### Measurements of the Neutron Magnetic Form Factor with CLAS

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## Overview

- Motivation
- Jefferson Lab Experiment 94-017
- CEBAF Large Acceptance Spectrometer
- Data Analysis
- Preliminary Results
- Summary



#### Motivation

- Nucleon elastic form factors have been calculated using a wide variety of models (e.g., VMD, CQM, RQM, pChQM, Skyrme/soliton, CBM, ...)
- Predictions have been made at high Q<sup>2</sup>using pQCD
- Progress is being made on lattice QCD calculations
- Form factors are related to the generalized parton distributions
- A definitive test of these theoretical efforts is to predict all four nucleon form factors simultaneously
- Limited range and quality of the existing data for  $G_M^n$  reduces the discriminating power of such a test
- High quality data for  $\mathbf{G}_M^n$  over a large  $\mathbf{Q}^2$  range are clearly needed

### **Experimental Technique of E94-017**

- Measure the ratio of quasi-elastic e-n to e-p scattering in deuterium with CLAS in Hall B at Jefferson Lab
- Extract  $\mathbf{G}_M^n$  from ratio using  $\mathbf{G}_E^p$ ,  $\mathbf{G}_M^p$ , and  $\mathbf{G}_E^n$
- Many of the systematic errors inherent in absolute cross section measurements cancel
- Neutron detection efficiency must be calibrated accurately
- Dual-cell deuterium-hydrogen target used in the experiment so that the ep  $\rightarrow$  e'  $\pi^+(n)$  reaction on the hydrogen target is used to measure the neutron detection efficiency

$$R_D = \frac{\frac{d\sigma}{d\Omega} \left[ D\left(e, e'n\right) \right]}{\frac{d\sigma}{d\Omega} \left[ D\left(e, e'p\right) \right]} \approx R_F = \frac{\left[ \frac{G_{En}^2 + \tau G_{Mn}^2}{1 + \tau} + 2\tau G_{Mn}^2 \tan^2\left(\frac{\theta}{2}\right) \right]}{\left[ \frac{G_{Ep}^2 + \tau G_{Mp}^2}{1 + \tau} + 2\tau G_{Mp}^2 \tan^2\left(\frac{\theta}{2}\right) \right]}$$







#### **Particle Detection and ID**





- Charged particle  $\theta = 8^{\circ} 144^{\circ}$
- Neutral particle  $\theta = 8^{\circ} 75^{\circ} (144^{\circ})$
- Charged particle momentum resolution  $\sim 0.5\%$
- Charged particle angular resolution  $\sim 0.5$  mrad
- Particle ID (e<sup>-</sup>, p,  $\pi^+$ ,  $\pi^-$ , K<sup>+</sup>,  $\gamma$ , n,...)

## The CLAS E5 Run

- Approximately 2.3 billion triggers were acquired in 30 days of operation
- $Q^2$  range 0.3 4.5  $(\text{GeV/c})^2$
- Three different running conditions provided overlapping  $Q^2$  coverage
  - 4.2-GeV beam with magnetic field at 90% of maximum reaches high  $Q^2$
  - 2.5-GeV beam with field at 60% covers intermediate  $Q^2$  range
  - 2.5-GeV beam with field at 60% but with reversed polarity to reach lowest  $Q^2$

## **Benefits of CLAS**

- Numerous cross-checks
  - Three independent, overlapping measurements of e-n
  - Three different beam energy and magnetic field combinations allow independent, overlapping measurements of e-p
  - Multiple, overlapping measurements of  $\mathbf{G}_M^n$
- Accommodates dual-cell target
  - In-situ neutron tagging
  - In-situ proton elastic scattering (proton detection efficiency, momentum corrections, and alignment)
- Inelastic background suppression using information from detected nucleons





# **Z-Vertex Distribution for e-p Events**



### **EC** Neutron Detection Efficiency







## Ratio of Quasi-Elastic e-n to e-p Scattering in Deuterium



#### **Corrections to the Ratio**

 $R_C\left(Q^2\right) = c_{nuc}\left(Q^2\right)c_{rad}\left(Q^2\right)c_{fermi}\left(Q^2\right)R_D\left(Q^2\right)$ 

- $c_{fermi}(Q^2)$  corrections for losses near the edge of the acceptance due to fermi motion in the target estimated with Monte-Carlo simulation
- $c_{rad}(Q^2)$  radiative corrections performed with a modified version of EXCLURAD (Afanasev *et al.*)
- $c_{nuc}(Q^2)$  nuclear corrections performed using calculations by Arenhovel for  $Q^2 \leq 1$  (GeV/c)<sup>2</sup> and Jeschonnek for  $Q^2 \geq 1$ (GeV/c)<sup>2</sup>

### Extraction of $\mathbf{G}_M^n$ from the Ratio

$$R_{C} = \frac{\sigma_{Mott}^{n} \left(G_{En}^{2} + \frac{\tau_{n}}{\epsilon_{n}}G_{Mn}^{2}\right) \left(\frac{1}{1+\tau_{n}}\right)}{\sigma_{Mott}^{p} \left(G_{Ep}^{2} + \frac{\tau_{p}}{\epsilon_{p}}G_{Mp}^{2}\right) \left(\frac{1}{1+\tau_{p}}\right)}$$

$$G_M^n = \sqrt{\left[R_C\left(\frac{\sigma_{Mott}^p}{\sigma_{Mott}^n}\right)\left(\frac{1+\tau_n}{1+\tau_p}\right)\left(G_{Ep}^2 + \frac{\tau_p}{\epsilon_p}G_{Mp}^2\right) - G_{En}^2\right]\frac{\epsilon_n}{\tau_n}}$$

- Arrington parametrization used for  $\mathbf{G}_E^p$  and  $\mathbf{G}_M^p$
- Galster parametrization used for  $\mathbf{G}_E^n$

## **Sources of Systematic Errors**

- Uncertainties in the other form factors
- Detection efficiencies
- Inelastic background
- Acceptance
- Fermi loss corrections
- Radiative corrections
- Nuclear corrections



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#### Summary

- $G_M^n$  has been extracted from the ratio of quasi-elastic e-n to e-p scattering from deuterium over a broad range of  $Q^2$
- A dual-cell H-D target was used allowing in-situ measurements of the neutron detection efficiency using the ep  $\rightarrow e' \pi^+(n)$  reaction
- The use of multiple beam energies, magnetic field settings, and neutron detectors provided independent measurements over overlapping ranges in  $Q^2$
- Consistency of overlapping measurements indicates that systematic errors are under control
- The dipole parametrization provides a good description of the data at  $Q^2>1~({\rm GeV/c})^2$
- Technique can be extended to higher  $Q^2$