

بسم الله الرحمن الرحيم

# Effects of ion-channel guiding on saturation mechanism of a single-pass FEL

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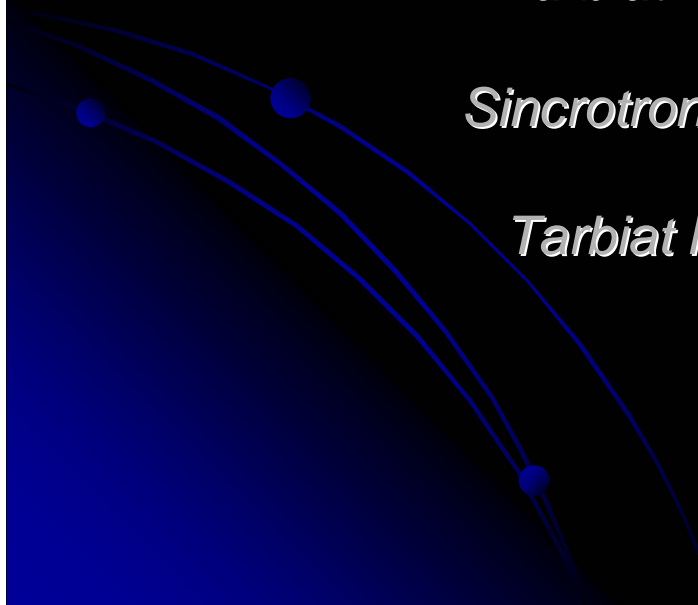
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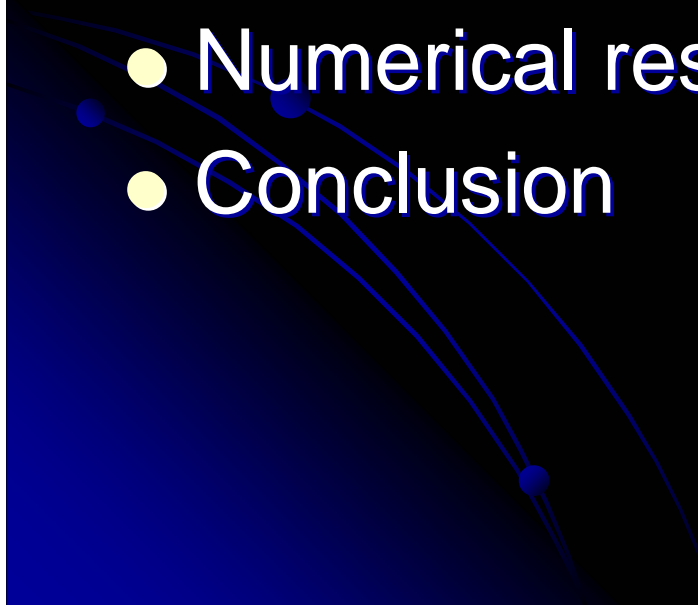
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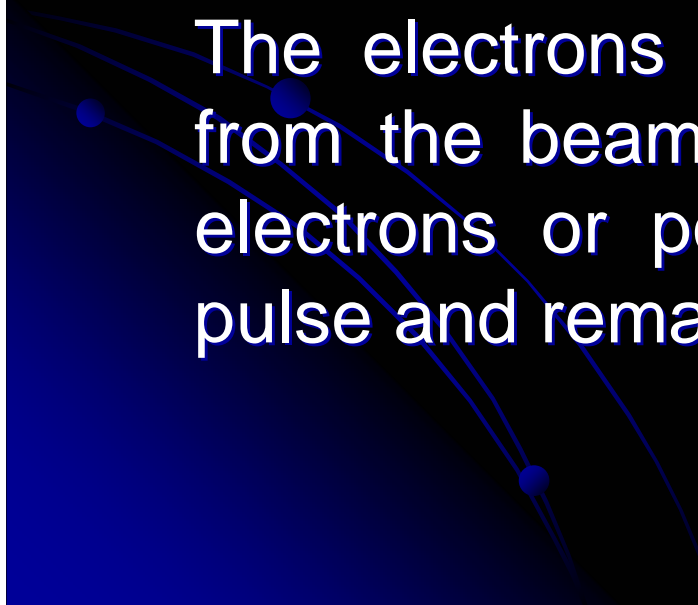
# Outline

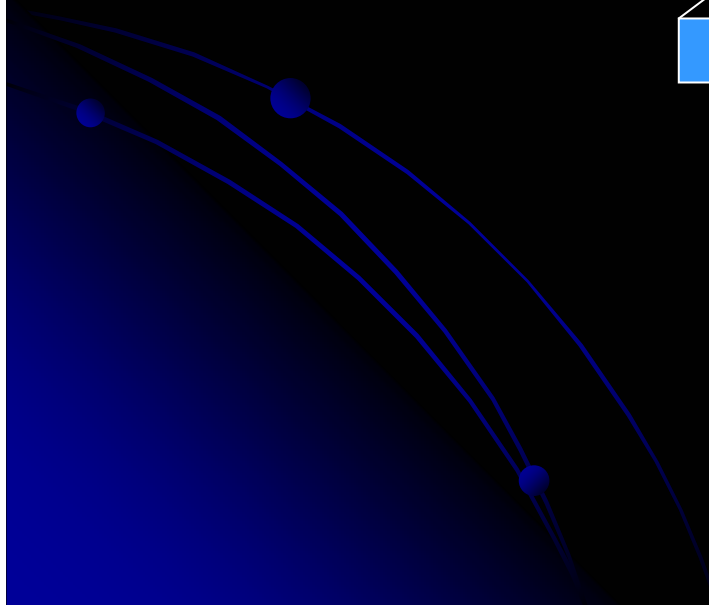
- Ion-channel guiding concept
  - Ion-channel benefits and detriments
  - The physical configuration
  - Theoretical model
  - Numerical results
  - Conclusion
- 

# Ion-channel guiding concept

An ion-channel can be produced inside the FEL structure by ionizing a neutral gas either by using the electron beam itself or a high power laser pulse.

The electrons of the ionized gas are expelled from the beam volume either by self-field of the electrons or ponderomotive force of the laser pulse and remains a channel of heavy ions.



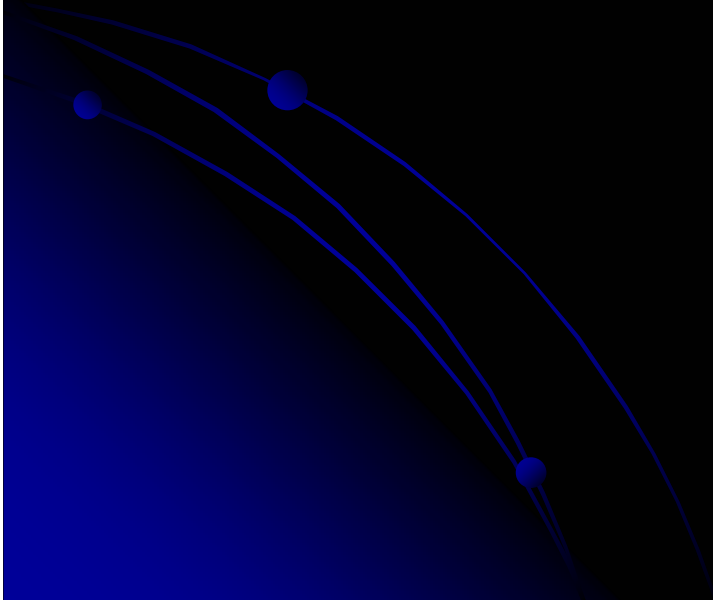


# Ion-channel guiding benefits

- Ions neutralize the space-charge of electrons and reduce the electron beam divergence.
- It permits beam currents higher than vacuum limit.
- It is a less expensive alternative to other guiding methods.
- It helps the radiation guiding.

# Detriments

- Beam head erosion
- Ion-hose instability
- Ion-channel collapse



# Physical configuration

- Helical Wiggler  $\vec{B}_w = B_w(\hat{e}_x \cos k_w z + \hat{e}_y \sin k_w z)$
- Ion-channel electrostatic field  $\vec{E}_i = 2\pi e n_i(\hat{e}_x x + \hat{e}_y y)$
- Radiation field  $\delta\vec{A}(z, t) = \delta\hat{A}(z)[\hat{e}_x \cos \alpha_+(z, t) - \hat{e}_y \sin \alpha_+(z, t)]$   
 $\delta\Phi(z, t) = \delta\hat{\Phi}(z) \cos \alpha(z, t)$
- Electron beam  $E = \gamma m_e c^2$

# Equations of motion in the wiggler frame

$$\frac{du_1}{dz} = u_2 + \frac{d\delta a}{dz} \cos \psi - \omega_p^2 \frac{x_1}{\beta_3}$$

$$\frac{du_2}{dz} = -u_1 - \frac{d\delta a}{dz} \sin \psi - \omega_p^2 \frac{x_2}{\beta_3} - a_w$$

$$\frac{du_3}{dz} = a_w \frac{u_2}{u_3} + k_+ \delta a \frac{u_1 \sin \psi + u_2 \cos \psi}{u_3} - \frac{1}{\beta_3} \left( k \delta \varphi \sin \psi_{sc} - \frac{d}{dz} \delta \varphi \cos \psi_{sc} \right)$$

$$\vec{u} \equiv \vec{p} / m_e c, \quad \omega_p^2 = 2\pi n_i e^2 / m_e k_w^2 c^2, \quad \omega_b^2 = 4\pi n_b e^2 / m_e k_w^2 c^2$$

$$a_w = e B_w / m_e k_w c^2, \quad \delta a = e \delta \hat{A}(z) / m_e c^2, \quad \delta \varphi = e \delta \hat{\varphi}(z) / m_e c$$

$$k \equiv k / k_w, \quad \omega \equiv \omega / k_w c, \quad \beta \equiv v / c$$

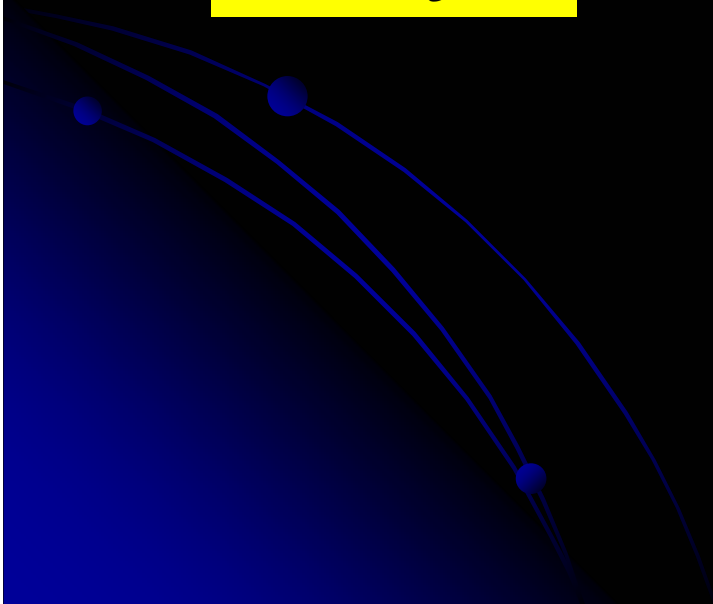


$$\frac{d\psi(z)}{dz} = k_+(z) + 1 - \frac{\omega}{\beta_3}$$

$$\frac{d\psi_{sc}(z)}{dz} = k(z) - \frac{\omega}{\beta_3}$$

$$\frac{dx_1}{dz} = \frac{u_1}{u_3} + x_2$$

$$\frac{dx_2}{dz} = \frac{u_2}{u_3} - x_1$$



# Field equations

- Maxwell equations in Coulomb gauge

$$\left( \frac{\partial^2}{\partial z^2} - \frac{\partial^2}{\partial t^2} \right) \delta \vec{a}(z, t) = -4\pi \delta \vec{J}_\perp(z, t)$$

$$\frac{\partial^2}{\partial z \partial t} \varphi(z, t) = 4\pi \delta J_z(z, t)$$

$$\delta \vec{J}(z, t) = -\frac{1}{4} \omega_b^2 \int_{-\infty}^{\infty} dt_0 \vec{\beta}(t, t_0) \frac{\delta[t - \tau(z, t_0)]}{|\beta_z(t, t_0)|}$$

$$\frac{d}{dz} \delta a = \Gamma_+ \delta a$$

$$\frac{d}{dz} \Gamma_+ = k_+^2 - \omega^2 - \Gamma_+^2 + \frac{\omega_b^2 \beta_{z0}}{\delta a} \left\langle \frac{u_1 \cos \psi - u_2 \sin \psi}{|u_3|} \right\rangle$$

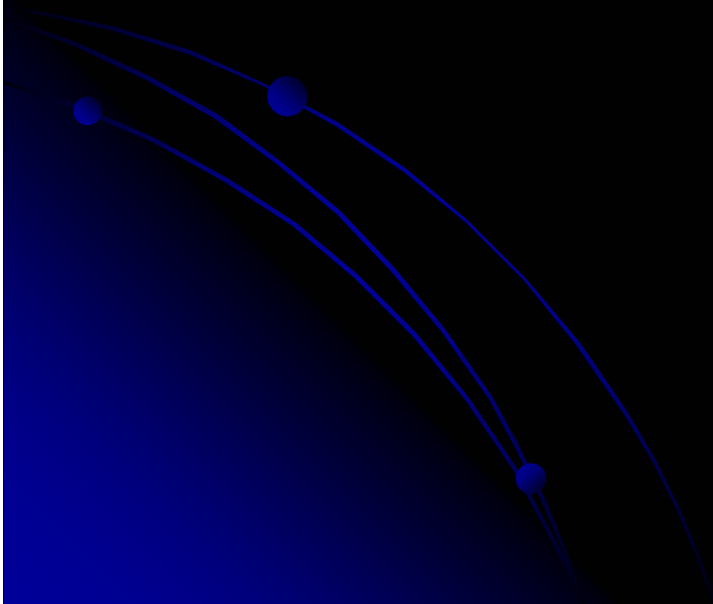
$$\frac{d}{dz} k_+ = -2k_+ \Gamma_+ - \frac{\omega_b^2 \beta_{z0}}{\delta a} \left\langle \frac{u_1 \sin \psi + u_2 \cos \psi}{|u_3|} \right\rangle$$

$$\frac{d}{dz} \delta \phi = -2 \frac{\omega_b^2 \beta_{z0}}{\omega} \langle \sin \psi_{sc} \rangle$$

$$k = -2 \frac{\omega_b^2 \beta_{z0}}{\delta \phi \omega} \langle \cos \psi_{sc} \rangle$$

$$\langle(\cdots)\rangle \equiv \frac{1}{2\pi} \int_{-\pi}^{\pi} d\psi_0 (\cdots)$$

$$\psi_0 = -\omega t_0$$



# Numerical analysis

The first example is corresponding to a 1.2GeV electron beam,  $a_w=4$  and without ion-channel

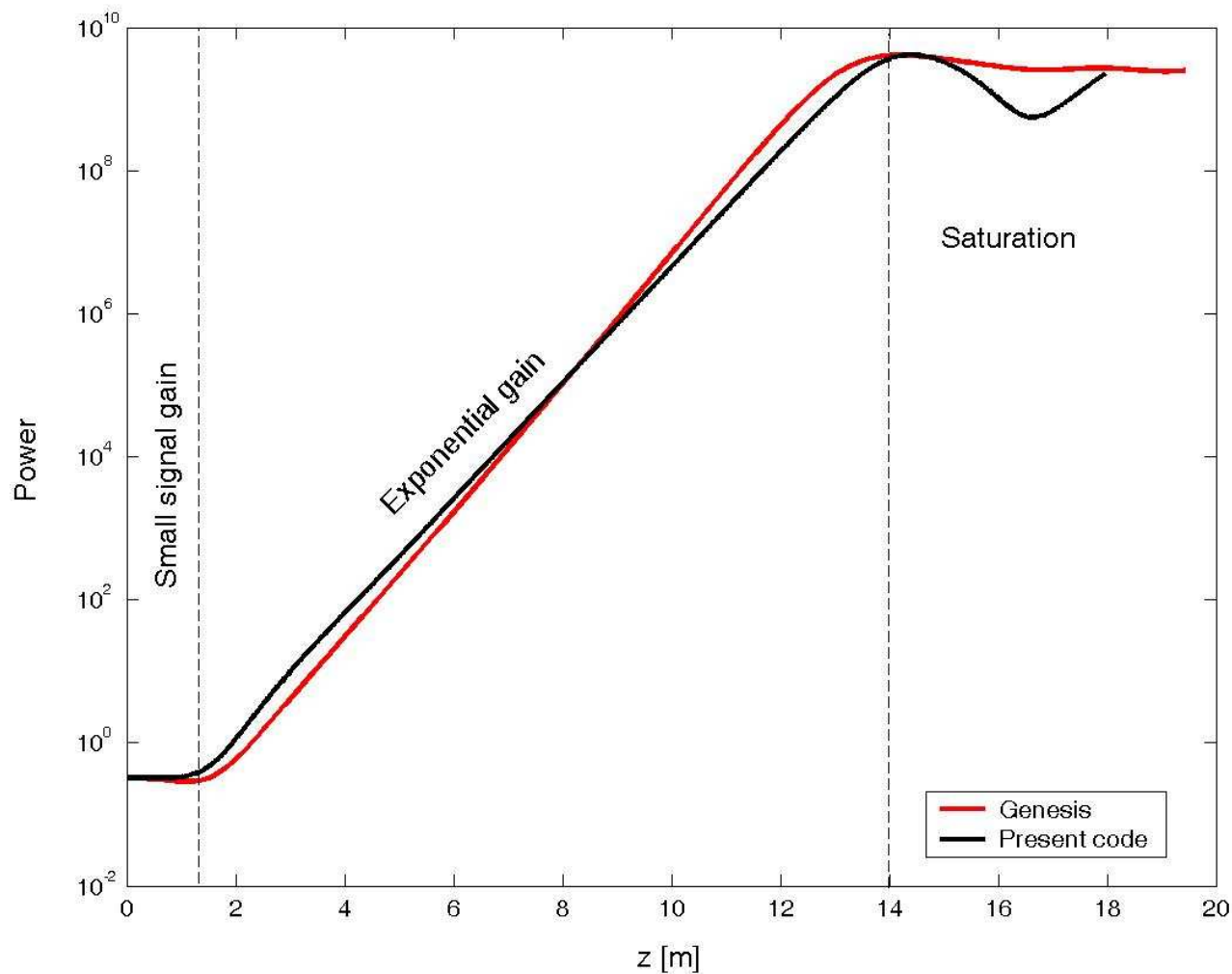
(Fermi@Elettra at 100 nm)

$$\gamma_0 = 1250, \quad \omega_b = 47.5$$

$$K = 4, \quad \omega_p = 0$$

There is  $4N+7$  first order differential equations which must to be solve self consistently

# The output power versus axial position without ion-channel



Ion-channel affects the ordinary FEL parameter. An effective FEL parameter could be defined as:

$$K = \left| \frac{\gamma_0 \beta_{0z}^2 a_w}{\gamma_0 \beta_{0z}^2 - \omega_p^2} \right|$$

The output wavelength has its ordinary formula

$$\lambda = \frac{\lambda_w}{2\gamma_0^2} [1 + K^2]$$

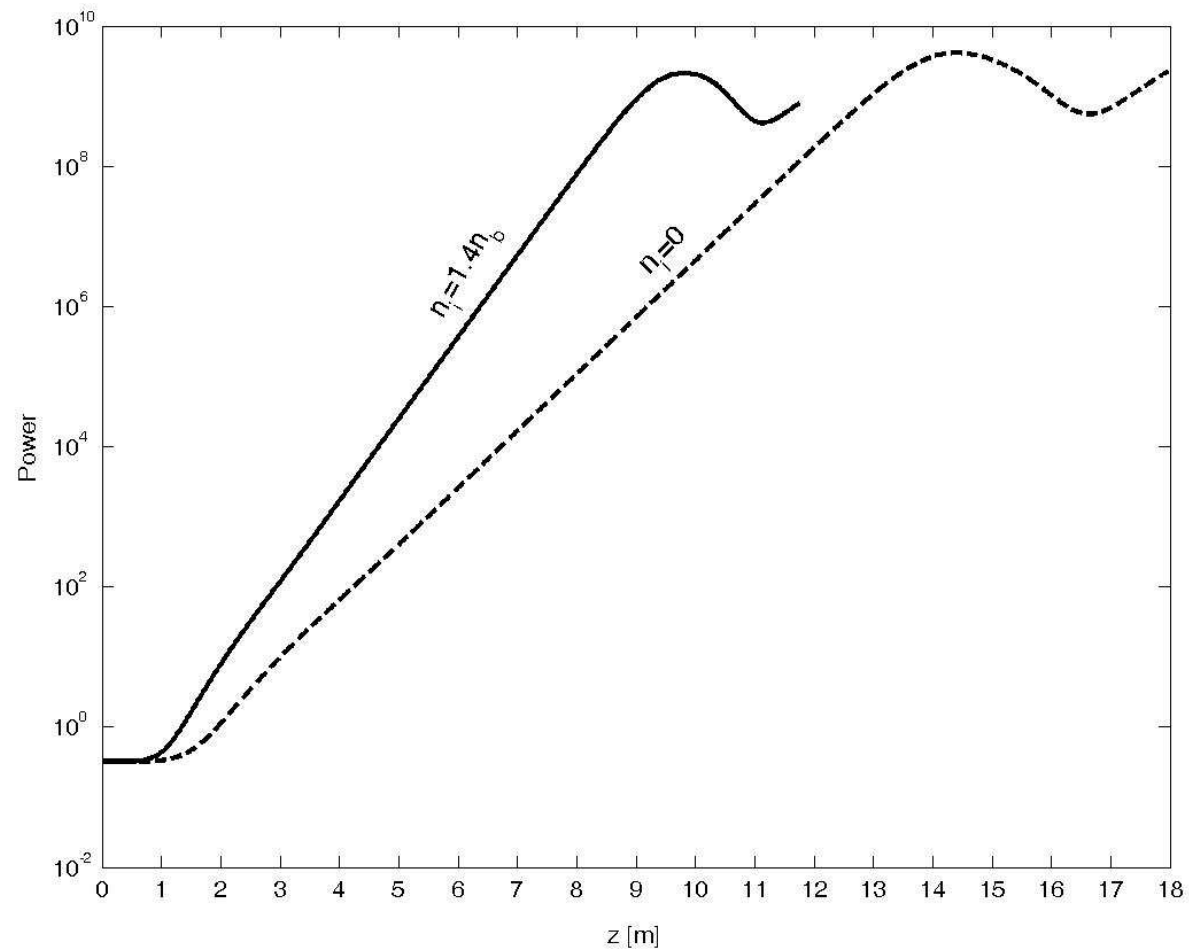
Consequently, ion-channel acts as a tuning parameter for the standard FELs.

It can be used in two different schemes:

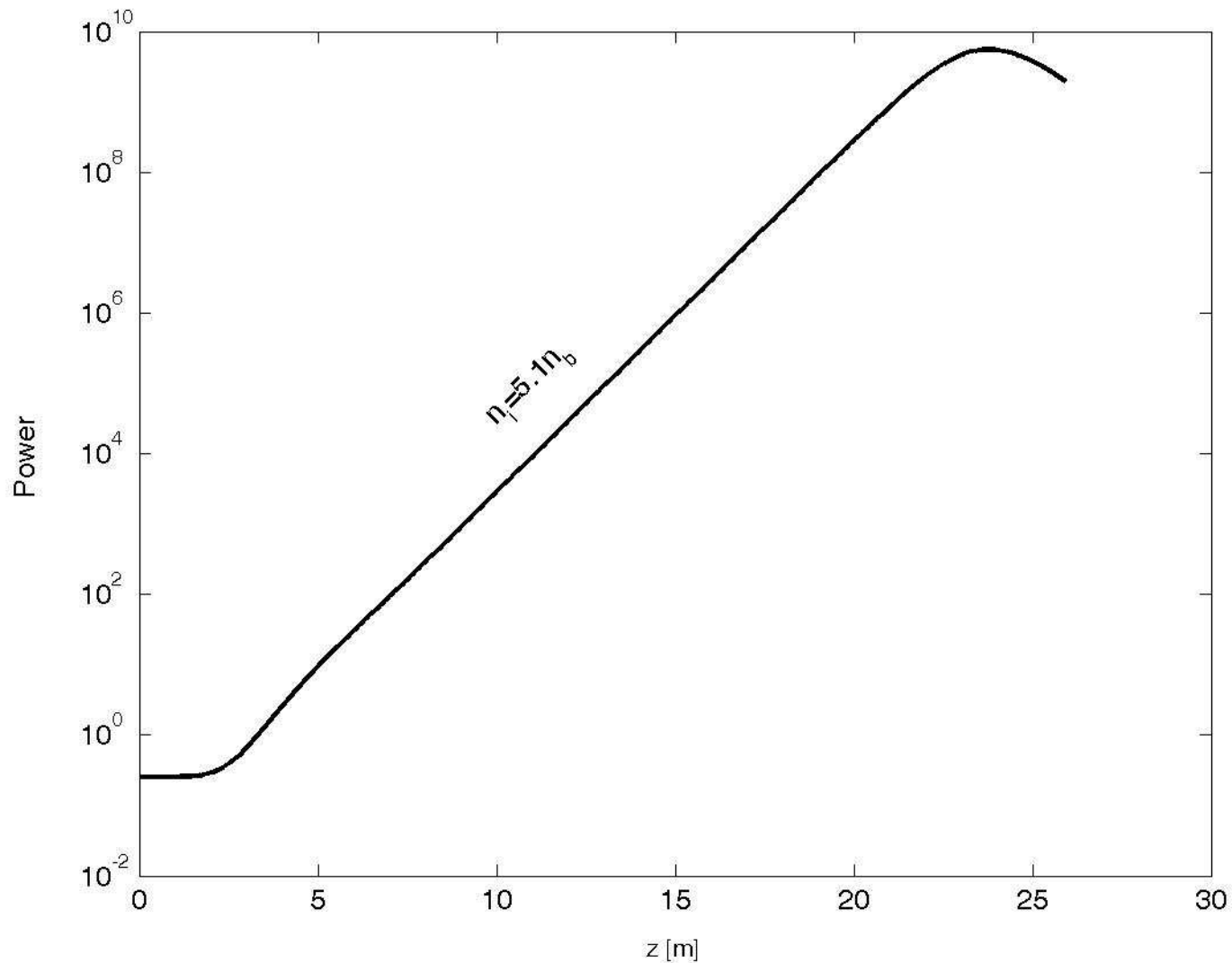
1. Changing the output wavelength for a fixed FEL structure
2. Getting to the same output wavelength with different structures



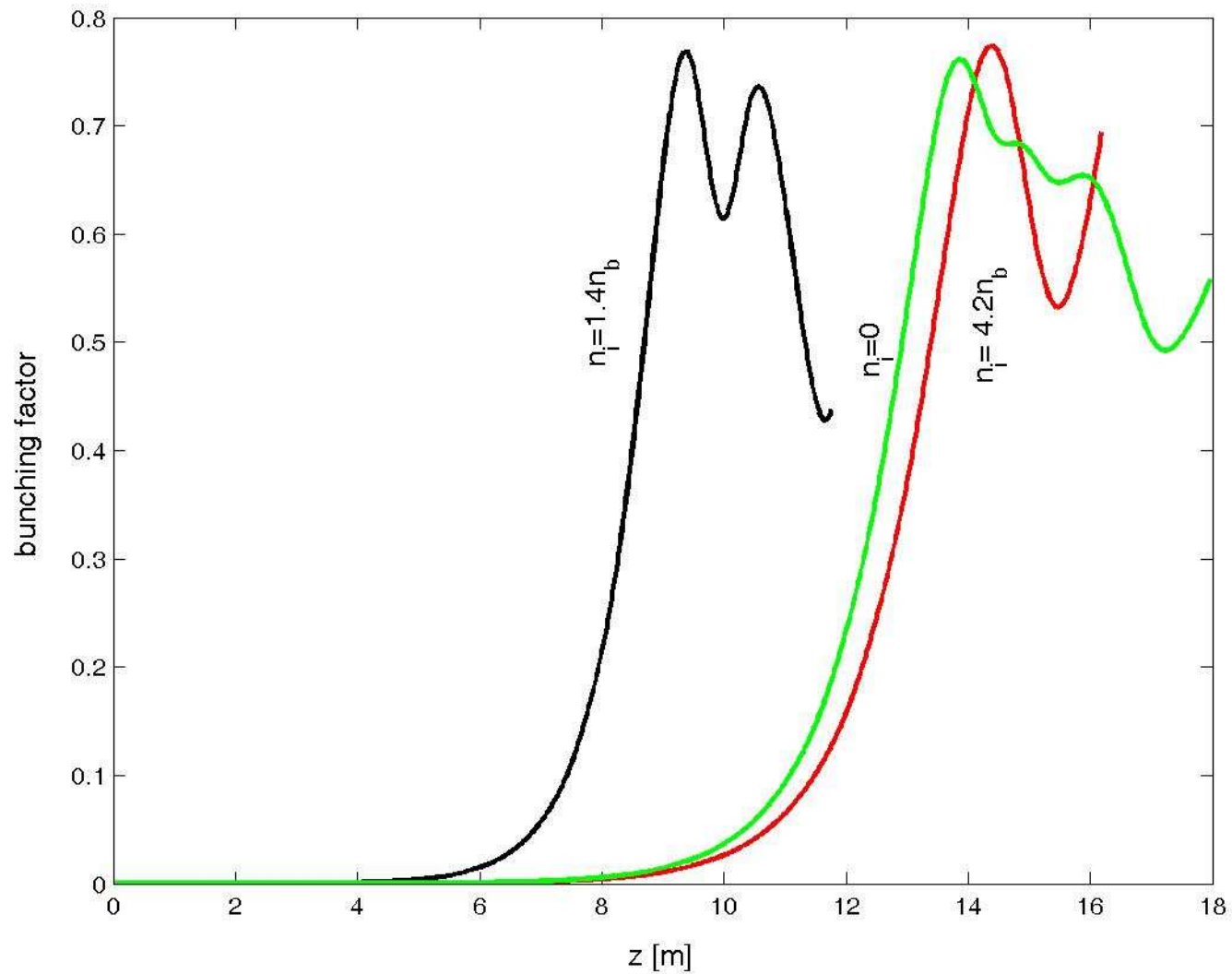
Output power for a FEL with  $a_w=1.27$ , ion-channel density  $n_i=1.4n_0$  and  $\lambda=100\text{nm}$  ( $K=4$ )



Output power for a FEL with  $a_w=4$  ion-channel density  
 $n_i=5.1n_b$  and  $\lambda=50\text{nm}$  ( $K=2.74$ )



# Bunching parameter for different ion-channel densities



# Conclusion

*In addition to this evident fact that ion-channel guiding allows higher electron currents and less divergence;*

- The ion-channel affects the ordinary FEL parameter in a way which can be used as a tuning parameter to control the output characteristics of the FEL.
- It is possible to reduce the saturation length for a fixed wavelength by reducing the ordinary FEL parameter and compensating it by an appropriate value of ion-channel.

Thank you for your attention

