EuroGDR Supersymmetry - 4th meeting – Frascati 25 November 2004 –

# Recent developments in Astroparticle physics

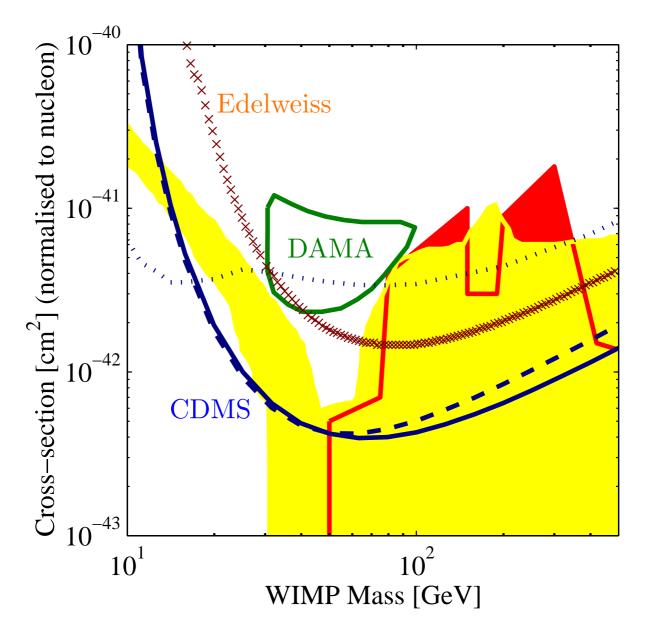
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# OUTLINE

- 1. News from the observational front:
- Dark matter: CDMS
- Some  $\gamma$  rays puzzles from our galaxy...
- CMB polarization: CAPMAP, CBI, DASI
- Large Scale Structure: SDSS Ly $\alpha$ data
- 2. News from the theory side:
- Running spectral index or not ?
- Cosmology meets String theory
- Beyond the classical neutralino DM scenarios
- 3. Outlook

# DARK MATTER:

New limit spin-independent cross-section from CDMS ! [astro-ph/0405033]



DAMA region excluded... BUT: a spin-dependent cross section could still explain all data, for mass window  $\simeq 5-13$  GeV [Savage, Gondolo & Freese '04]

 $\rightarrow$  most of the SUSY parameter space still available more in the talks tomorrow afternoon !

# Galaxy $\gamma$ ray puzzles:

SPI/INTEGRAL has (again) evidence for a 0.511 MeV emission line from the centre of the galaxy: e<sup>+</sup>e<sup>-</sup> annihilation at rest !
??? Annihilating/decaying CDM ???
→ see P. Fayet talk tomorrow

• EGRET excess of  $\gamma$  rays at about 1-10 GeV  $\rightarrow$  see W. de Boer and A. Morselli talks

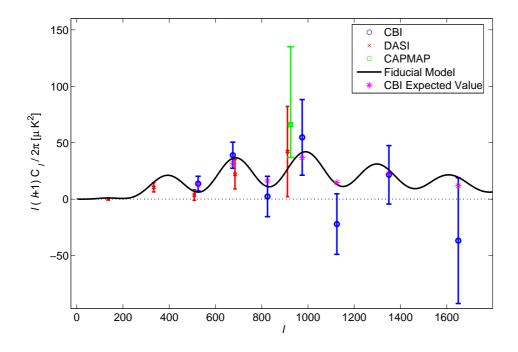
 also some γ ray emission seen around TeV by VERITAS, CANGAROO and H.E.S.S., but the position and energy spectra do not match...: different sources ???

Possible explanation from annihilating DM requires very heavy masses  $M_{\chi} > 12$  TeV

Astrophysical explanations probably more plausible... more data needed !

#### CMB polarization

New measurements of the E-mode by DASI, CAPMAP and CBI [astro-ph/0409357,380,569]



#### Figure from CBI [astro-ph/0409569]

 $7\sigma$  detection, consistent with WMAP and scalar perturbations, but error bars still very large ! No evidence for the B-mode.

WMAP 2nd release and other polarization measurements (Boomerang...) still to come...

#### Large Scale Structure: the Sloan Digital Sky Survey

New determination of the power spectrum at small scales from Ly  $\alpha$  data [astro-ph/0405013,07377]

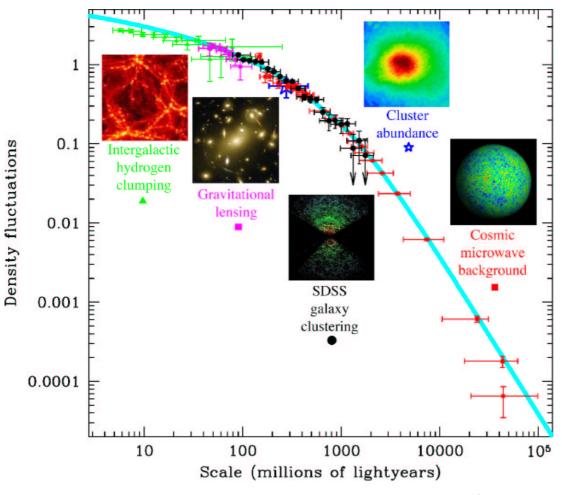


Figure by M. Tegmark

- $\longrightarrow$  better control of systematics
- $\longrightarrow$  longer lever arm in scale

**NEW** determination of the cosmological parameters from WMAP data, SDSS galaxy clustering, bias and Lyman  $\alpha$  data, SN Ia data. New analysis by Seljak *et al* [astro-ph/0407372] What has changed ???

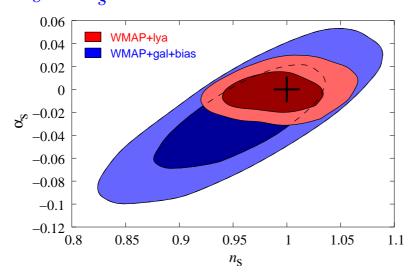
#### Improved accuracy on all parameters, e.g. 0.78 WMAP+gal+SN+bias+lya 0.76 WMAP+gal+SN+bias 0.74 0.72 Ч 0.7 0.68 0.66 0.64 0.62 0.24 0.26 0.28 0.32 0.34 0.36 0.38 0.3

More stringent bound on neutrino masses

 $\Omega_{\rm m}$ 

$$\sum m_{\nu} \le 0.66 \text{ eV} (3\nu) \text{ or } m_{\nu} \le 0.79 \text{ eV} (3+1\nu)$$

Spectral index  $n_s = 0.977^{+0.025}_{-0.021}$  and  $\alpha_s = n'_s = -0.003 \pm 0.010$ : NO RUNNING !



Also bound on the tensors as  $r \leq 0.45$  at  $1\sigma$ .

### What can we say about inflation then ?

In single field inflation with  $V(\phi)$  the scalar power spectrum is

$$\mathcal{P}_{\mathcal{R}}(k) = \left. \frac{1}{12\pi^2 M_P^6} \frac{V^3}{V'^2} \right|_{k=aH} \propto k^{n-1}$$

and the spectral index:

$$n(k) - 1 = \left. \frac{d \log(\mathcal{P}_{\mathcal{R}})}{d \log(k)} \right|_{k=aH} = 2\eta - 6\epsilon + \dots$$

So  $n'(k_0)$  arises only at second order:

$$n'(k) = \frac{2}{3} \left( (n-1)^2 - 4\eta^2 \right) + 2\xi$$

where

$$\epsilon = \frac{M_P^2}{16\pi} \frac{(V')^2}{V^2} \qquad \eta = \frac{M_P^2}{8\pi} \frac{V''}{V} \qquad \xi = \frac{M_P^4}{64\pi^2} \frac{V'V'''}{V^2}$$

so we expect  $n' \propto (n-1)^2$  (or  $\xi$  must be large).

The new result is consistent with this !

Lyman  $\alpha$  data give stronger constraints for models with substantial running, e.g.  $\rightarrow$  running mass models !

# Running mass model(s) [Stewart '96,'97] $\phi \rightarrow \text{flat direction of the} \qquad V'_{SUSY}(\phi) = 0$ SUSY potential

Break supersymmetry explicitly in a hidden sector and obtain a soft mass for the inflaton field:

$$V(\phi) = V_0 \left( 1 + \frac{\mu^2 \phi^2}{2M_P^2} \right) + \dots \text{ for } \phi < M_P.$$

At tree level, for a generic scalar field one has naturally  $|\mu^2| \simeq 1$   $\eta$  problem !

 $\rightarrow V(\phi)$  is NOT flat at high scale

Assume that the inflaton field interacts **not** so weakly and add one loop corrections to the potential by substituting

 $\mu^2 \to \mu^2 (Q = \phi)$  running mass

The running of the mass can flatten the potential somewhere in the region  $\phi < M_P$ .

Slow roll inflation

#### Parameterize n(k) and $\mathcal{P}_{\mathcal{R}}(k)$ :

 $\begin{array}{l}
n(k) - 1 \ll 1 \text{ on} \\
\text{cosmological scales} \\
\text{So take the running mass as}
\end{array} \Rightarrow \begin{array}{linear expansion} \\
\text{around pivot } \phi_0 (\leftrightarrow k_0) \\
\end{array}$ 

$$m^2(\phi) \simeq m^2(\phi_0) + c * \log\left(\frac{\phi}{\phi_0}\right)$$

where  $c \propto \beta_m = \frac{dm^2}{d \log(Q)}(\phi_0) = coupling \times m_{loop}^2$ .

Then defining  $\phi_*$  by  $V'_{lin}(\phi_*) = 0$  and introducing the parameter  $s = c \log(\phi_*/\phi_0)$ , we have

$$\frac{n(k)-1}{2} = s\left(\frac{k}{k_0}\right)^c - c$$

and

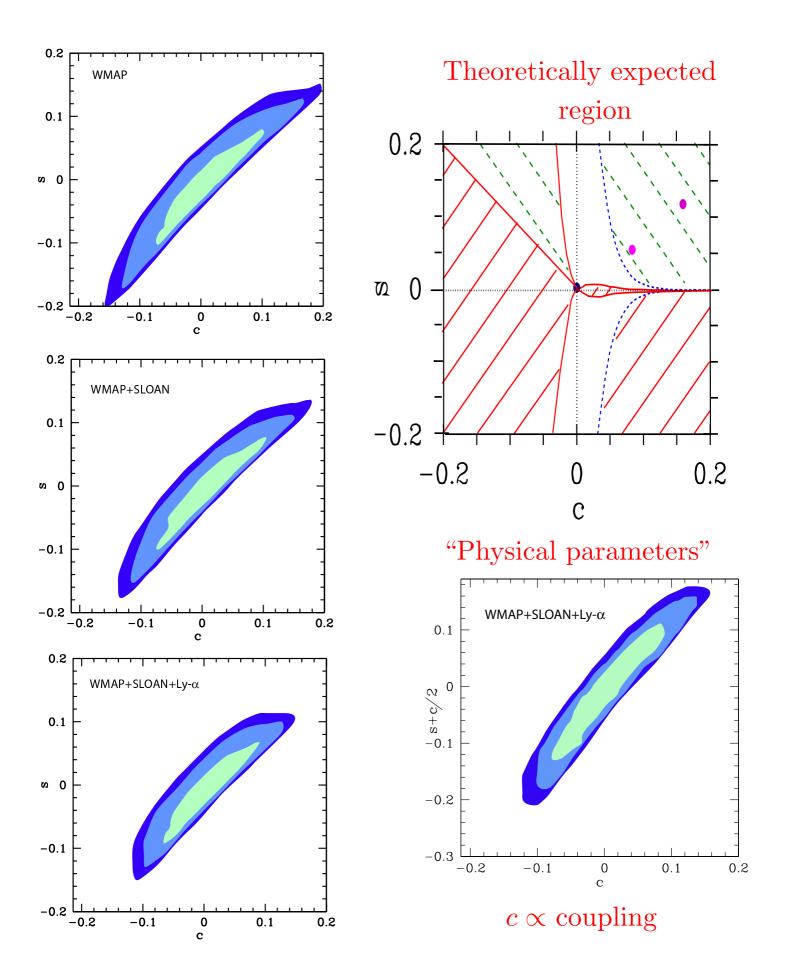
$$n'(k) = 2sc \left(\frac{k}{k_0}\right)^c$$

"Strong (exponential !)" scale dependence !! NOTE: s, c are related to physical parameters rescaled by the inflationary Hubble scale  $H_I^2$ :

$$c \equiv -\frac{\beta_m(\phi_0)}{3H_I^2} \qquad s + \frac{1}{2}c \equiv \frac{m^2(\phi_0)}{3H_I^2}$$

c suppressed by a coupling, s also to have slow roll...

What are the constraint from the new data for s, c in such models ? [LC, Lyth, Melchiorri & Odman astro-ph/0408129]



### String theory meets inflationary cosmology

Lots of activity on inflationary model building in string models, too many models to review here, but a couple of general results are:

• it is possible to have "unstable", but sufficiently long-lived de Sitter minima in string theory

 $\rightarrow$  KKLT

• the string inflationary potentials are pretty complex, they must include non-perturbative contributions (difficult to compute...) and in general a fine-tuning of the order  $\mathcal{O}(10^{-2})$  is necessary to have slow roll...

• in certain realizations with colliding branes cosmological "superstrings" could form at the end of inflation

 $\rightarrow$  similar to the old local strings, but with different scale and smaller interconnection probability P < 1 [Copeland et al, Dvali at al... '04 ]

# Beyond the classical CMSSM supersymmetric DM

The CMSSM neutralino scenario appears nowadays to be a bit fine-tuned... Most of the activities lately are concentrated on less constrained scenarios:

- More general SUSY breaking schemes, e.g.  $\rightarrow$  Y. Mambrini
- NMSSM, e.g. [Cerdeño *et al*, Menon *et al* '04]
- Split SUSY  $\rightarrow$  A. Romanino, A. Mazumdar
- $scalar/Kaluza-Klein DM \rightarrow P. Fayet, G. Servant$
- Super Weakly Interacting Massive Particles: gravitinos (see e.g. [W. Buchmüller *et al...*, J. Ellis *et al*, J. Feng *et al...*, Roszkowski & Ruiz de Austri, K. Hamaguchi *et al* '04]) or axinos (e.g. [LC *et al...*, Brandenburg & Steffen '04]) or ....

SuperWIMPS: such particles reach thermal equilibrium at very high temperatures and "freeze out" when relativistic with high number density  $\Omega < 1 \rightarrow$  very light masses  $\rightarrow$  Hot or Warm DM ! They can be Cold DM if  $T_{RH} < T_f$ . Their yield is given (at least) by two mechanisms:

– thermal scattering and decays in the plasma

$$\frac{dY_{\rm SW}}{dT} = \frac{-1}{HTs(T)} \left[ \sum_{ij} \langle \sigma(i+j \to {\rm SW} + ...)v_{rel} \rangle n_i n_j \right]$$
scatterings
$$\sum_i \langle \Gamma(i \to {\rm SW} + ...) \rangle n_i \right]$$
decays

strongly dependent on  $T_{RH}$  !

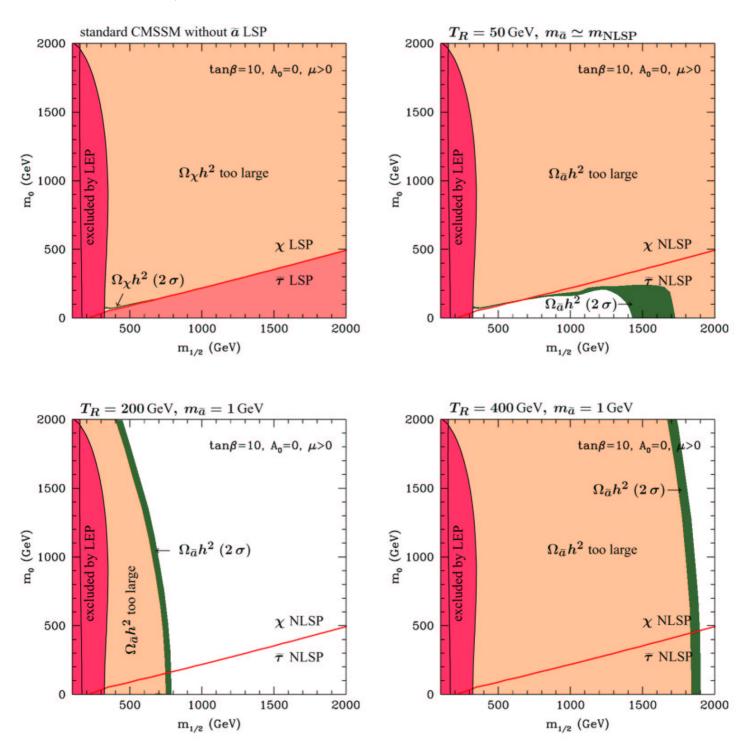
- decay out of equilibrium of the NLSP:

$$\Omega_{\rm sw}^{NT} = \frac{m_{\rm sw}}{m_{\rm NLSP}} \ \Omega_{\rm NLSP}$$

BEWARE of the other decay products ( $\gamma$ s or hadrons) not spoiling Nucleosynthesis or distort the CMB !

What are the consequences for the supersymmetric parameters from SWIMP CDM ?

More parameter space allowed, especially the  $\tilde{\tau}$ NLSP region. See e.g. for the axino (less constrained by BBN), [LC, Roszkowski, Ruiz de Austri & Small 04]



# Collider signature: long-lived $\tilde{\tau}$ s ???

Then it will be necessary to study the  $\tilde{\tau}$  decay to distinguish between the different scenarios: gravitino, axino, singlino, R-parity violation, other...

Note measuring the  $\tilde{\tau}$  lifetime and mass is not enough, the main decay channel and possibly a radiative one are probably needed !

⇒ need to block the  $\tilde{\tau}$  s and store them for a sufficient time..., but the lifetime could range from minutes to years ! [Feng *et al*, Hamaguchi *et al* 04] If the decay is seen, the main signals are:

• SuperWIMP: main  $\tilde{\tau}$  decay into SW+ $\tau$ , then SW+ $\tau + \gamma$ ...

Perhaps the angular distribution in the radiative decay could give information on the spin of the invisible particle and distinguish the gravitino if not mainly goldstino [Buchmüller *et al* 04]

• R-parity breaking: mainly 3-body decay to visible particles [Allanach *et al* 04]

• NMSSM "singlino": different topologies and shorter lifetimes due to non-negligible mixing in the neutralino sector. [Ellwanger & Hugonie 98, Martin 00]

# Conclusions and Outlook

The era of precision cosmology continues:

- the new LSS data from Lyman  $\alpha$  allow to put better constraints on the models of structure formation, e.g.  $\Omega_{\nu}h^2$ , and the scale dependence of the spectral index
- the simple single field inflationary paradigm with negligible tensor perturbation and running is sufficient to describe the data
- Dark matter bounds are becoming more stringent and starting to probe the SUSY parameter space
- still there are important open puzzles:
  - galaxy  $\gamma$  emissions ?
  - reionization ?
  - UHECR ?
  - **-**Λ?
  - ?

We are looking forward to the next year with more data from WMAP, SDSS, DM experiments, etc...