

# The Quantum Free Electron Laser (QFEL)

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# **OUTLINE**

**Classical SASE and spiking**

**Quantum theory and QFEL parameter**

**Linear quantum theory**

**Quantum purification and narrow linewidth**

**Why a laser wiggler? Emittance criteria,  
possible experimental set up and parameters**

**Conclusions**

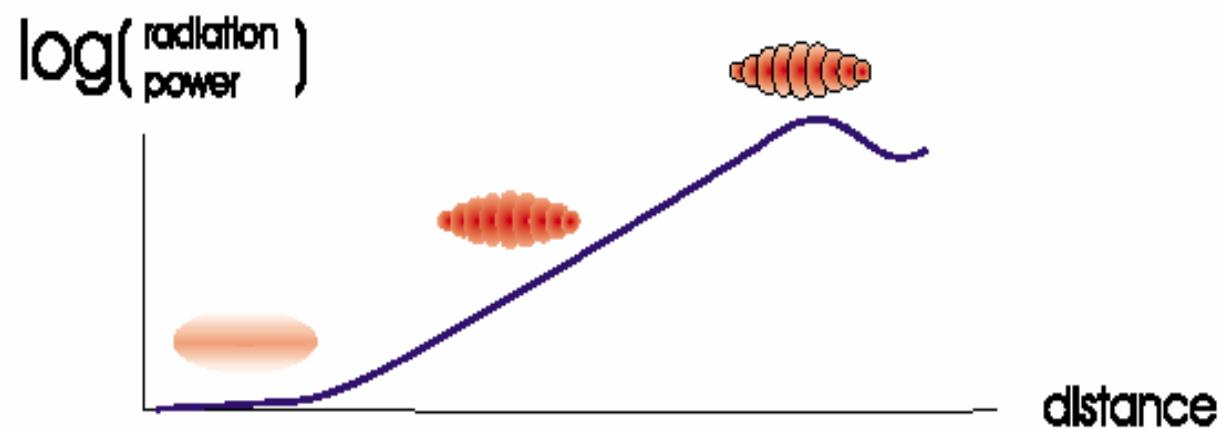
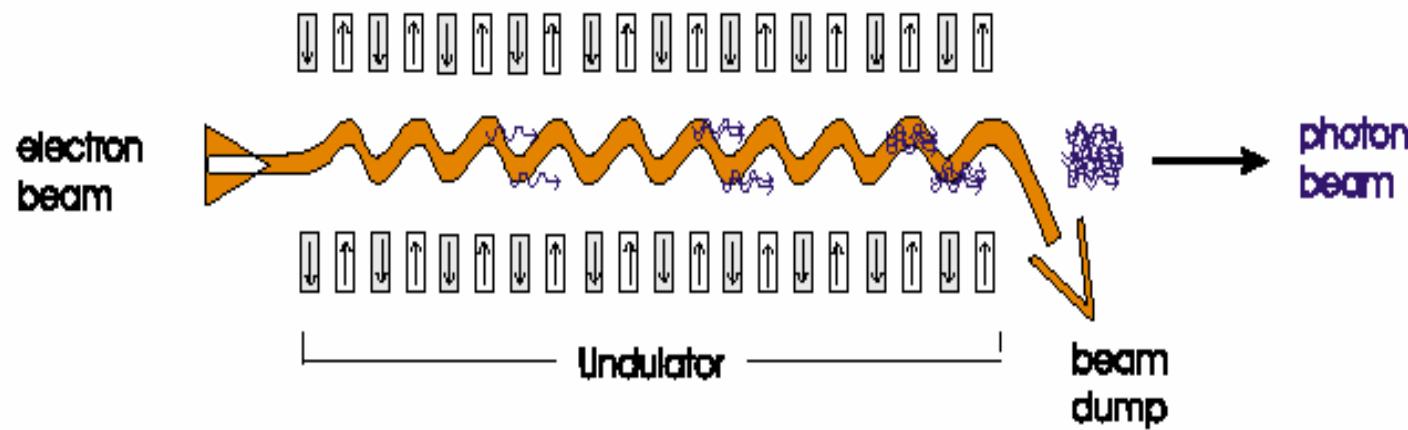
# what is QFEL?

QFEL is a novel macroscopic quantum coherent effect:

collective Compton backscattering of a high-power laser wiggler by a low-energy electron beam.

The QFEL linewidth can be four orders of magnitude smaller than that of the classical SASE FEL

# SASE high-gain regime



# SELF-BUNCHING

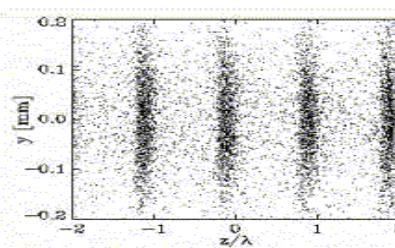
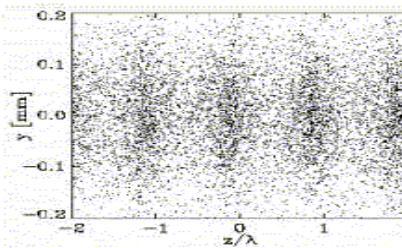
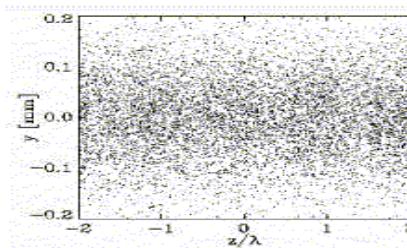
- start up from noise
- exponential growth of intensity and bunching
- saturation ( $P_{\text{rad}} \sim P_{\text{beam}}$ ) after several  $L_g$

$b \sim 0$



$b \sim 0.8$

bunching:



$$b = \left| \frac{1}{N} \sum_{j=1}^N e^{-i\theta_j} \right|$$

→

wiggler length (several  $L_g$ )

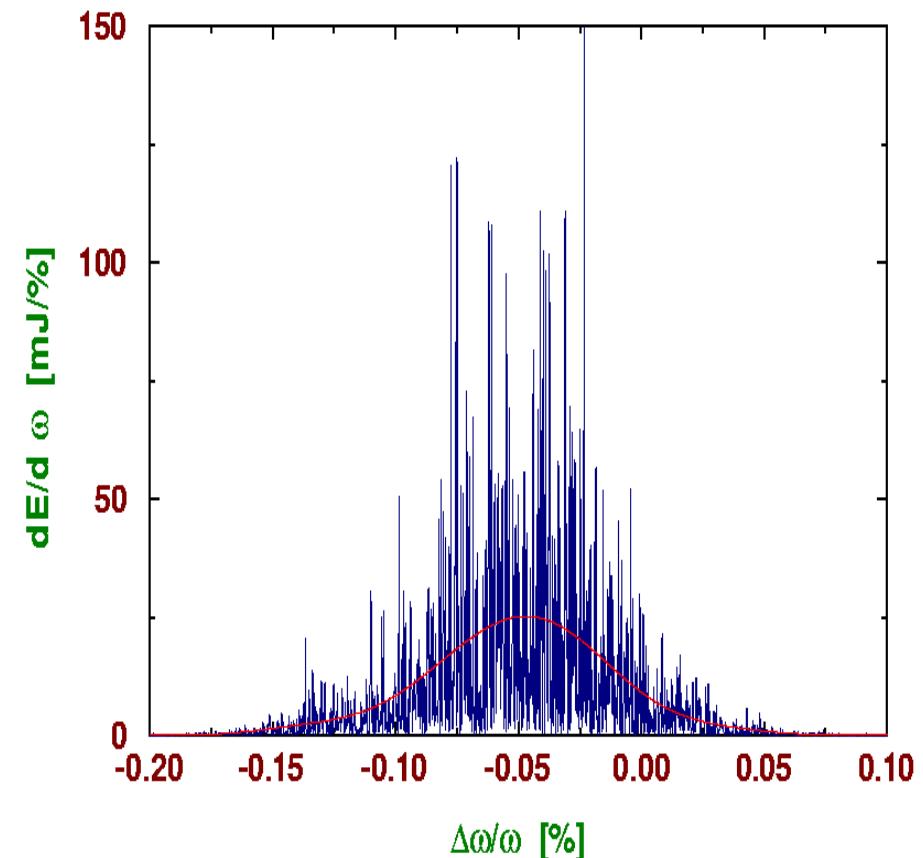
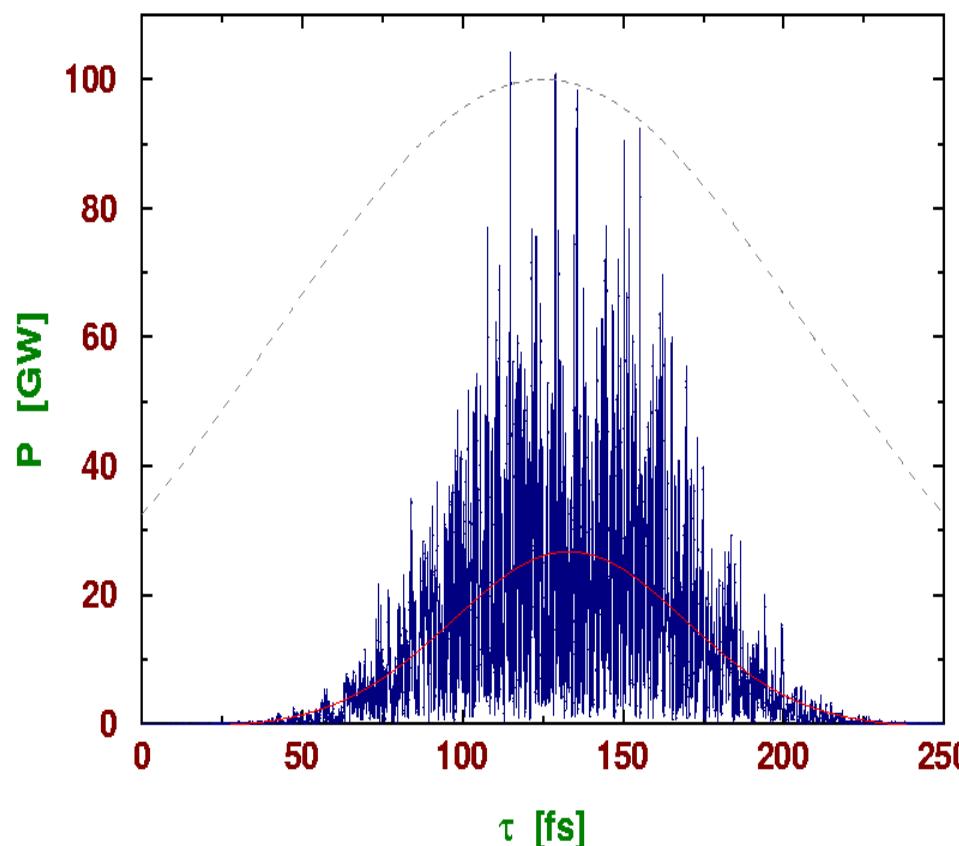
# DRAWBACKS OF 'CLASSICAL' SASE

R.Bonifacio, L. De Salvo, P.Pierini, N.Piovella, C. Pellegrini, PRL (1994)

Time profile has many random spikes ( $n = L_b/L_c$ )

Broad and noisy spectrum at short wavelengths (x-ray FELs)

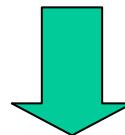
simulations from DESY for the SASE experiment



# QUANTUM FEL MODEL

## Procedure :

Describe N particle system as a **Quantum Mechanical ensemble**



Write a **Schrödinger equation** for macroscopic wavefunction:  $\Psi$

or equivalently ..

the equation for the **Wigner function** (quantum distribution):  $W$

Include **slippage**  $z_1$   
(i.e. propagation)



$\Psi(\theta, \bar{z}, z_1)$   
or  
 $W(\theta, p, \bar{z}, z_1)$

# 1D QUANTUM FEL MODEL

$$i \frac{\partial \Psi}{\partial \bar{z}} = -\frac{1}{2 \bar{\rho}^{3/2}} \frac{\partial^2 \Psi}{\partial \theta^2} - i \left\{ A(z_1, \bar{z}) e^{i\theta} - c.c. \right\} \Psi$$

$$\frac{\partial A}{\partial \bar{z}} + \frac{\partial A}{\partial z_1} = \int_0^{2\pi} d\theta | \Psi(\theta, z_1, \bar{z}) |^2 e^{-i\theta}$$

$$\begin{aligned}\bar{z} &= \frac{z}{L_g} \\ z_1 &= \frac{z - v_z t}{L_c} \\ L_c &\approx \left( \frac{\lambda}{\lambda_L} \right) L_g\end{aligned}$$

$A$  : normalized FEL amplitude

$\bar{\rho} = \rho_{FEL} \left( \frac{mc\gamma}{\hbar k} \right)$  : **QUANTUM FEL** parameter

the **classical** model is valid when  $\bar{\rho} \gg 1$

$\frac{\partial A}{\partial z_1} = 0$   **G. Preparata** from QFT, PRA (1988)

# LINEAR THEORY

Classical theory

$$(A \propto e^{i\lambda\bar{z}})$$

1)  $(\lambda - \Delta)\lambda^2 + 1 = 0$

R.Bonifacio, C.Pellegrini, L.Narducci, Opt. Commun. (1985)

Quantum theory

Previous approach

2)  $(\lambda - \Delta')\lambda^2 + 1 = 0 \quad \Delta' = \Delta - \frac{1}{2\bar{\rho}} \quad \Delta_{\max} = \frac{1}{2\bar{\rho}}$

C.B.Schroeder, C.Pellegrini, P.Chen,  
PRE (2001)

Mistake of equation (2): incorrect quantization  $pe^{-i\theta} \rightarrow \hat{p}e^{-i\hat{\theta}}$

Correct quantization

$$[e^{-i\hat{\theta}}, \hat{p}] = e^{-i\hat{\theta}}$$

$$pe^{-i\theta} \rightarrow \frac{1}{2}(\hat{p}e^{-i\hat{\theta}} + e^{-i\hat{\theta}}\hat{p})$$

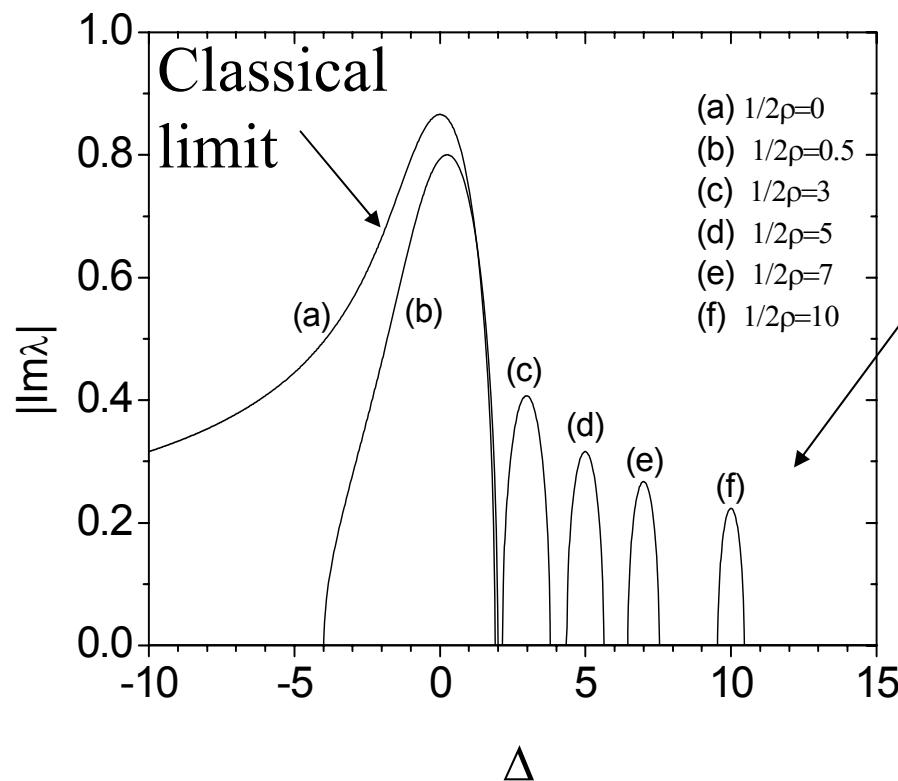
**Weyl symmetrization rule**

3)  $(\lambda - \Delta)\left(\lambda^2 - \frac{1}{4\bar{\rho}^2}\right) + 1 = 0$

R.Bonifacio, N.Piovella, G.Robb,  
A. Schiavi, PRST-AB (2006)

# Quantum Linear Theory $(A \propto e^{i\lambda\bar{z}})$

$$(\lambda - \Delta) \left( \lambda^2 - \frac{1}{4\bar{\rho}^2} \right) + 1 = 0$$



Quantum regime for  $\bar{\rho} < 1$

Resonance:

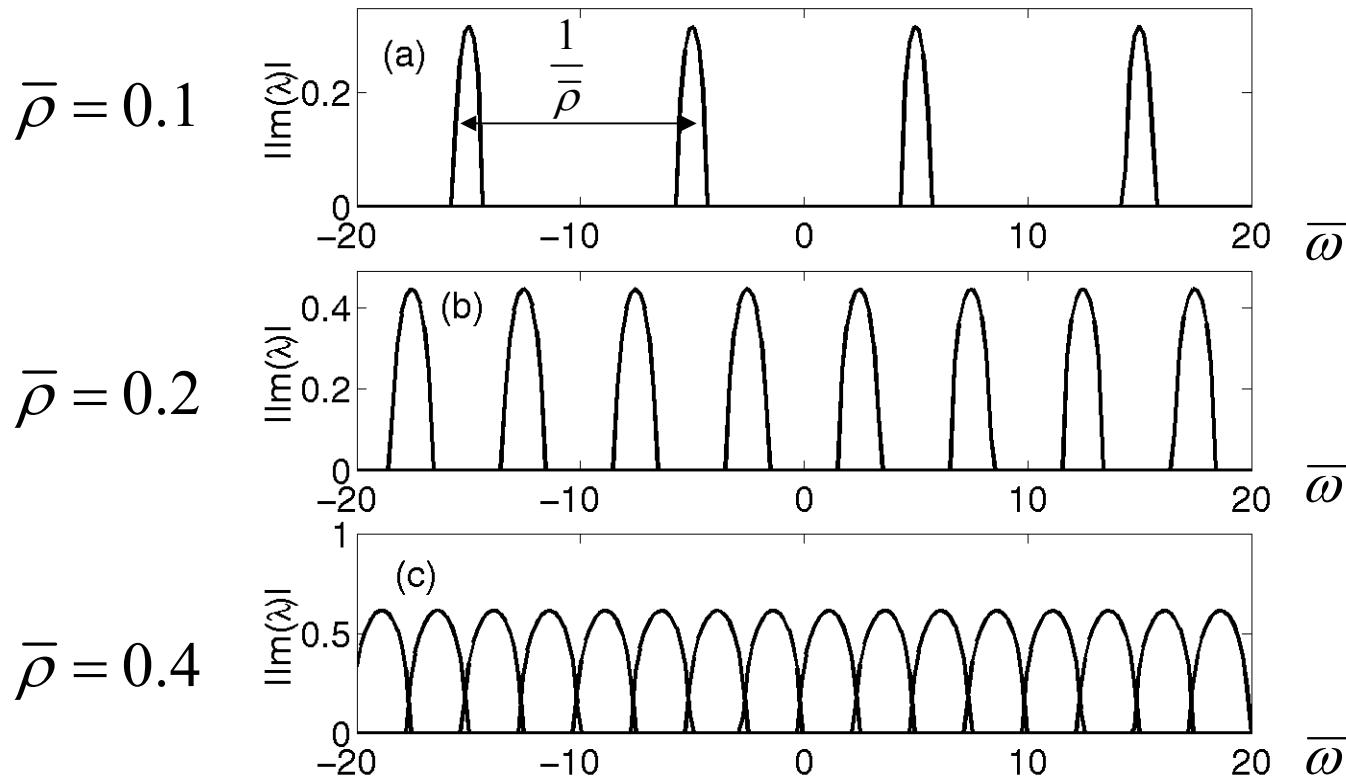
$$\Delta = \frac{1}{2\bar{\rho}}$$

width  $\propto \sqrt{\bar{\rho}}$

discrete frequencies as in a cavity

$$(\lambda - \Delta) \left( \lambda^2 - \frac{1}{4\bar{\rho}^2} \right) + 1 = 0 \quad \left( \Delta = \frac{n}{2\bar{\rho}} - \bar{\omega} \right)$$

$$\left( \bar{\omega} = \frac{\omega - \omega_{sp}}{2\rho\omega_{sp}} \right)$$



max for  $\Delta = 1/2\bar{\rho}$

$$\bar{\omega}_n = \frac{1}{2\bar{\rho}}(2n-1)$$

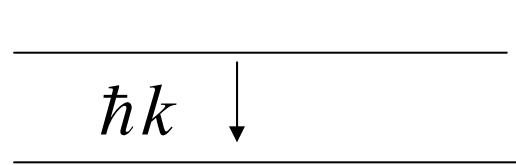
width  $4\sqrt{\bar{\rho}}$

Continuous limit

$$4\sqrt{\bar{\rho}} \geq 1/\bar{\rho} \rightarrow \bar{\rho} \geq 0.4$$

# The physics of Quantum FEL

$$\bar{\rho} = \rho \frac{mc\gamma}{\hbar k} = \frac{\sigma(p_z)}{\hbar k}$$

Momentum-energy levels: 

$$\omega_n = E_n - E_{n-1} \propto [n^2 - (n-1)^2] \propto 2n-1 \quad (n = 0, -1, \dots)$$

**Equally spaced frequencies as in a cavity**

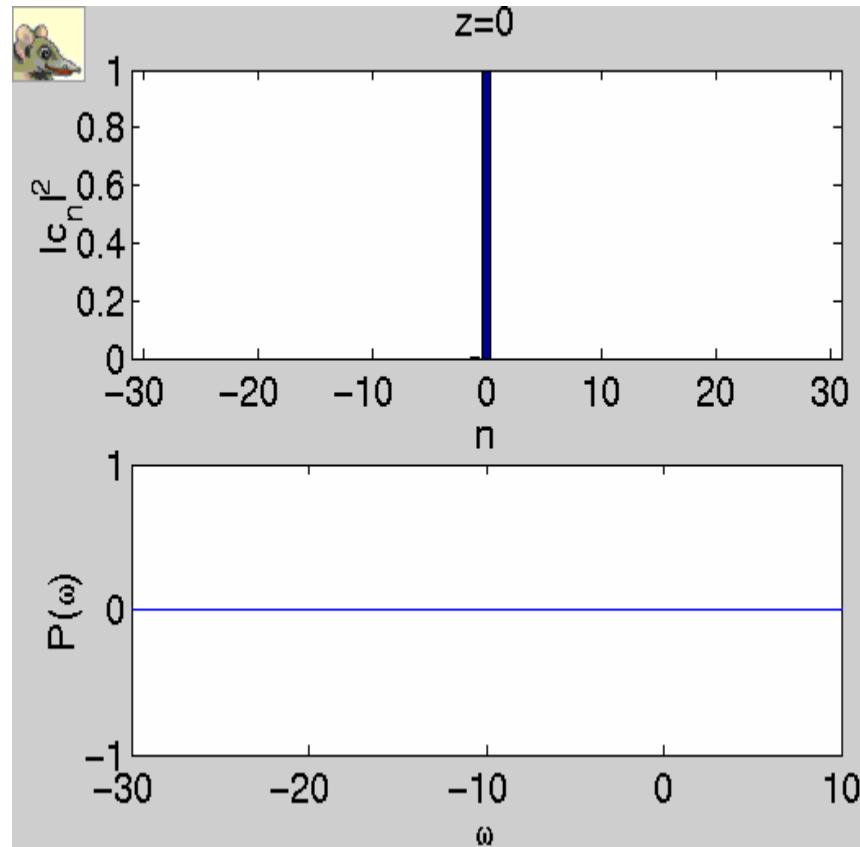
**CLASSICAL REGIME:**  $\bar{\rho} \gg 1$   
many momentum level  
transitions  
 $\Rightarrow$  **many spikes**

**QUANTUM REGIME:**  $\bar{\rho} \leq 1$   
a single momentum level  
transition  
 $\Rightarrow$  **single spike**

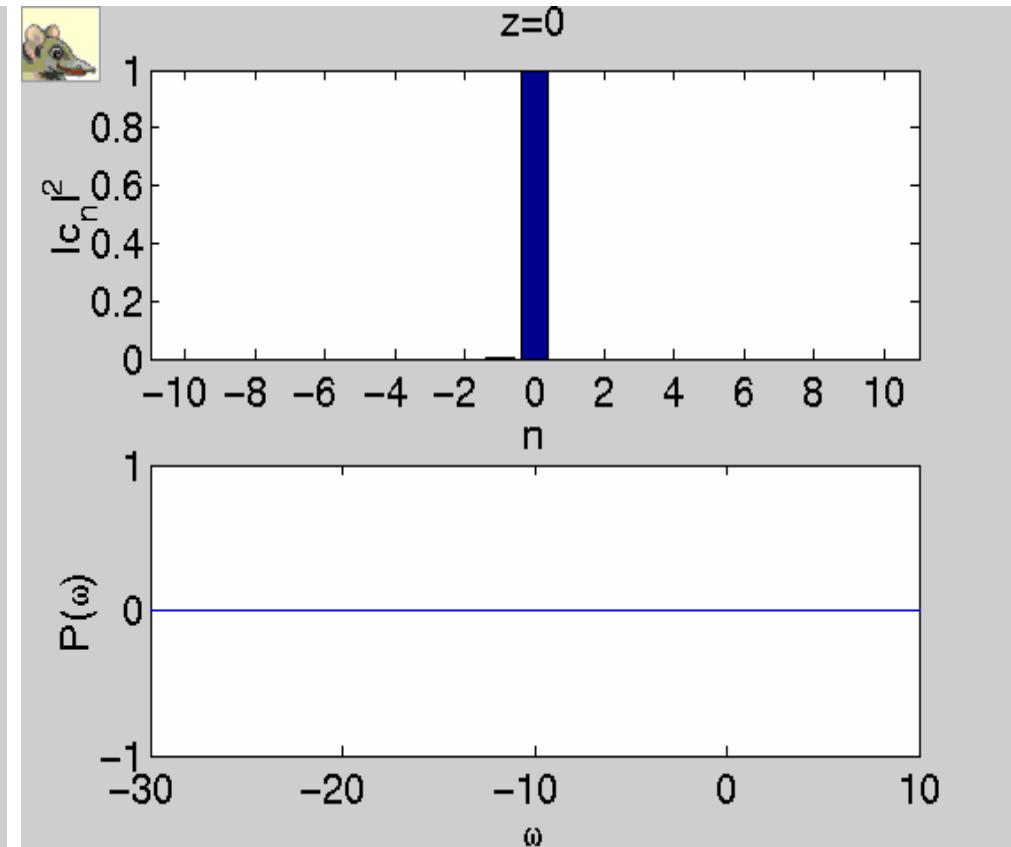
In classical regime with universal scaling no dependence on  $\bar{\rho}$

# momentum distribution for SASE

CLASSICAL REGIME:  $\bar{\rho}=5$



QUANTUM REGIME:  $\bar{\rho}=0.1$



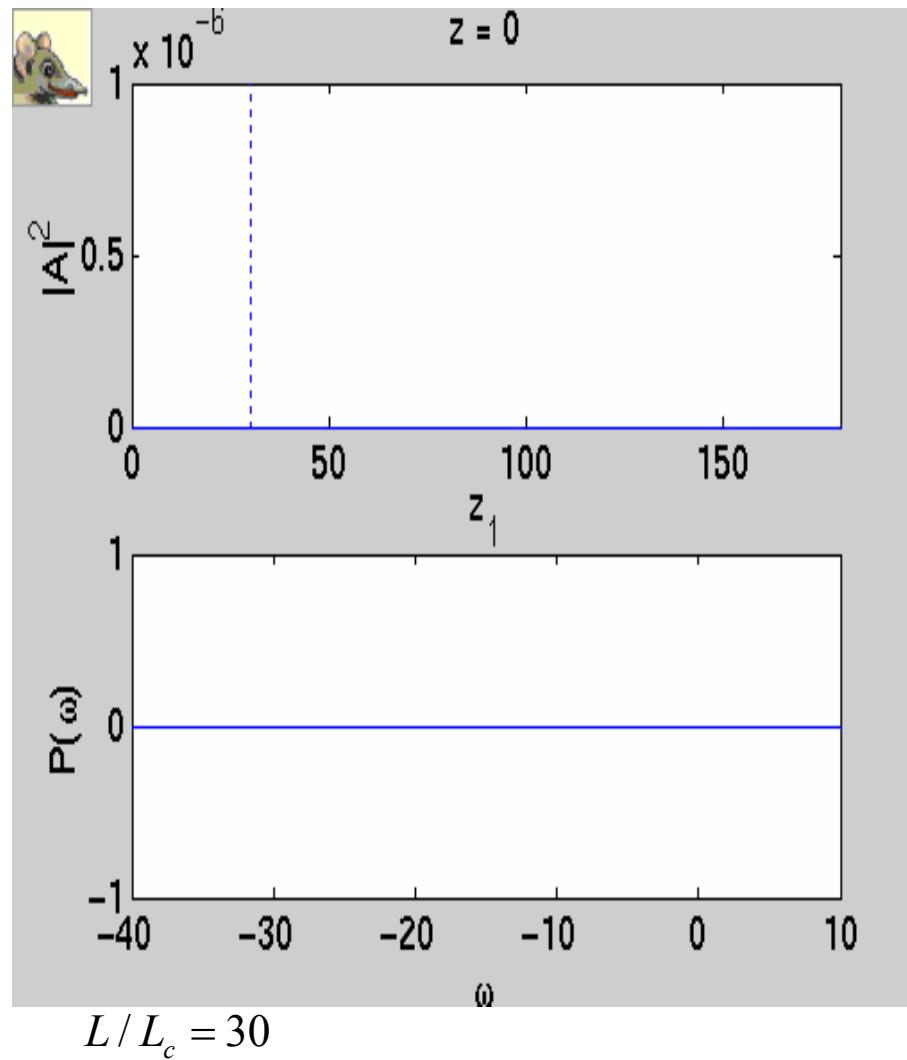
Classical regime:  
both  $n < 0$  and  $n > 0$  occupied

Quantum regime:  
sequential SR decay, only  $n < 0$

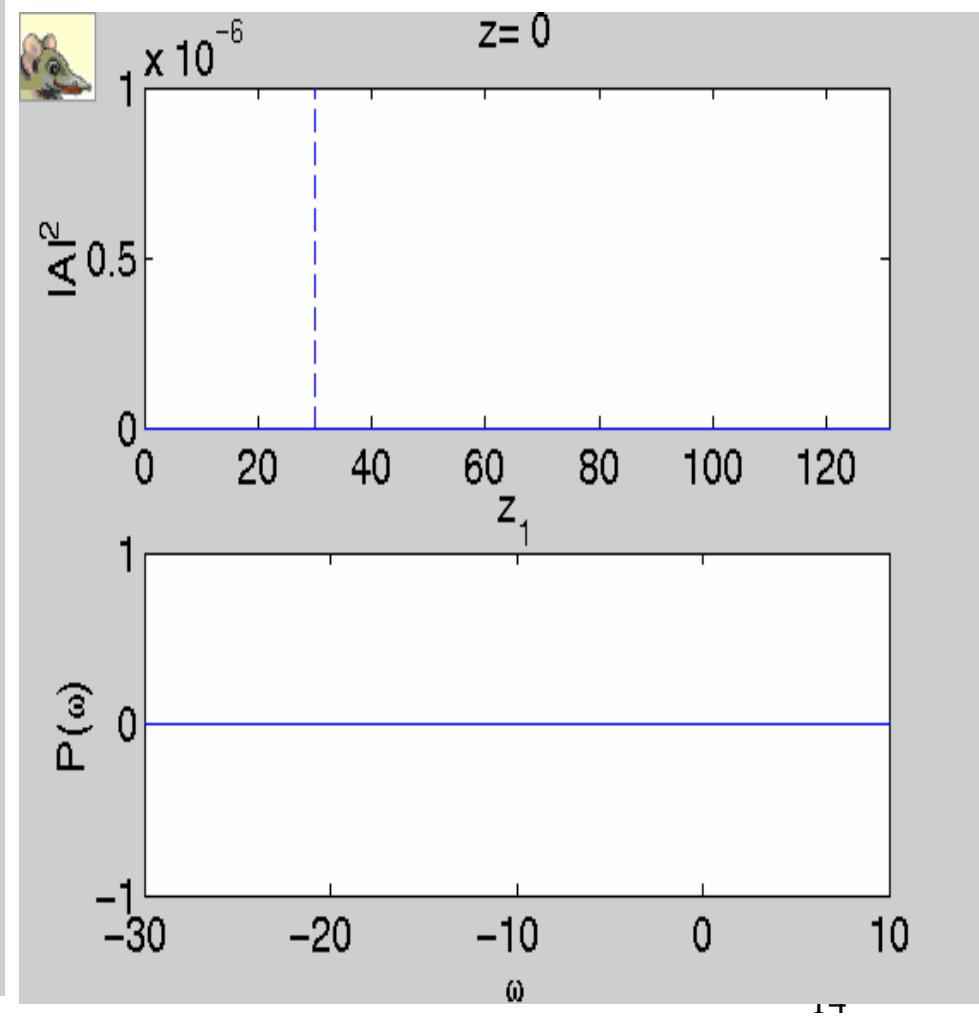
# SASE mode operation

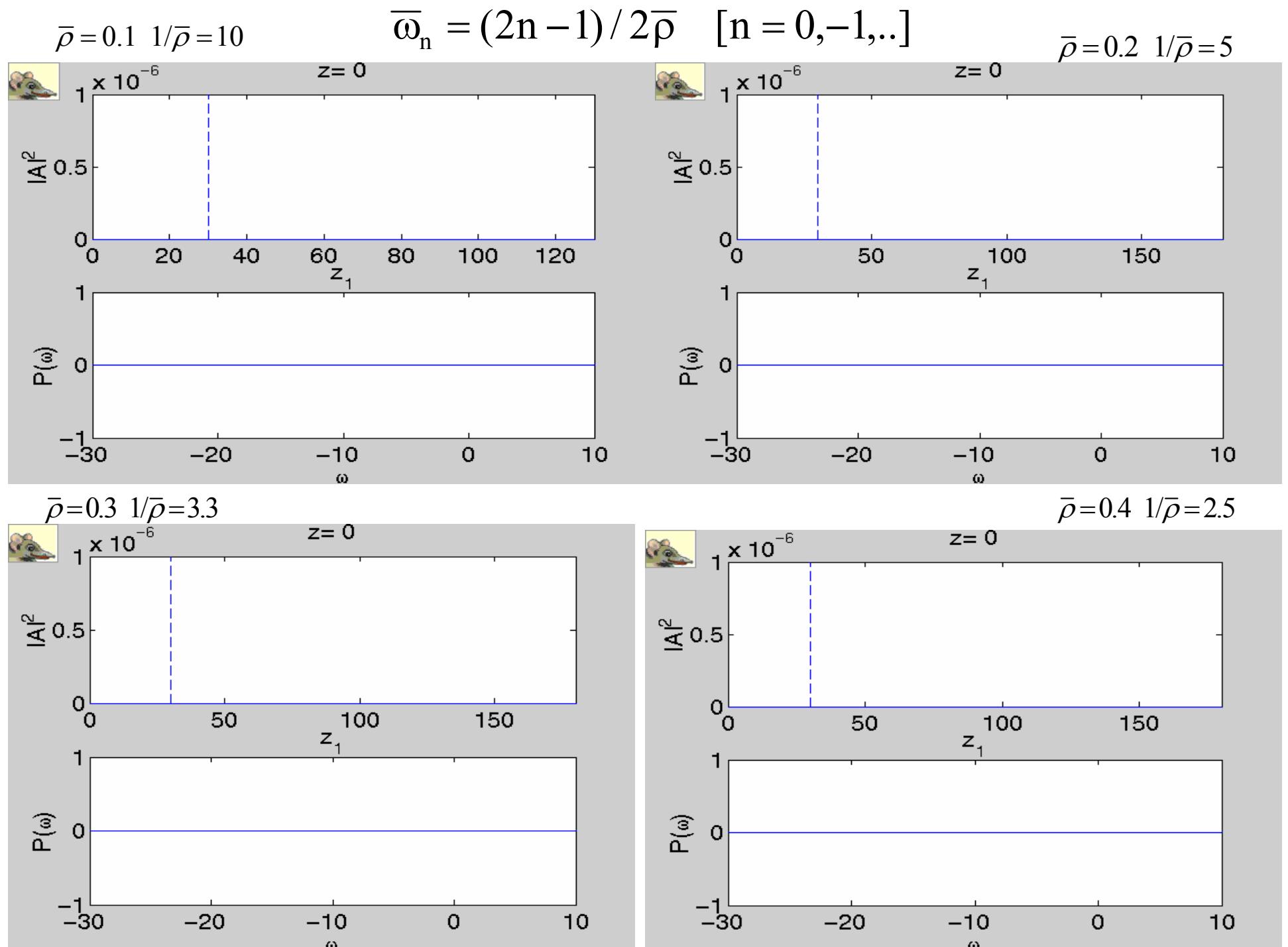
R.Bonifacio, N.Piovella, G.Robb, NIMA(2005)

quantum regime ( $\bar{\rho} = 0.05$ )

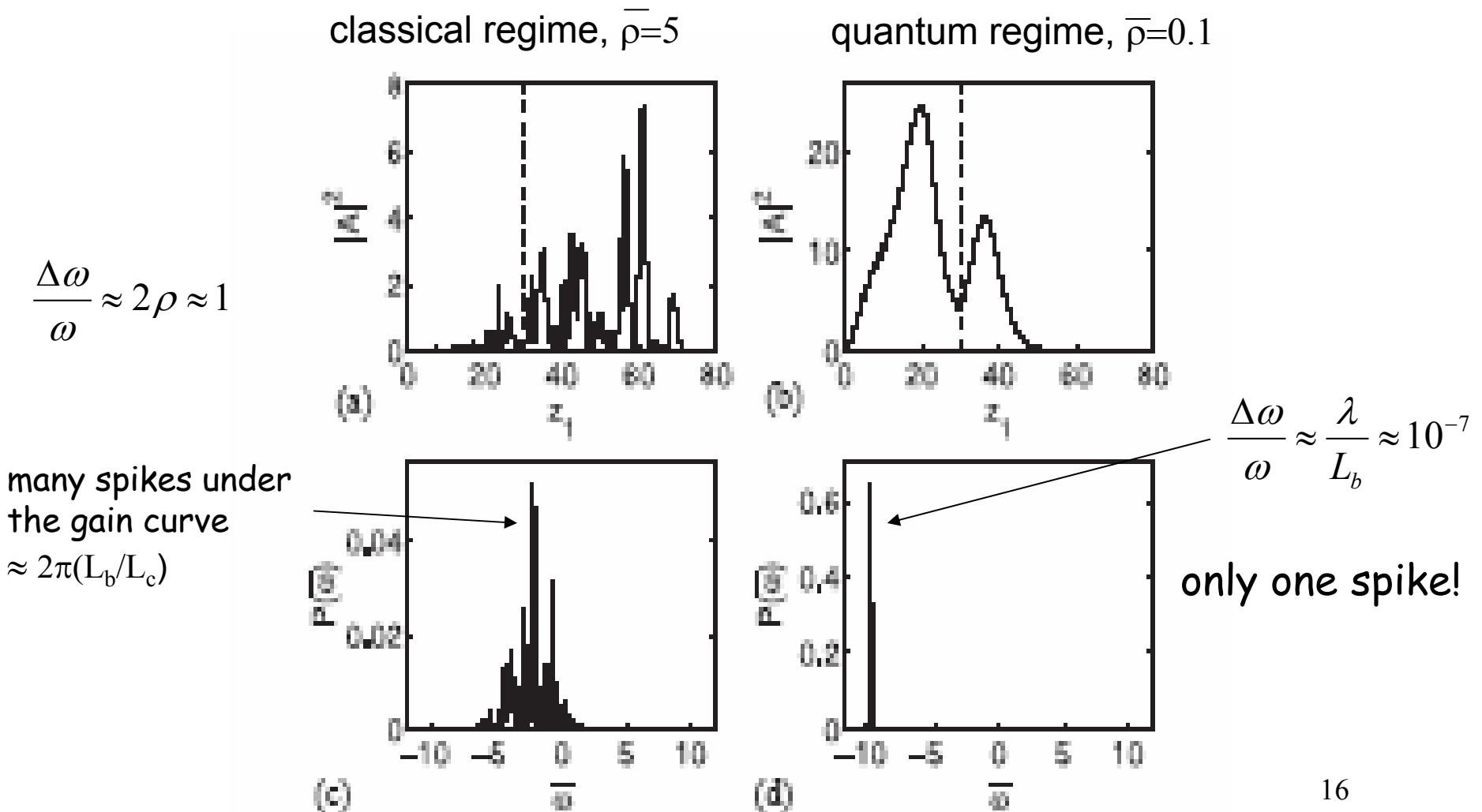


classical regime ( $\bar{\rho} = 5$ )

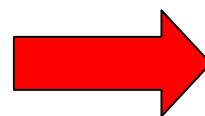
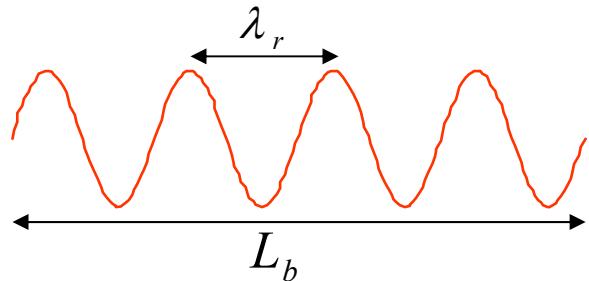




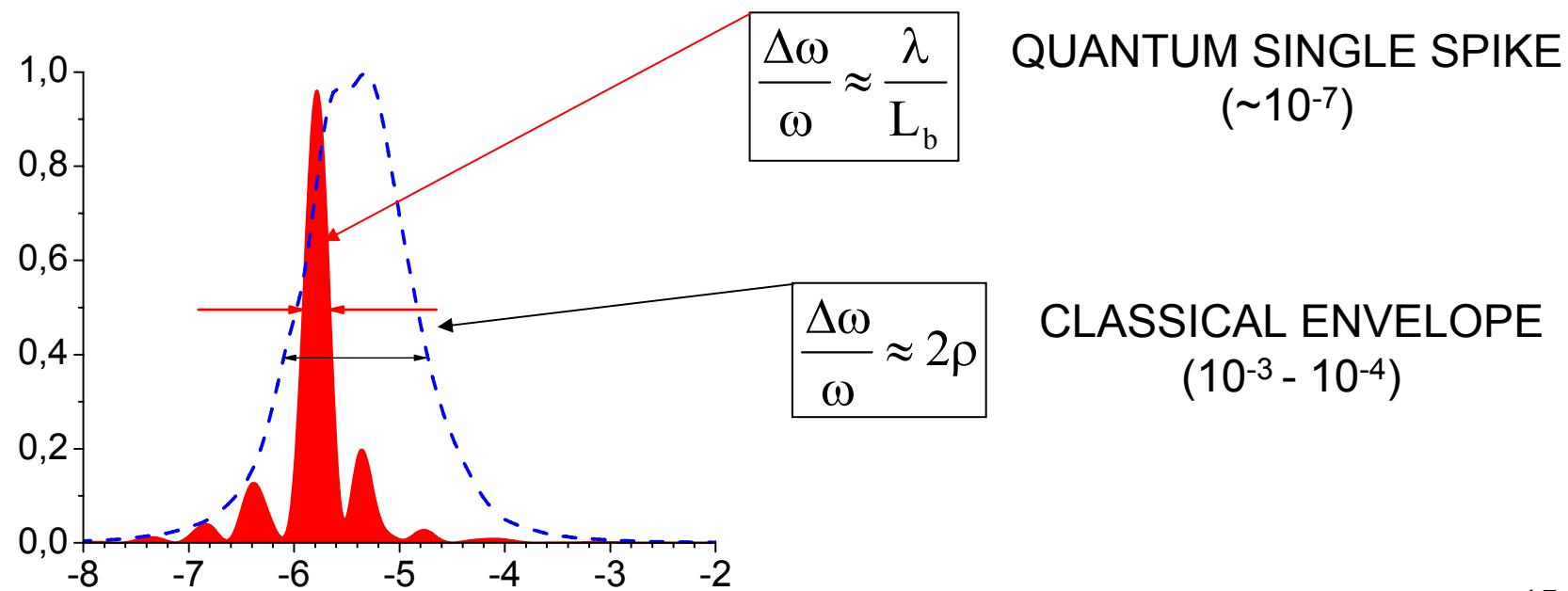
# CLASSICAL AND QUANTUM SASE



# LINEWIDTH OF THE SPIKE IN THE QUANTUM REGIME



$$\left( \frac{\Delta\omega}{\omega} \right)_{\text{QFEL}} \approx \frac{\lambda_r}{L_b}$$



# why QFEL requires a LASER WIGGLER?

$$\bar{\rho} = \rho \frac{mc\gamma}{\hbar k_r} = \rho \gamma \frac{\lambda_r}{\lambda_c} \quad \left( \lambda_c = \frac{h}{mc} \right)$$

$$\gamma = \sqrt{\frac{\lambda_w(1+a_w^2)}{2\lambda_r}}$$

$$\bar{\rho} \leq 1 \Rightarrow \rho \leq \frac{\sqrt{2}\lambda_c}{\sqrt{\lambda_r \lambda_w(1+a_w^2)}} \quad \text{and}$$

$$L_w \approx \frac{\lambda_w}{\rho} \geq \frac{\sqrt{\lambda_r \lambda_w^3(1+a_w^2)}}{\sqrt{2}\lambda_c}$$

for a laser wiggler  $\lambda_w \rightarrow \lambda_L / 2$

to lase at  $\lambda_r=1 \text{ \AA}$ :

## MAGNETIC WIGGLER:

$\lambda_w \sim 1 \text{ cm}, E \sim 3.5 \text{ GeV}$

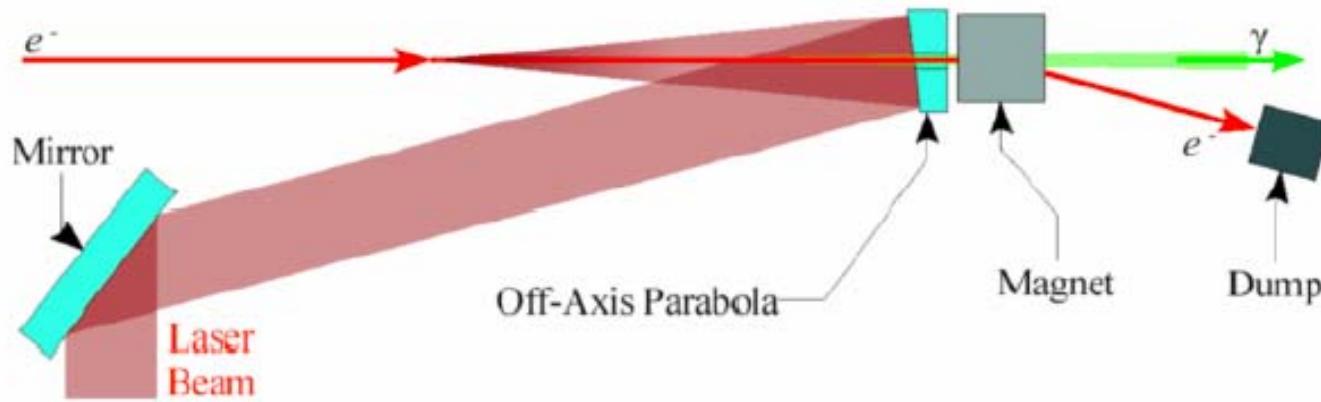
$\rho \sim 10^{-6}, L_w \sim 3 \text{ Km}$

## LASER WIGGLER

$\lambda_w \sim 1 \mu\text{m}, E \sim 25 \text{ MeV}$

$\rho \sim 10^{-4}, L_w \sim 2 \text{ mm}$  18

# typical parameters for QFEL



**Electron beam**

E [MeV]	21
I [A]	420
$\varepsilon_n$ [mm mrad]	0.1-1
$\delta\gamma/\gamma$ [%]	0.01

**Laser beam**

$\lambda_L$ [ $\mu\text{m}$ ]	1
$P_L$ [TW]	2
$\tau$ [ps]	40
$w_0$ [ $\mu\text{m}$ ]	30
$Z_r$ [mm]	3

**QFEL beam**

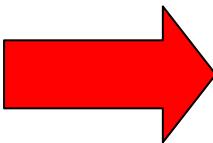
$\lambda_r$ [Å]	1.5
$P_r$ [MW]	3.5
$\Delta\omega/\omega$	$10^{-7}$
Brilliance	$10^{28}$

# EMITTANCE CRITERIA FOR A LASER WIGGLER

**1) e-beam contained in the laser beam:**

$$Z_R = \frac{4\pi R^2}{\lambda_L} , \quad \beta^* = \frac{\sigma_b^2 \gamma}{\varepsilon_n}$$

$$Z_R \leq \beta^*$$



$$\varepsilon_n \leq \frac{\gamma \lambda_L}{4\pi} \left( \frac{\sigma}{R} \right)^2$$

for instance  $\gamma=50$ ,  $\lambda_L=1 \text{ }\mu\text{m}$ ,  $\varepsilon_n \leq 4 \left( \frac{\sigma}{R} \right)^2$

If  $R \sim 2\sigma$ ,  $\varepsilon_n < 1 \text{ mm mrad} !!$

**2) Resonance broadening  
due to off-axis emission smaller  
than the natural linewidth**



$$\begin{aligned} \varepsilon_n &\leq \frac{\gamma \lambda_r}{2\pi} \sqrt{\frac{Z_{rad}}{L_{gain}}} \\ &\approx 0.06 \lambda_L \frac{\sigma}{R} \sqrt{(1 + a_w^2)(2Z_L / L_g)} \end{aligned}$$

$$Z_r = \frac{4\pi\sigma^2}{\lambda_r}$$

# SCALING LAWS FOR A LASER WIGGLER

independent parameters:

$$\lambda_r(\text{\AA}), \lambda_L(\mu\text{m}), a_w, \bar{\rho}$$

$$\rho \approx 5 \cdot 10^{-4} \frac{\bar{\rho}}{\sqrt{\lambda_r \lambda_L (1 + a_w^2)}}$$

$$L_g(\text{mm}) \approx 0.1 \sqrt{\lambda_r \lambda_L^3 (1 + a_w^2) \frac{1 + \bar{\rho}}{\bar{\rho}^3}}$$

e-beam:

$$\gamma \approx 50 \sqrt{\frac{\lambda_L}{\lambda_r} (1 + a_w^2)}$$

$$I(A) \approx \frac{5 \cdot 10^3}{a_w^2} \sqrt{\frac{\lambda_L}{\lambda_r^5} (1 + a_w^2) \bar{\rho}^3 (1 + \bar{\rho})}$$

$$\varepsilon_n(\text{mm rad}) \approx 0.13 \cdot \lambda_L \sqrt{1 + a_w^2}$$

$$\sigma(\mu\text{m}) \approx 4 \left[ \lambda_r \lambda_L^5 (1 + a_w^2) \frac{1 + \bar{\rho}}{\bar{\rho}^3} \right]^{1/4}$$

laser beam:

$$P_L(\text{TW}) \approx 3 \cdot a_w^2 \sqrt{\lambda_r \lambda_L (1 + a_w^2) \frac{1 + \bar{\rho}}{\bar{\rho}^3}}$$

$$\tau(\text{ps}) \approx 1.35 \cdot \sqrt{\lambda_r \lambda_L^3 (1 + a_w^2) \frac{1 + \bar{\rho}}{\bar{\rho}^3}}$$

R. Bonifacio, N.Piovella, M.Cola,  
L.Volpe, NIMA (2007)

$$2Z_L = 5L_g, R = 1.5\sigma$$

<b>QFEL parameter</b>	$\bar{p}$	0.2	5	0.2
<b>Laser wave length</b>	$\lambda_L$ ( $\mu\text{m}$ )	1	1	10
<b>Radiation length</b>	$\lambda_r$ ( $\text{\AA}$ )	1.5	1.5	1.5
<b>Wiggler parameter</b>	$a_0$	0.3	0.8	0.3
<b>FEL parameter</b>	$\rho$	$7.5 \cdot 10^{-5}$	$1.5 \cdot 10^{-3}$	$2.4 \cdot 10^{-5}$
<b>Gain length</b>	$L_g$ (mm)	1.2	0.03	37.5
<b>Laser Rayleigh range</b>	$Z_L$ (mm)	3	0.07	94
<b>Laser radius</b>	$R$ ( $\mu\text{m}$ )	15.4	2.4	274
<b>Laser power</b>	$P_L$ (TW)	2	0.3	6
<b>Laser duration</b>	$\tau_L$ (ps)	40	1	1200
<b>Interaction length</b>	$L_{\text{int}}$ (mm)	6	0.14	187
<b>E-beam energy</b>	$\gamma$	42.6	52.3	135
<b>E-beam radius</b>	$\sigma$ ( $\mu\text{m}$ )	10.3	1.6	183
<b>Peak current</b>	I (kA)	0.4	22	1.3
<b>Emittance limit</b>	$\epsilon_n$ (mm mrad)	0.1	0.1	0.9
<b>Gain band width</b>	$\Delta\gamma/\gamma$	$3.4 \cdot 10^{-5}$	$1.5 \cdot 10^{-3}$	$1.1 \cdot 10^{-5}$
<b>FEL line width</b>	$\Delta\omega/\omega$	$2.1 \cdot 10^{-7}$	$1.5 \cdot 10^{-3}$	$6.6 \cdot 10^{-7}$
<b>Number of spikes</b>	$N_s$	1	278	1
<b>FEL power</b>	(MW)	3.5	900	11
<b>Photons' number</b>	$N_{\text{ph}}$	$6 \cdot 10^9$	$3 \cdot 10^{10}$	$6 \cdot 10^9$
<b>Peak Brilliance</b>	$B^{(*)}$	$10^{28}$	$1.6 \cdot 10^{26}$	$1.2 \cdot 10^{23}$

## Classical versus quantum SASE

Classical SASE FEL X-ray experiments (DESY,LCLS):

- require very long Linac (~GeV, Km) and undulators (~100 m)
- Generate chaotic radiation with broad and spiky spectrum ( $\Delta\omega/\omega \sim 10^{-3}$ ).
- Have very high cost ( $10^9$  U\$) and large size

a QFEL experiment

- will generates a single spike almost monochromatic X-ray radiation ( $\Delta\omega/\omega \sim 10^{-7}$ ).
- Needs a laser wiggler
- Reduces cost ( $\sim 10^6$  U\$)
- Very compact apparatus ( $\sim$  m)

## Summary

**Classical regime :  $\bar{\rho} \gg 1$**

**Quantum regime:  $\bar{\rho} \leq 1$  : discreteness of momentum exchange  
relevant=> quantum effects.**

**The system is prepared in a defined momentum state  $p_0$ , making  
transition to the lower state**

**The system radiates a monocromatic train wave  $\lambda$ , whose  
length is  $L_b$ . Hence one has a single line with linewidth**

$$\lambda_r / L_b \approx 10^{-7}$$

**In the opposite case random transition from many momentum  
states. Each transition gives a spike with different frequency.**

**Total bandwidth:**  $\rho \approx 10^{-3}$

**QFEL has a linewidth 4 orders of magnitude smaller than the  
classical**

**The dimensions and cost are three order of magnitude  
smaller**