



Euridice Network

**Quantum Loops in the
Resonance Chiral Theory:
The Vector Form Factor**

Ignasi Rosell
IFIC, València (Spain)
May 2004

1. Introduction

MOTIVATION

- Confinement: We do NOT know how to work with the \mathcal{L}_{QCD} at low energies \longrightarrow effective field theories
- At intermediate energies ($M_\rho \lesssim E \lesssim 2 \text{ GeV}$) there is no natural expansion parameter \longrightarrow expansion in $1/N_C$
- The Effective Theory for QCD at intermediate energies which uses $1/N_C$ as a expansion parameter \longrightarrow Resonance Chiral Theory (R χ T)

1. Introduction

MOTIVATION

- Confinement: We do NOT know how to work with the \mathcal{L}_{QCD} at low energies \longrightarrow effective field theories
- At intermediate energies ($M_\rho \lesssim E \lesssim 2 \text{ GeV}$) there is no natural expansion parameter \longrightarrow expansion in $1/N_C$
- The Effective Theory for QCD at intermediate energies which uses $1/N_C$ as a expansion parameter \longrightarrow Resonance Chiral Theory ($\text{R}\chi\text{T}$)

PROBLEM!

The renormalization is not clear: the theory has been used only at tree level.

SOLUTION

A study of any function at NLO is needed \rightarrow VFF of the pion.

EFFECTIVE FIELD THEORIES (EFT)

The only remnant of the high-energy dynamics are in the low-energy couplings and in the symmetries of the EFT.

- Theoretical interpretation: integration of the heavy particles

$$e^{i\Gamma_{eff}[\Phi_l]} = \int [d\Phi_h] e^{iS[\Phi_l, \Phi_h]}$$

- Example: the Fermi Theory of Weak Interactions



- The choice of the suitable **degrees of freedom** is fundamental.
- **Interchanges** with heavy particles \longrightarrow local operators with light particles.
- They can be worked as **renormalizables** if one stops at certain precision.

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.
- Massless quark limit \rightarrow chiral symmetry:

$$\begin{aligned}q_L &\longrightarrow q_L' = g_L q_L, \\q_R &\longrightarrow q_R' = g_R q_R,\end{aligned}$$

where $g_{L,R} \in SU(n_f)_{L,R}$. As the symmetry is global, one expects an influence into the hadronic spectroscopy:

$$\left. \begin{array}{l}SU(3)_V \text{ Representations} \\ M_{\pi,K} \ll M_\rho\end{array} \right\} \text{SCSB: } SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.
- Massless quark limit \rightarrow chiral symmetry:

$$\begin{aligned}q_L &\longrightarrow q_L' = g_L q_L, \\q_R &\longrightarrow q_R' = g_R q_R,\end{aligned}$$

where $g_{L,R} \in SU(n_f)_{L,R}$. As the symmetry is global, one expects an influence into the hadronic spectroscopy:

$$\left. \begin{array}{l}SU(3)_V \text{ Representations} \\ M_{\pi,K} \ll M_\rho\end{array} \right\} \text{SCSB: } SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

- The CCWZ formalism allows to obtain effective Lagrangians for SSB theories \rightarrow an EFT with Goldstone bosons can be developed: χ PT.

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.
- Massless quark limit \rightarrow chiral symmetry:

$$\begin{aligned}q_L &\longrightarrow q_L' = g_L q_L, \\q_R &\longrightarrow q_R' = g_R q_R,\end{aligned}$$

where $g_{L,R} \in SU(n_f)_{L,R}$. As the symmetry is global, one expects an influence into the hadronic spectroscopy:

$$\left. \begin{array}{l}SU(3)_V \text{ Representations} \\ M_{\pi,K} \ll M_\rho\end{array} \right\} \text{SCSB: } SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

- The CCWZ formalism allows to obtain effective Lagrangians for SSB theories \rightarrow an EFT with Goldstone bosons can be developed: χ PT.
- Although $\Lambda_\chi \sim 1$ GeV, χ PT only for regions $E \ll M_\rho$.

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.
- Massless quark limit \rightarrow chiral symmetry:

$$\begin{aligned} q_L &\longrightarrow q_L' = g_L q_L, \\ q_R &\longrightarrow q_R' = g_R q_R, \end{aligned}$$

where $g_{L,R} \in SU(n_f)_{L,R}$. As the symmetry is global, one expects an influence into the hadronic spectroscopy:

$$\left. \begin{array}{l} SU(3)_V \text{ Representations} \\ M_{\pi,K} \ll M_\rho \end{array} \right\} \text{SCSB: } SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

- The CCWZ formalism allows to obtain effective Lagrangians for SSB theories \rightarrow an EFT with Goldstone bosons can be developed: χ PT.
- Although $\Lambda_\chi \sim 1$ GeV, χ PT only for regions $E \ll M_\rho$.
- Organization in terms of increasing powers of momentum, E/Λ_χ .

CHIRAL PERTURBATION THEORY (χ PT)

- The running of α_s drives naively to the confinement \rightarrow EFT's.
- Massless quark limit \rightarrow chiral symmetry:

$$\begin{aligned}q_L &\longrightarrow q_L' = g_L q_L, \\q_R &\longrightarrow q_R' = g_R q_R,\end{aligned}$$

where $g_{L,R} \in SU(n_f)_{L,R}$. As the symmetry is global, one expects an influence into the hadronic spectroscopy:

$$\left. \begin{array}{l}SU(3)_V \text{ Representations} \\ M_{\pi,K} \ll M_\rho\end{array} \right\} \text{SCSB: } SU(3)_L \otimes SU(3)_R \rightarrow SU(3)_V$$

- The CCWZ formalism allows to obtain effective Lagrangians for SSB theories \rightarrow an EFT with Goldstone bosons can be developed: χ PT.
- Although $\Lambda_\chi \sim 1$ GeV, χ PT only for regions $E \ll M_\rho$.
- Organization in terms of increasing powers of momentum, E/Λ_χ .
- With dimensional regularization the counterterms are in the Lagrangian.

THE $1/N_C$ EXPANSION

Which expansions can be used for QCD EFT's?

- $E \ll M_\rho$: $E/\Lambda_\chi \longrightarrow \chi$ PT
- $M_\rho \lesssim E \lesssim 2\text{ GeV}$: ???

THE $1/N_C$ EXPANSION

Which expansions can be used for QCD EFT's?

- $E \ll M_\rho$: $E/\Lambda_\chi \longrightarrow \chi\text{PT}$
- $M_\rho \lesssim E \lesssim 2\text{GeV}$: ???

\hookrightarrow **A SOLUTION: $1/N_C$ expansion**

't Hooft proposed the QCD generalization to a gauge theory with N_C colours, waiting simplifications in the case of $N_C \rightarrow \infty$.

THE $1/N_C$ EXPANSION

Which expansions can be used for QCD EFT's?

- $E \ll M_\rho$: $E/\Lambda_\chi \longrightarrow \chi\text{PT}$
- $M_\rho \lesssim E \lesssim 2\text{GeV}$: ???

\hookrightarrow A SOLUTION: $1/N_C$ expansion

't Hooft proposed the QCD generalization to a gauge theory with N_C colours, waiting simplifications in the case of $N_C \rightarrow \infty$.

- Main features in the limit $N_C \rightarrow \infty$
 1. Mesons and glue states are free and stable.
 2. The number of meson states is infinity.
 3. The hadronic physics can be studied through an effective Lagrangian with mesons \longrightarrow EFT's.

THE RESONANCE CHIRAL THEORY ($R_\chi T$)

- Basic features

1. EFT for $E \sim M_R$ ($M_\rho \lesssim E \lesssim 2 \text{ GeV}$).
2. Expansion in powers of momentum does not work $\longrightarrow 1/N_C$ expansion.
3. Goldstone bosons and resonances as degrees of freedom:

$$\mathcal{L}_{R_\chi T}(u, V, A, S, P) = \mathcal{L}_u(u) + \mathcal{L}_R(u, V, A, S, P).$$

4. Construction of the Lagrangian *à la* CCWZ, using the chiral symmetry.
5. The antisymmetric formalism is used for the massive spin-1 fields.
6. The matching with QCD bans the introduction of $\mathcal{L}_{\chi PT}$ local terms of $\mathcal{O}(p^4)$ or higher at LO in $\mathcal{L}_{R_\chi T}$ and fixes the values of the couplings.

2. VFF in $R_\chi T$ at NLO

[I. ROSELL, J.J. SANZ-CILLERO, A. PICH, in preparation]

- The VFF of the pion for the neutral current, $\mathcal{F}(q^2)$, is defined through the matrix element:

$$\langle \pi^+(p_1) \pi^-(p_2) | \left(\frac{1}{2} \bar{u} \gamma_\mu u - \frac{1}{2} \bar{d} \gamma_\mu d \right) | 0 \rangle = \mathcal{F}(q^2) (p_1 - p_2)^\mu .$$

- The VFF at tree level has the contributions



which give the result

$$\mathcal{F}_{LO}(q^2) = 1 + \frac{F_V G_V}{F^2} \frac{q^2}{M_V^2 - q^2} .$$

- In our work this observable is studied at **NLO**.

GREEN FUNCTIONS

1. Pion self-energy



2. ρ self-energy



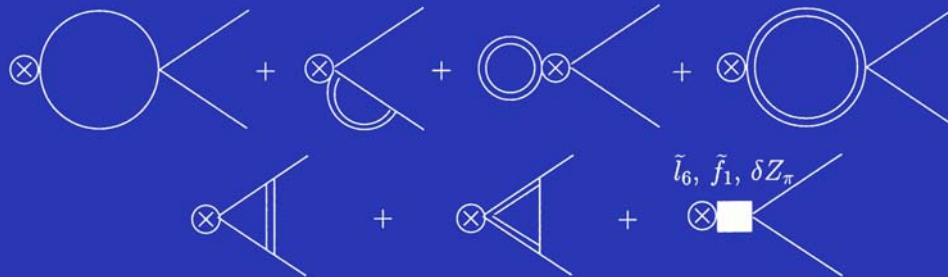
3. $\bar{q}\gamma^\mu q \rightarrow V_{\rho\sigma}$ vertex



4. $V_{\mu\nu} \rightarrow \pi\pi$ vertex



5. $\bar{q}\gamma^\mu q \rightarrow \pi\pi$ vertex



3. Conclusions

Intermediate energies
($M_\rho \lesssim E \lesssim 2 \text{ GeV}$)



expansion in $1/N_C$

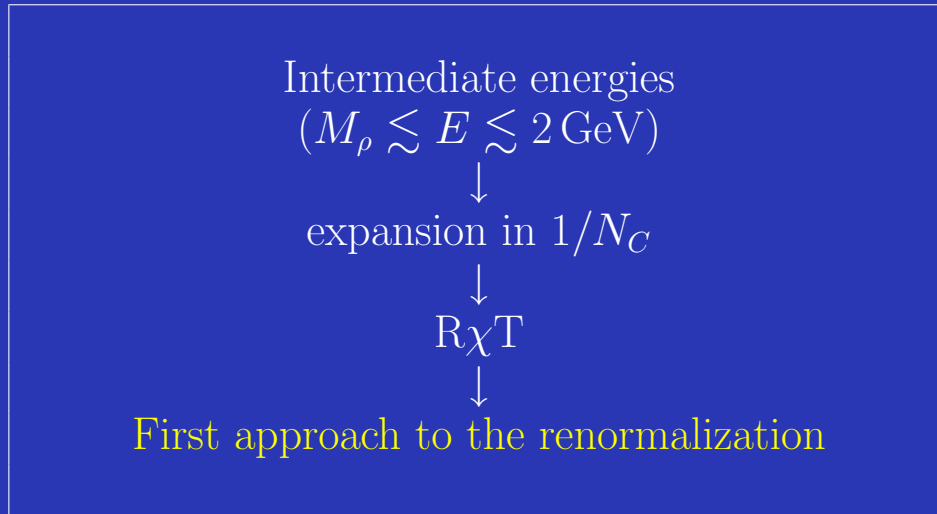


R χ T



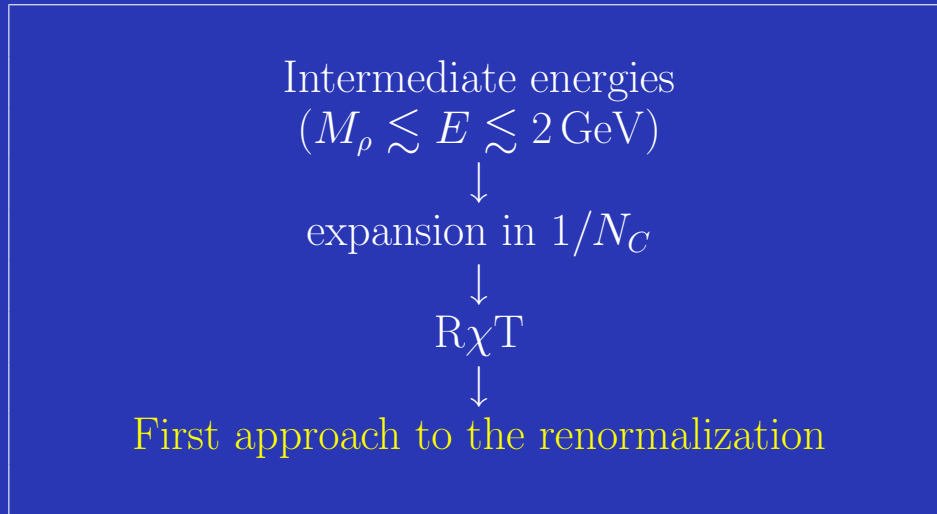
First approach to the renormalization

3. Conclusions



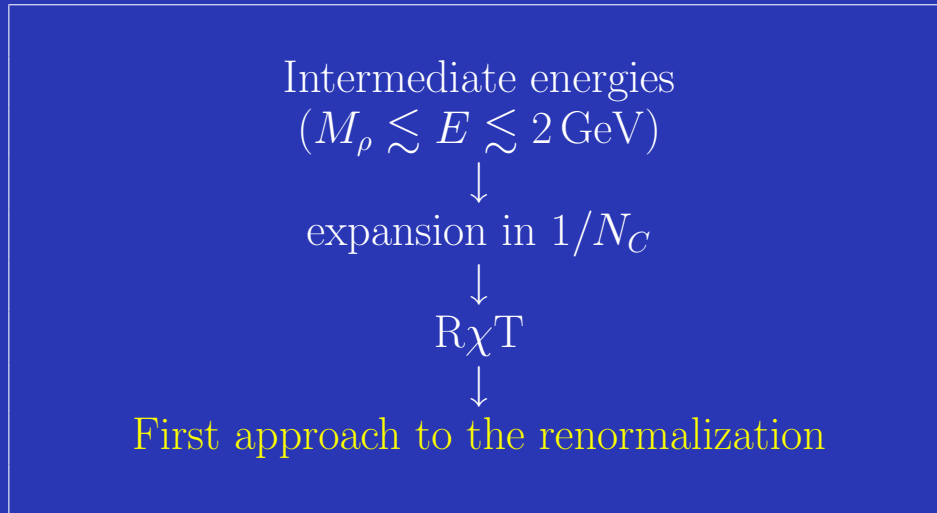
- Using the EOM, terms with resonances of $\mathcal{O}(p^4)$ are not needed for the VFF, at least out of the pole.

3. Conclusions



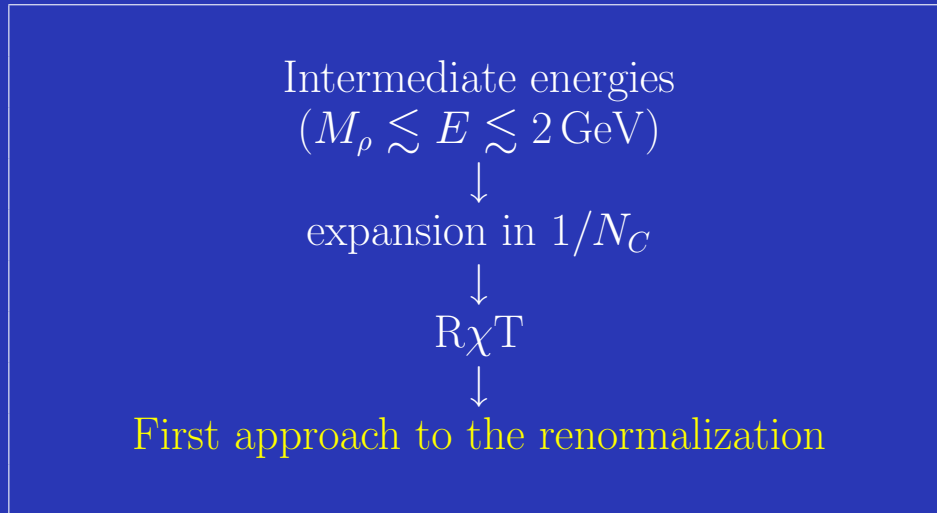
- Using the EOM, terms with resonances of $\mathcal{O}(p^4)$ are not needed for the VFF, at least out of the pole.
- **Generalization?** \longrightarrow only necessary for the NLO in $\text{R}\chi\text{T}$ terms with two or more resonances of $\mathcal{O}(p^2)$ + matching.

3. Conclusions



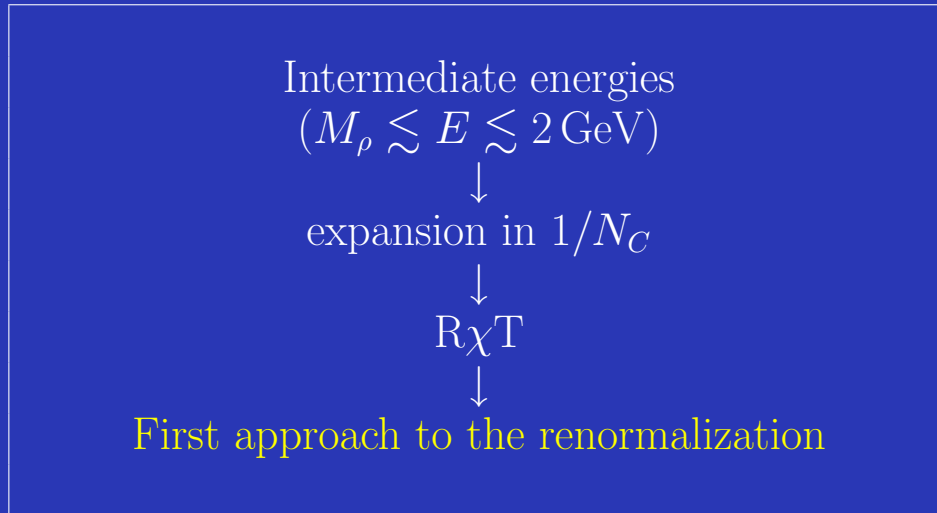
- Using the EOM, terms with resonances of $\mathcal{O}(p^4)$ are not needed for the VFF, at least out of the pole.
- Generalization? \longrightarrow only necessary for the NLO in $R\chi T$ terms with two or more resonances of $\mathcal{O}(p^2)$ + matching.
- χ PT decoupling from $R\chi T$.

3. Conclusions



- Using the EOM, terms with resonances of $\mathcal{O}(p^4)$ are not needed for the VFF, at least out of the pole.
- Generalization? \longrightarrow only necessary for the NLO in $R\chi T$ terms with two or more resonances of $\mathcal{O}(p^2)$ + matching.
- χ PT decoupling from $R\chi T$.
- Short-distance behaviour problem \longrightarrow we are working now on this item: introduction of new operators, effective resummation?

3. Conclusions



- Using the EOM, terms with resonances of $\mathcal{O}(p^4)$ are not needed for the VFF, at least out of the pole.
- Generalization? \longrightarrow only necessary for the NLO in R χ T terms with two or more resonances of $\mathcal{O}(p^2)$ + matching.
- χ PT decoupling from R χ T.
- Short-distance behaviour problem \longrightarrow we are working now on this item: introduction of new operators, effective resummation?
- Future work: functional integration \longrightarrow a fundamental step to work with R χ T at NLO.