

# Quantum Loops in Resonance Theory : The Vector Form-Factor

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# Outline

I. ) The  $1/N_c$  counting and  $R_{\chi T}$  at LO

II. )  $R_{\chi T}$  lagrangian at NLO in  $1/N_c$

III.) Low/High energy:       $\rightarrow$  From  $R_{\chi T}$  to  $\chi PT$   
    $\rightarrow$  From  $R_{\chi T}$  to pQCD

# The Problem

Perturbative

Expansion in  $\alpha_s$

$Q^2 \gg 1 \text{ GeV}^2$

pQCD  
+  
OPE

Non-perturbative

$Q^2 \sim 1 \text{ GeV}^2$

RESONANCES

QFT ?

Non-perturbative

$\chi$  Expansion

$Q^2 \ll 1 \text{ GeV}^2$

$\chi$ Pt

Perturbative

*QCD at  $N_c \rightarrow \infty$*   
*(Minimal Hadronical Ansatz)*

# $1/N_c$ Counting and $R\chi T$

[ 't Hooft 74 ]  
[ Witten 79 ]

$1/N_c$  Suppression  
of  $q\bar{q}$  loops

$1/N_c$  Suppression of  
meson loops

**LO** given by tree-level

LO in  $1/N_c$

$R\chi T$

[Ecker et al '89]

Resonances (  $R=V,A,S,P$  ) + Goldstones (  $pNGB=\pi,K,\eta_8,\eta_1$  )

Chiral symmetry

General lagrangian for  $R + pNGB$ :  $\langle R \mathcal{O}(p^2) \rangle$

Minimal Hadronic Ansatz: *Long/short distances* ✓

# RχT Lagrangian

at LO in  $1/N_c$

$$\mathcal{L}_{2\chi}[\pi] = \frac{F^2}{4} \langle \mathbf{u}^\mu \mathbf{u}_\mu + \chi_+ \rangle$$

$$\mathcal{L}_{2V}[\pi, \mathbf{V}] = -\frac{1}{2} \langle \nabla_\lambda \mathbf{V}^{\lambda\mu} \nabla^\nu \mathbf{V}_{\nu\mu} - \frac{1}{2} \mathbf{M}_V^2 \mathbf{V}_{\mu\nu} \mathbf{V}^{\mu\nu} \rangle$$

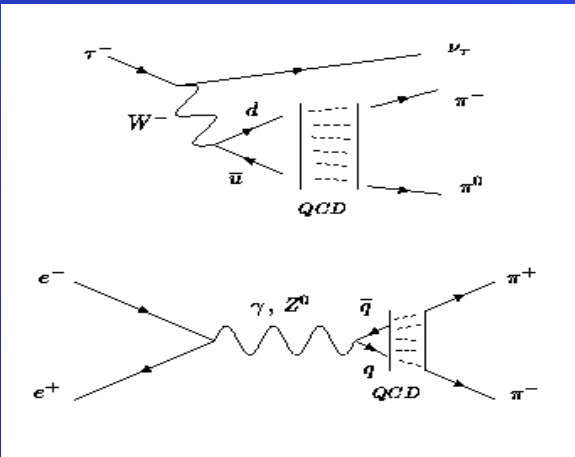
$$+ \frac{\mathbf{F}_V}{2\sqrt{2}} \langle \mathbf{V}_{\mu\nu} \mathbf{f}_+^{\mu\nu} \rangle + \frac{\mathbf{iG}_V}{2\sqrt{2}} \langle \mathbf{V}_{\mu\nu} [\mathbf{u}^\mu, \mathbf{u}^\nu] \rangle$$

$$\mathcal{L}_{2S}[\pi, \mathbf{S}] = c_d \langle \mathbf{S} \mathbf{u}^\mu \mathbf{u}_\mu \rangle + \dots$$

$$\mathcal{L}_{2A}[\pi, \mathbf{A}], \quad \mathcal{L}_{2P}[\pi, \mathbf{P}]$$

## QCD - Long/Short Distance

- Vector FF
- Axial FF
- Scalar FF
- 2 point Green-Funcs.
- Scattering Ampls.



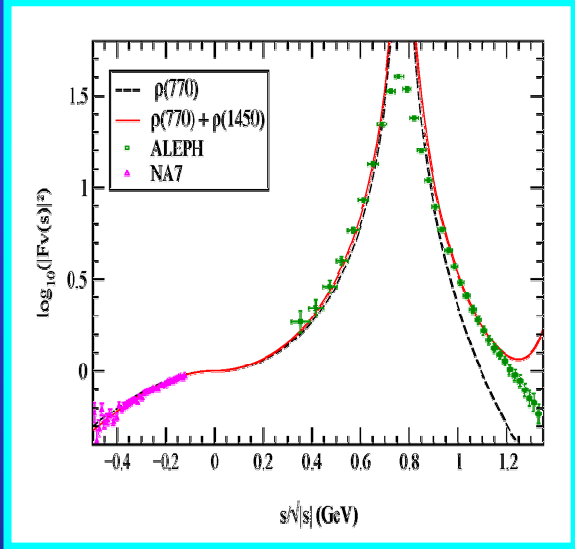
# VFF at LO in $1/N_c$ : Tree-level amplitudes

$$\langle \pi^+ \pi^- | \left( \frac{1}{2} \bar{u} \gamma^\mu u - \frac{1}{2} \bar{d} \gamma^\mu d \right) | 0 \rangle = \mathcal{F}(q^2) (p_{\pi^+} - p_{\pi^-})^\mu$$

**LO in  $1/N_c$**

$$\mathcal{F}(q^2) = 1 + \underbrace{\frac{F_V G_V}{F^2}}_{\text{QCD Short Dist.}} \frac{q^2}{M_V^2 - q^2} \xrightarrow{q^2 \rightarrow \infty} 0$$

$F_V G_V = F^2$



*QCD at NLO in  $1/N_c$ :*  
*Quantum loops*



# $\mathcal{L}_{R\chi T}$ at LO in $1/N_C$ + QCD short distance

\* Couplings fixed by QCD-short distances:

$$\frac{F_V}{\sqrt{2}} = \sqrt{2} G_V = 2 c_d = F \sim \mathcal{O}(\sqrt{N_C})$$

\* There are also mass relations

$$M_A^2 = 2 M_V^2, \quad M_P^2 = 2 M_S^2 (1 - \delta) \sim \mathcal{O}(1)$$

$$\delta \approx 3\pi\alpha_S F^2 / M_S^2 \sim 0.08\alpha_S$$

WELL DEFINED

$1/N_C$  counting

with LO vertices

$$\sum_{n_{\text{loops}}} \left( \frac{1}{(4\pi)^2} \frac{p^2}{F^2} \right)^{n_{\text{loops}}}$$

# NLO Lagrangian

$R\chi T$  at LO [Ecker et al'89]

One Loop Diagrams  $\longrightarrow$  NLO Contributions

However, Loops=Divergencies  $\mathcal{O}(p^4)$ !!

New NLO pieces (NLO couplings)

**Minimal Extension:** Pieces to renormalize

# Perturbative calculation:

Next-to-leading order in  $1/N_C$

Simplifications:

- $m_q=0$
- $U(2)_L \times U(2)_R$

## NLO in $1/N_C$ :

- One loop  $\rightarrow$  LO vertices
- Tree-level  $\rightarrow$  LO vertices and one NLO vertex

## Renormalization procedure:

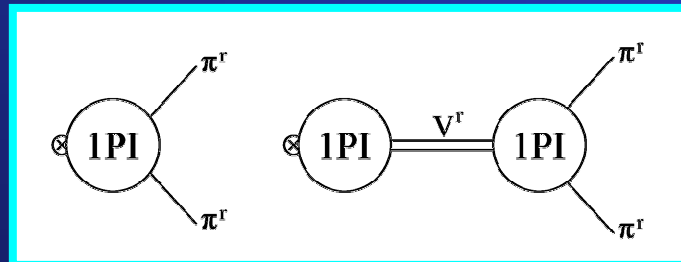
Renormalized  
propagators  
(self-energy)



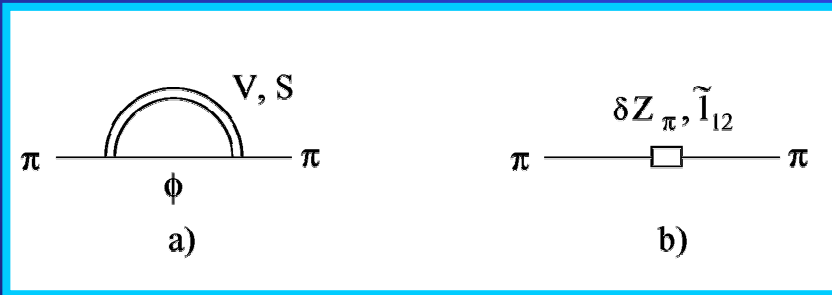
Renormalized  
vertex-functions  
(1PI)



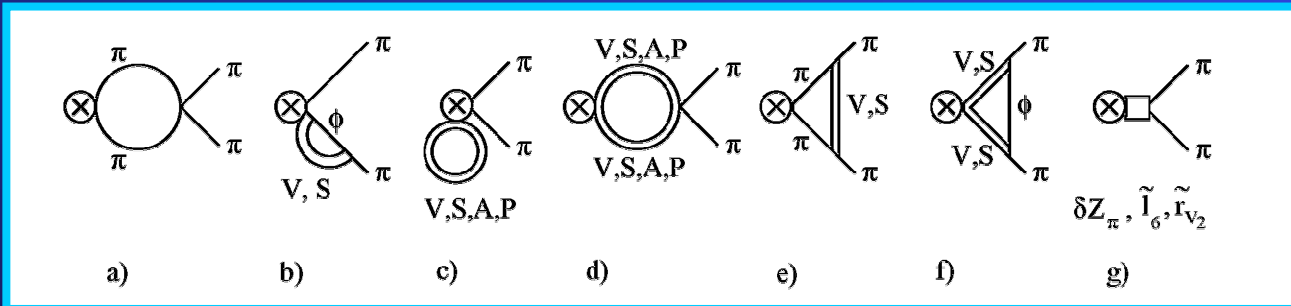
Renormalized  
VFF



# • Only External Pions



$$-\tilde{l}_{12} \langle \nabla^\mu \mathbf{u}_\mu \nabla^\nu \mathbf{u}_\nu \rangle$$

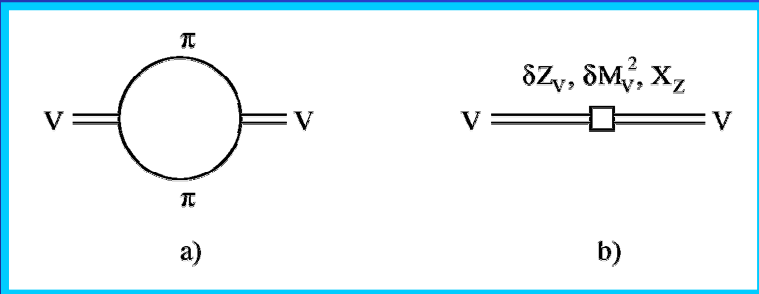


$$\frac{i \tilde{l}_6}{4} \langle \mathbf{f}_+^{\mu\nu} [\mathbf{u}_\mu, \mathbf{u}_\nu] \rangle + i \tilde{c}_{51} \langle \nabla^\rho \mathbf{f}_+^{\mu\nu} [\mathbf{h}_{\mu\rho}, \mathbf{u}_\nu] \rangle + \dots$$

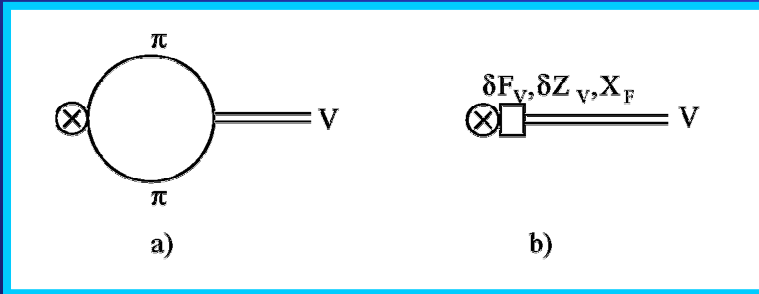
$\mathcal{L}_{2\chi}$   
 +  
 $\tilde{\mathcal{L}}_{4\chi}$   
 +  
 $\tilde{\mathcal{L}}_{6\chi}$

Other processes  $\longrightarrow$  Remaining operators  
 and  $\tilde{l}_i$  and  $\tilde{c}_j$  couplings

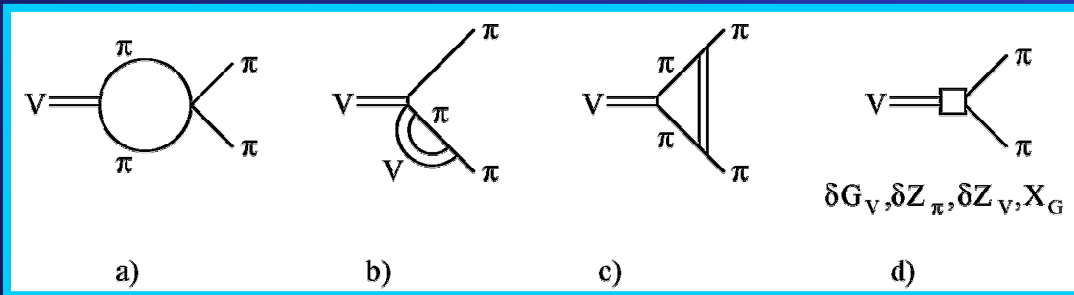
# External Pions and Vectors



$$\frac{X_{Z_1}}{2} \langle \nabla^2 V^{\mu\nu} \{ \nabla_\nu, \nabla^\rho \} V_{\mu\rho} \rangle + \dots$$



$$X_{F_1} \langle V_{\mu\nu} \nabla^2 f_+^{\mu\nu} \rangle + \dots$$



$$i X_{G_1} \langle \{ \nabla^\rho, \nabla_\mu \} V^{\mu\nu} [u_\nu, u_\rho] \rangle + \dots$$

$\mathcal{L}_{2V}$

+

$\mathcal{L}_{4V}$

# Equations of Motion

$$\ell_6^{\text{eff}} = \ell_6 + 2X_Z F_V G_V - 2\sqrt{2} X_F G_V - 4\sqrt{2} X_G F_V$$

$$\mathbf{r}_{V_1}^{\text{eff}} = \mathbf{r}_{V_1} = 4F^2(c_{51} - c_{53})$$

$$\mathbf{F}_V^{\text{eff}} = \mathbf{F}_V + 2X_Z F_V M_V^2 - 2\sqrt{2} X_F M_V^2$$

$$\mathbf{G}_V^{\text{eff}} = \mathbf{G}_V + 2X_Z G_V M_V^2 - 4\sqrt{2} X_G M_V^2$$

$$M_V^{\text{eff}2} = M_V^2 + 2X_Z M_V^4$$

• EOM from the LO lagrangian

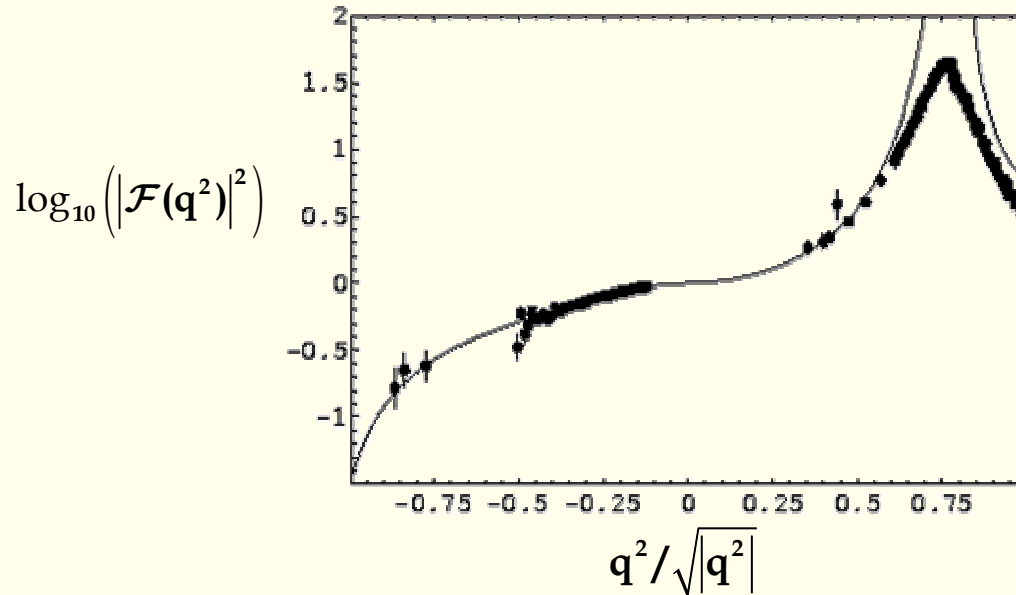


• Full reduction of  $\mathcal{L}_{4V}$  to  $(\mathcal{L}_\chi + \mathcal{L}_{2V})^{\text{eff}}$

$$\mathcal{F}(q^2) = \mathcal{F}(\ell_6^{\text{eff}}, \mathbf{r}_{V_1}^{\text{eff}}, \mathbf{F}_V^{\text{eff}}, \mathbf{G}_V^{\text{eff}}, M_V^{\text{eff}})$$

*Experimental Results*  
*and*  
*High/Low Energy Limits*

# Space-like form-factor



$$\mu = 770 \text{ MeV}$$

$$\frac{F_V^{\text{eff}} G_V^{\text{eff}}}{F^2} = 1.1$$

$$\frac{2G_V^{\text{eff}2}}{F^2} = 1$$

$$M_V^{\text{eff}} = 776 \text{ MeV}$$

Radiative corrections:

**UNDER CONTROL !**

$$M_A^2 = 2M_V^2$$

$$2c_d = F$$

$$M_S = 1 \text{ GeV}$$

$$M_P^2 = 2M_S^2$$

$$\ell_6^{\text{eff}} = 0, \quad r_{V_1}^{\text{eff}} = 0$$

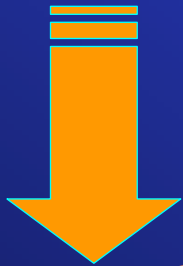


# From $R\chi T$ to $\chi PT$

$R\chi T$  ( $q^2 \ll M_R^2, \Lambda_\chi^2$ )

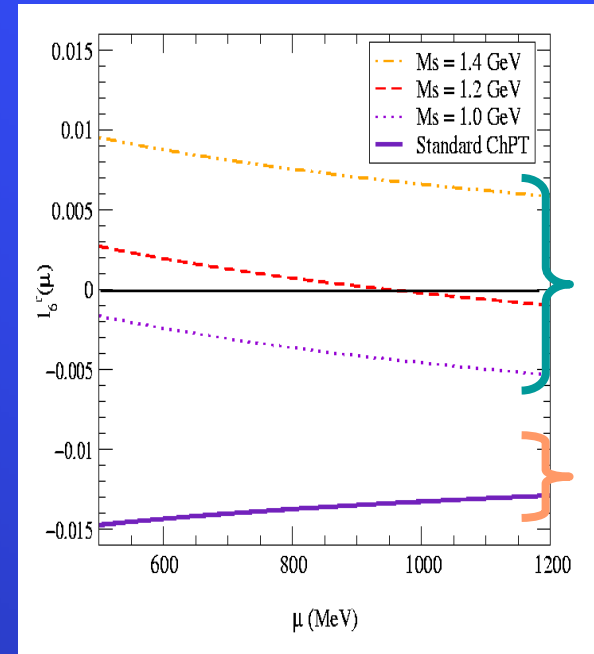
$$\mathcal{F}(q^2) = 1 + \left[ \frac{F_V^{\text{eff}}(\mu) G_V^{\text{eff}}(\mu)}{M_V^{\text{eff}2}(\mu)} - \frac{\tilde{\ell}_6^{\text{eff}}(\mu)}{F^2} - \frac{1}{12\pi^2} \ln\left(\frac{M_V^2}{\mu^2}\right) + \dots \right] q^2 +$$

$$-\frac{1}{96\pi^2 F^2} q^2 \ln\left(-\frac{q^2}{\mu^2}\right) + \dots$$



$\mathcal{O}(p^4)$   $\chi PT$

$$\mathcal{F}(q^2) = 1 - \frac{\ell_6^{\chi PT}(\mu)}{F^2} q^2 - \frac{1}{96\pi^2 F^2} q^2 \ln\left(-\frac{q^2}{\mu^2}\right) + \dots$$



$\tilde{\ell}_6^{\text{eff}}(\mu)$

$\ell_6^{\chi PT}(\mu)$

$\left( \begin{array}{l} \ell_6 \text{ in } n_f=2 \\ L_9 \text{ in } n_f=3 \end{array} \right)$

\* Phenomenological suppression

$$\tilde{\ell}_6^{\text{eff}}(\mu) \ll \ell_6^{\chi PT}(\mu)$$

$\uparrow$   
 $\mathcal{O}(1)$

$\uparrow$   
 $\mathcal{O}(N_D)$

\* Similar with  $\mathcal{O}(p^6)$   $\chi PT$

# From $R\chi T$ to pQCD

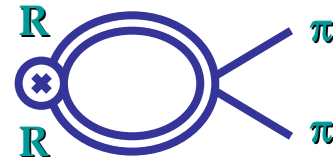
**pQCD**  $Q^2 \rightarrow \infty$  :

$$\mathcal{F}(Q^2) \xrightarrow{Q^2 \rightarrow \infty} 0$$

**$R\chi T$**   $Q^2 \rightarrow \infty$  :

Intermediate Resonance states:

**Constant Resonance form-factors !**



$$\mathcal{F}(Q^2) = a_{(4)} \frac{Q^4}{F^2} + a_{(2)} \frac{Q^2}{F^2} + a_{(0)} \frac{M_R^2}{F^2} + \mathcal{O}\left(\frac{M_R^4}{Q^2 F^2}\right)$$

constants  
&  $\ln(Q^2)$

If  $a_{(j)}=0$  then  
**RENORMALIZATION**  
through LO lagrangian

# *Conclusions*

- QCD at large  $N_c$  : *Minimal Hadronical Ansatz*

- $R\chi T$  at NLO in  $1/N_c$  :

- Quantum loops: Well defined counting and renormalization

- Radiative corrections: Small within the range  $|q^2| \leq 1 \text{ GeV}^2$

- Low energy limit:  $R\chi T \xrightarrow{q^2 \ll \Lambda_\chi^2} \chi PT$

- High energy limit: - Wrong behaviour  
(wrong resonance FF)

- Need for extra operators

- with 2, 3... Resonance fields, e.g.  $\langle V A \mathcal{O}(p^2) \rangle$