

# Chiral Perturbation Theory and the $1/N_c$ Expansion

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- The large  $N_c$  limit of QCD
- The effective Lagrangian and the  $\eta'$
- The  $\delta$  expansion  
(as opposed to the  $p$  expansion)
- Kaplan-Manohar transformation
- Results
- Summary and outlook

Work done in collaboration with H. Leutwyler

## The large $N_c$ limit of QCD

- $N_c \rightarrow \infty$ ,  $g^2 N_c$  fixed G.'t Hooft, 74
- As a consequence the  $U(1)_A$  anomaly disappears in this limit
- Spontaneous symmetry breakdown  
 $U(3)_L \times U(3)_R \rightarrow U(3)_V$
- Nonet of Goldstone bosons:  $\pi, K, \eta, \eta'$   
E.Witten, 79
- Effective prescription : need to include the degree of freedom of the  $\eta'$
- Perturbative analysis yields counting rules for correlation functions  
The  $\eta'$  mass is a special case

## Effective Lagrangian

- Leading order ( $\sim$  current algebra):
  - take Lagrangian of SU(3):  
 $U(x) \in \text{SU}(3) \rightarrow U(x) \in \text{U}(3)$
  - $\det U = e^{i\psi}$ ,  $\psi$  describes the  $\eta'$
  - add a mass term for the  $\psi$  field
  - chiral symmetry:  $\psi^2 \rightarrow (\psi + \theta)^2$

$$\mathcal{L} = \frac{F^2}{4} \langle D_\mu U^\dagger D^\mu U + U^\dagger \chi + \chi^\dagger U \rangle - \frac{\tau_{GD}}{2} (\psi + \theta)^2$$

Di Vecchia & al, Rosenzweig & al, Witten, 80

- $\theta(x)$ : source field for the winding number density  $\omega(x)$

$$\omega = \frac{1}{16\pi^2} \text{tr}_c G_{\mu\nu} \tilde{G}^{\mu\nu}$$

- $\eta'$  mass ( $m_q = 0$ ) :  $M_{\eta'}^2 = 6 \tau_{GD}/F^2$
- $F^2 = O(N_c)$ ,  $\tau_{GD} = O(1)$  :  $M_{\eta'}^2 = O(1/N_c)$

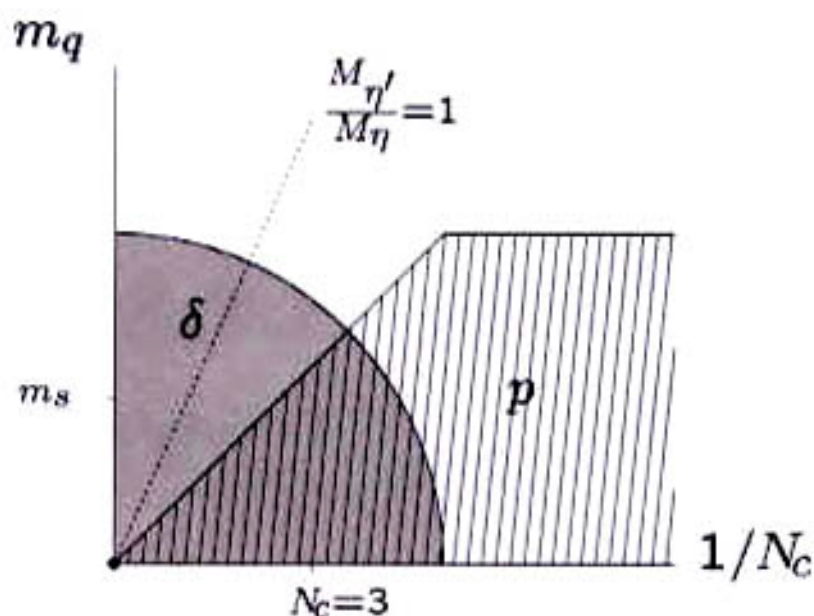
## Effective Lagrangian, cont.

- Leading order

$$\mathcal{L} = \frac{F^2}{4} \langle D_\mu U^\dagger D^\mu U + U^\dagger \chi + \chi^\dagger U \rangle - \frac{\tau_{GD}}{2} (\psi + \theta)^2$$

- Low energy constants are independent of the light quark masses  $m_u, m_d, m_s$
- Low energy constants do know about  $N_c$ :  
 $F^2 = O(N_c), B = O(1), \tau_{GD} = O(1), \dots$
- Introduce counting parameter  $\delta$ :  
 $p = O(\sqrt{\delta}), m = O(\delta), 1/N_c = O(\delta)$
- Leading order Lagrangian =  $O(1)$
- $\mathcal{L}_{eff} = \mathcal{L}^{(0)} + \mathcal{L}^{(1)} + \mathcal{L}^{(2)} + \dots$

## Domain of validity



- Equation of Motion for the  $\eta'$   

$$(\square + M_{\eta'}^2)\psi(x) = f_{\text{ext}}(x) + \dots$$
- $N_c = \infty$  : the  $\eta'$  is massless
- $p^2 \ll M_{\eta'}^2$  :  $\psi(x) = M_{\eta'}^{-2} f_{\text{ext}}(x) + \dots$
- $p^2, M_{\eta'}^2$  both small :  $\delta$  expansion



## $\delta$ expansion

- Chiral symmetry
- Large  $N_c$  counting rules

$$\langle 0 | T j_1 \dots j_n \omega_1 \dots \omega_m | 0 \rangle = O(N_c^{2-l-m})$$

$j_i = \bar{q} \Gamma_i q$ ,  $l = \#$  of quark loops

In particular: correlation functions involving  $\omega$  are suppressed

- Translates into counting rules for the effective coupling constants...  
There are, however, ambiguities  $\leftrightarrow$  freedom in choice of the effective variables
- Nevertheless, it can be shown that the input is sufficient for the construction of an effective Lagrangian which  
(a) describes QCD correctly and  
(b) has the desired properties, namely...

R.K. & H.Leutwyler, to appear

## $\delta$ expansion, cont.

- To any given order in  $\delta$  only a finite number of terms is needed
- The coupling constants in the Lagrangian do obey the 'canonical counting rules' of large  $N_c$  J.Gasser & H.Leutwyler, 1985
- In particular, the coupling constants  $c_n$  are bounded,  $c_n \leq O(N_c), \forall n$
- The Lagrangian of order  $\delta^n$  is a polynomial of degree  $n + 2$  in  $\theta$
- Note that it is easy to find a Lagrangian which does not have these properties: a suitable redefinition of the effective fields,  $\psi \rightarrow \psi + f(\psi + \theta)$ ,  $f = O(N_c^2)$ , e.g., generates one

## Kaplan-Manohar transformations

- SU(3) effective Lagrangian up to and including  $O(p^4)$  is invariant under

$$m_u \rightarrow m_u + \lambda m_d m_s, \text{ cyclic}$$

$$L_8 \rightarrow L_8 + \lambda F^2 / (16B), \quad L_6, L_7 \text{ similar}$$

D.Kaplan & A.Manohar, 86

- In order to preserve the U(1)-transformation properties of  $m$  we must include  $\theta$ :

$$m \rightarrow m + \lambda e^{-i\theta} m^{\dagger-1} \det m^{\dagger}$$

- This implies the change

$$\begin{aligned} \langle U^{\dagger} \chi \rangle &\rightarrow \\ &\langle U^{\dagger} \chi \rangle + \frac{\lambda}{4B} e^{-i(\psi+\theta)} \left( \langle \chi^{\dagger} U \rangle^2 - \langle \chi^{\dagger} U \chi^{\dagger} U \rangle \right) \end{aligned}$$

- Introduces a non-polynomial dependence on  $\theta$

Conclusion:  $\lambda$  must vanish to all orders in  $1/N_c$



## Summary and outlook

- The occurrence of a new low energy scale ( $M_{\eta'}$ ) leads to a more complicated low energy structure
- The formulation of the effective theory at large  $N_c$  involves a simultaneous expansion in powers of  $p$ ,  $m_q$  and  $1/N_c$ :  $\delta$  expansion
- The Kaplan-Manohar transformation is forbidden in this framework
- Promising results at order  $\delta$  - extension to order  $\delta^2$  (loops)