Wilson Lines in Warped Space Dynamical Symmetry Breaking and Restoration

Kin-ya Oda

University of Bonn

26.11.2004, Euro GDR workshop, Frascati

in collaboration with Andreas Weiler (TU München) hep-ph/0410061

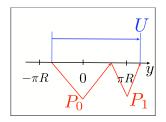
Outline



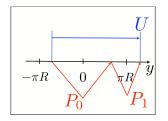
- 2 Setup: 5-Dim SU(N) with Warped Compactific'n on S^1/Z_2
- 3 Analysis: KK Expansions and Effective Potential

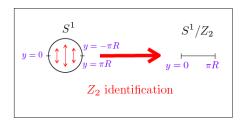
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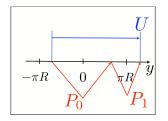


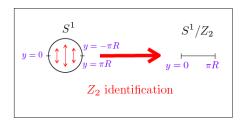
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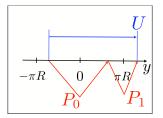
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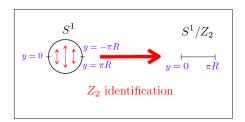




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Orbifold Compactification on S^1/Z_2





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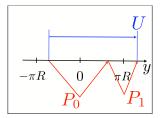
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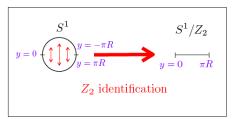
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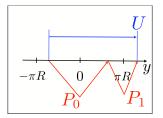
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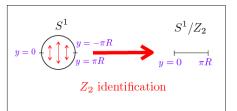
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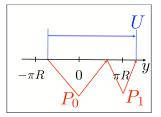
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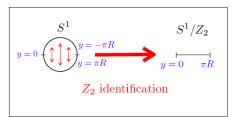
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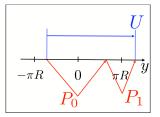


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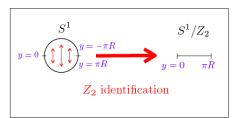
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$$P_i: \partial_y \to -\partial_y$$

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5-dim SU(5) GUT on S^1/Z_2 Kawamura's realization of doublet-triplet (2-3) splitting

[Kawamura '00, '01; Altarelli, Feruglio '01; Hall, Nomura '01, '02; Hebecker, March-Russell '01]

$$P_0 = I_5, \quad P_1 = \begin{pmatrix} -I_3 \\ I_2 \end{pmatrix}, \quad U = P_1 P_0 = \begin{pmatrix} -I_3 \\ I_2 \end{pmatrix}.$$

Then, (P_0, P_1) parities are

Take

$$5_H, \overline{5}_H : \begin{pmatrix} +-\\ ++ \end{pmatrix}, \quad A_\mu : \begin{pmatrix} ++ & +-\\ +- & ++ \end{pmatrix}, \quad A_y : \begin{pmatrix} -- & -+\\ -+ & -- \end{pmatrix}$$

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$$\Phi(++) \sim \cos rac{n}{R} y, \quad \Phi(--) \sim \sin rac{n}{R} y, \ \Phi(\pm \mp) \sim iggl\{ \cos s \ \sin
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Comments:

• $1/R \sim M_{GUT}$

(*n* = 0, 1, 2, . . .)

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- Zero mode resides only in ++ part
- 2-3 splitting realized!

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Gauge Twist and Scherk-Schwarz SUSY Breaking

Similarly, for $SU(2)_R$ auxiliary gauge field, continuous twists

$$\mathcal{U} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}, \quad \mathcal{P}_0 = \begin{pmatrix} 1 \\ & -1 \end{pmatrix}, \quad \mathcal{P}_1 = \begin{pmatrix} \cos \theta & \sin \theta \\ \sin \theta & -\cos \theta \end{pmatrix}$$

yield SUSY breaking: $m_{\text{SUSY}} \simeq \theta / R$. Recall: $u = \mathcal{P}_1 \mathcal{P}_0, \mathcal{P}_0^2 = \mathcal{P}_1^2 = 1$

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- One might think of putting

$$heta \sim m_{
m SUSY}/M_{
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at classical level. [Barbieri, Hall, Nomura '01]

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 However, quantum correction leads [von Gersdorff, Quiros, Riotto '02]

$$heta = 0$$
 or $heta \sim 1$

w/o extra boundary term.

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w/o extra boundary term.

• SUSY unbroken or broken at GUT scale!

Motivation

5-dim SU(5) GUT on S¹ /Z₂ Scherk-Schwarz SUSY Breaking Motivation: Warped Space, Two Hierarchical Scales

We can use large hierarchy in Randall-Sundrum warped compactification, which is also on S^1/Z_2 , if we learn how to treat continuous Wilson lines associated with continuous twists!

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Randall-Sundrum Geometry

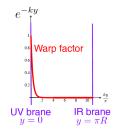
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Randall-Sundrum Geometry

AdS Metric for $-\pi R < y \leq \pi R$:

$$e^{-2k|y|}\eta_{\mu
u}dx^{\mu}dx^{
u}+dy^{2}$$

(k: [AdS radius]⁻¹ $\lesssim M_{\text{Planck}}$)



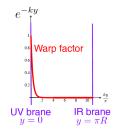
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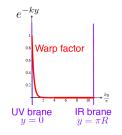
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*M*_{UV} in this coordinate frame (normalized at UV brane) will be observed as the physical mass:

$$m_{IR} = e^{-k\pi R} M_{UV},$$

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by an observer at IR brane.

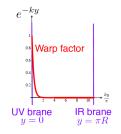
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 If, say, kR ≃ 11, large hierarchy generated

$$a \equiv e^{-k\pi R} \simeq 10^{-15}.$$

Gauge Twists and $\langle A_y \rangle$ Are Mixed Hosotani Mechanism

Winson Line in Orbifold S^1/Z_2

• In compact space, there's non-contractable loop.

Gauge Twists and $\langle A_y \rangle$ Are Mixed Hosotani Mechanism

Winson Line in Orbifold S^1/Z_2

- In compact space, there's non-contractable loop.
- Integral along the loop cannot be gauged away.

Gauge Twists and $\langle A_y \rangle$ Are Mixed Hosotani Mechanism

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Slogan

Make use of zero modes of A_y as adjoint Higgs fields.

Gauge Twists and $\langle A_y \rangle$ Are Mixed Hosotani Mechanism

General Twists, Zero Modes

 Most general SU(N) gauge twists can be block-diagonalized into SU(2) subspaces:

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- Zero modes of $\langle A_y \rangle$ is living only inside each SU(2).
- **Kawamura's twist** doesn't contain any *SU*(2) blocks and hence any continuous d.o.f.
- Hereafter concentrate on a SU(2) block with above twists.

Gauge Twists and $\langle A_y \rangle$ Are Mixed Hosotani Mechanism

Hosotani Mechanism

Large gauge transformation

$$\Omega = \mathbf{e}^{i\varphi \mathbf{y}\mathbf{T}}: \quad \mathbf{g}\langle \mathbf{A}_{\mathbf{y}} \rangle \to \mathbf{g}\langle \mathbf{A}_{\mathbf{y}} \rangle - \frac{\varphi}{\pi \mathbf{R}}, \quad \theta \to \theta - \varphi, \quad \mathbf{m}_{\mathbf{n}} \to \mathbf{m}_{\mathbf{n}}$$

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T: gauge direction of the zero-mode ($T \propto \sigma_2$ in this case)

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- Symmetry breaking pattern determined by dynamics rather than by hand!

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KK Expansions 1-Loop Determination of Wilson Lines

Kaluza-Klein Expansions, Simple when Even or Odd

 When field has definite odd or even parity, KK expansions, obtained straightforwardly [Gherghetta, Pomarol '00]

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- When field has definite odd or even parity, KK expansions, obtained straightforwardly [Gherghetta, Pomarol '00]
- As in flat case, there is zero mode of A_y contributing to Wilson lines
- This time, it is **not flat** but exponentially localized on IR brane $\propto z(y)^2$, where

$$z(y)=e^{k|y|}.$$

(Recall: *k* is [AdS curvature radius]⁻¹ $\lesssim M_{\text{Planck}}$)

KK Expansions 1-Loop Determination of Wilson Lines

KK Expansions for Most General Twist and B.G.

• Difficulty: Non-diagonal twist $\theta \neq 0$



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- → We've found a gauge-transf. to **remove** higher modes!
- \rightarrow Then we can **freely** take $\theta = 0$ gauge (or whatever)

KK Expansions 1-Loop Determination of Wilson Lines

KK Exp'ns for Most General Twist and B.G. (Cont'd)

KK expansions in $\theta = 0$ gauge:

$$\begin{aligned} A_{\mu}^{\pm} &\sim \left(\cos\frac{\mathbf{v}z^{2}}{2} \pm i\epsilon(\mathbf{y})\sin\frac{\mathbf{v}z^{2}}{2}\right) z \left[\alpha_{n}J_{1}(\mu_{n}z) + \beta_{n}Y_{1}(\mu_{n}z)\right], \\ A_{y}^{\pm} &\sim \left(\epsilon(\mathbf{y})\cos\frac{\mathbf{v}z^{2}}{2} \pm i\sin\frac{\mathbf{v}z^{2}}{2}\right) z^{2} \left[\alpha_{n}J_{0}(\mu_{n}z) + \beta_{n}Y_{0}(\mu_{n}z)\right], \\ \text{where } z &= e^{k|\mathbf{y}|}, \ \mathbf{v} \sim \langle A_{y} \rangle \text{ and } \epsilon(\mathbf{y}) = \begin{cases} +1 & \text{for } \mathbf{y} > 0, \\ -1 & \text{for } \mathbf{y} < 0. \end{cases} \\ \text{Note:} \end{aligned}$$

• α_n, β_n , coefficients of two independent solutions to KK eq

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where $z = e^{k|\mathbf{y}|}, \ \mathbf{v} \sim \langle A_{y} \rangle$ and $\epsilon(\mathbf{y}) = \begin{cases} +1 & \text{for } \mathbf{y} > 0, \\ -1 & \text{for } \mathbf{y} < 0. \end{cases}$
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Note:

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- μ_n , dimensionless KK masses
- Wave functions continuous, even if accompanied by $\epsilon(y)$

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KK Expansions 1-Loop Determination of Wilson Lines

Calculation of Effective Potential for $\langle A_{y} \rangle$

Effective potential for $v = \langle A_y \rangle / k$: (Recall: $k = [AdS curvature radius]^{-1} \leq M_{Planck}$)

$$V_{\text{eff}}(\mathbf{v}) = a^4 I_{IR} + I_{UV} \\ + \frac{(ka)^4}{(4\pi)^2} \int_0^\infty dx \, x^3 \log\left[1 - \frac{I_0(x)K_1(x) - K_0(x)I_1(x) - \frac{1}{x}\cos a^{-2}\mathbf{v}}{2I_0(x)I_1(x)\left(\gamma + \log\frac{ax}{2}\right)}\right],$$

where $a = e^{-k\pi R} \sim 10^{-15}$ is warp factor

Dimensional reduction in UV frame

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- Dimensional reduction in UV frame
- Dim-reg in 4-dim with infinite number of KK towers
- Zeta-function regularization, replacing infinite KK sum by contour integral
- I_{IR}, I_{UV}: independent of v and a → Can be renormalized by brane tensions

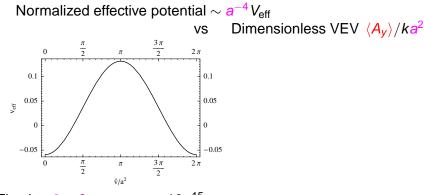
KK Expansions 1-Loop Determination of Wilson Lines

Effective Potential for $\langle A_y \rangle$

Normalized effective potential $\sim a^{-4} V_{eff}$ vs Dimensionless VEV $\langle A_y \rangle / k a^2$

KK Expansions 1-Loop Determination of Wilson Lines

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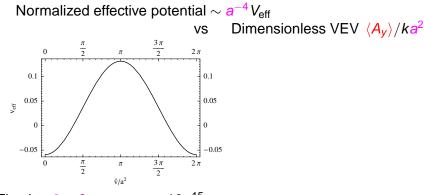


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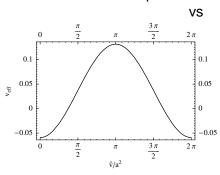
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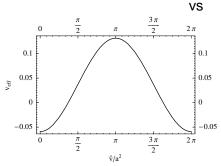
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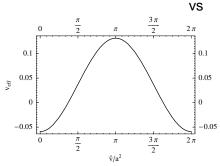
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- DSB scale is ka ~ TeV!

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 Studied 5-dim SU(N) pure gauge theory with warped compactification on S¹/Z₂

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• Dynamical Schwerk-Schwarz SUSY breaking...

Necessary to find a supergravity configuration proposed by Abe and Sakamura

[Recently appeared toy models: Eto, Maru, Sakai '03; Correia, Schmidt, Tavartkiladze '04]

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