

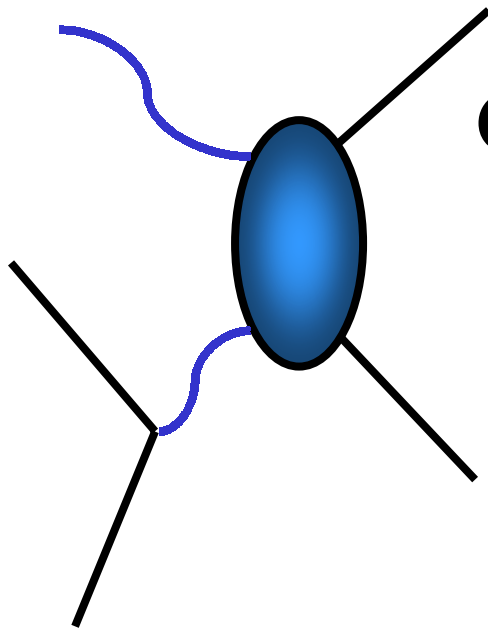
# N05

Workshop on Nucleon Form Factors

INFN-Frascati, Italy

12-14 October 2005

## Recent Results in Real and Virtual Compton Scattering at Jefferson Lab

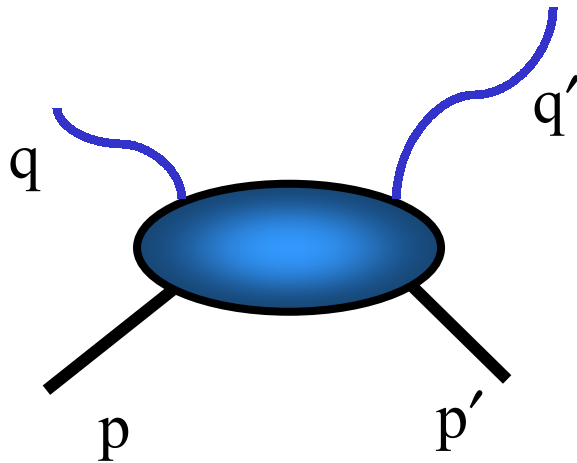


Charles E. Hyde-Wright  
*Old Dominion University*  
*Norfolk VA 23529 USA*  
chyde@odu.edu



# Compton Scattering

Different physics highlighted in different kinematic domains.



$$s = (q+p)^2,$$

$$u = (p-q')^2$$

$$t = (q-q')^2$$

## Real Compton Scattering

Forward limit

Dispersion Relations,

Low Energy Theorem:

GDH Sum Rule

Electric, Magnetic, & Spin

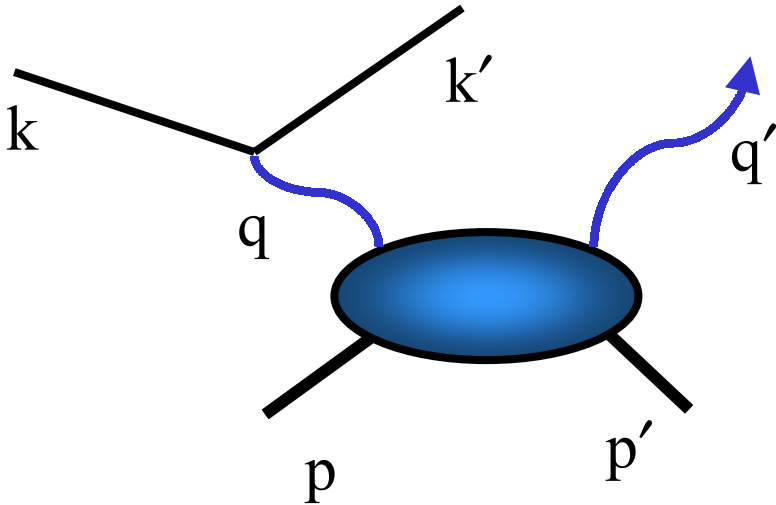
Polarizabilities

High Energy,

Wide Angle Compton Scattering

$s-M^2$ ,  $M^2-u$ ,  $-t \gg \Lambda_{\text{QCD}}^2$

# Virtual Compton Scattering



$$s = (q+p)^2,$$

$$Q^2 = -q^2,$$

$$u = (p-q')^2$$

$$t = (q-q')^2$$

Low Energy Theorem:

(P. Guichon et al)

Generalized Polarizabilities

Dispersion Relations,

(B. Pasquini et al)

Deeply Virtual Compton Scattering

$s \sim M^2$ ,  $Q^2 \gg \Lambda_{\text{QCD}}^2$

$-t \ll Q^2$

Generalized Parton Distributions

Spatial imaging of quarks & gluons

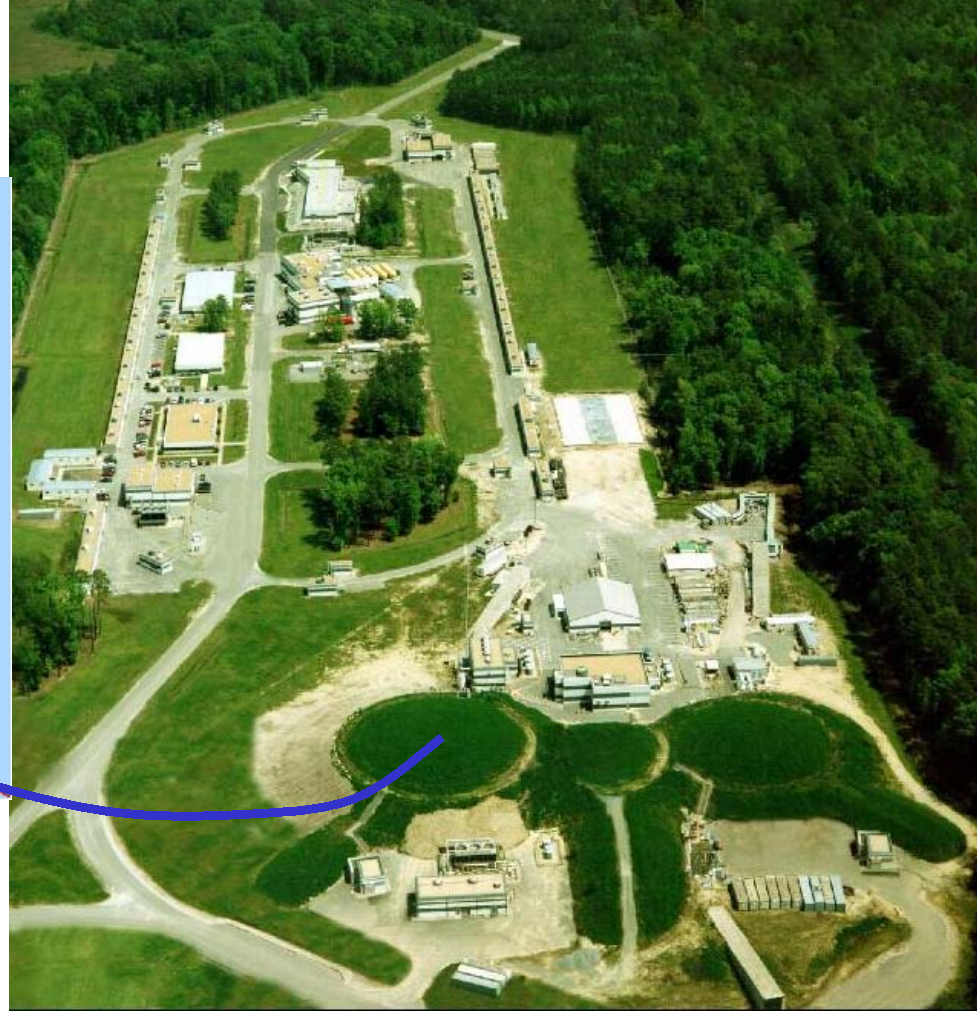
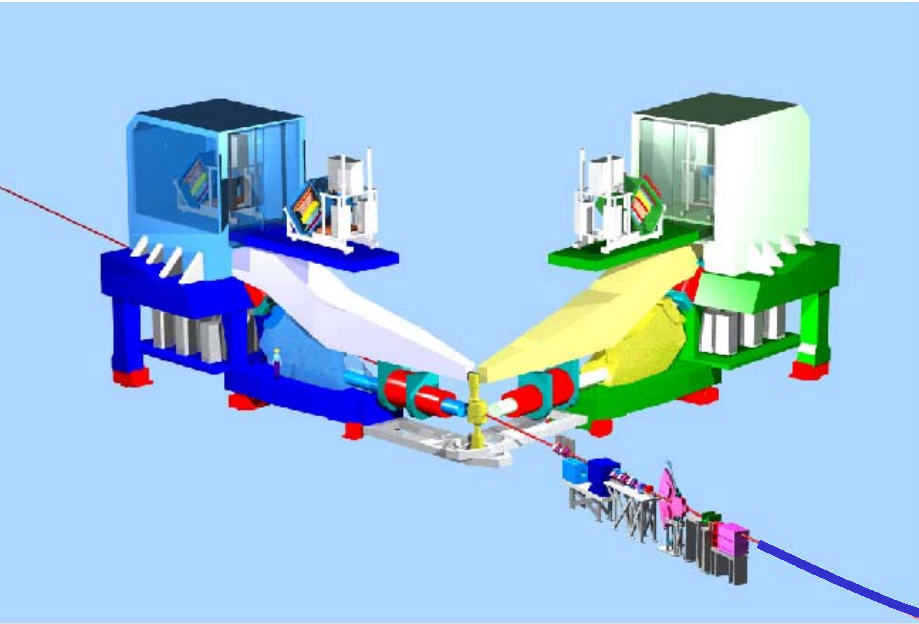
Wide angle Compton Scattering

$s, Q^2, -t$  large, no constraint on  $u$ .

# Jefferson Lab Continuous Electron Beam Accelerator Facility

Hall A

High Resolution Spectrometers (HRS<sup>2</sup>)



VCS:  $H(e,e'p)X$

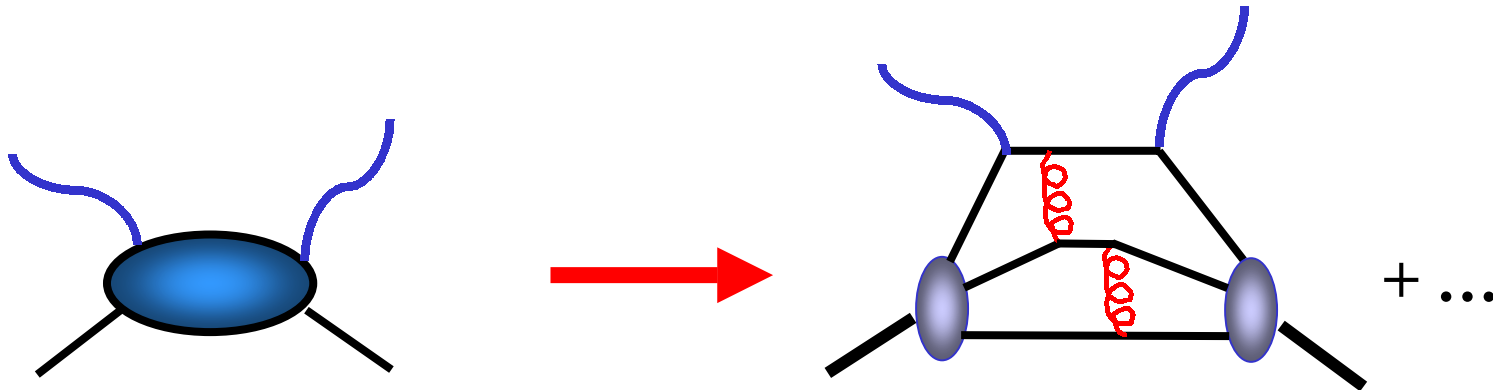
$100\mu\text{A } e^-$  on 15 cm liquid  $\text{H}_2$  target

RCS:  $H(\gamma,\gamma'p)$

$\leq 40\mu\text{A } e^-$  on Cu radiator + 15 cm liquid  $\text{H}_2$  target

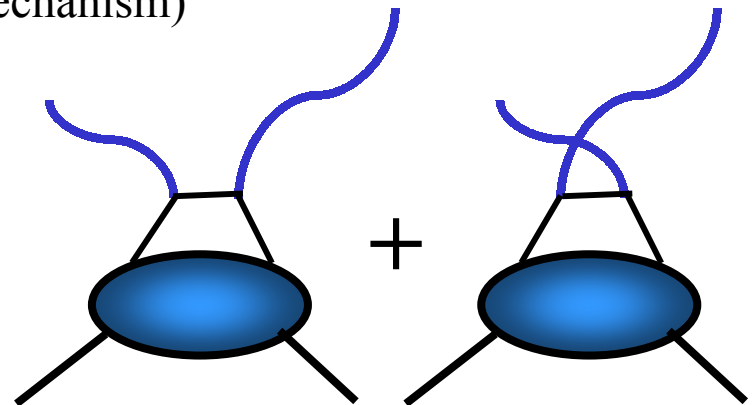
# Wide Angle Compton Scattering

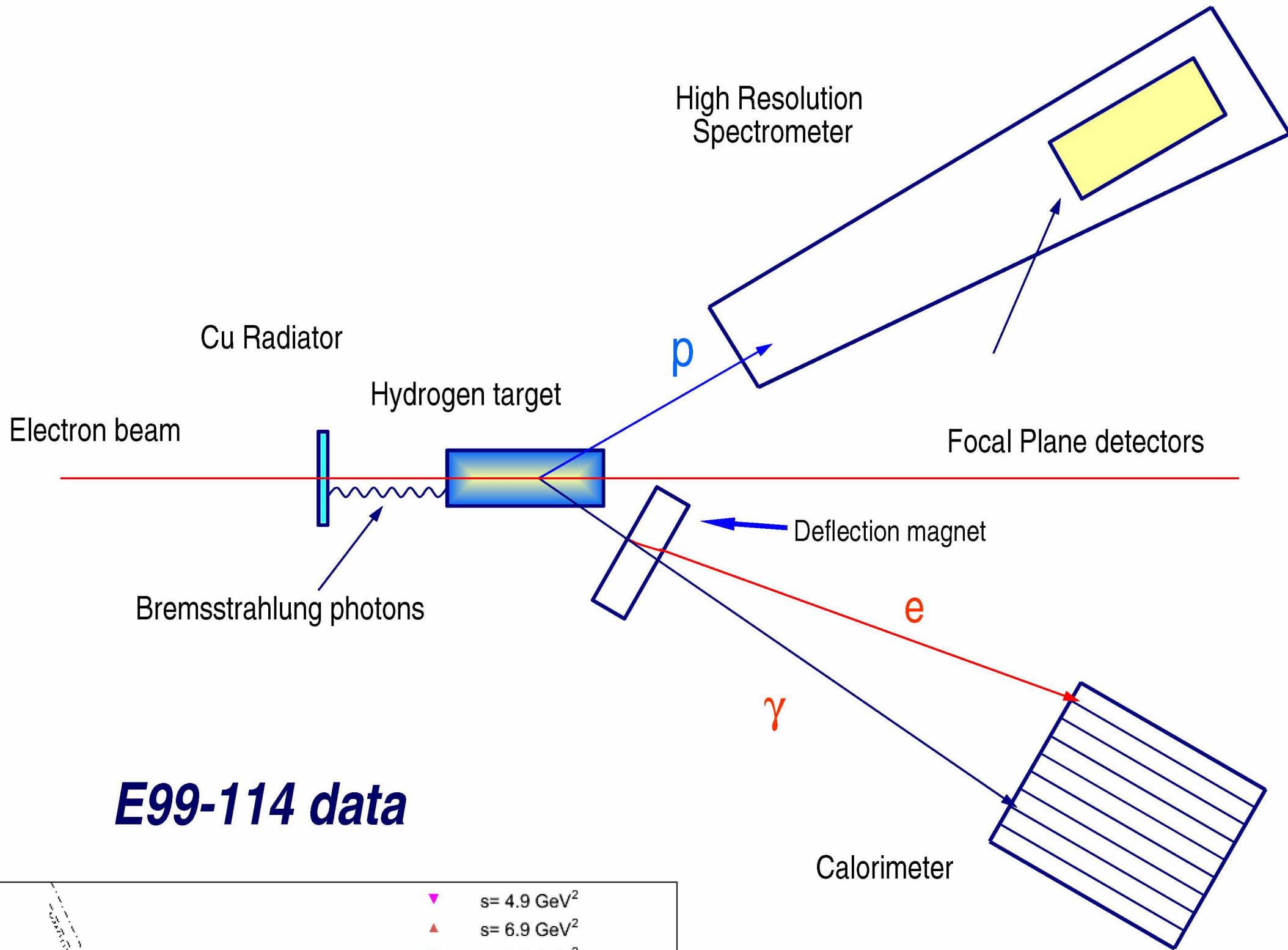
- **Asymptotic limit:**  $s \rightarrow M^2, -t, M^2-u \rightarrow \infty \dots$ 
  - “pure” Perturbative QCD (G.Farrar & S.J.Brodsky...)



- **Sub-asymptotic:**  $s \rightarrow M^2, -t, M^2-u \leq 10 \text{ GeV}^2$

- Conjecture: Handbag dominance (Feynman mechanism)
- Radyushkin, Kroll: Factorized ansatz (GPD).
  - Quark propagator  $(xP+k_{\perp}+q)^2 \rightarrow xP \cdot q$
- G.Miller, unfactorized
  - Exact Klein-Nishina amplitude
    - Constituent Quarks
  - Wave function ansatz



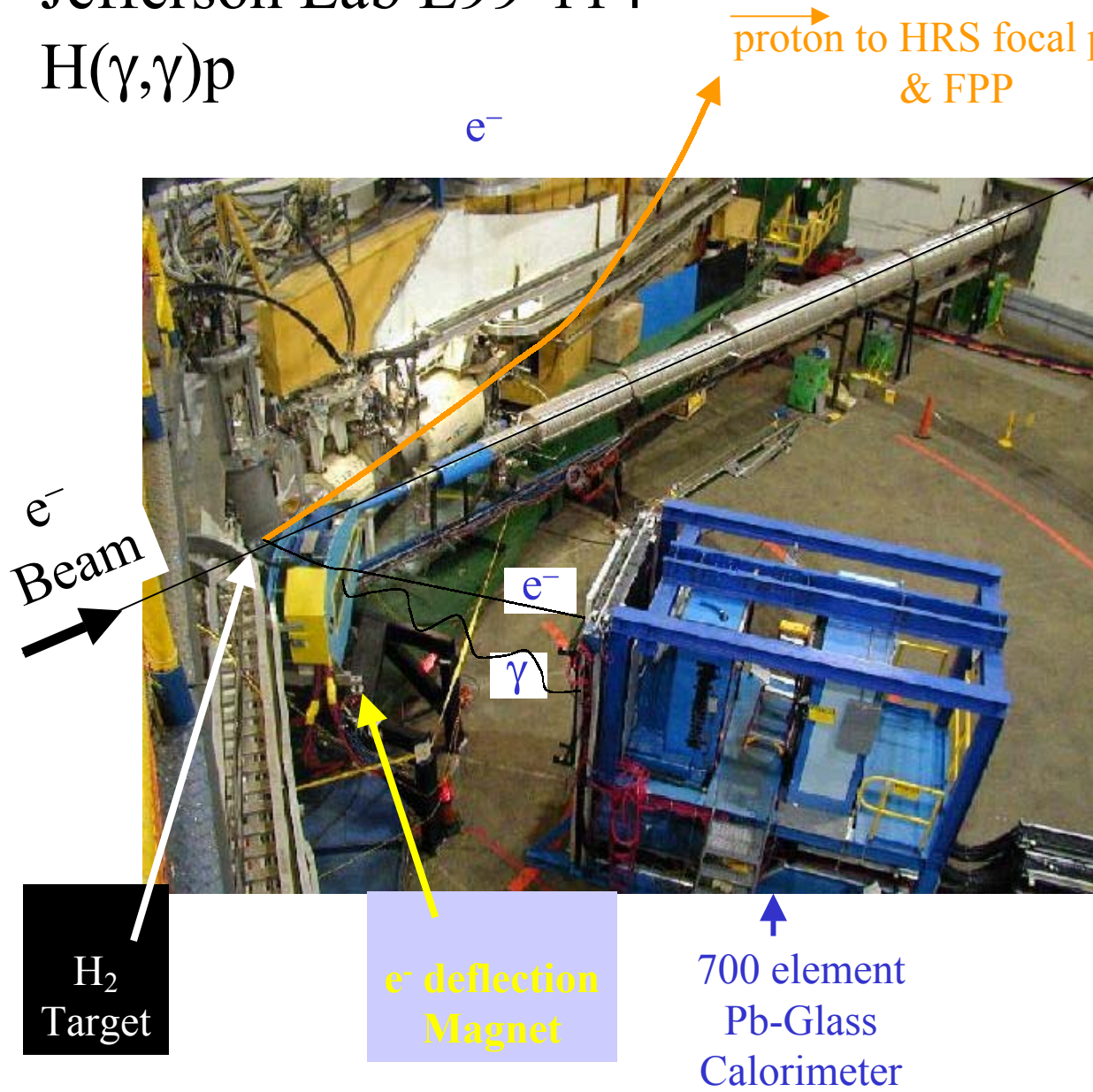


## ***E99-114 data***

- ▼  $s = 4.9 \text{ GeV}^2$
- ▲  $s = 6.9 \text{ GeV}^2$
- $s = 8.9 \text{ GeV}^2$

# Jefferson Lab E99-114

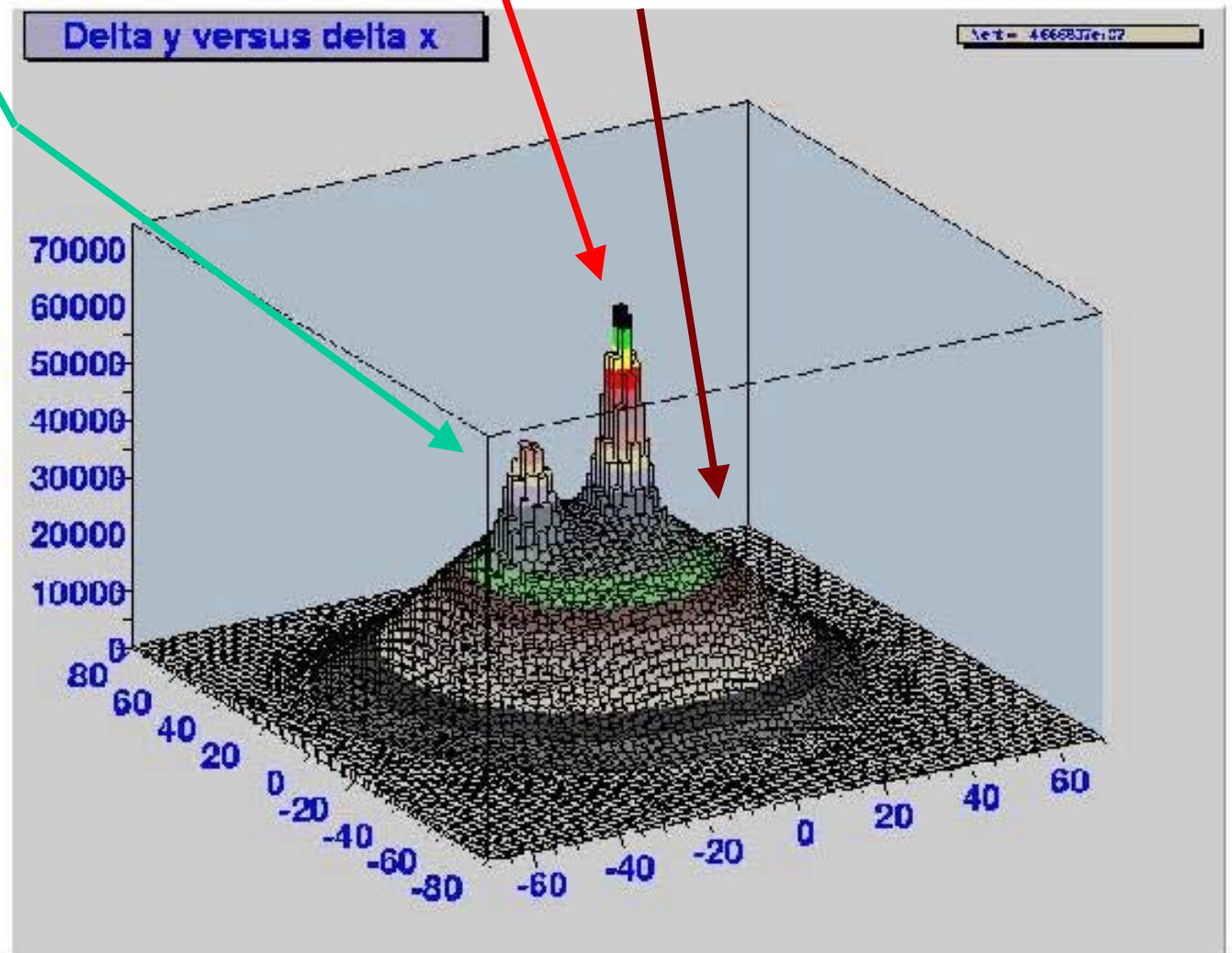
## $H(\gamma,\gamma)p$



- 10-40 $\mu$ A beam on 6% radiator and 1g/cm<sup>2</sup> H<sub>2</sub> (mixed  $\gamma$ +e beam)
- Recoil Proton detected in HRS
  - Focal Plane Polarimeter
- Scattered photon detected in 700 element Pb-Glass array
- Scattered electron deflected by magnet and detected in Pb-Glass.

Electron,

Compton,  $\pi^0$  events



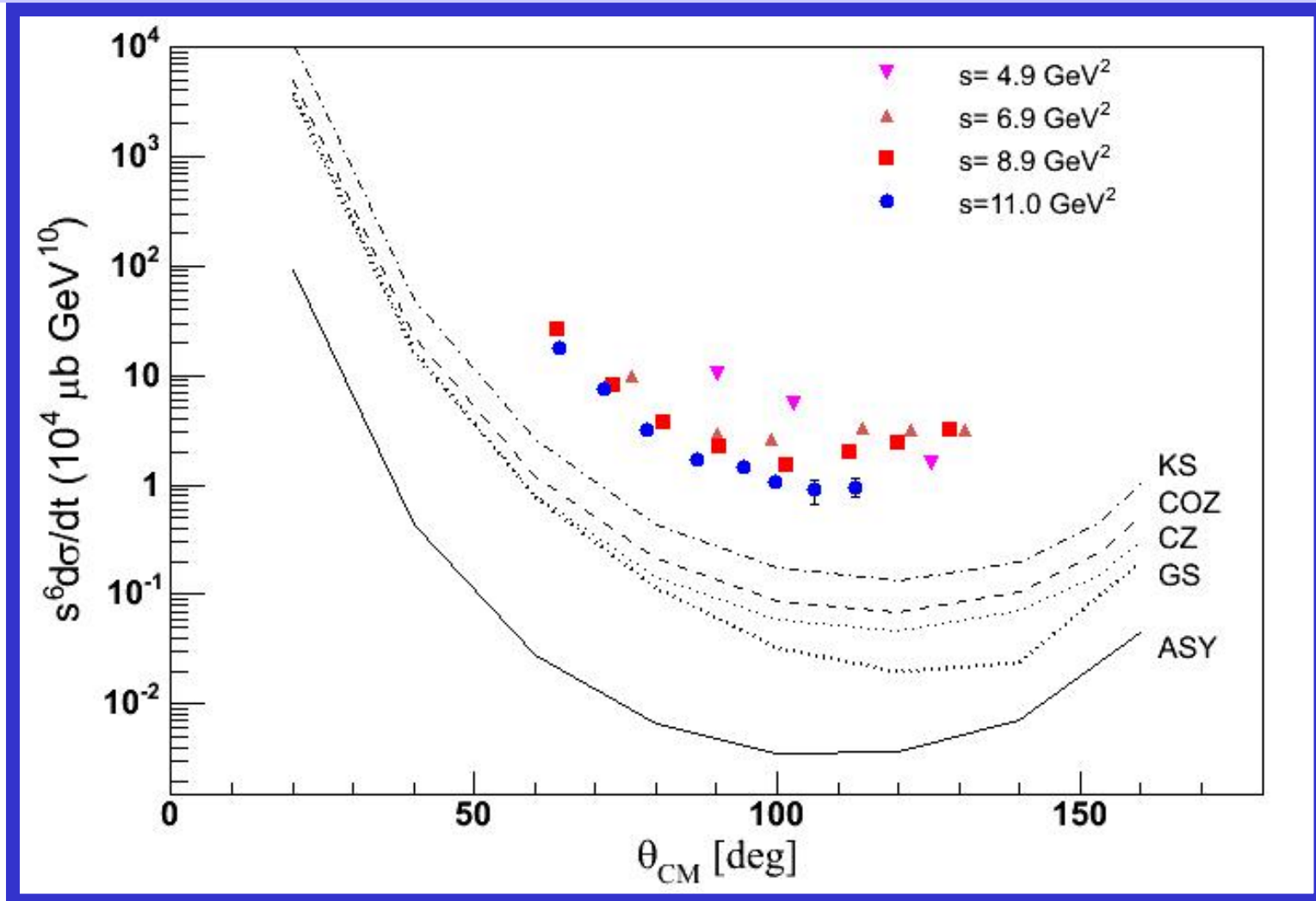
- 2-D plot of position of EM shower, relative to elastic prediction from recoil proton kinematics



# RCS Differential Cross Sections

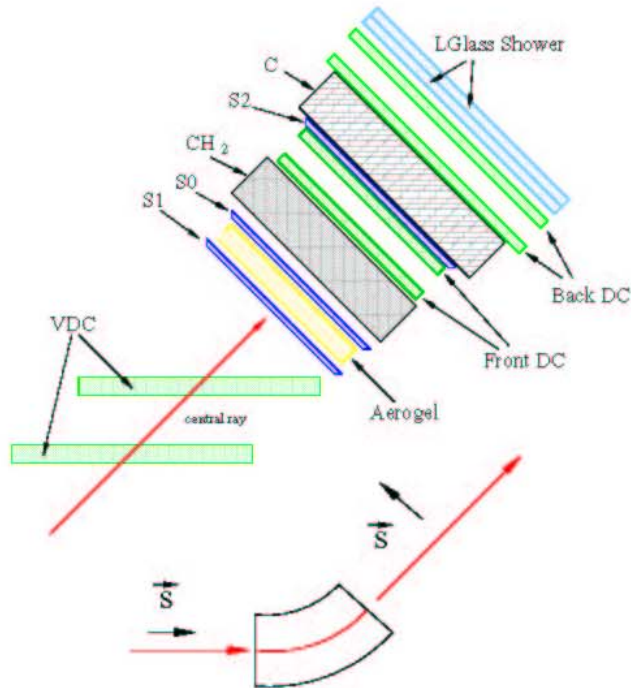
JLab E99-114, 25 kinematic points w/ statistical errors.

2 gluon exchange pQCD calculations only generate a small piece of the scattering amplitude, even with highly asymmetric Distribution Amplitudes



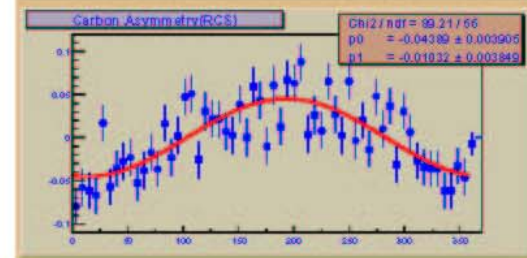
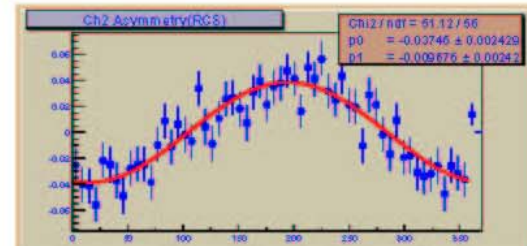
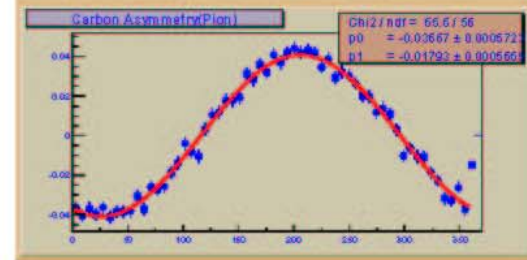
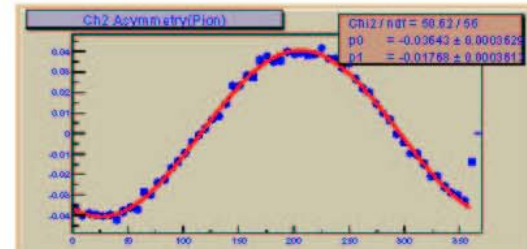
# Asymmetries for Pion and Compton processes

## Proton Spectrometer and Focal Plane Polarimeter



$$\Theta_{sp} = 86^\circ \cdot E_p(\text{GeV}) \cdot \frac{\theta_{\text{deflect}}}{45^\circ}$$

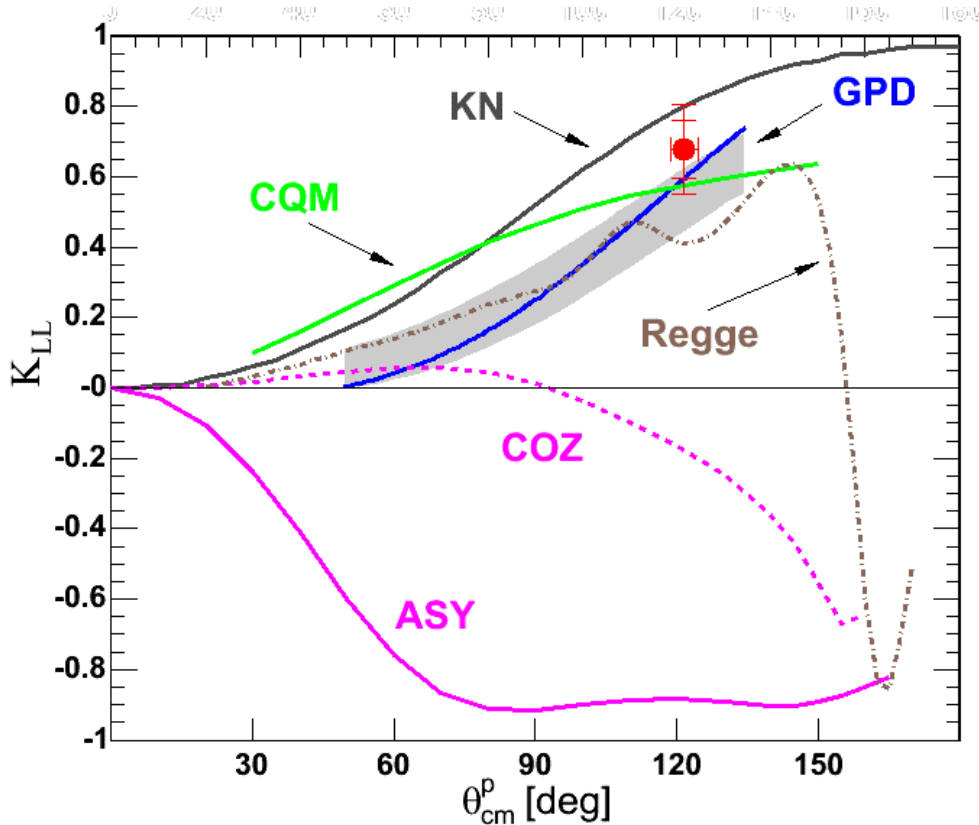
in E99-114  $\Theta_{sp}$  at focal plane is 270 degree. Asymmetries are proportional to  $p_L$  and  $p_S$  of the proton polarization.



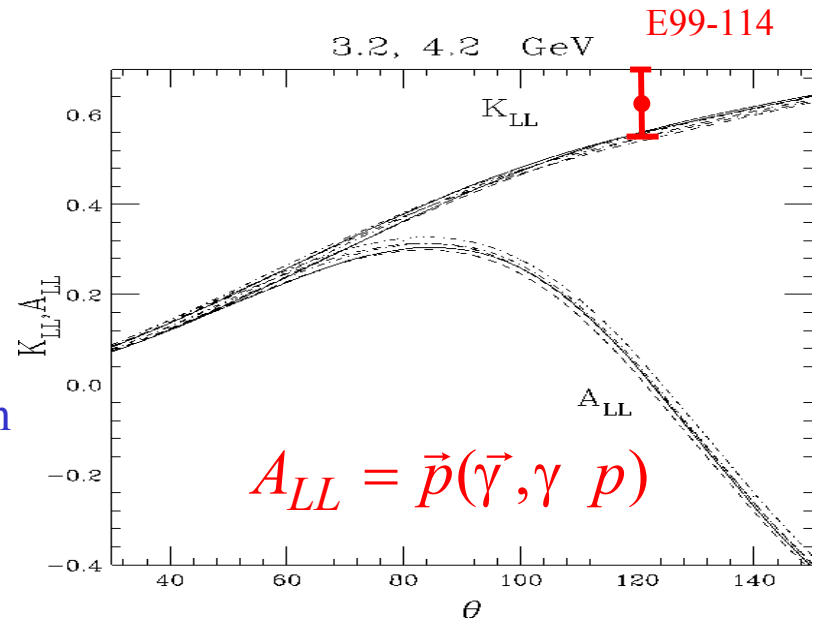
# Longitudinal Polarization Transfer in RCS

$$K_{LL} = p(\vec{\gamma}, \gamma \vec{p})$$

$$E_\gamma = 3 \text{ GeV}$$

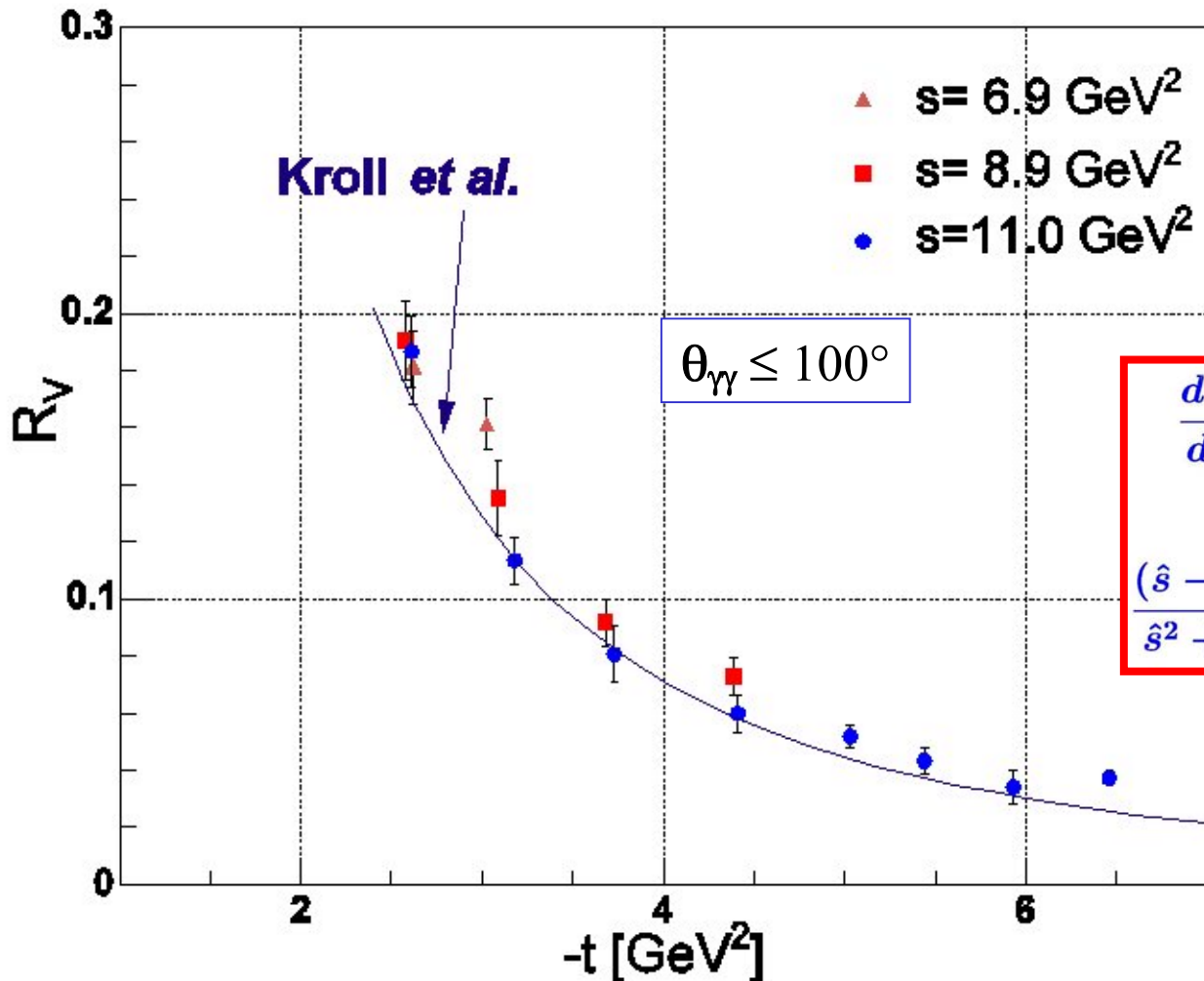


CQM=G.A.Miller PRC 69(2004)



- Experimental evidence for Handbag mechanism
  - CQM & GPD Calculations
- pQCD (2Gluon exchange) strongly ruled out

# Handbag Amplitude: Klein-Nishina Scaling at fixed $-t$ ?



$$\frac{d\sigma_{\text{RCS}}}{d\sigma_{\text{KN}}} = \frac{2\hat{s}\hat{u}}{\hat{s}^2 + \hat{u}^2} R_A^2(t) +$$

$$\frac{(\hat{s} - \hat{u})^2}{\hat{s}^2 + \hat{u}^2} \left[ R_V^2(t) + \frac{-t}{4m^2} R_T^2(t) \right]$$

$\gamma p \rightarrow \gamma p$

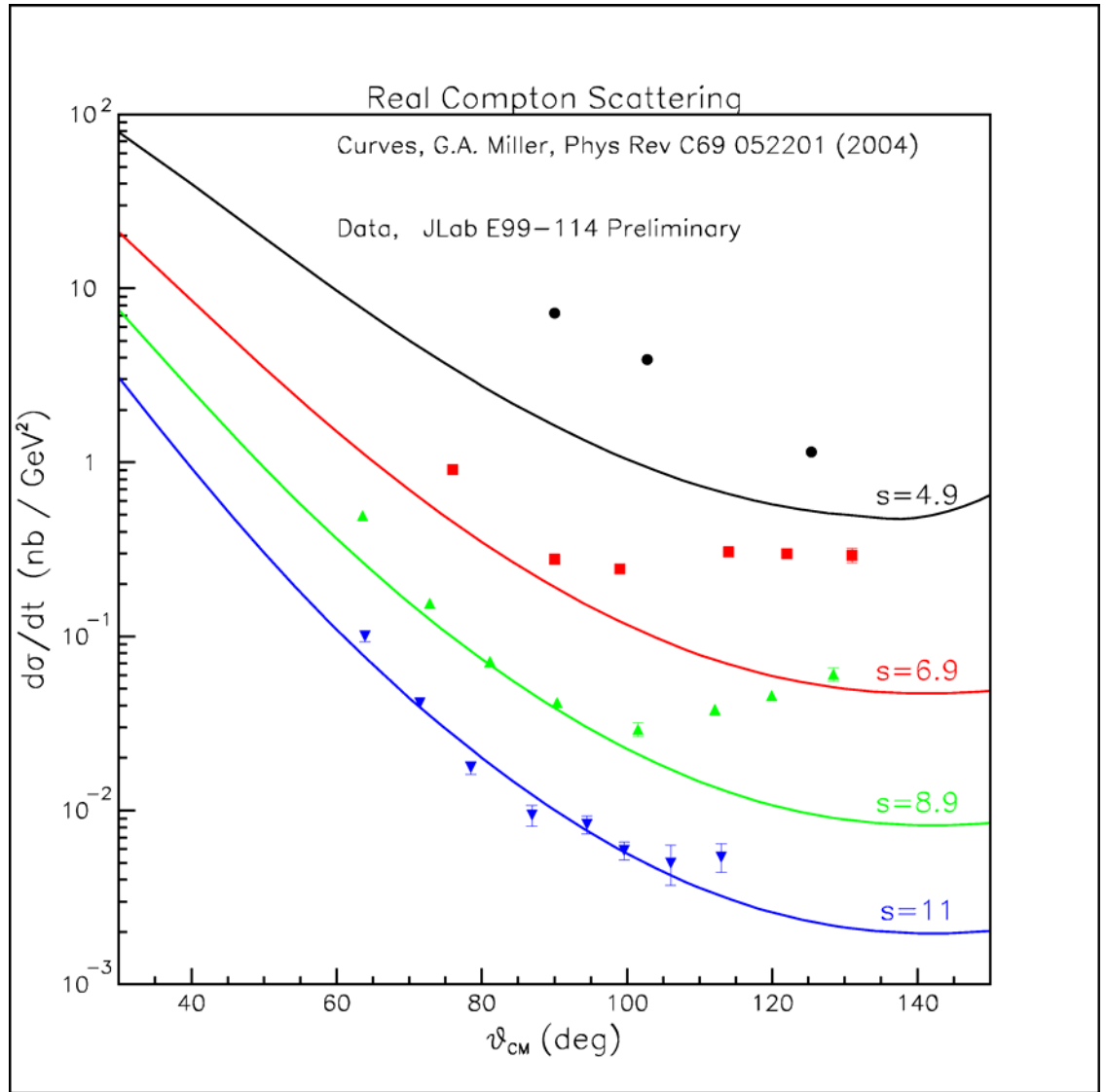
E99-114 Statistical errors only

Light Front Cloudy Bag Model  
(only 3-quark content included  
at large x for RCS)

Handbag amplitude &  
Wavefunction ansatz

$$\Phi(k_1, k_2, k_3) = \frac{(\text{spin})N}{\left[ M_{123}^2 + \beta^2 \right]^\gamma}$$

Parameters  $\beta, \gamma, m_q$   
fitted to  $H(e, e')p$

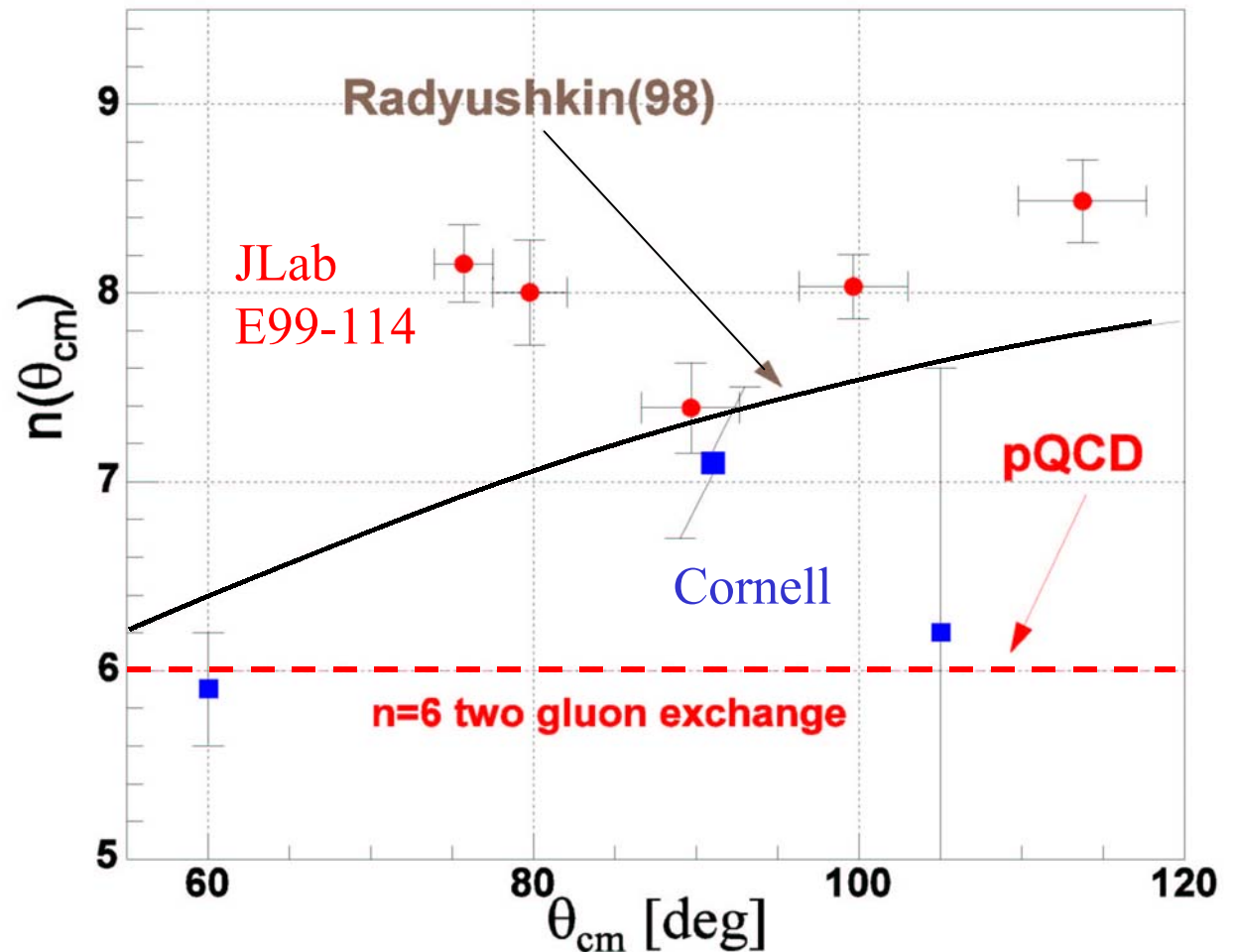


# Cross section scaling at fixed $\theta_{\gamma\gamma}^{\text{CM}}$

$$\frac{d\sigma}{dt} \rightarrow \frac{f(\theta)}{s^n}$$

$n+2 =$  Sum of # of elementary fields in initial & final state

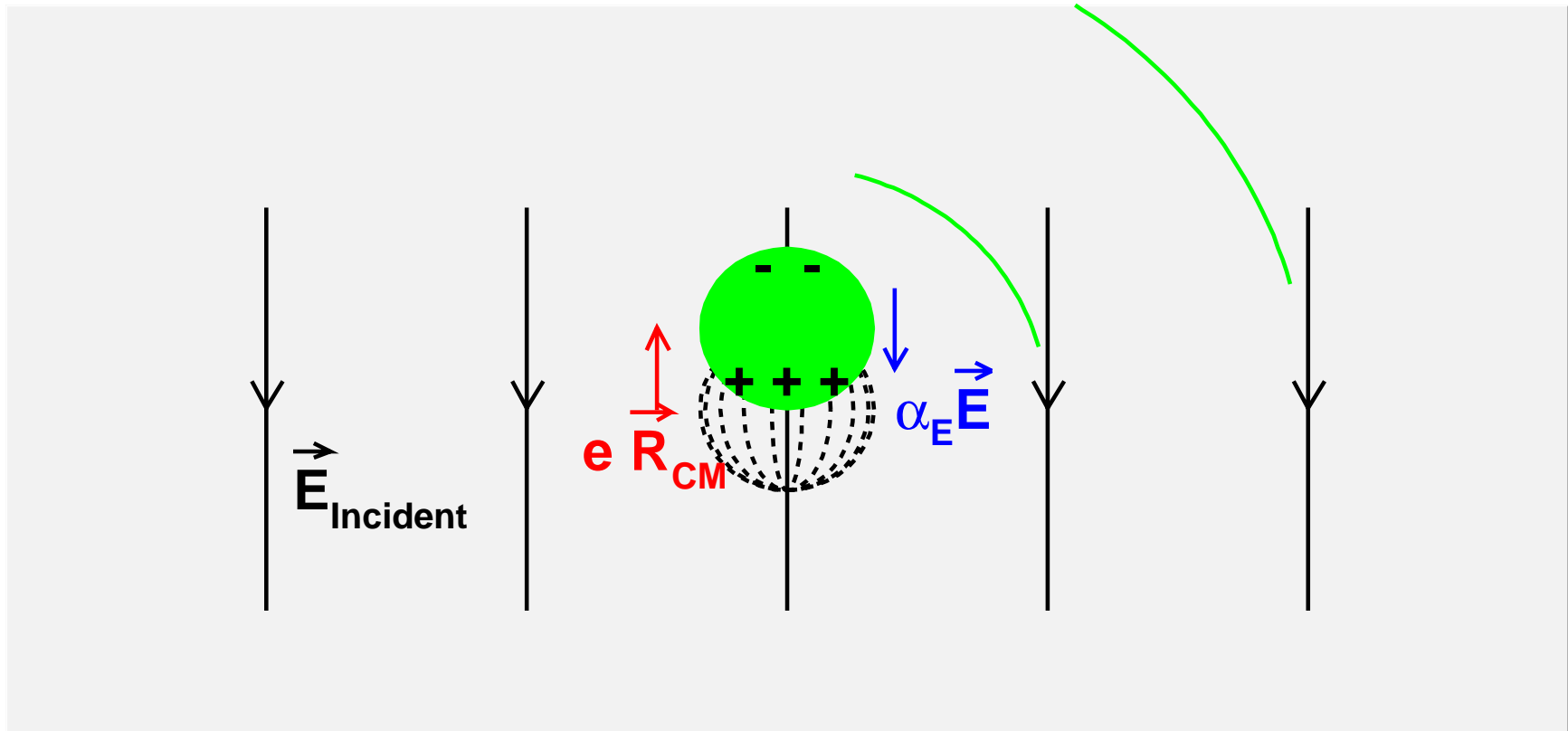
$p_{\perp}$  largest at  $\theta^{\text{CM}}=90^{\circ}$



# Future Prospects for RCS

- Ales Psaker, A. Radyushkin (ODU):
  - Improved treatment of  $k_{\perp}$  in handbag amplitude
- C. Weiss: Massless quarks in constituent quark wavefunction of proton.
- Measurement of  $A_{LL}$  at 4 GeV and  $120^{\circ}$  (Hall C)
- Double kinematic range in  $s$  and  $t$  with JLab @ 12 GeV

# Low Energy Compton Scattering



Power radiated: 
$$P = \frac{2}{3c^3} \left| \frac{\partial^2 \vec{d}}{\partial t^2} \right|^2$$

Center of Mass motion and induced dipole: 
$$\frac{\partial^2 \vec{d}}{\partial t^2} = \frac{e^2}{m} \mathbf{E} - \alpha_E \omega^2 \mathbf{E}$$



## Forward and Backward Low Energy Scattering Amplitudes

D. Babusci, *et al.*, *Phys. Rev.* **C58**, 1013 (1998):

$$\begin{aligned}
 \frac{1}{8\pi M} T_{fi}(0^\circ) &= \vec{\epsilon}_f^* \cdot \vec{\epsilon}_i \left[ -\frac{e^2}{4\pi M} + \omega^2(\alpha_E + \beta_M) + \mathcal{O}(\omega^4) \right] \\
 &+ i\sigma \cdot \vec{\epsilon}_f^* \times \vec{\epsilon}_i \left[ -\omega \frac{e^2 \kappa^2}{4\pi 2M^2} + \omega^3 \gamma_0 + \mathcal{O}(\omega^5) \right] \\
 -\gamma_0 &= \gamma_{E1} + \gamma_{M1} + \gamma_{E2} + \gamma_{M2} \\
 \frac{1}{8\pi M} T_{fi}(\pi) &= \vec{\epsilon}_f^* \cdot \vec{\epsilon}_i \left[ -\frac{e^2}{4\pi M} + \omega\omega'(\alpha_E - \beta_M) + \mathcal{O}(\omega^4) \right] \\
 &+ i\sqrt{\omega\omega'}\sigma \cdot \vec{\epsilon}_f^* \times \vec{\epsilon}_i \left[ -\omega \frac{e^2 \kappa^2 + 4\kappa + 2}{4\pi 2M^2} + \omega\omega' \gamma_0 + \mathcal{O}(\omega^4) \right] \\
 \gamma_\pi &= -\gamma_{E1} + \gamma_{M1} + \gamma_{E2} - \gamma_{M2}
 \end{aligned}$$

## Compton Scattering Cross Section

$$\frac{d\sigma}{d\Omega_{\gamma\gamma}^{\text{Lab}}} = \left| \frac{1}{8\pi M} \frac{\omega_f}{\omega_i} T_{fi} \right|^2$$

## Real Compton Scattering & Proton Polarizabilities

World Data: (1960 - 2001, Moscow, Saskatoon, Illinois, Mainz)

RCS Experiment: V. Olmos de Leon *et al.*, *Eur. Phys. J A***10**, 207 (2001):

$\mathcal{O}(p^3)\chi$ PT: V. Bernard, *et al.*, *Phys. Lett. B* **319** 269 (1993).

$\mathcal{O}(p^4)\chi$ PT: V. Bernard, *et al.*, *Z. Phys.* **348** 317 (1993):

	Experiment	$\mathcal{O}(p^3)\chi$ PT	
$\alpha_E$	$(12.1 \pm 0.3_{\text{stat}} \mp 0.4_{\text{syst}})$	12.5	$10^{-4} \text{ fm}^3$
$\beta_M$	$(1.6 \pm 0.4_{\text{stat}} \pm 0.4_{\text{syst}})$	1.25	$10^{-4} \text{ fm}^3$
$\gamma_\pi$	$-(36.1 \pm 2.1_{\text{stat}} \pm 0.9_{\text{syst}})$	-38.3	$10^{-4} \text{ fm}^4$
$\gamma_0$	$(1.02 \pm 0.08_{\text{stat}} \pm 0.10_{\text{syst}})$	4.5	$10^{-4} \text{ fm}^4$

$\mathcal{O}(p^4)\chi$ PT introduces uncertainties  $\approx \pm 3$  in each quoted value from phenomenological constants, esp.  $\Delta$  and  $N^*$  terms.

Harmonic Oscillator model:  $\alpha_E \approx \alpha_{\text{QED}} \cdot \text{Volume} \cdot (b/\lambda_C)$

$\alpha_E \ll \text{Volume}$ : proton is very stiff, intrinsically relativistic ( $\lambda_C \approx \text{size}$ )

$\beta_M \ll \alpha_E$ : Strong cancellation of para- and dia-magnetism.

$N \rightarrow \Delta$  transition over saturates  $\beta_M$ .

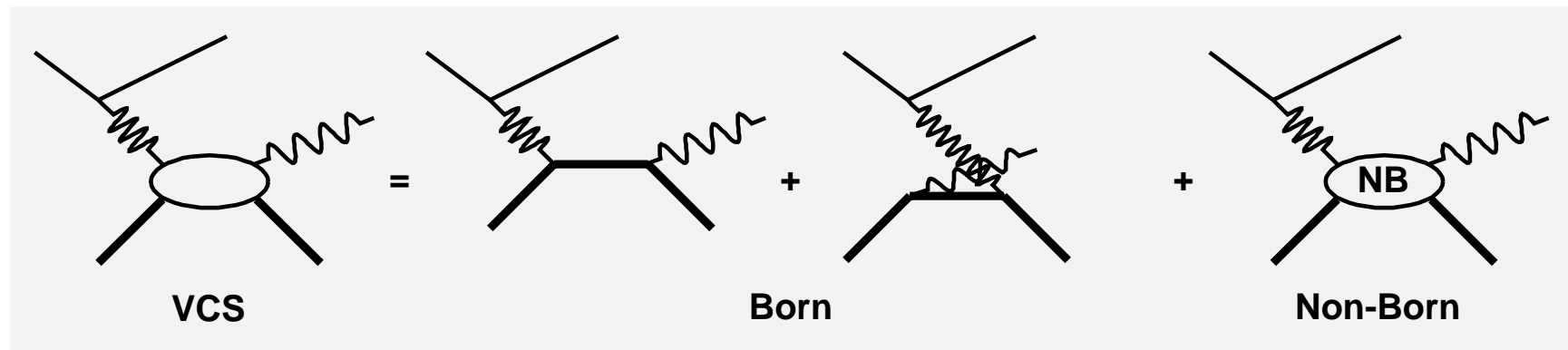
Idea of Virtual Compton Scattering for Polarizabilities:

$\gamma + p$  invariant mass  $\leq M + m_\pi$ .

$2\pi/q =$  virtual-photon wavelength, controlled by experiment.

Measure spatial variation of polarization inside proton:

$q =$  fourier transform variable.



P.A.M. Guichon, G.Q. Liu, A.W. Thomas, *Nucl. Phys. A*, **591** 606 (1995).

P.A.M. Guichon, M. Vanderhaeghen, *Prog. Part. Nucl. Phys.* **41** 125 (1998).

Lvov...

## Low Energy Expansion (LEX) of VCS

$q'$  = final photon energy in  $\gamma p$  CM frame.

$$\begin{aligned}
 d\sigma &= d\sigma^{\text{BH+Born}} + v_{LL} [P_{LL} - P_{TT}/\epsilon] + v_{LT} P_{LT} + \mathcal{O}(q'), \\
 P_{LL} &= -\sqrt{24}M G_{E,p}(\tilde{Q}^2) P^{(C1,C1)0}(\tilde{Q}^2) = \frac{4M}{\alpha_{\text{QED}}} G_{E,p}(\tilde{Q}^2) \alpha_E(\tilde{Q}^2) \\
 P_{TT} &= 6M(1 + \tilde{\tau}) G_{M,p}(\tilde{Q}^2) \left[ P^{(M1,M1)1}(\tilde{Q}^2) + \sqrt{8}\tilde{\tau} P^{(C1,M2)1}(\tilde{Q}^2) \right] \\
 P_{LT} &= \sqrt{\frac{3}{2}}M \sqrt{1 + \tilde{\tau}} \left[ G_{E,p}(\tilde{Q}^2) P^{(M1,M1)0}(\tilde{Q}^2) - \sqrt{6}G_{M,p}(\tilde{Q}^2) P^{(C1,C1)1}(\tilde{Q}^2) \right] \\
 &= -\frac{2M}{\alpha_{\text{QED}}} \sqrt{1 + \tilde{\tau}} G_{M,p}(\tilde{Q}^2) \beta_M(\tilde{Q}^2) - \text{spin G.P.}
 \end{aligned}$$

where  $v_{LL}$ ,  $v_{TT}$ , and  $\epsilon$  are kinematic factors and  $\tilde{Q}^2 = Q^2$  in  $q' \rightarrow 0$  limit.

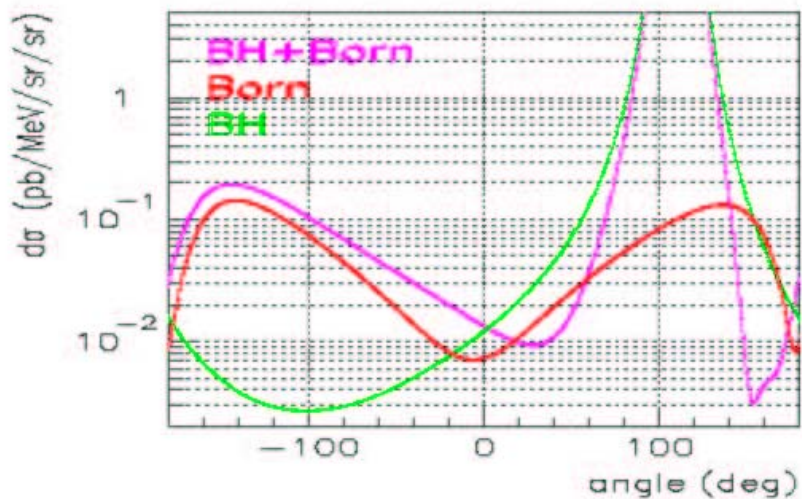
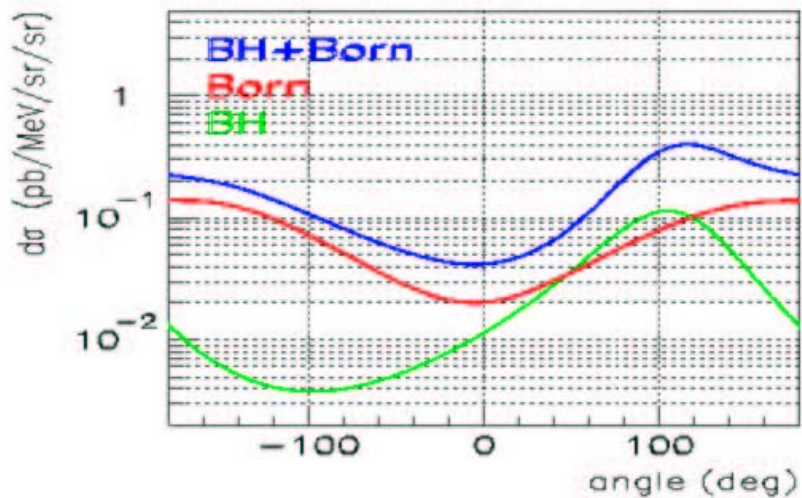
Generalized Polarizabilities (GP)  $P^{(\Lambda_f, \Lambda_i) \Delta S}(\tilde{Q}^2)$ ,  
 $(\Lambda_f, \Lambda_i) = (\text{final}, \text{initial})$  Multipolarity;  $\Delta S = 0, 1$ : Proton spin flip.

In  $\tilde{Q}^2 \rightarrow 0$  limit:

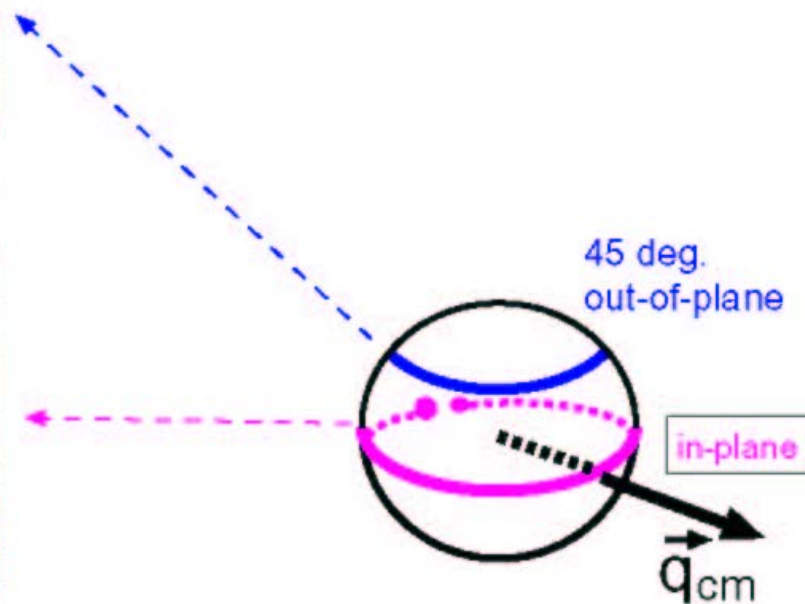
$$\begin{aligned}
 P^{(C1,C1)0} &\longrightarrow -\sqrt{2/3} \bar{\alpha}_E / \alpha_{\text{QED}} \\
 P^{(M1,M1)0} &\longrightarrow -\sqrt{8/3} \bar{\beta}_M / \alpha_{\text{QED}} \\
 P_{TT} &\longrightarrow 0 \\
 P^{(C1,C1)1} &\longrightarrow 0
 \end{aligned}$$

# Components of BETHE-HEITLER + BORN cross section

$$d^5\sigma(ep \rightarrow e\pi\gamma)$$



$$|T_{\text{BH}} + T_{\text{Born}}|^2$$

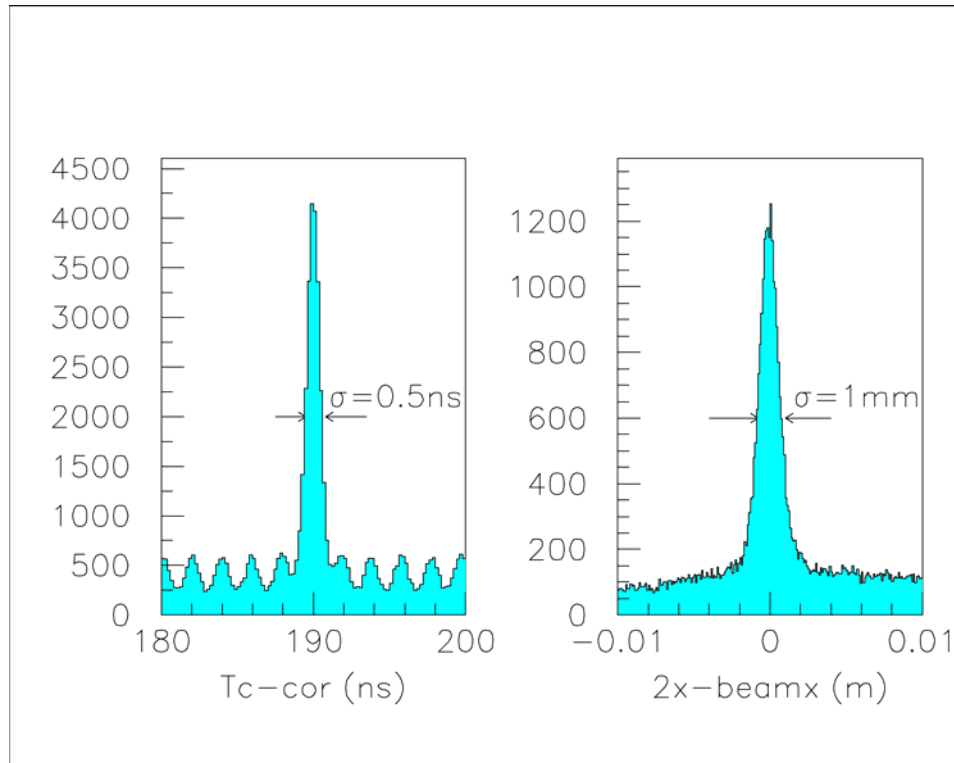


## Virtual Compton Scattering Experiments:

- MAMI:  $k \leq 0.8$  GeV       $q_{cm} = 0.6$  GeV/c  
J. Roche *et al.*, *Phys. Rev. Lett.*, **85**, 708 (2000)
- Jefferson Lab:  $k = 4$  GeV       $Q^2 = 1.0, 1.9$  GeV<sup>2</sup>.  
[hallaweb.jlab.org/physics/experiments/E93-050](http://hallaweb.jlab.org/physics/experiments/E93-050)  
G. Laveissiere *et al.*, *Phys. Rev. Lett.*, **93**, 122001 (2004).
- Bates-Linac:  $k = 0.6$  GeV       $Q^2 = 0.05$  GeV<sup>2</sup>.  
R. Miskimen, UMass-Amherst, spokesperson

# Reconstruct $H(e,e'p)X$ coincidence at target

500 MHz Beam  
Structure



Time Coincidence

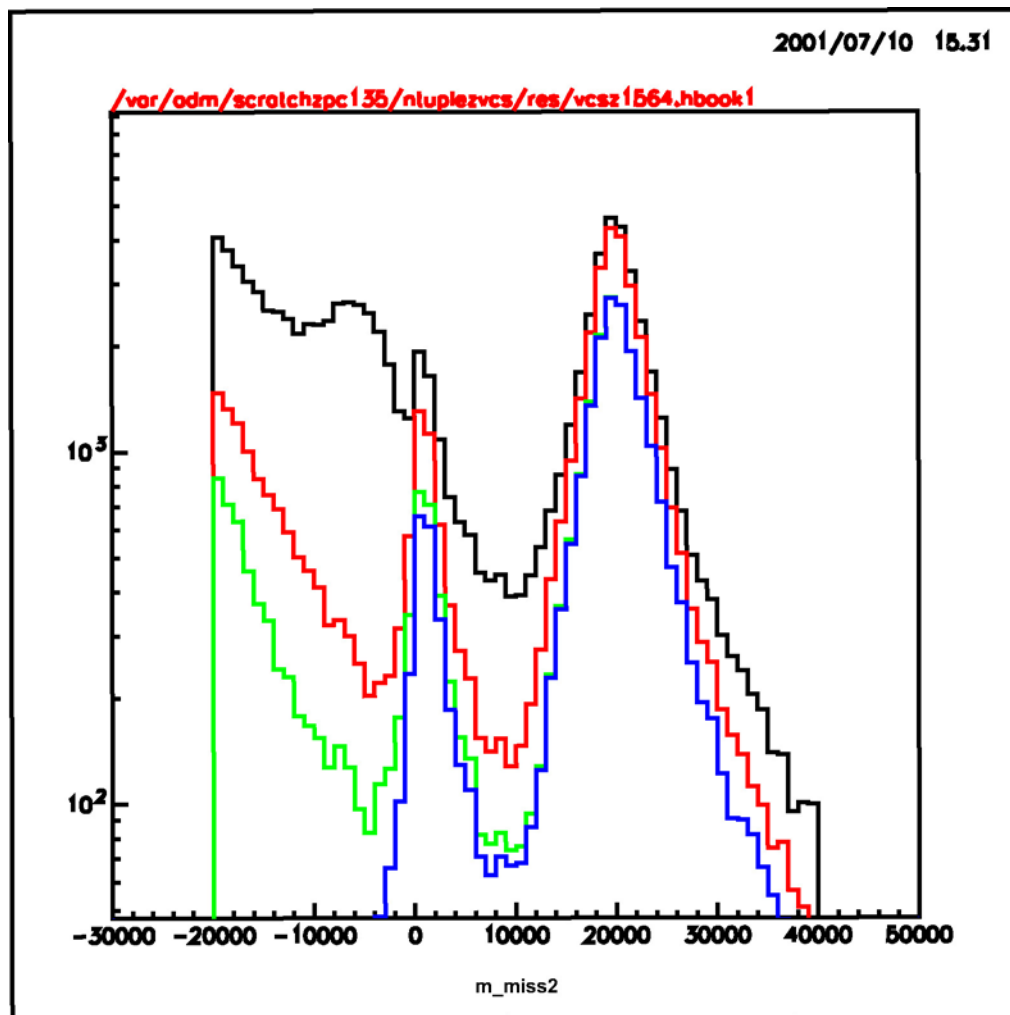
Spatial Coincidence

Reconstruct 1mm  
precision from  
 $4\text{mm} \otimes 20\text{KHz}$  raster,  
with 30 KHz  
bandwidth Beam  
Position Monitor

# Raw H(e,e'p)X Missing Mass squared ( $M_X^2$ ) Spectrum

Prominent H(e,e'p) $\pi^0$  peak at  $M_X^2 = m_{\pi^0}^2 = 20,000 \text{ MeV}^2$

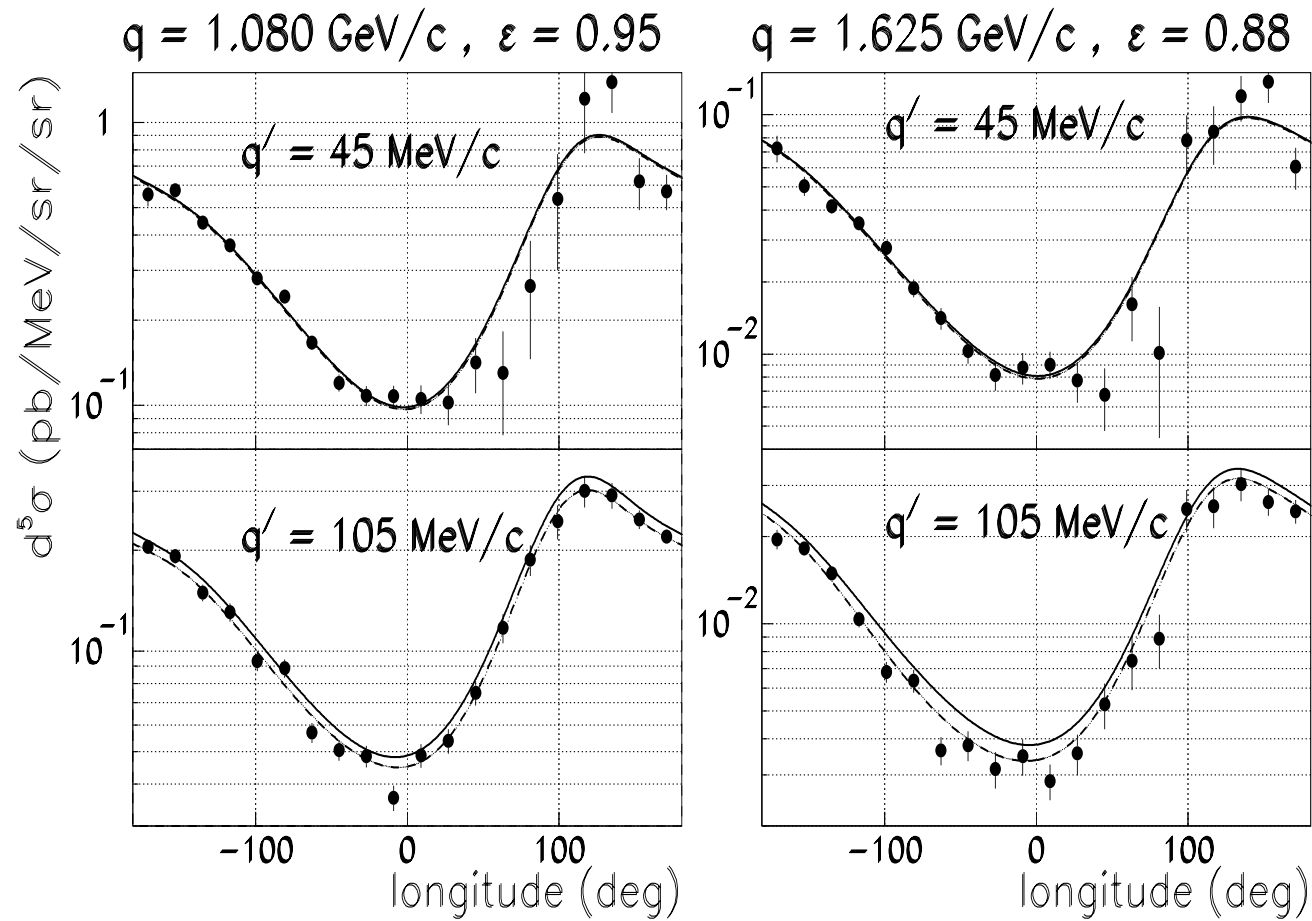
Large unphysical background obscures for  $M_X^2 = 0$  VCS events



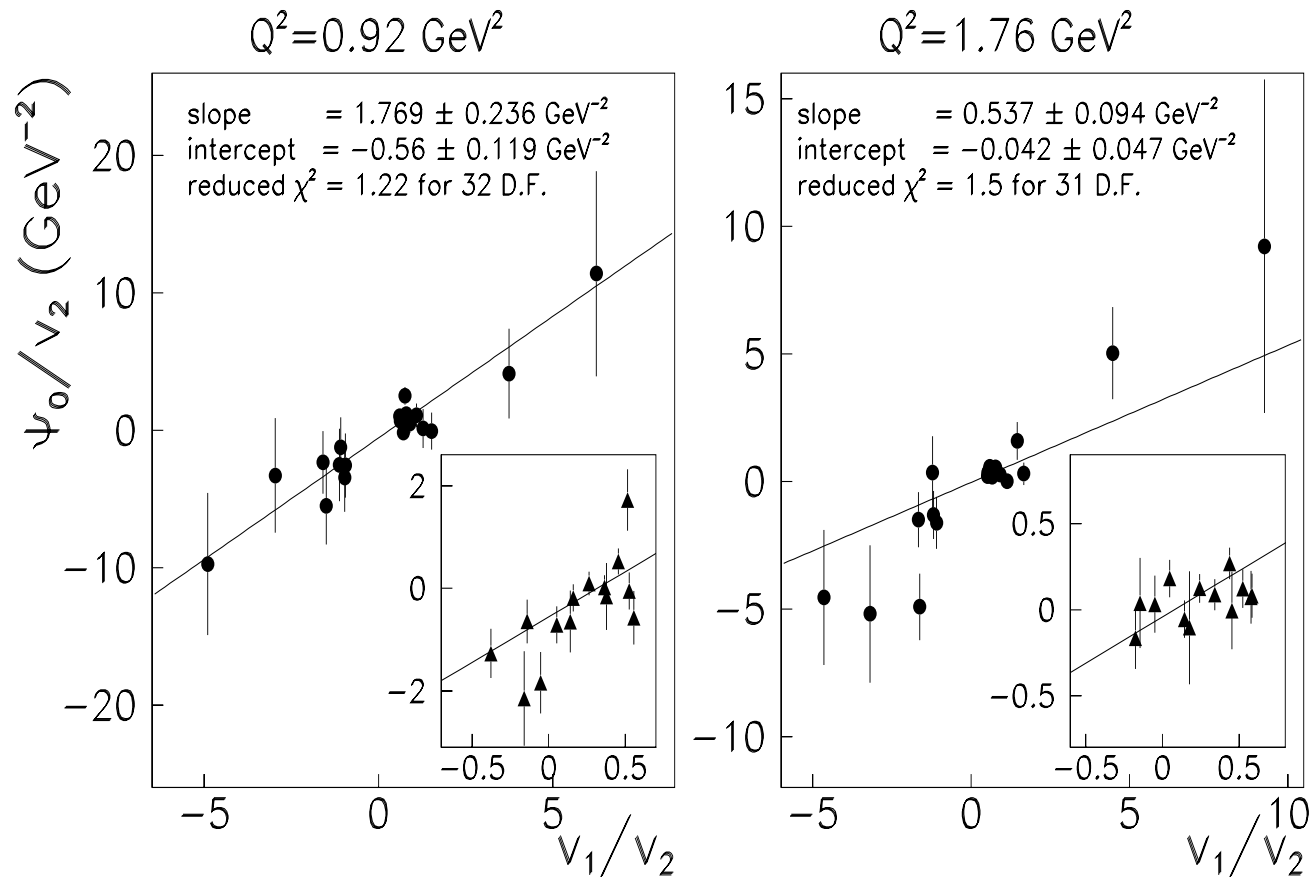
Cuts to remove:

- Elastic  $ep \rightarrow ep$  events. Cut on electron removes mis-identified elastic events with proton punching through collimator
- H(e,e'p)X events with proton reconstructed to punch through collimator
- H(e,e'p) $\gamma$  events, if assumption that  $\gamma \parallel$  beam puts proton in collimator



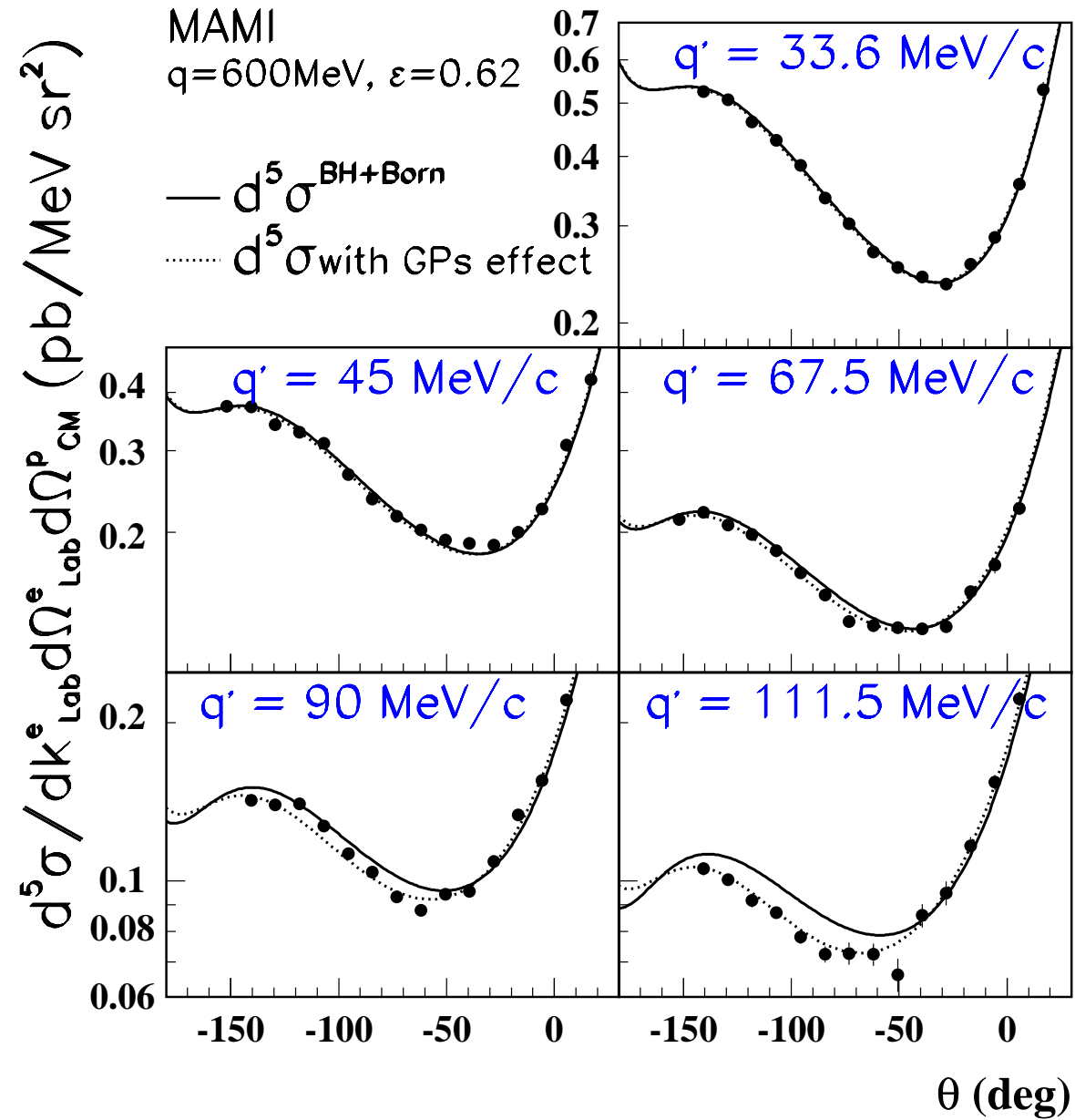


JLab ( $ep \rightarrow ep\gamma$ ) cross section for the lowest and highest  $q'$  bin, at  $40^\circ$  out-of-plane (latitude). Only statistical errors are shown. The abscissa is the azimuthal angle (or longitude). The full curve is the (BH+Born) cross section, the dashed curve includes the first-order GP effect fitted in this analysis.



LEX fit to VCS data below threshold (straight line) for each data set of JLab data. Black circles correspond to out-of-plane data, and the inner plot is a zoom on the lepton plane data (triangles).

J. Roche *et al.*  
 Phys. Rev. Lett.  
 85, 708 (2000)



# Dispersion Relations

B. Pasquini, *et al.*, Eur. Phys. J. A **11**, 185 (2001)

I. Complete formalism for VCS cross section up to  $N\pi\pi$  threshold.

Input from

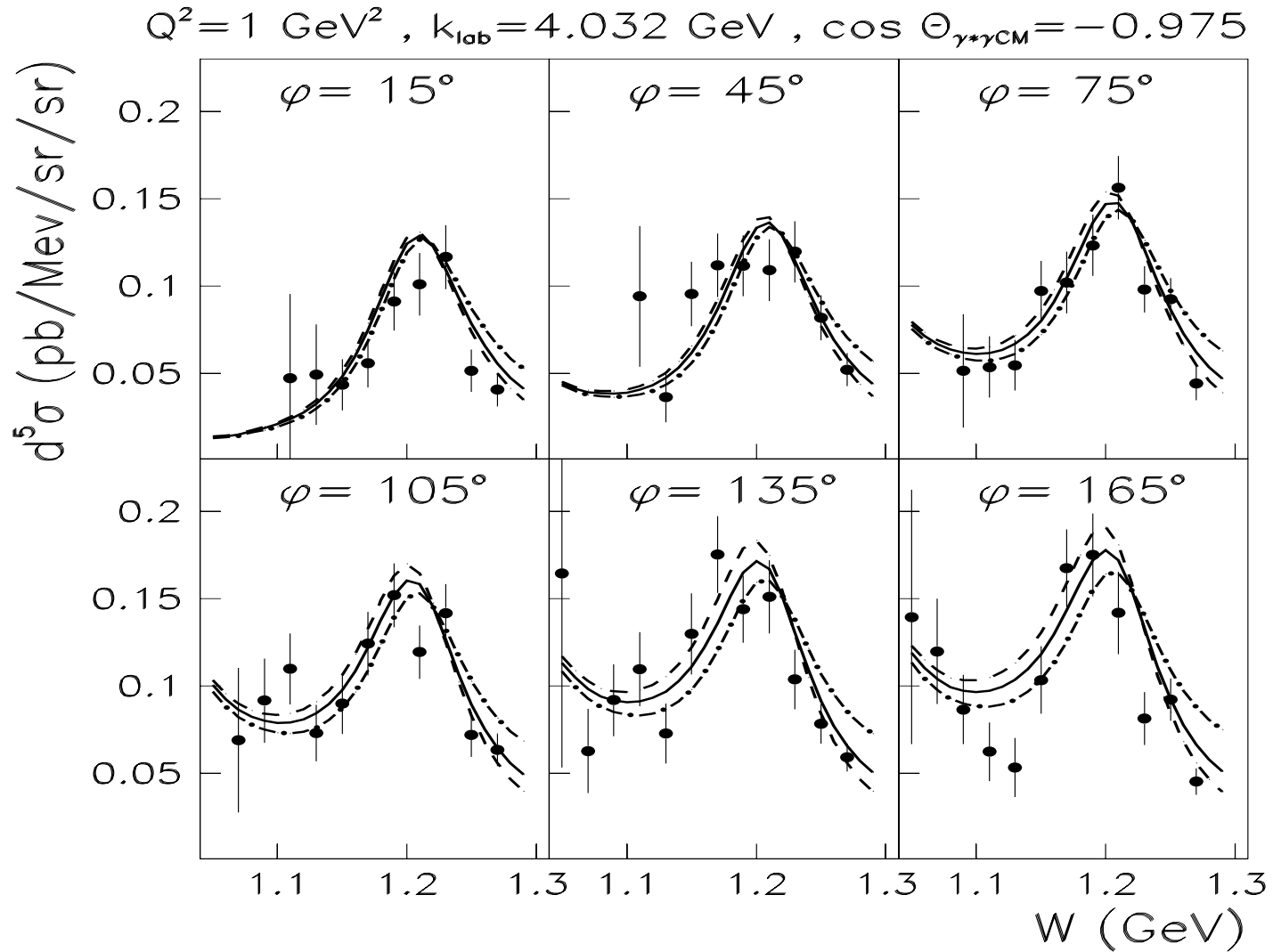
- $\gamma N \rightarrow N\pi$  multipoles (MAID)
- $t$ -channel  $\pi^0$  exchange
- Two low-energy subtraction “constants” (functions of  $Q^2$ ):
  - $\Delta\beta(Q^2)$ :  $t$ -channel  $\sigma$ -meson exchange
  - $\Delta[\alpha + \beta](Q^2)$   $s$ -channel  $N\pi\pi$  and  $N\eta$  resonances not included in MAID.
- $\Delta\beta$  and  $\Delta\alpha$  fitted independently to data at  $Q^2 = 1$  and  $Q^2 = 1.9$  GeV<sup>2</sup> with dipole ansatz:

$$\Delta\beta = \frac{\Delta\beta(0)}{\left[1 + Q^2/\Lambda_\beta^2\right]^2} \quad \Delta\alpha = \frac{\Delta\alpha(0)}{\left[1 + Q^2/\Lambda_\alpha^2\right]^2}$$

II. Predictions of spin polarizabilities.

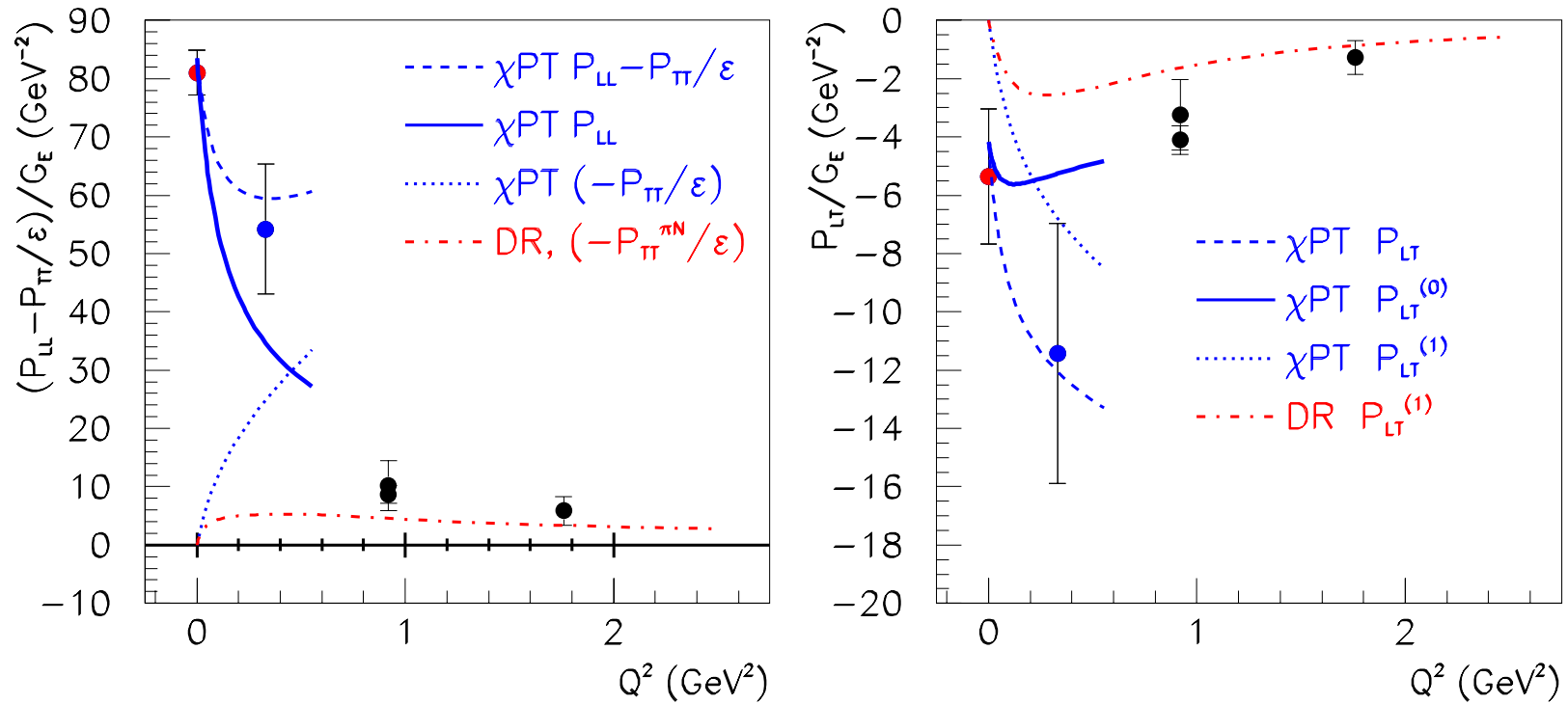
- Separation of  $\alpha_E$  from  $[P_{LL} - P_{TT}/\epsilon]$  and  $\beta_M$  from  $P_{LT}$

III. Interpretation of generalized polarizabilities.



$(ep \rightarrow ep\gamma)$  cross section for JLab data set I-b in six intervals of the azimuthal angle  $\varphi$  (angle between lepton and hadron planes) as a function of  $W$ .

The curves are  $1\text{-}\sigma$  DR fits to total of 700 data points.



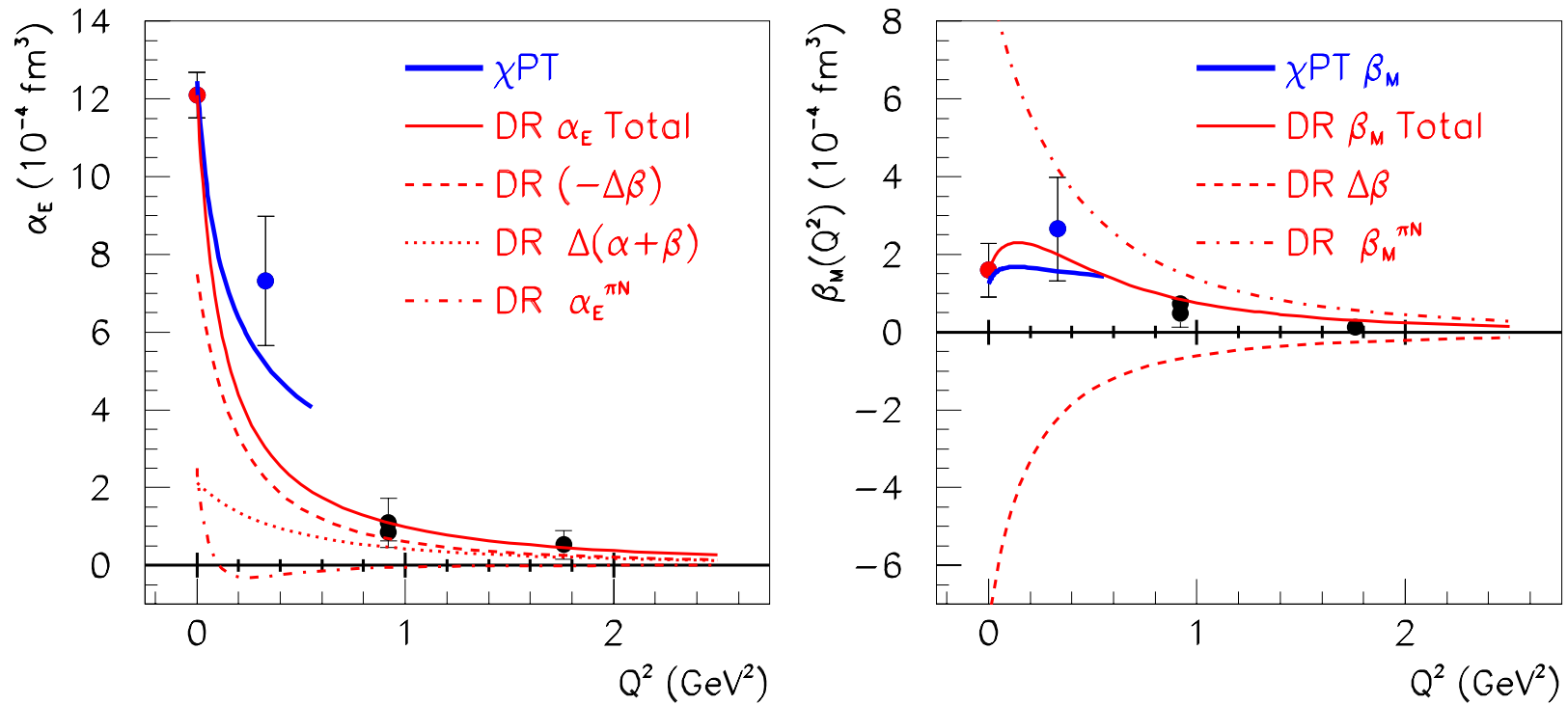
RCS: Olmos de Leon et al [MAMI],,

VCS: J. Roche et al [MAMI-A1], G. Laveissiere et al [JLab E93050].

Dashed Curves are total  $\chi\text{PT}$ , Solid curves are  $\chi\text{PT}$  for  $P_{LL}$  and scalar part of  $P_{LT}$  only.

Dotted curves are  $\chi\text{PT}$  for  $P_{TT}$  and spin-flip part of  $P_{LT}$ .

DotDashed Curves are Dispersion Relation  $\pi N$  predictions for  $P_{TT}$  and spin-flip part of  $P_{LT}$ .



Generalized Polarizabilities:

Obtained by subtracting DR predictions for spin-flip polarizabilities from data.

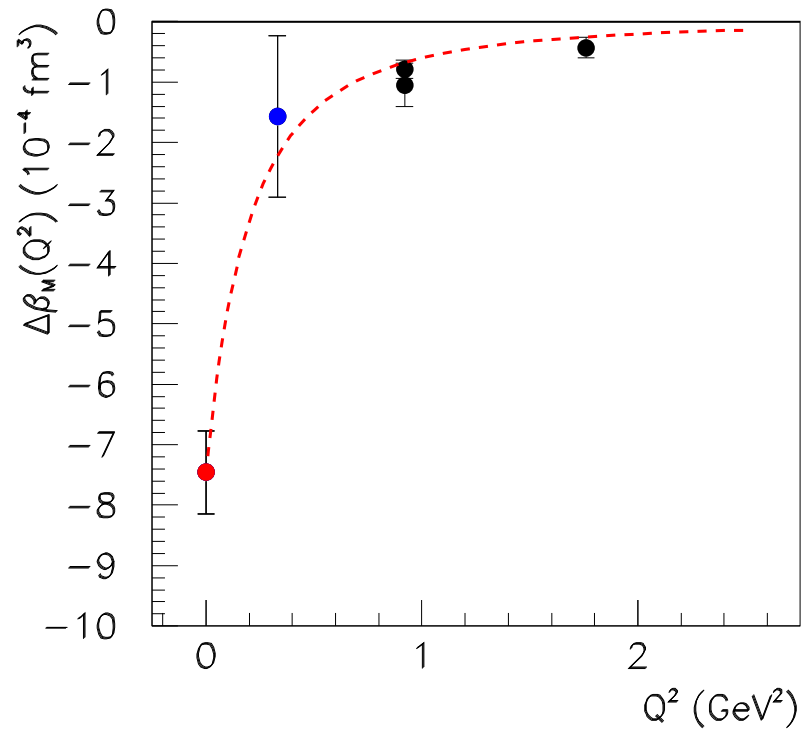
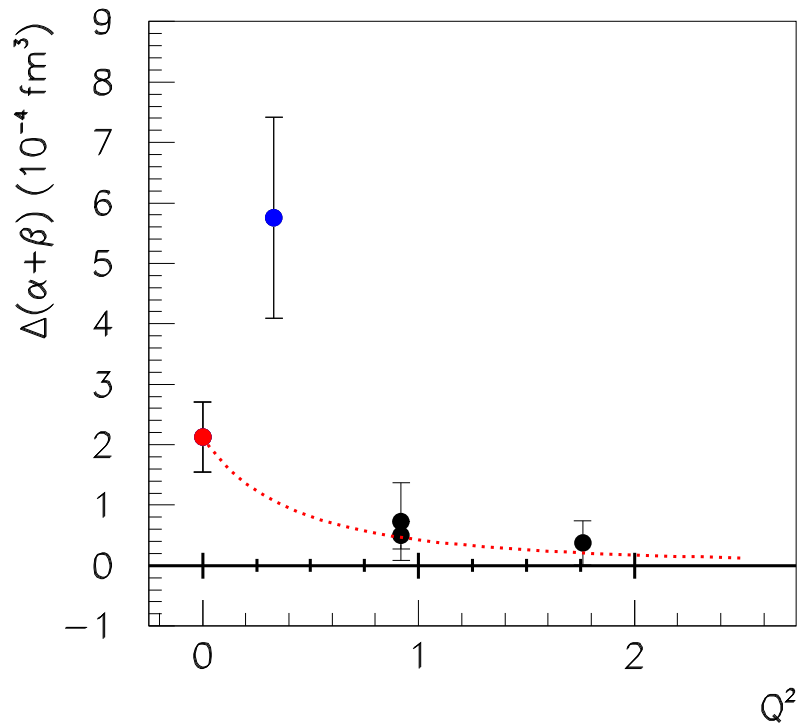
Solid curves are  $\chi\text{PT}$  for  $\alpha_E$  and  $\beta_M$ .

DotDashed curves are  $\pi N$  contributions to  $\alpha_E$  and  $\beta_M$ .

Dashed curves are  $\sigma$ -meson (pion-cloud) exchange term  $\Delta\beta$  (fitted)

Dipole parameter  $\Lambda_\beta < \Lambda(\text{elastic})$ : Indicative of large size of pion cloud.

Dotted curve is  $\Delta[\alpha + \beta]$  term, fitted with dipole ansatz to  $Q^2 = 0 \& 1 \text{ GeV}^2$  data.



$$\Delta[\alpha + \beta] = \alpha_E - \alpha_E^{\pi N} + \Delta\beta$$

$$\Delta[\alpha + \beta] = F_2^{asy} \text{ amplitude in DR.}$$

Spatial distribution not the same as  $G_E$ .

Dashed Curve is  $\sigma$ -exchange term  $\Delta\beta$ ,

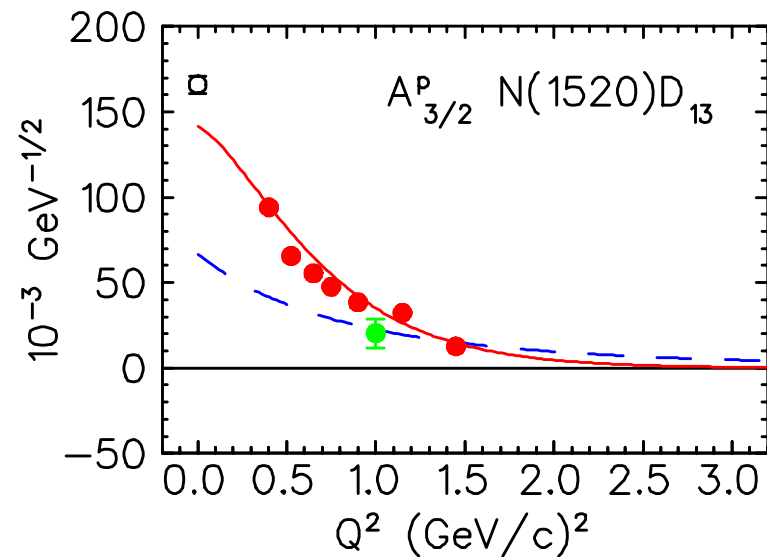
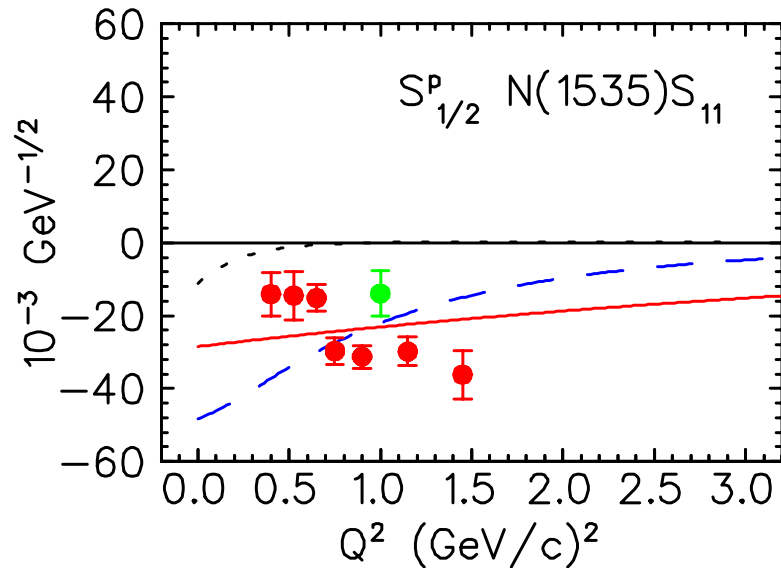
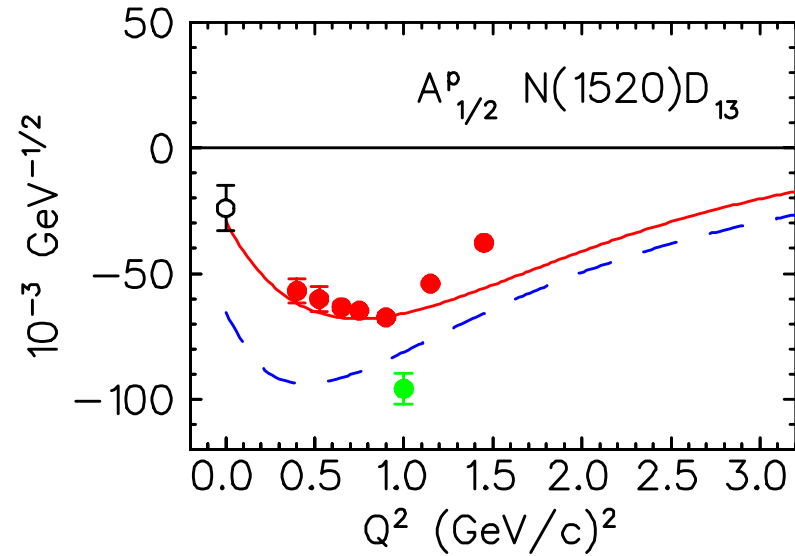
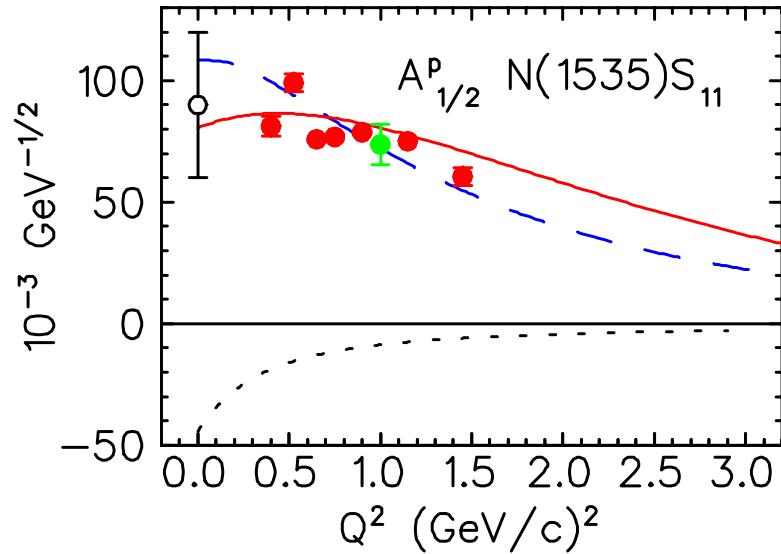
Dipole parameter  $\Lambda_\beta < \Lambda(\text{elastic})$ .

Indicative of large size of pion cloud.



# Contributions to $\Delta[\alpha + \beta](Q^2)$

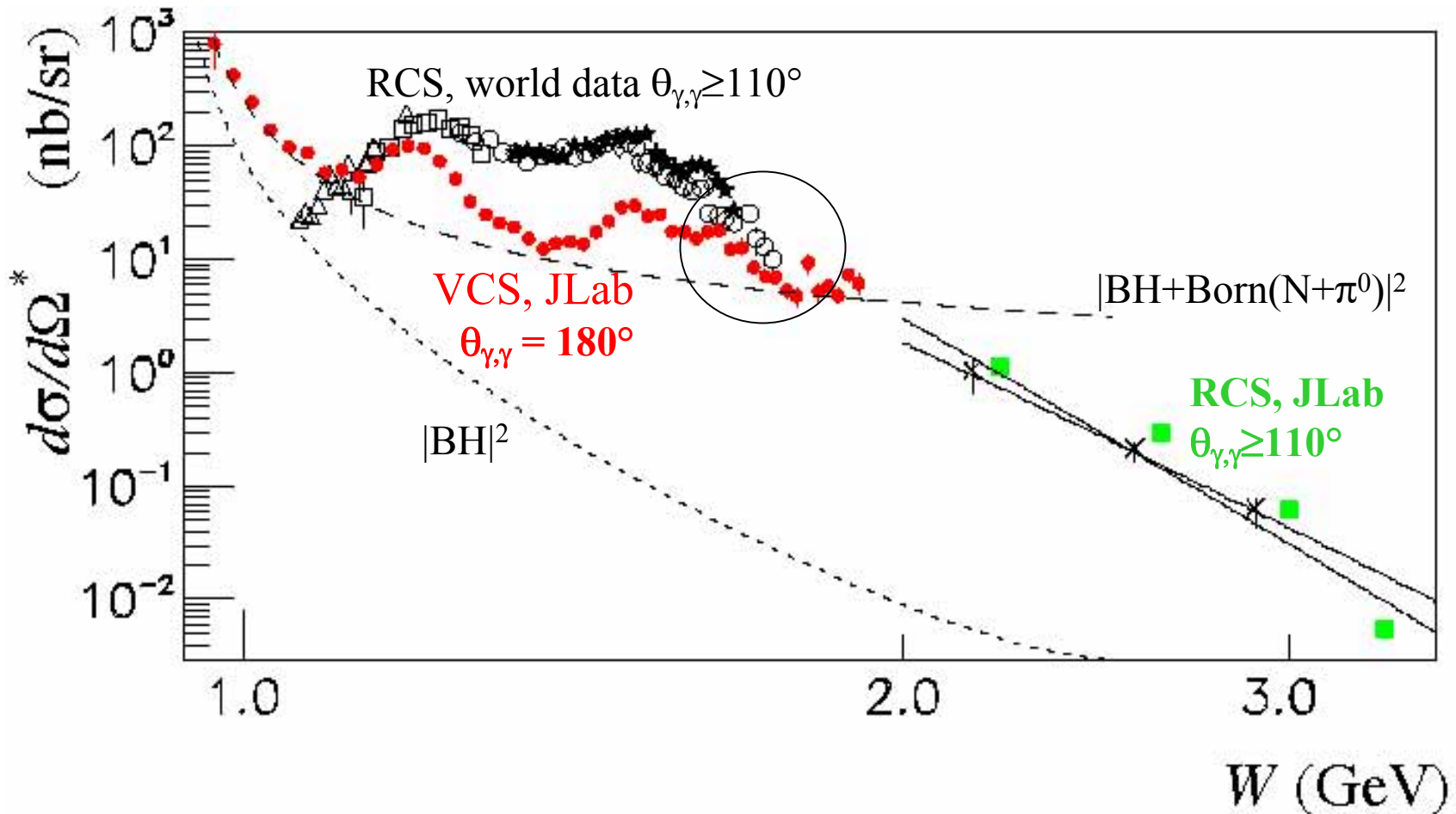
$\gamma^*p \rightarrow N^*$  Amplitudes: L. Tiator, *et al.*, nucl-th 0310041



# Conclusions and Prospects for VCS/Generalized-Polarizabilities

- Low Energy Expansion analysis of data up to  $N\pi$  threshold and Dispersion Relation Analysis of data up through  $N\pi\pi$  threshold give consistent results.
- DR interpretation of Generalized Polarizabilities:
  - Diamagnetism shows large spatial size, as expected from pion cloud.
  - Large contribution to Electric Polarizability from excitations beyond  $N\pi$ , with non-trivial  $Q^2$  dependence.
  - Two photon amplitude is not dominated by low energy excitations.
- New Results soon from
  - Low  $Q^2$  Bates experiment;
  - Single and Double Spin asymmetries at MAMI
- Future experiments with MAMI upgrade.

# VCS( $Q^2=1\text{ GeV}^2$ ) $d\sigma(\gamma^*,\gamma)$ at large $-t$ : (for $\theta_{\gamma\gamma}\approx 180^\circ$ , $-t \approx W^2 = s$ )



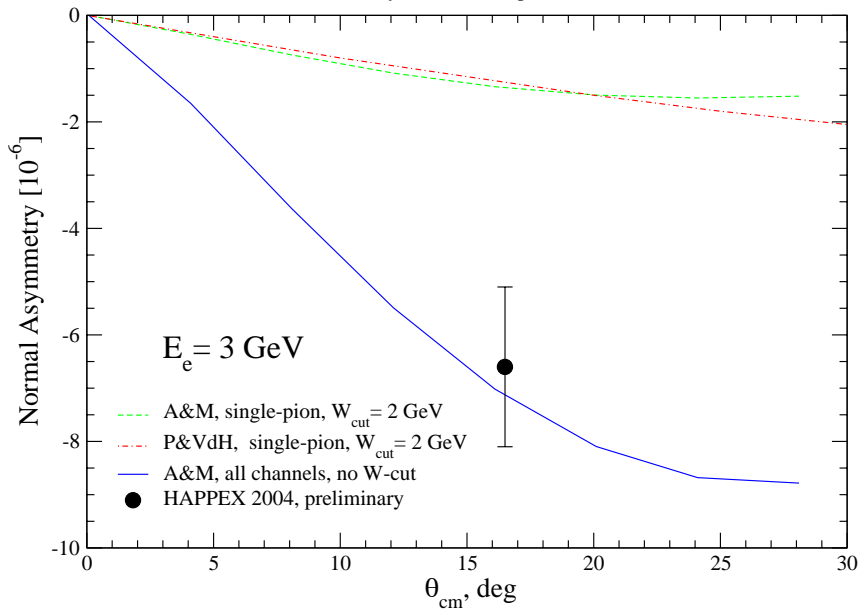
- Resonance form-factors at  $\Delta$  and S11, D13
- Evidence for Compton scattering from point-like objects for  $W \approx 2$  GeV?; or
- Evidence for Nucleon pole dominance of u-channel Regge:  $\mathcal{T} \propto s^{a(u)}$ ?

# Normal Beam Asymmetry from $2\gamma$ -exchange

- Preliminary data from HAPPEX on the proton and He-4 targets
- Measures absorptive part of Compton scattering amplitude, integrated over photon virtualities and  $W$
- Calculations by Afanasev&Merenkov

Normal beam asymmetry for elastic ep-scattering

Unitarity-based model predictions



Normal beam asymmetry for elastic  $e^{-4}\text{He}$  scattering

Unitarity-based model predictions

