

K_{eff} and Low Energy Constants *

G.A., J. BIJNENS

Lund University

P. TALAVERA

Orsay

- Definitions. Form-Factors
- CHPT. Calculation to 2-loops
- Checks.
- LEC fit. Errors.
- Masses, decay constants, form-factors to 2-loops.
- π - π scattering parameters.
- Notes and Summary.

* "K_{eff} Form-Factors and π - π Scattering"
hep-ph/0003258

G T Amotoss

before starting ...

the 2-loop K_{π} calculation is one of
the next steps in CHPT

for selfconsistent checks

to compare with $\pi\pi$ scatt.

to obtain a first
3 flavour LEC fit

(how do higher orders modify
their values?)

to be updated with the
new measurements
(form-factors, ...)

$$K^+(p) \longrightarrow \pi^+(p_+) \pi^-(p_-) \ell^+(p_\ell) \nu_\ell(p_\nu)$$

$$\begin{array}{c} K^+ \\ K^0 \end{array} \longrightarrow \begin{array}{ccccc} \pi^0 & \pi^0 & \ell^+ & \nu_\ell \\ \pi^- & \pi^0 & \ell^+ & \nu_\ell \end{array}$$

$$T^{+-} = \frac{G_F}{\sqrt{2}} V_{us}^* \bar{u} \gamma_\mu (1 - \gamma_5) v (V^\mu - A^\mu)$$

$$V_\mu = - \frac{H}{m_K^3} E_\mu \nu \rho \sigma (p_e + p_\nu)^\nu (p_+ + p_-)^\rho (p_+ - p_-)^\sigma$$

$$A_\mu = - \frac{i}{m_K} [(p_+ + p_-)_\mu F + (p_+ - p_-)_\mu G + (p_e + p_\nu)_\mu R]$$

- $m_e^2 R \approx 0$
- H to $O(p^6)$: 1-loop
- F and G to $O(p^6)$: 2-loops

$$\mathcal{L} = \mathcal{L}_2 + \mathcal{L}_4 + \mathcal{L}_6 + \dots$$

$$F = \frac{m_K}{\sqrt{2} F_\pi} \left[1 + \frac{F_{1\text{loop}}}{F_\pi^2} + \frac{F_{LL} + F_V + F_{CT}}{F_\pi^4} + \dots \right]$$

$$G = \frac{m_K}{\sqrt{2} F_\pi} \left[1 + \frac{G_{1\text{loop}}}{F_\pi^2} + \frac{G_{LL} + G_V + G_{CT}}{F_\pi^4} + \dots \right]$$

$(F, G)_{1\text{loop}}$: 1-loop diagrams + $\mathcal{O}(p^4)$ CT

$(F, G)_V$: $\mathcal{L}_2 \times \mathcal{L}_2 \times \mathcal{L}_2$ $\mathcal{O}(p^6)$ 

$(F, G)_{CT}$: \mathcal{L}_6 CT $\mathcal{O}(p^6)$ 

$(F, G)_{LL}$: $\mathcal{L}_4 \times \mathcal{L}_2, \mathcal{L}_2 \times \mathcal{L}_2^*$ $\mathcal{O}(p^6)$



* $\mathcal{O}(p^6)$ contribution from
the rewriting with
physical quantities

What to worry about ?

- to reproduce the 1-loop result. ✓
- isospin relations: $T^{+-} = \frac{T^{-0}}{\sqrt{2}} + T^{00}$ ✓
- to renormalize the wave function ✓
- to consider physical [2-loop] master and decay constants.* ✓
- to cancel the divergences
 - polynomial (with CT) ✓
 - spurious, non-local ✓
- to use or not to use the GMO relation (the difference is $O(p^8)$). WE USE. ✓
- values for the CT:
 - $O(p^4)$ CT \rightarrow FIT * [OUR PROJECT]
 - $O(p^6)$ \rightarrow saturation from the nearest resonances:
vector, axial-vector, scalars
and the η' (η - η' mixing).✓

We have a huge 2-loop result.
What and how do we fit?

Theoretical side:

$B_0 \hat{m}$, $B_0 m_s$, F_0 , $L_1 - L_{10}$ parameters from
the CHPT Lagrangian : 13

- L_{10} gives no contribution
- L_9 appears with $S_0 = (P_e + P_\nu)^2$. The fit is done with $S_0 = 0$.
- L_4 and L_6 from the large- N_c limit. (≈ 0).
- Note that

$$\begin{aligned} F_0 &= F_0(F_\pi, L_i, m_i) \\ (B_0 \hat{m}) &= (B_0 \hat{m})(m_\pi^2, L_i, \dots)^+ \end{aligned} \quad \left. \begin{array}{l} \text{THEY ARE OBTAINED} \\ \text{AFTER THE FIT!} \end{array} \right\}$$

Finally,

$$\underbrace{\frac{m_s}{\hat{m}}, \frac{F_K}{F_\pi}, L_1, L_2, L_3, L_5, L_7, L_8}_{\text{written orarser and } L_i.} \quad \left. \begin{array}{l} L_1, L_2, L_3, L_5, L_7, L_8 \\ \text{parameters to fit} \end{array} \right\}$$

$$* m_\pi^2 = 2B_0 \hat{m} + (m^2)^{(4)} + (m^2)^{(6)}$$

with 'physical' quantities

Experimental side:

- masses:

$$F_K / F_\pi^*$$

$$\left. \frac{m_s}{\hat{m}} \right|_1 = \frac{2m_{\phi K}^2 - m_{\phi \pi}^2}{m_{\phi \pi}^2} *$$

||

$$\left. \frac{m_s}{\hat{m}} \right|_2 = \frac{3m_{\phi \pi}^2 - m_{\phi \eta}^2}{2m_{\phi \pi}^2} *$$

||

24 FIXED

$$*(m^2)_{\text{phys}} = m_0^2 + (m^2)^{(4)} + (m^2)^{(6)}.$$

- $f_s(0), g_s(0), d_f \cdot [K_{e4} \text{ FORM-FACTORS}]$

↳ NEXT SLIDE

$F_K / F_\pi, \left. \frac{m_s}{\hat{m}} \right|_1, \left. \frac{m_s}{\hat{m}} \right|_2, f_s(0), g_s(0), d_f$

Input

Rosselet et al. (77)

$$f_s(q^2) = f_s(0)(1 + \lambda_f q^2)$$

$$q^2 = \frac{s_n - 4m_n^2}{4m_n^2}$$

$$g_p(q^2) = g_p(0)(1 + \lambda_g q^2)$$

$$f_s(0) = 5.59 \pm 0.14$$

$$g_p(0) = 4.77 \pm 0.27$$

$$\lambda_f = \lambda_g = 0.08 \pm 0.02$$

Relation with the F and G form-factors:

$$F = f_s e^{is_s} + f_p e^{is_p} \cos \theta_n + d\text{-wave} + \dots$$

$$G = g_p e^{is_p} + d\text{-wave} + \dots$$

$$f_s(0) = F(s_n \approx 4m_n^2, s_p = 0, \cos \theta_n = 0)$$

$$g_p(0) = G(s_n \approx 4m_n^2, s_p = 0, \cos \theta_n = 0)$$

Table 1: $L_4^r = 0$, $L_6^r = 0$ and $L_8^r = 6.90 \cdot 10^{-3}$. All $L_i^r(\mu)$ values quoted have been brought to the scale $\mu = 0.77$ GeV. The standard values are $m_s/\bar{m} = 24$, $\sqrt{s_\pi} = 0.336$ GeV and $s_f = 0$. [1] Gasser and Leutwyler, Nucl. Phys. B250 (1985) 465; [2] Bijnen, Colangelo and Gasser, Nucl. Phys. B427 (1994) 427

	Main Fit	[1,2]	$\mathcal{O}(p^4)$	fit 2	fit 3	fit 4	fit 5	fit 6	fit 7	fit 8	fit 9
$10^3 L_1^r$	0.52 ± 0.23	$0.37 ± 0.23$	0.46	0.53	0.50	0.49	0.52	0.45	0.44	0.63	0.65
$10^3 L_2^r$	0.72 ± 0.24	$1.35 ± 0.23$	1.49	0.73	0.67	0.74	0.80	0.51	1.04	0.73	0.85
$10^3 L_5^r$	-2.70 ± 0.99	$-3.5 ± 0.85$	-3.18	-2.71	-2.58	-2.73	-2.75	-2.16	-2.95	-2.67	-3.27
$10^3 L_6^r$	0.65 ± 0.12	$1.4 ± 0.5$	1.46	0.62	0.65	0.64	0.82	0.67	1.00	0.51	0.60
$10^3 L_7^r$	-0.26 ± 0.15	$-0.4 ± 0.2$	-0.49	-0.20	-0.26	-0.26	-0.30	-0.26	-0.23	-0.25	-0.26
$10^3 L_8^r$	0.47 ± 0.18	$0.9 ± 0.3$	1.08	0.35	0.47	0.47	0.58	0.46	0.52	0.44	0.48
changed quantity		$\mathcal{O}(p^4)$	m_s/\bar{m}	$\sqrt{s_\pi}$	$\sqrt{s_\pi}$	L_4^r, L_6^r	V_B, η'	μ	μ	$g^{(0)}$	
Unit			26	0.1	0.293	-0.3; -0.2	only	0.5	1.0	4.93	
				GeV	GeV	10^{-3}		GeV	GeV		

Figure 2: The F and G form-factor at $s_t = 0$ and $\cos \theta_s = 0$. The shaded band is the result of Rosselet et al. Shown are the full result for the absolute value, and the real part to lowest-order, $\mathcal{O}(p^4)$ and $\mathcal{O}(p^6)$. We also show the contribution from the $\mathcal{O}(p^6)$ Lagrangian; this is dominated by the vector contribution.

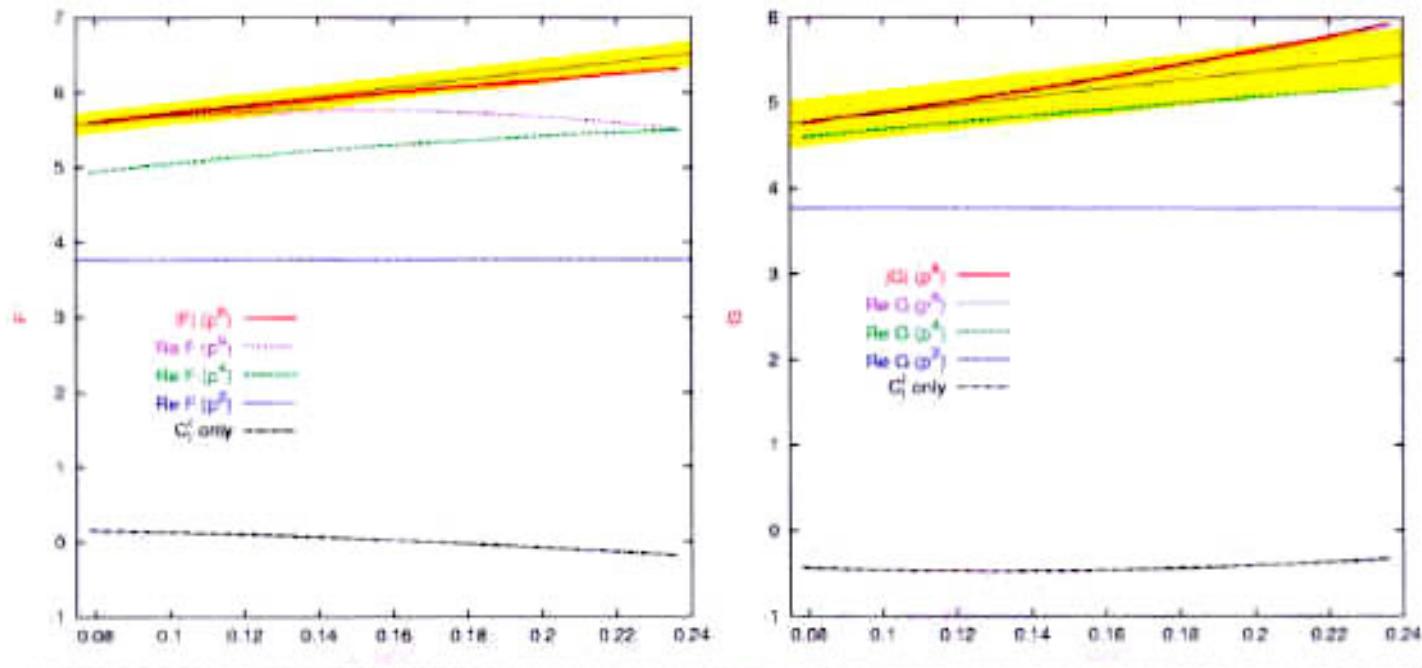
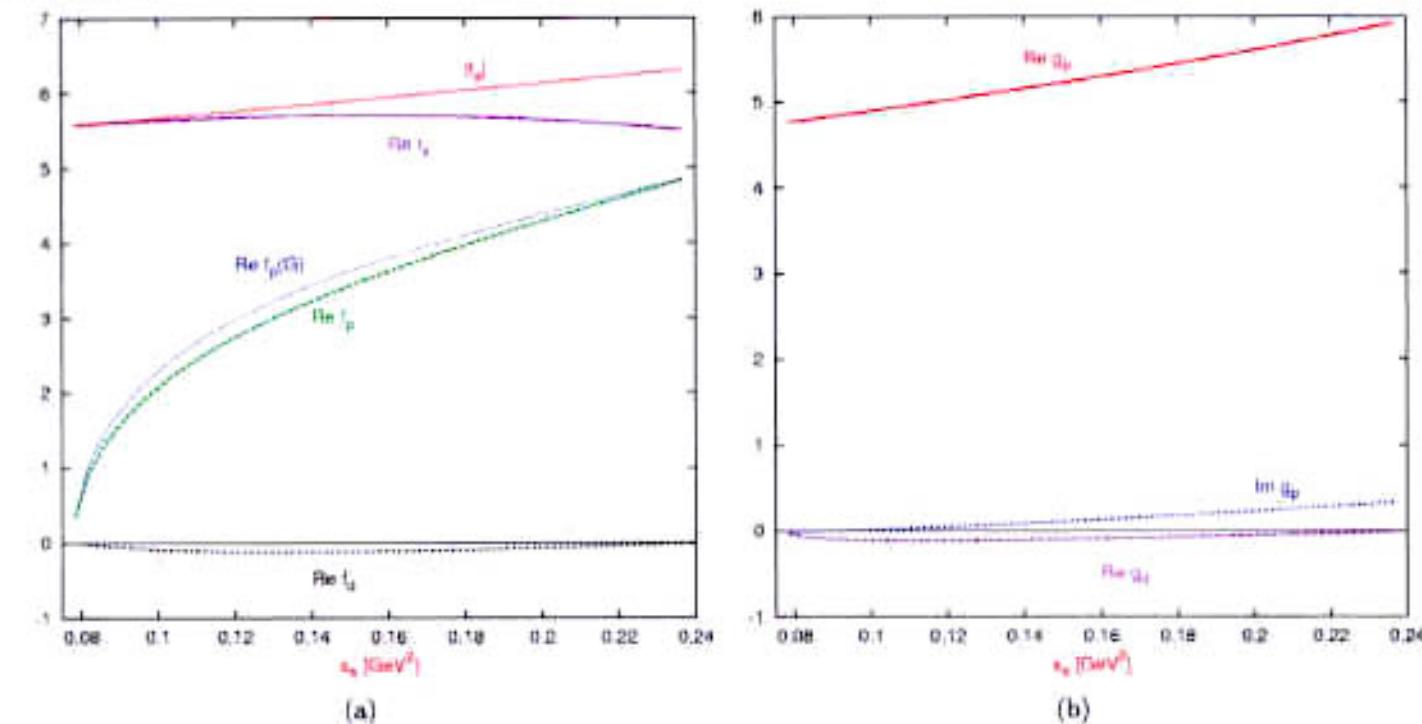


Figure 3: (a) The partial wave expansion of the combination $\tilde{F} = F + \sigma_\pi PL \cos \theta_s G/X$. Plotted are the absolute value of f_s and the real part of f_s , f_p and f_d . We also show the f_p contribution from the second term in \tilde{F} . (b) The partial wave expansion of G . Plotted are the real part of g_p and g_d and the imaginary part of g_p .



$\pi\pi$ scattering parameters.

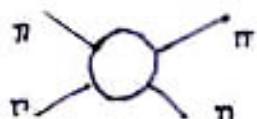
- $\pi\pi$ is known to 2-loops
in two-flavour CHPT : l_i
- Key: THREE-flavour CHPT: L_i

The matching $l_i(\text{SU}(2)) \leftrightarrow L_i(\text{SU}(3))$
is known to 1-loop, then
the l_i can have $O(p^4)$ corrections.

A 2-loops matching has to be done!?

Table 2: Threshold parameters in units of m_{π^+} . Set I is $\tilde{l}_1 = -1.7$ and $\tilde{l}_2 = 6.1$ while Set II is $\tilde{l}_1 = -1.5$ and $\tilde{l}_2 = 4.5$. Both sets have $\tilde{l}_3 = 2.9$ and $\tilde{l}_4 = 4.3$. Experimental values are taken from Dumbrajs et al. Errors are from the 68% confidence level ranges only.

	Main Fit	$\mathcal{O}(p^2)$	$\mathcal{O}(p^4)$	$\mathcal{O}(p^6)$	Set I	$\mathcal{O}(p^6)$	Set II	experiment
a_0^0	0.219 ± 0.005	0.159	0.203	0.222		0.212		0.26 ± 0.05
b_0^0	0.279 ± 0.011	0.182	0.254	0.282		0.256		0.25 ± 0.03
$-10a_0^2$	0.420 ± 0.010	0.454	0.423	0.420		0.450		0.28 ± 0.12
$-10b_0^2$	0.756 ± 0.021	0.908	0.755	0.729		0.81		0.82 ± 0.08
$10a_1^1$	0.378 ± 0.021	0.303	0.356	0.404		0.38		0.38 ± 0.02
$10^2 b_1^1$	0.59 ± 0.12	0	0.31	0.83		0.56		—
$10^2 a_2^0$	0.22 ± 0.04	0	0.16	0.28	input		0.17 ± 0.03	
$-10^3 b_2^0$	0.32 ± 0.10	0	0.40	0.24		0.27		—
$10^3 a_2^2$	0.29 ± 0.10	0	0.32	0.24	input		0.13 ± 0.30	
$-10^3 b_2^2$	0.36 ± 0.9	0	0.23	0.33		0.22		—
$10^4 a_3^1$	0.62 ± 0.11	0	0.20	0.65		0.48		—
$-10^4 b_3^1$	0.36 ± 0.02	0	0.15	0.38		0.37		—



Isospin amplitudes:

$$T^I(s, t) = 32 \pi \sum_{l=0}^{\infty} (2l+1) P_l(\cos \theta) t_e^I(s) .$$

$$t_e^I(s) = \sqrt{\frac{s}{s - 4m_\pi^2}} \frac{1}{2i} \left[e^{2i\delta_e^I} - 1 \right] .$$

$$\text{Re}[t_e^I(s)] = q^{2l} \left[a_e^I + q^2 b_e^I + O(q^4) \right] .$$

$$s = 4(m_\pi^2 + q^2)$$

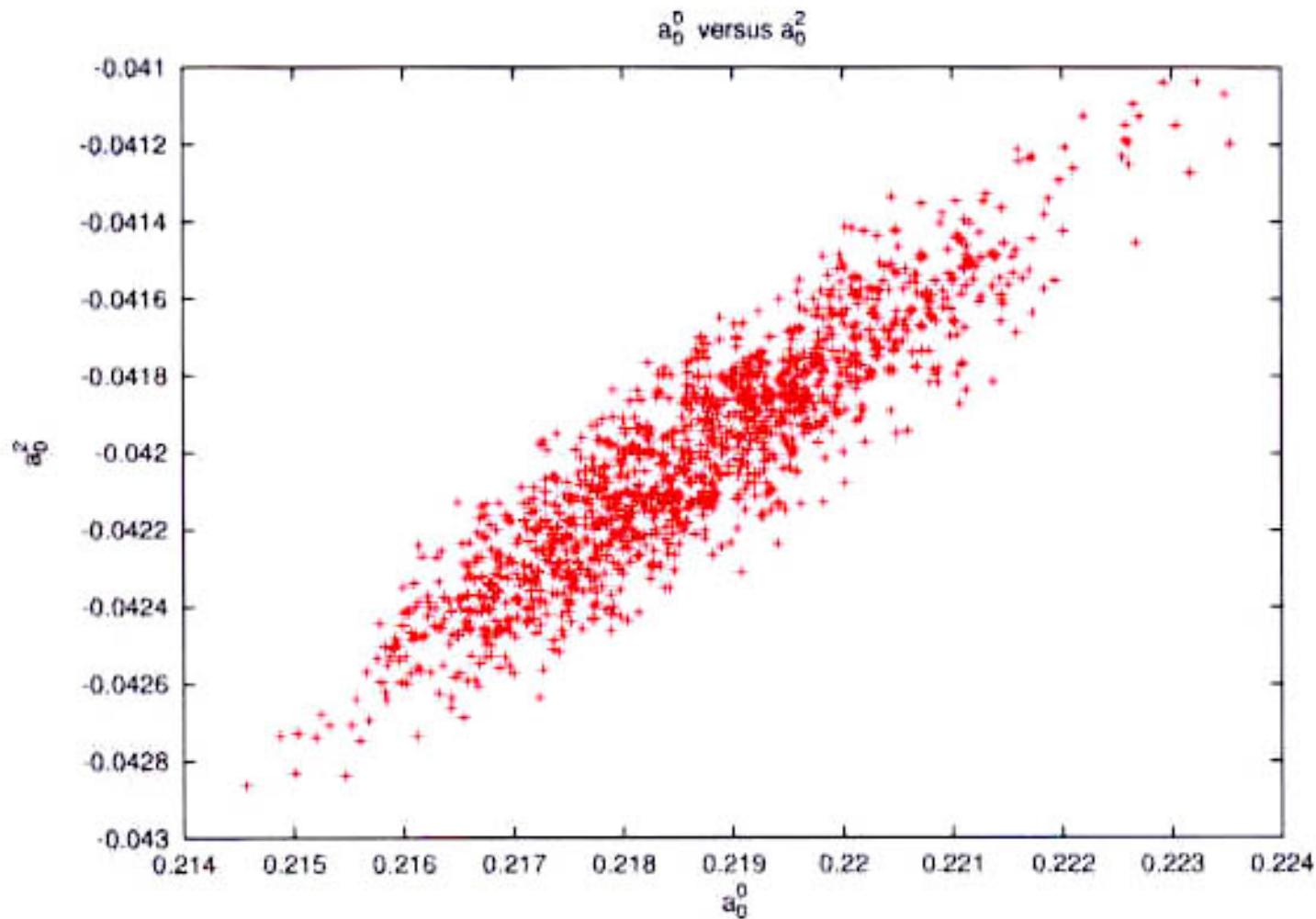


Figure 4: The prediction for the a_0^0 and a_0^2 scattering lengths for the set of L_i^* distributed according to the correlations of our main fit. Only points within the 68% confidence level domain are plotted.

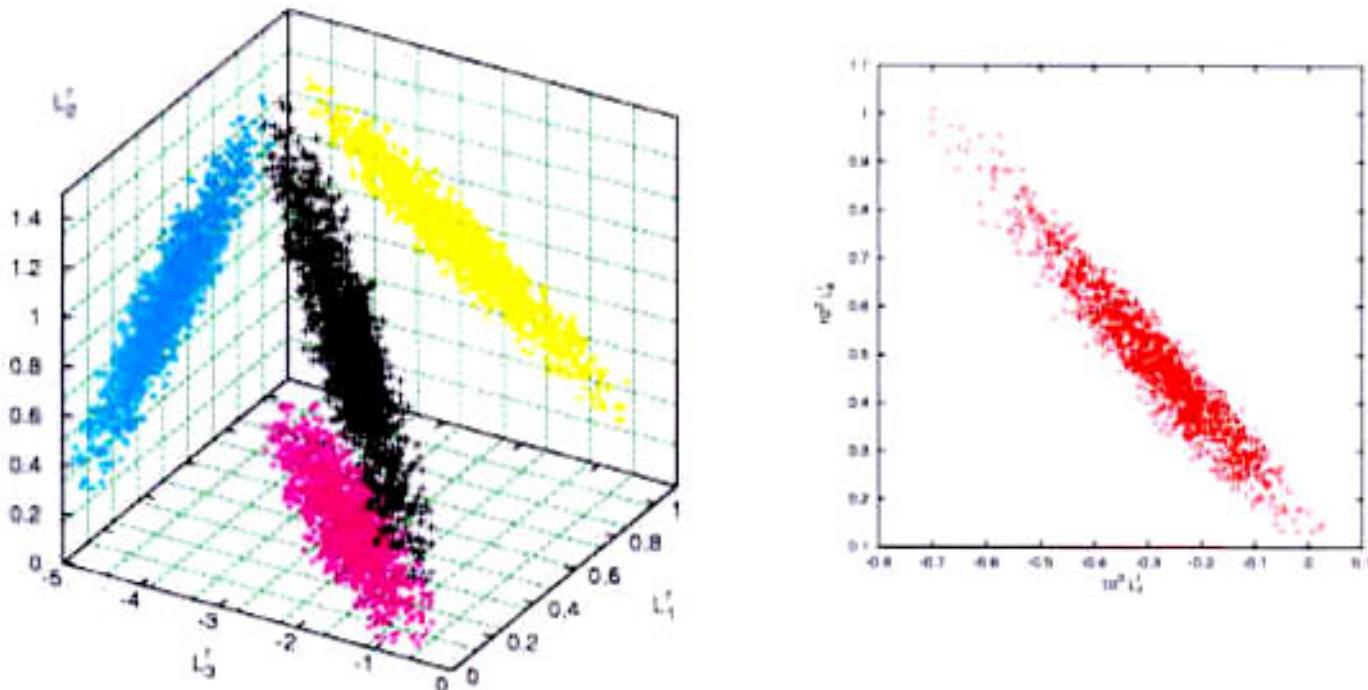
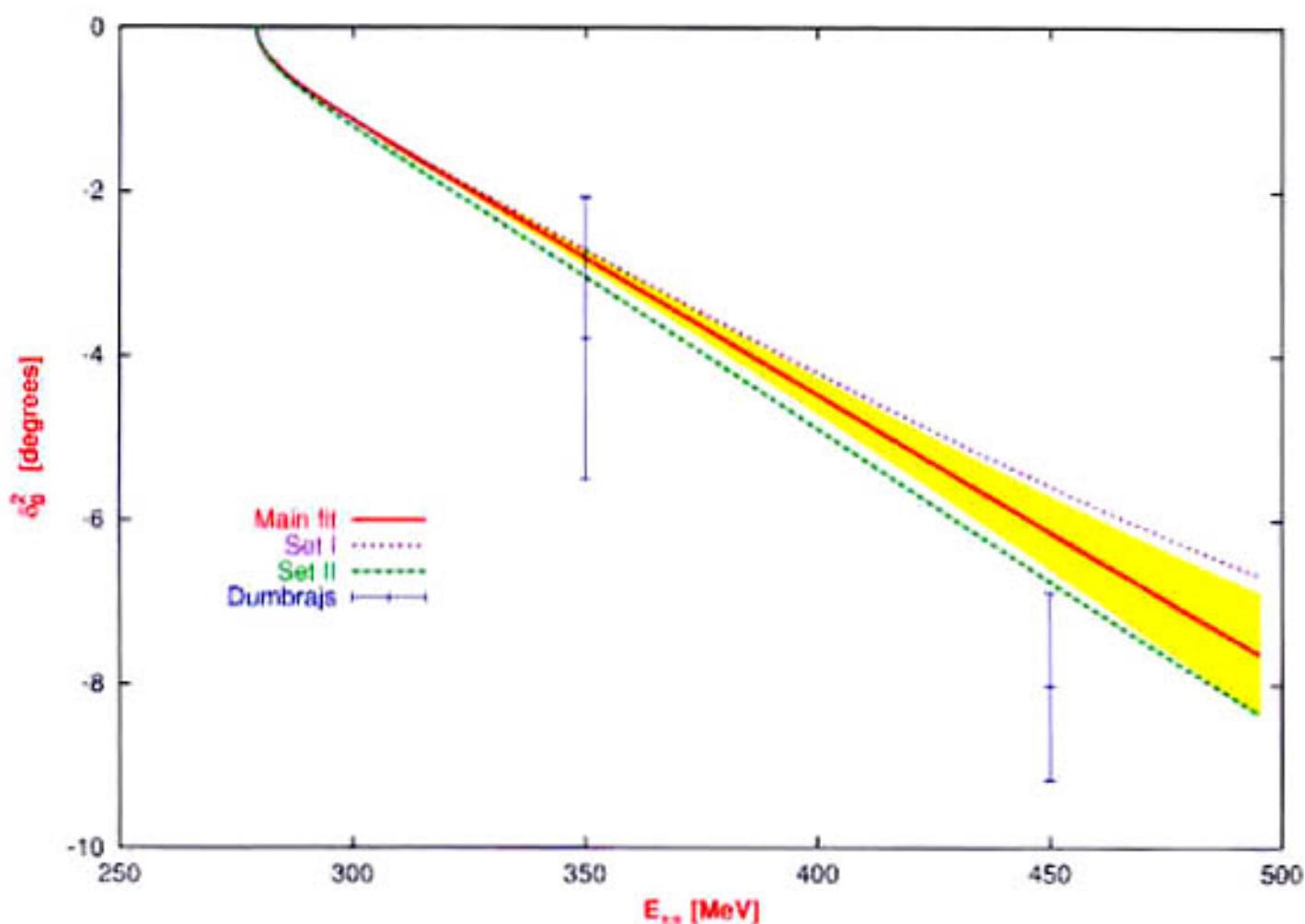
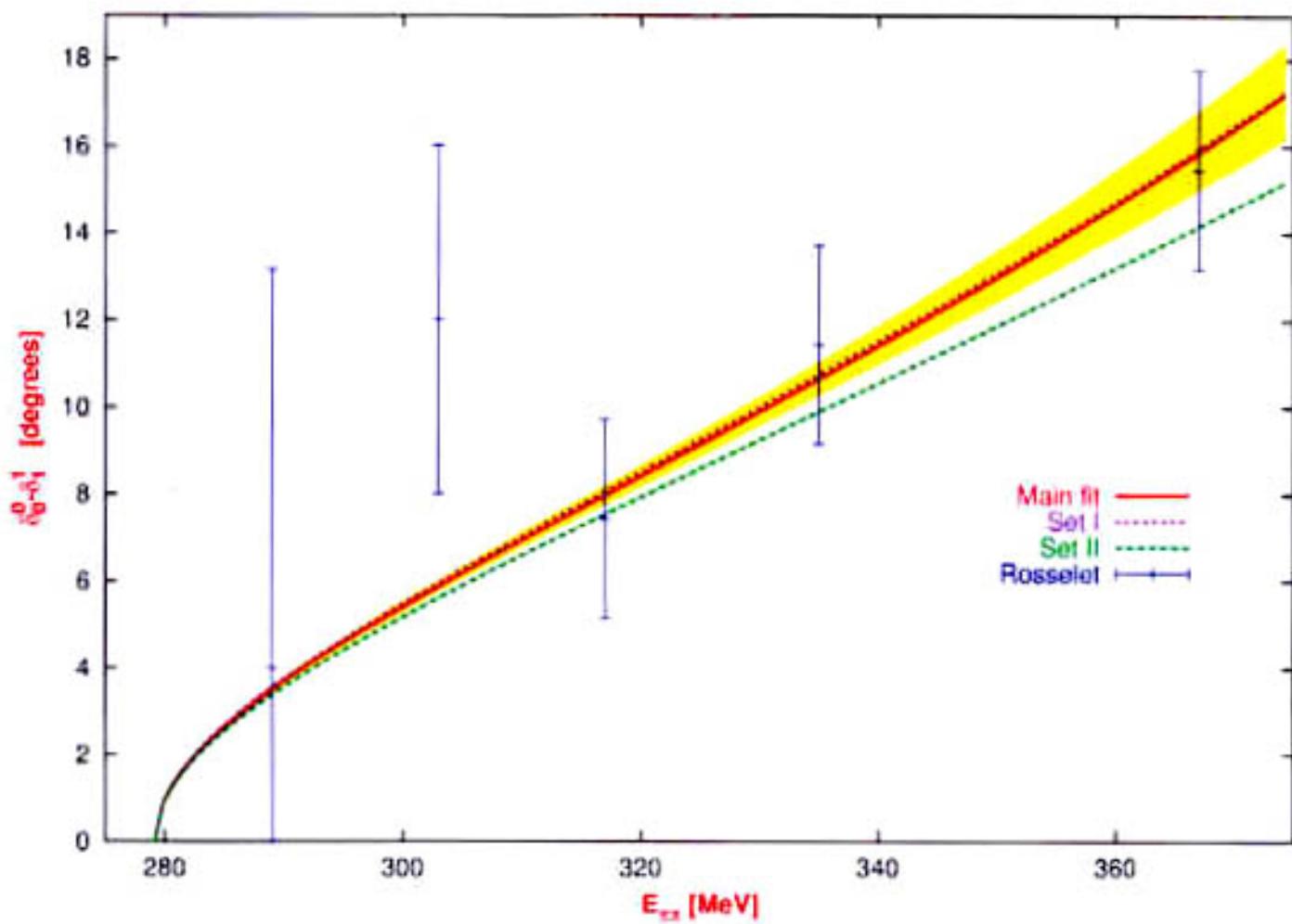


Figure 1: The sets of L'_i within a 68% confidence level range of χ^2 . In the first plot we show + the values of the first three L'_i and \times the projections on the coordinate planes. The second plot shows the L'_7 - L'_8 correlations.



Phase-shift δ_0^2 as function of the center of mass energy $E_{\pi\pi}$. Data points are taken from Hoogland et al. The shaded band indicates the 68% confidence level range.



Difference of phase-shifts $\delta_0^0 - \delta_1^1$ plotted as a function of the center of mass energy $E_{\pi\pi}$. Data points are taken from Rosselet et al. The shaded band is the 68% confidence level range.

masses and decay constants

$$\frac{m_n^2}{m_n^2 \text{ phys}} = 0.746 + 0.007 + 0.247$$

$$\frac{m_K^2}{m_K^2 \text{ phys}} = 0.695 + 0.019 + 0.286$$

$$\frac{m_\eta^2}{m_\eta^2 \text{ phys}} = 0.742 - 0.040 + 0.298$$

+ there is a cancellation to $O(p^4)$

(contributions with different sign)

+ large (?) $O(p^6)$

(contributions with the same sign)

$$F_n/F_0 = 1 + 0.136 - 0.076$$

$$\bar{F}_K/F_n = 1 + 0.134 + 0.086$$

$$F_\eta/F_n = 1 + 0.202 + 0.069$$

Why is $O(p^6)/O(p^4)$ so big?

- $m_s/\hat{m} = 20, 30$.
- resonance saturation.
- $L_4 \neq 0, L_6 \neq 0$.
- Is it an accident?

Summary

- NEW FIT FOR L_i . SOME OF THEM WITH APPRECIABLE CHANGE.
- THE FORM-FACTORS IN AGREEMENT WITH THE EXPERIMENT. Rosselet
- 'STRANGE' BEHAVIOUR (?) FOR THE MASSES.
- RECALCULATION OF THE $\pi\pi$ Scattering param. [WARNING: matching to one-loop!].