A numerical study of partially twisted boundary conditions

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Orientation - lattice phenomenology

• explore flavour sector of the SM



- QCD contribution → non-perturbative techniques necessary
- lattice QCD
- often theoretical uncertainty dominant

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$$\Gamma(D \to K l \nu) = \frac{G_F^2 |V_{cs}|^2}{192 \pi^3 m_D^3} \int_0^{(m_D - m_K)^2} dq^2 [\lambda(q^2)]^{3/2} |f_K^+(q^2)|^2$$

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e.g. semi-leptonic meson decay



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$$\underbrace{\langle K|O_{\mu}|D\rangle}_{LQCD} = f_{K}^{+}(q^{2}) \left(p_{D} + p_{K} - \frac{m_{D}^{2} - m_{K}^{2}}{q^{2}}\right)_{\mu} + f^{0}(q^{2})q_{\mu} \frac{m_{D}^{2} - m_{K}^{2}}{q^{2}}$$

There is a problem



The lattice approach



$$\begin{aligned} \mathbf{x}_{\mu} &= \mathbf{an}_{\mu}, \, \mathbf{n}_{\mu} = 1, 2, \dots, \mathbf{L/a} \\ \text{Fields:} & \overline{\psi}(\mathbf{x}), \, \psi(\mathbf{x}), \, \mathbf{U}_{\mu}(\mathbf{x}) \in \mathrm{SU}(3) \\ \text{momentum} \\ \text{cut-off:} & \Lambda = \frac{\pi}{\mathbf{a}} \end{aligned}$$

$$\begin{split} \langle \psi(\mathbf{x}) \overline{\psi}(0) \rangle &= \int \prod_{\mathbf{x}} dU \frac{1}{\not{p} + \mathbf{m}} e^{-\mathbf{S}_{\mathrm{eff}}[\mathbf{U}]} \\ & \stackrel{\mathrm{MC}}{\approx} \frac{1}{N_{\mathrm{meas}}} \sum_{i=1}^{N_{\mathrm{meas}}} \{\mathbf{S}(\mathbf{x}, 0)\}[\{\mathbf{U}_i\}] \\ & \text{sea quarks} \\ & \text{valence quarks} \end{split}$$

Momentum on the lattice



[Bedaque], [de Divitiis, Petronzio, Tantalo], [Sachrajda, Villadoro]

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[Sachrajda, Villadoro] - Twisted chiral perturbation theory

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$$E_{\pi,\rho} = \sqrt{m_{\pi,\rho}^2 + (\vec{p}_{lat} - \frac{1}{L}(\vec{\theta}_1 - \vec{\theta}_2))^2}$$

• partial twisting

$$\langle \bar{\psi}(\mathbf{x})\psi(0)\rangle \approx \frac{1}{N_{\text{meas}}} \sum_{i=1}^{N_{\text{meas}}} \{\mathbf{S}(\vec{\theta};\mathbf{x},0)\}[\{\mathbf{U}_i\}_{|\vec{\theta}=0}]$$

- finite volume errors exponentially supressed
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Our numerical study

- partial twisting of valence quarks $(\vec{\theta} = (0, 0, 0), (2, 0, 0), (0, \pi, 0), (3, 3, 3))$
- $\bullet\,$ combination with Fourier momenta $|\vec{\rm pL}|=0,2\pi$

• Observables:

$$\begin{split} \mathbf{E}_{\pi,\rho}^{2} &= \mathbf{m}_{\pi,\rho}^{2} + \left(\vec{p}_{lat} - \frac{\vec{\theta}_{1} - \vec{\theta}_{2}}{\mathbf{L}}\right)^{2} \\ \mathbf{E}_{\pi}\mathbf{f}_{\pi} &= \langle 0|\mathbf{A}_{0}(0)|\pi(\mathbf{p})\rangle \\ \mathbf{Z}_{\mathbf{P}} &= \langle 0|\mathbf{P}(0)|\pi(\mathbf{p})\rangle \\ \boldsymbol{\epsilon}_{\lambda}(\lambda,\mathbf{p})\mathbf{m}_{\rho}\mathbf{f}_{\rho} &= \langle 0|\mathbf{V}_{i}(0)|\rho(\lambda,\mathbf{p})\rangle \end{split}$$

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- Observables:

$\mathrm{E}_{\pi, ho}^2$	=	$m_{\pi,\rho}^2 + \left(\vec{p}_{lat} - \frac{\vec{\theta}_1 - \vec{\theta}_2}{L}\right)^2$
$E_{\pi} f_{\pi}$	=	$\langle 0 \mathcal{A}_0(0) \pi(\mathbf{p}) \rangle$
Z_P	=	$\langle 0 \mathbf{P}(0) \pi(\mathbf{p})\rangle$
$\epsilon_{\lambda}(\lambda, \mathbf{p})\mathbf{m}_{\rho}\mathbf{f}_{\rho}$	=	$\langle 0 V_i(0) \rho(\lambda, p) \rangle$

Numerical results - Dispersion relation



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Conclusions and outlook

Results:

- partial twisting → physical quantities at any momentum (cf. meson dispersion relation)
- matrix elements momentum-independent
- method does not introduce apreciable noise

Outlook:

- check for lighter quark masses (Finite volume effects $\propto e^{-Lm_{\pi,\rho}}$)
- phenomenological application (semi-leptonic decays, ...)