bar{a}^2} + 4s_{13}^2 c_{23}^2 \frac{\sin^2(\bar{a} - 1) \Delta}{(\bar{a} - 1)^2}$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}, \quad \bar{a} = \frac{2Ea}{\Delta m_{31}^2}$$

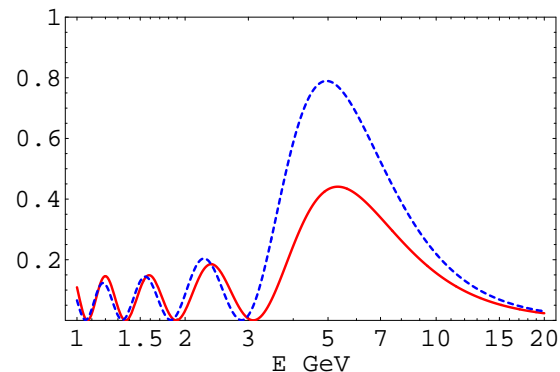
★ Perturbation vs exact calculation

Reference values:  $\delta = 90^\circ$ ,  $\theta_{13} = 0.23$

$L = 4500$  km



$L = 6000$  km



dotted line=perturbation, solid line=exact calculation

Remark

Perturbative calculation is significantly different from numerical calculation **MSW resonance region**.

## Comparison in Two Generations

★ Hamiltonian in two generations

$$H = O \text{diag}(0, \Delta) O^T + \text{diag}(a, 0)$$

★ Second order perturbation due to small mixing

$$P = \frac{4\Delta^2 \sin^2 \theta}{(\Delta - a)^2} \sin^2 \frac{(\Delta - a)L}{2}$$

$$= O(\sin^2 \theta)$$

★ Exact calculation

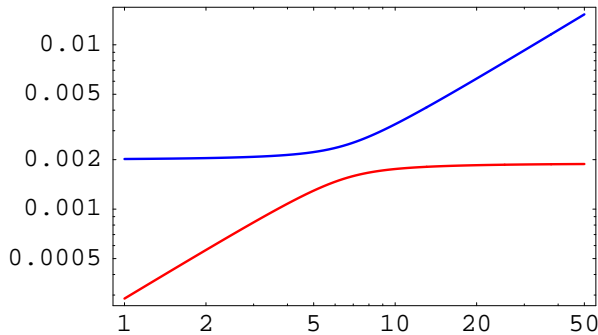
$$P = \frac{\Delta^2}{\tilde{\Delta}^2} \sin^2 \frac{\tilde{\Delta}L}{2}$$

$$= O(\sin^2 \theta) + O(\sin^4 \theta) + \dots$$

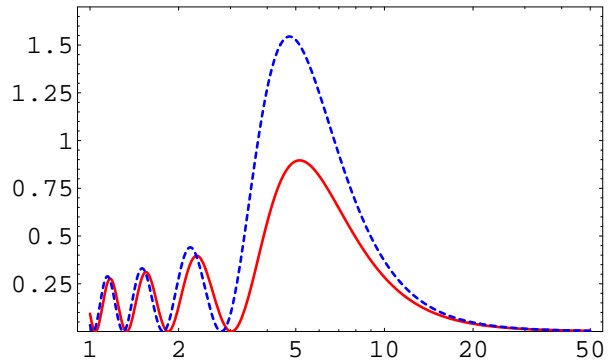
★ Perturbation vs exact calculation

$$\Delta m^2 = 2.0 \times 10^{-3} \text{eV}^2, \quad \sin \theta = 0.23$$

Level Crossing



Probability



dotted line=perturbation, solid line=exact

# Two Energy Regions

## ★ Standard Parametrization

$$U = O_{23}\Gamma_\delta O_{13}\Gamma_\delta^\dagger O_{12}$$

## ★ General relation

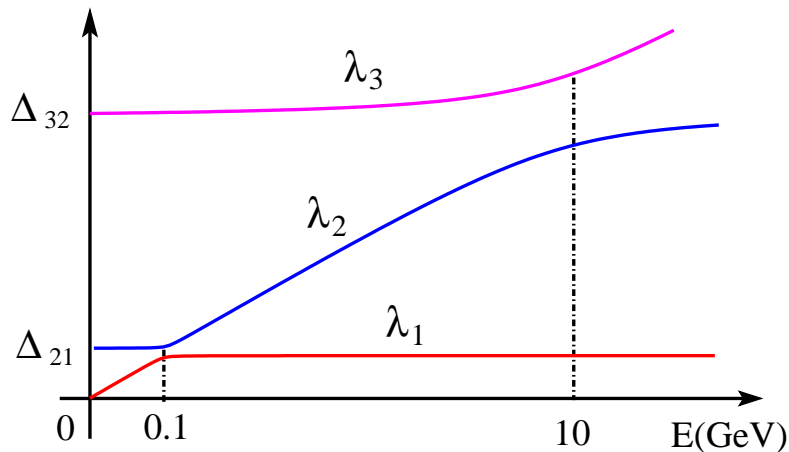
$$H(t) = [O_{23}\Gamma_\delta]H'(t)[O_{23}\Gamma_\delta]^\dagger$$

$$H' = O_{13}O_{12}\text{diag}(0, \Delta_{21}, \Delta_{31})O_{12}^T O_{13}^T + \text{diag}(a(t), 0, 0)$$

## ★ Two energy regions

$(\theta_{12}, \Delta m_{21}^2)$  : Low Energy Phenomenon

$(\theta_{13}, \Delta m_{31}^2)$  : High Energy Phenomenon



### Remark

Parameters  $\alpha$  and  $s_{13}$  control low and high energy MSW resonance, respectively.

## New Proposal

### ★ Double expansion

$$S' = \begin{array}{ccccccc} O(1) & + & O(\alpha) & + & O(\alpha^2) & \dots & \\ O(s_{13}) & + & O(\alpha s_{13}) & + & O(\alpha^2 s_{13}) & \dots & \\ O(s_{13}^2) & + & O(\alpha s_{13}^2) & + & O(\alpha^2 s_{13}^2) & \dots & \\ O(s_{13}^3) & + & O(\alpha s_{13}^3) & + & O(\alpha^2 s_{13}^3) & \dots & \\ \dots & \dots & \dots & \dots & \dots & \dots & \end{array}$$

### ★ Second order perturbation

$$S' = \begin{array}{cccc} O(1) & + & O(\alpha) & + & O(\alpha^2) \\ O(s_{13}) & + & O(\alpha s_{13}) & & \\ O(s_{13}^2) & & & & \end{array}$$

### ★ Our method

$$S' = \begin{array}{ccccccc} O(1) & + & O(\alpha) & + & O(\alpha^2) & + & O(\alpha^3) \dots \\ O(s_{13}) & & & & & & \\ O(s_{13}^2) & & & & & & \\ O(s_{13}^3) & & & & & & \\ O(s_{13}^4) & & & & & & \\ \dots & & & & & & \end{array}$$

### Remark

Approximate formula in our method **contains the higher order terms** of small  $\alpha$  and  $s_{13}$ .

# Two Generation Approximations

## ★ Double expansion

$$\begin{aligned}
 S' &= \begin{array}{cccccc}
 O(1) & + & O(\alpha) & + & O(\alpha^2) & \dots \\
 O(s_{13}) & + & O(\alpha s_{13}) & + & O(\alpha^2 s_{13}) & \dots \\
 O(s_{13}^2) & + & O(\alpha s_{13}^2) & + & O(\alpha^2 s_{13}^2) & \dots \\
 O(s_{13}^3) & + & O(\alpha s_{13}^3) & + & O(\alpha^2 s_{13}^3) & \dots \\
 \dots & & \dots & & \dots & \dots
 \end{array} \\
 &= \lim_{s_{13} \rightarrow 0} S' + \lim_{\alpha \rightarrow 0} S' - \lim_{s_{13}, \alpha \rightarrow 0} S' \\
 &\quad + O(\alpha s_{13}) + O(\alpha^2 s_{13}) + O(\alpha s_{13}^2) + \dots
 \end{aligned}$$

## ★ Our method

$$S' \sim \lim_{s_{13} \rightarrow 0} S' + \lim_{\alpha \rightarrow 0} S' - \lim_{s_{13}, \alpha \rightarrow 0} S'$$

Where

$$\begin{aligned}
 \lim_{s_{13} \rightarrow 0} S' &= \text{T exp} \left( -i \int \lim_{s_{13} \rightarrow 0} H' dt \right) \\
 \lim_{\alpha \rightarrow 0} S' &= \text{T exp} \left( -i \int \lim_{\alpha \rightarrow 0} H' dt \right) \\
 \lim_{s_{13}, \alpha \rightarrow 0} S' &= \text{T exp} \left( -i \int \lim_{s_{13}, \alpha \rightarrow 0} H' dt \right)
 \end{aligned}$$

Remark

We need to calculate Hamiltonian in the **vanishing limit of small parameters**.

# Hamiltonian in Two Generations

## ★ Hamiltonian

$$H' = O_{13}O_{12}\text{diag}(0, \Delta_{21}, \Delta_{31})O_{12}^T O_{13}^T + \text{diag}(a(t), 0, 0)$$

### 1. Vanishing limit of $s_{13}$

$$H^\ell = \lim_{s_{13} \rightarrow 0} H'$$

$$= \begin{pmatrix} \Delta_{21}s_{12}^2 + a(t) & \Delta_{21}s_{12}c_{12} & 0 \\ \Delta_{21}s_{12}c_{12} & \Delta_{21}c_{12}^2 & 0 \\ 0 & 0 & \Delta_{31} \end{pmatrix}$$

### 2. Vanishing limit of $\alpha$

$$H^h = \lim_{\alpha \rightarrow 0} H'$$

$$= \begin{pmatrix} \Delta_{31}s_{13}^2 + a(t) & 0 & \Delta_{31}s_{13}c_{13} \\ 0 & 0 & 0 \\ \Delta_{31}s_{13}c_{13} & 0 & \Delta_{31}c_{13}^2 \end{pmatrix}$$

### 3. Vanishing limit of $s_{13}$ and $\alpha$

$$H^d = \lim_{s_{13}, \alpha \rightarrow 0} H'$$

$$= \begin{pmatrix} \Delta_{31}s_{13}^2 + a(t) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Delta_{31}c_{13}^2 \end{pmatrix}$$



# Main Results

## ★ Zeros of Amplitudes

$$S^l = \begin{pmatrix} S_{ee}^l & S_{e\mu}^l & 0 \\ S_{\mu e}^l & S_{\mu\mu}^l & 0 \\ 0 & 0 & S_{\tau\tau}^l \end{pmatrix} \quad S^h = \begin{pmatrix} S_{ee}^h & 0 & S_{e\tau}^h \\ 0 & S_{\mu\mu}^h & 0 \\ S_{\tau e}^h & 0 & S_{\tau\tau}^h \end{pmatrix}$$
$$S^d = \begin{pmatrix} S_{ee}^d & 0 & 0 \\ 0 & S_{\mu\mu}^d & 0 \\ 0 & 0 & S_{\tau\tau}^d \end{pmatrix} \quad (S_{\mu\mu}^h = S_{\mu\mu}^d, S_{\tau\tau}^l = S_{\tau\tau}^d)$$

## ★ Our method

$$S' \sim S^l + S^h - S^d$$

## ★ Main Results

$$\begin{aligned} S'_{\mu e} &\sim S_{\mu e}^l \\ S'_{\tau e} &\sim S_{\tau e}^h \\ S'_{\tau\mu} &\sim 0 \\ S'_{ee} &\sim S_{ee}^l + S_{ee}^h - S_{ee}^d \\ S'_{\mu\mu} &\sim S_{\mu\mu}^l \\ S'_{\tau\tau} &\sim S_{\tau\tau}^h \end{aligned}$$

### Remark

We obtain **very simple results** except for  $S'_{ee}$ .

# Large Non-perturbative Effects

## ★ Simple CP dependence

$$S(t) = [O_{23}\Gamma_\delta]S'(t)[O_{23}\Gamma_\delta]^\dagger$$

$$P(\nu_e \rightarrow \nu_\mu) = |S(t)_{\mu e}|^2 = A \cos \delta + B \sin \delta + C$$

## ★ Our formula for $P(\nu_e \rightarrow \nu_\mu)$

$$A \simeq \sin 2\theta_{12}^\ell \sin 2\theta_{23} \sin 2\theta_{13}^h \sin \frac{\Delta_{21}^\ell L}{2} \sin \frac{\Delta_{31}^h L}{2} \cos \frac{\Delta_{32} L}{2}$$

$$B \simeq \sin 2\theta_{12}^\ell \sin 2\theta_{23} \sin 2\theta_{13}^h \sin \frac{\Delta_{21}^\ell L}{2} \sin \frac{\Delta_{31}^h L}{2} \sin \frac{\Delta_{32} L}{2}$$

$$C \simeq c_{23}^2 \sin^2 2\theta_{12}^\ell \sin^2 \frac{\Delta_{21}^\ell L}{2} + s_{23}^2 \sin^2 2\theta_{13}^h \sin^2 \frac{\Delta_{31}^h L}{2}$$

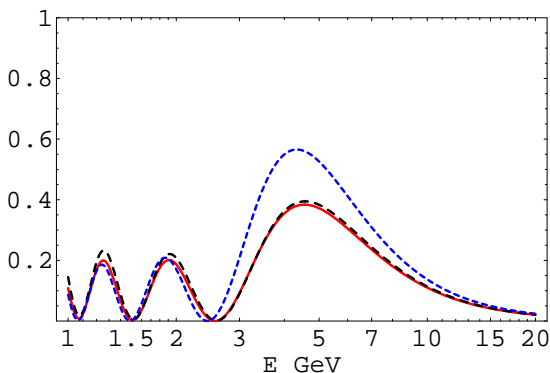
$$\sin 2\theta_{13}^h = O(s_{13}) + O(s_{13}^2) + O(s_{13}^3) + \dots$$

$$\sin \frac{\Delta_{21}^\ell L}{2} = O(1) + O(\alpha) + O(\alpha^2) + O(\alpha^3) + \dots$$

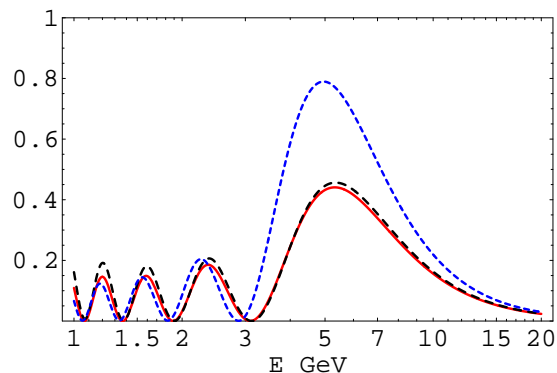
## ★ Our method vs perturbation

Reference values:  $\delta = 90^\circ$ ,  $\theta_{13} = 0.23$

$L = 4500$  km



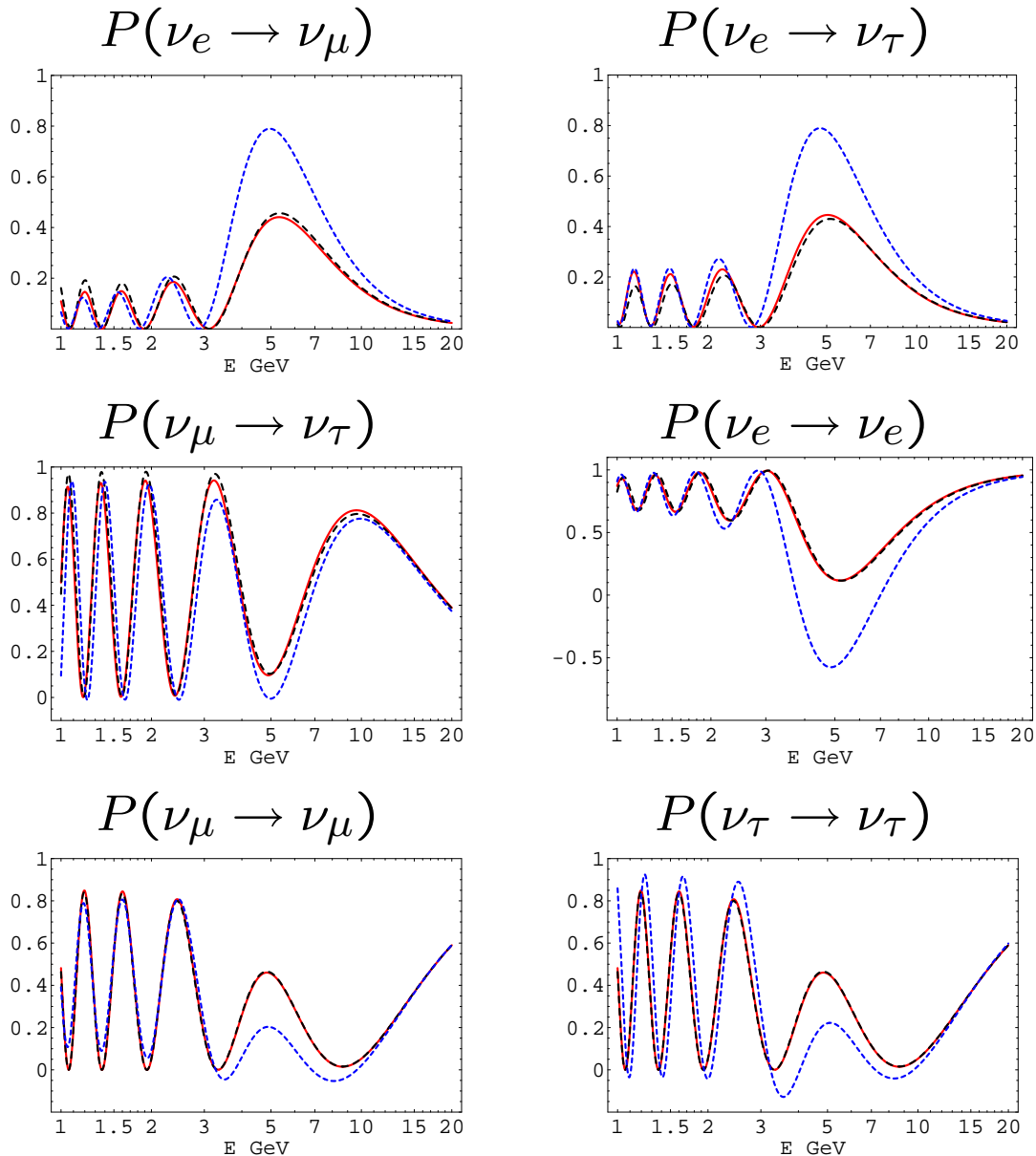
$L = 6000$  km



dotted line=perturbation, line=ours, dashed line=numerical

# Comparison in All Channels

$$L = 6000 \text{ km}, \quad \delta = 90^\circ, \quad \theta_{13} = 0.23$$



dotted line=perturbation, line=ours, dashed line=numerical

Remark

Our formulas are a good agreement with numerical calculation.

# Summary

We have studied non-perturbative effects due to small parameters  $\Delta m_{21}^2 / \Delta m_{31}^2$  and  $\sin \theta_{13}$ .

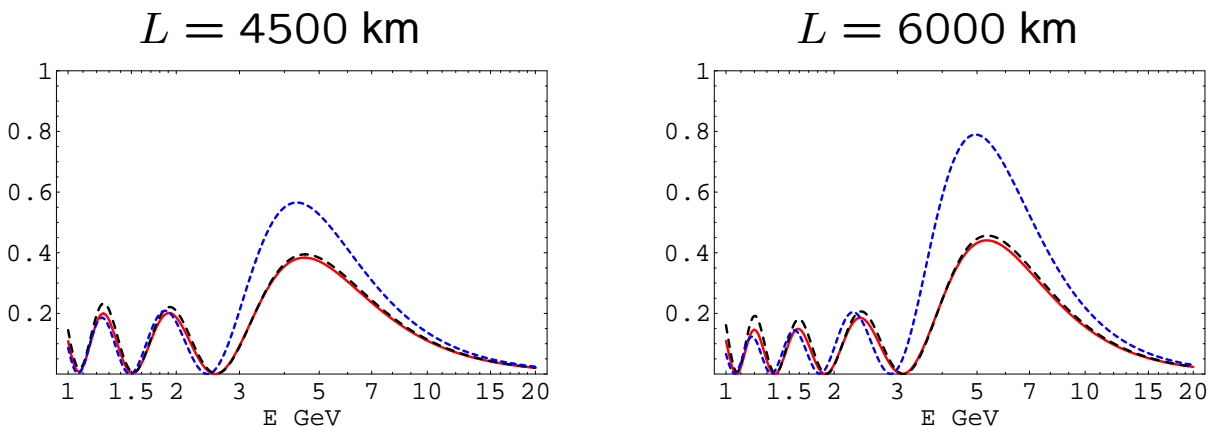
★ Second order perturbation

$$S' = \begin{matrix} O(1) & + & O(\alpha) & + & O(\alpha^2) \\ O(s_{13}) & + & O(\alpha s_{13}) & & \\ O(s_{13}^2) & & & & \end{matrix}$$

★ Our method

$$S' = \begin{matrix} O(1) & + & O(\alpha) & + & O(\alpha^2) & + & O(\alpha^3) & \dots \\ O(s_{13}) & & & & & & & \\ O(s_{13}^2) & & & & & & & \\ O(s_{13}^3) & & & & & & & \\ O(s_{13}^4) & & & & & & & \\ \dots & & & & & & & \end{matrix}$$

★ Our method vs perturbation



dotted line=perturbation, line=ours, dashed line=numerical