

# *Controlling chaos in the wave/particle interaction*

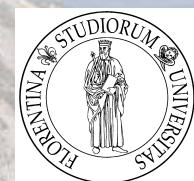
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# *Outline*

- Chaotization of the dynamics to stabilize a wave
  - Using a monochromatic Hamiltonian model of the FEL
  - Control based upon a linear stability analysis of periodic orbits
  - A control technique allowing to find the apt values of the parameters
- Regularization of the dynamics to reduce energy spread
  - Using a model of electrons interacting with test-waves
  - Method providing a control term which creates barriers preventing diffusion in phase-space
  - Experimental check

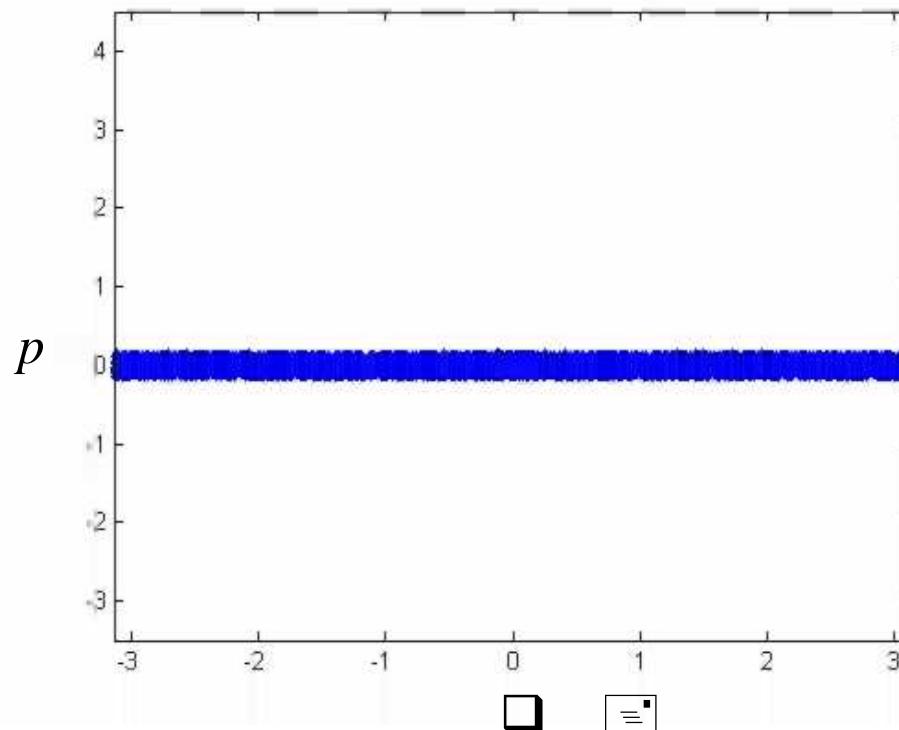
## *Hamiltonian model with $N$ bodies and one wave*

$$H_N = \sum_{j=1}^N \frac{p_j^2}{2} - \frac{2\sqrt{I}}{M} \sum_{j=1}^N \cos(\theta_j - \phi)$$

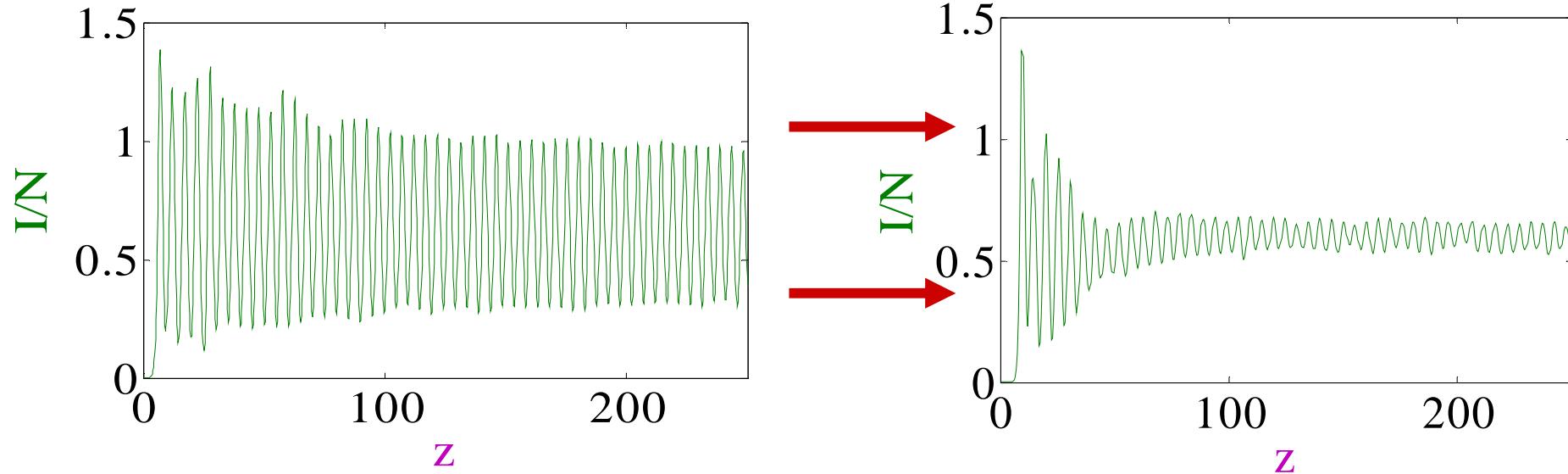
kinetic part      interaction term

$(q_j, p_j)$  longitudinal position/ momentum of the particles

$(\phi, I)$  phase/ intensity of the wave



# *Reducing oscillations*



A key structure : the macro-particle

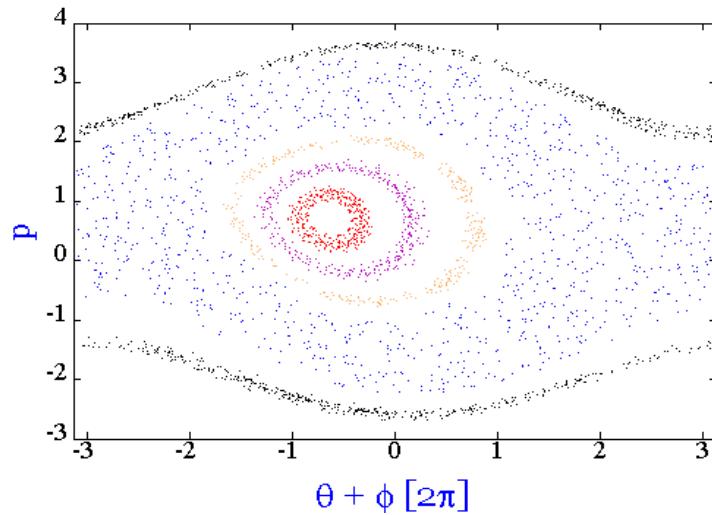
- An aggregate of particles (bunching)
- Its oscillations in phase-space is responsible for the wave's fluctuations

# *A mean-field Hamiltonian*

Self-consistent

$$H_N = \sum_{j=1}^N \frac{\dot{q}_j^2}{2} - 2\sqrt{N} \sum_{j=1}^N \cos(q_j - \mathbf{j})$$

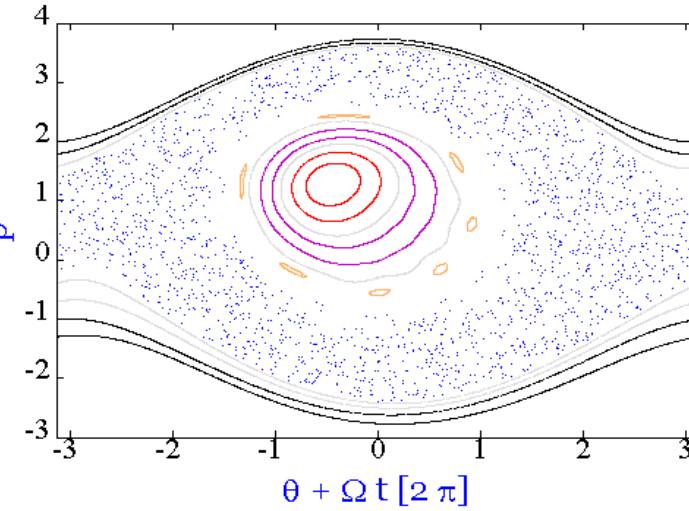
Poincaré sections at  $d\mathbf{I}/dt = 0$



Mean-field

$$H_{mf} = \frac{p^2}{2} - 2\sqrt{\frac{p\phi(t)}{N}} \cos(q - \phi(t))$$

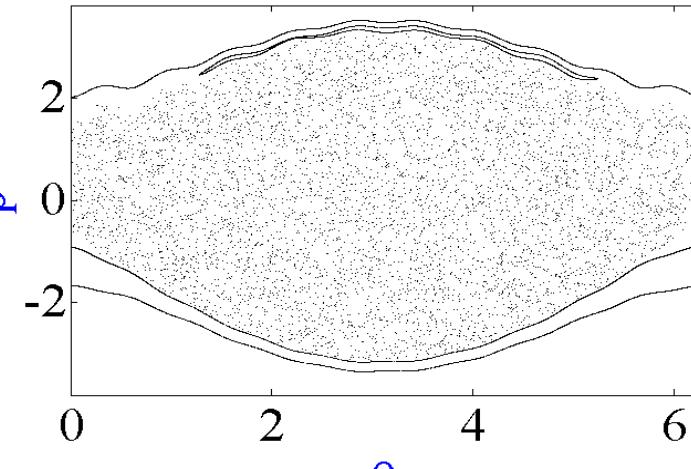
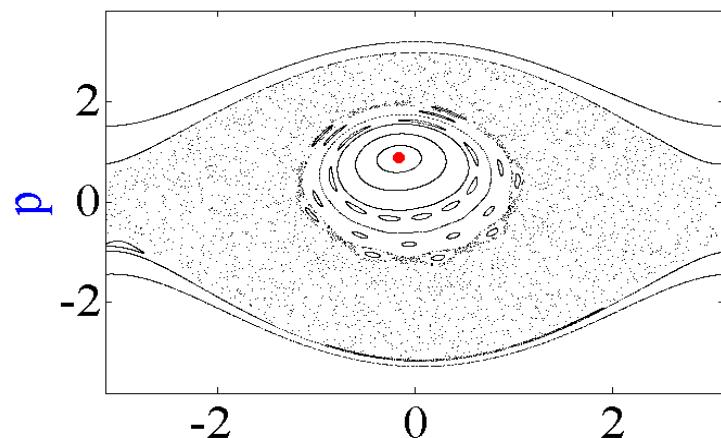
Poincaré sections at  $d\mathbf{Y}(t)/dt = 0$



Mean-field model :

- The macro-particle seen as an invariant structure of phase-space
- Methods of « control of chaos » to reshape this invariant structure

## The residue method

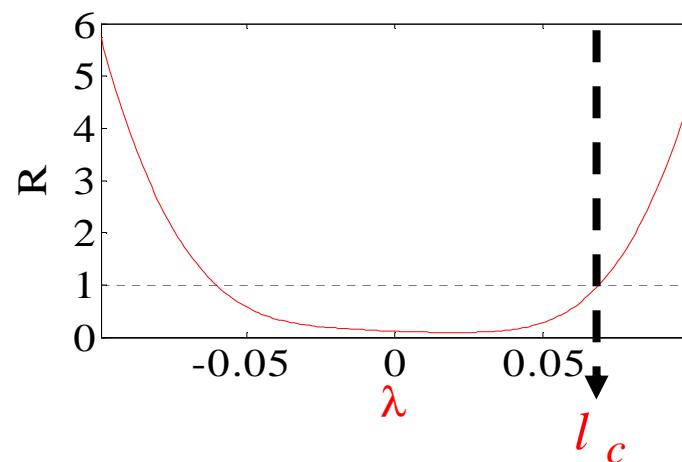


$$H_{mf} = \frac{p^2}{2} - 2\sqrt{\frac{\dot{I}(t)}{N}} \cos(q - \dot{\phi}(t))$$

Mean-field model

$$- 2l \cos(k(q - w_1 t))$$

+ test-wave

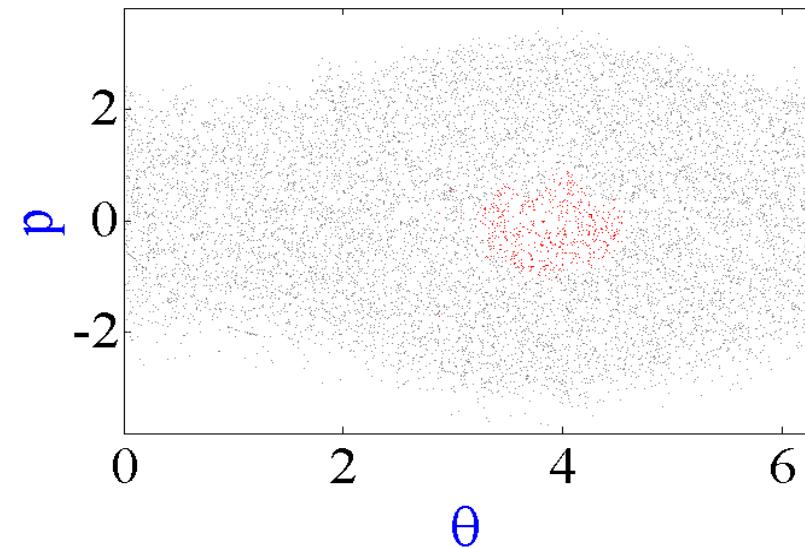
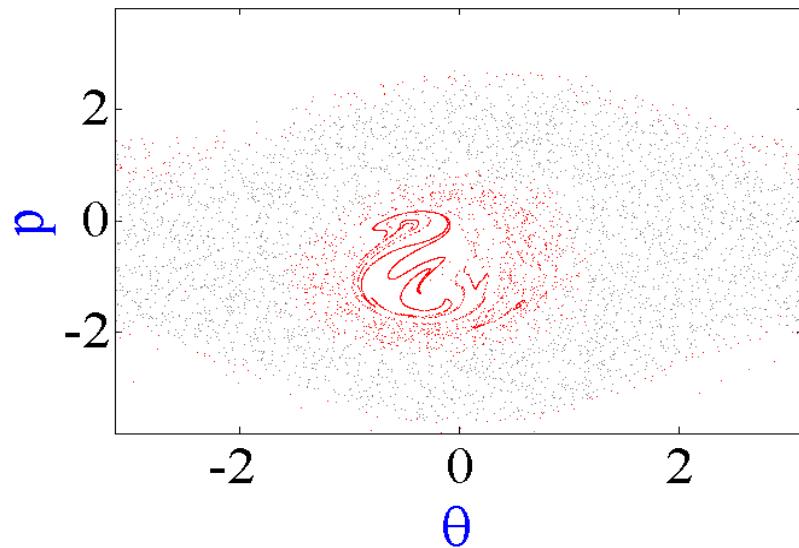


Linear Stability Analysis  
of the Periodic Orbit

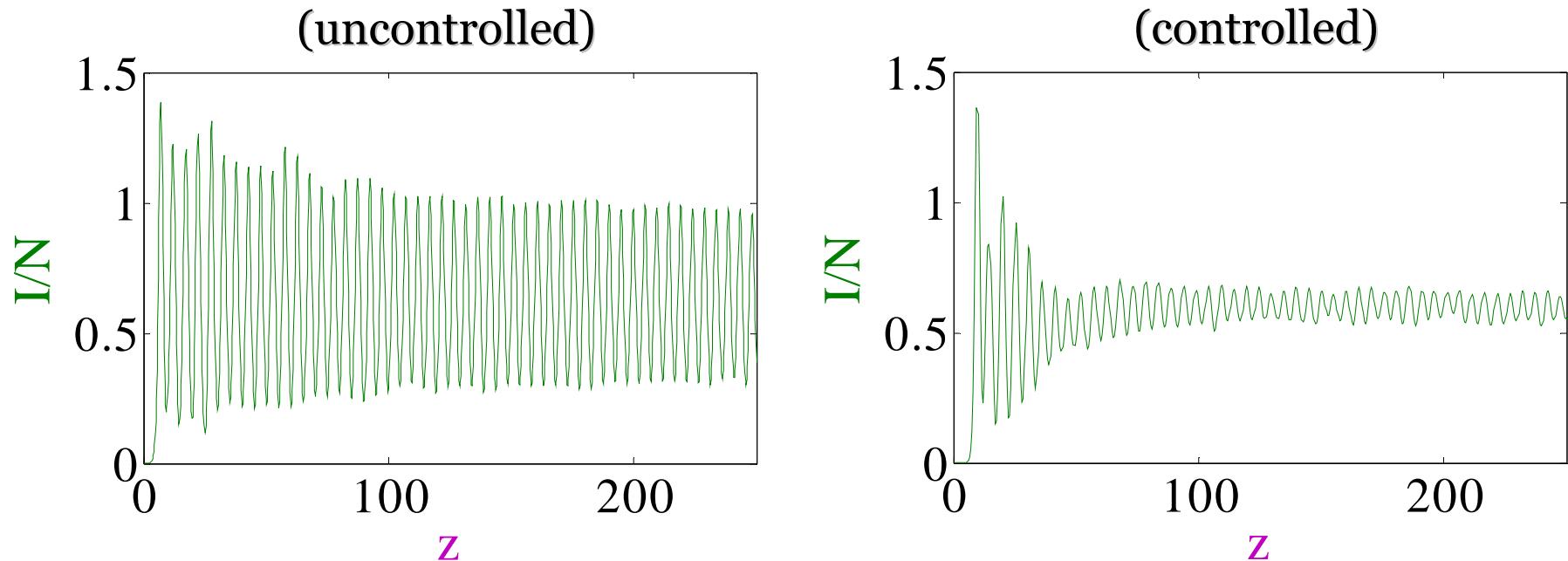
$\Rightarrow$  Destruction of the macro-particle

# Control of the self-consistent dynamics

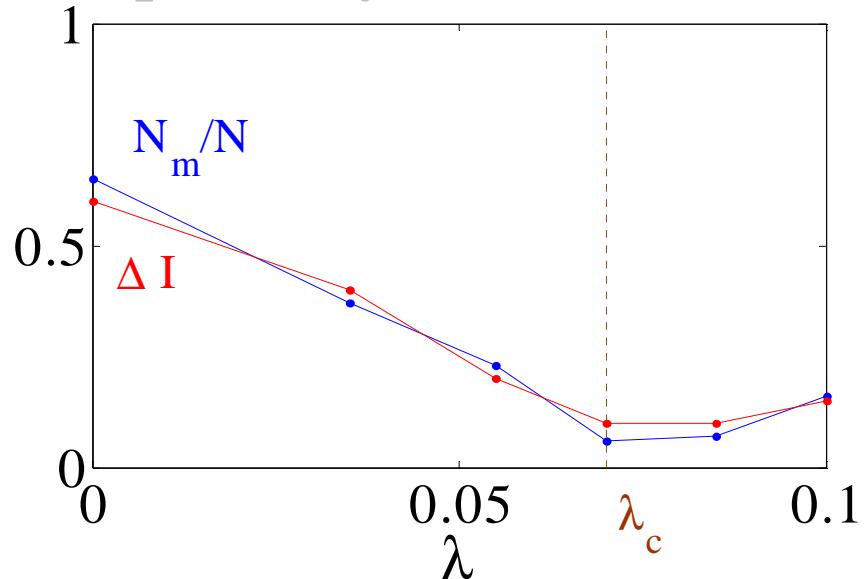
$$H_N = \sum_{j=1}^N \frac{\mathbf{p}_j^2}{2} - 2\sqrt{\frac{I}{N}} \sum_{j=1}^N \mathbf{a} \cos(\mathbf{q}_j + \mathbf{j}) - 2l_c \mathbf{a} \sum_j \cos(k(q_j - w_1 t))$$



As for the intensity...

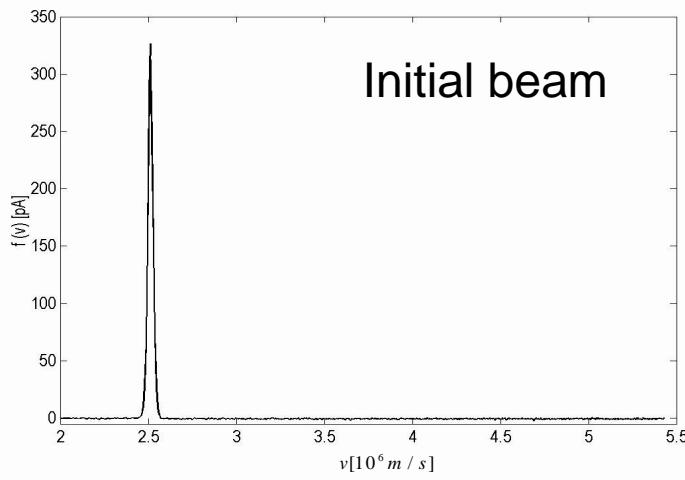


Optimality of the control

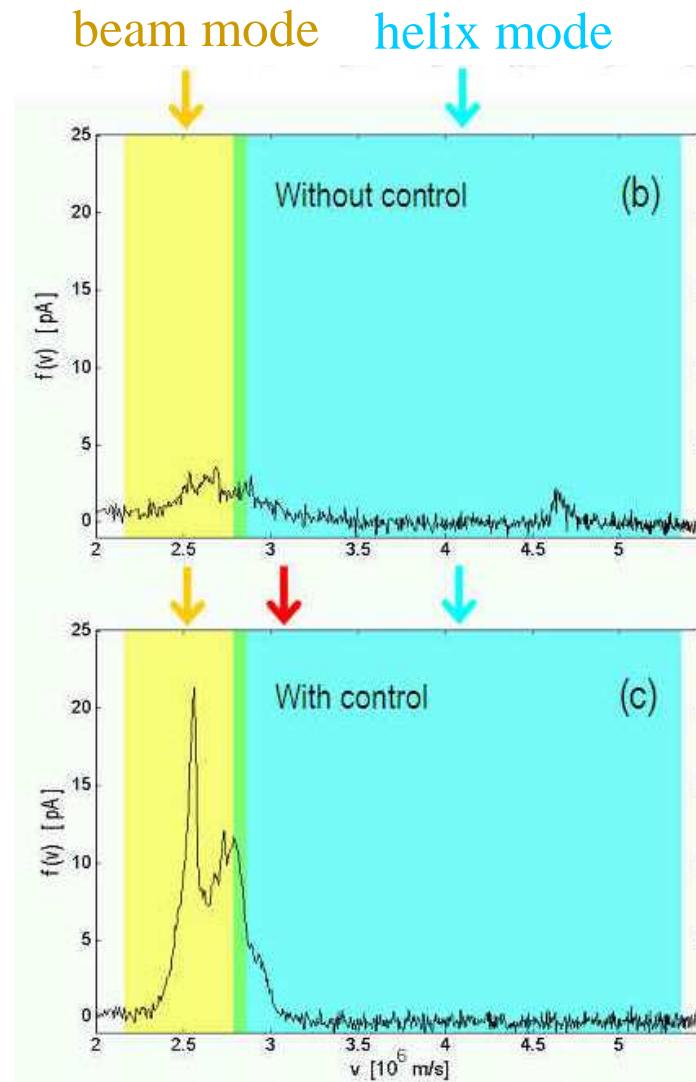


# *Experiment on a Travelling Wave Tube*

Beam velocity distribution

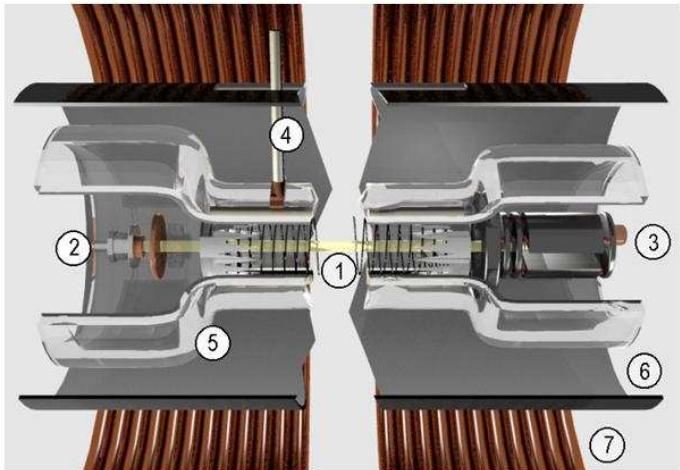


Initial beam

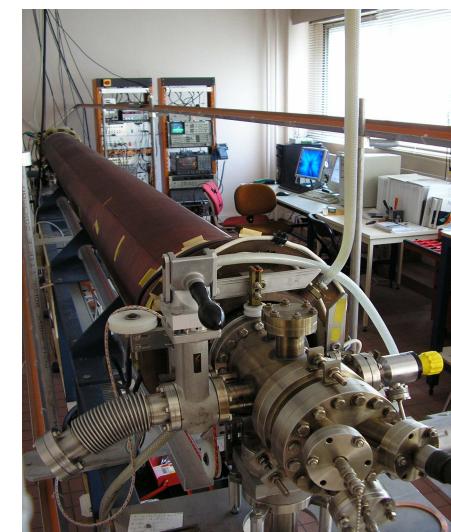
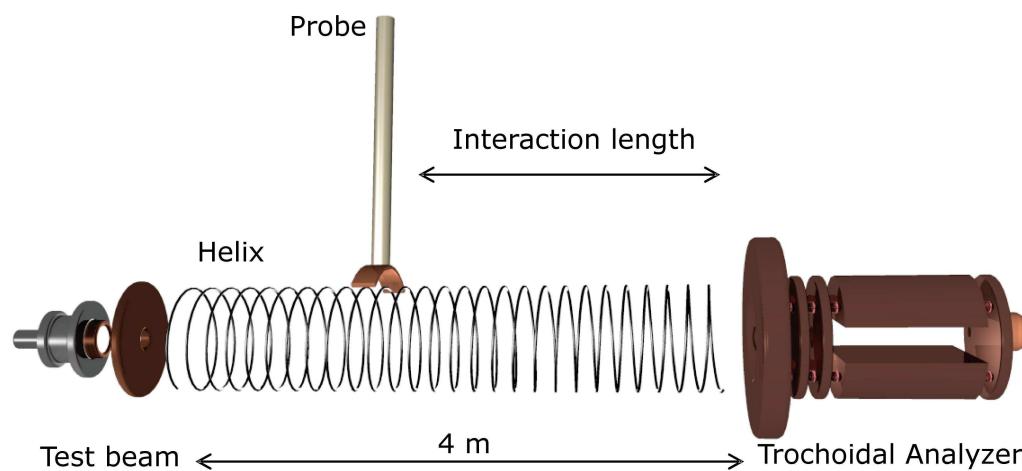


C. Chandre *et al.*, Phys. Rev. Lett. **94**, 074101 (2005)

# *Experiment on a Travelling Wave Tube*

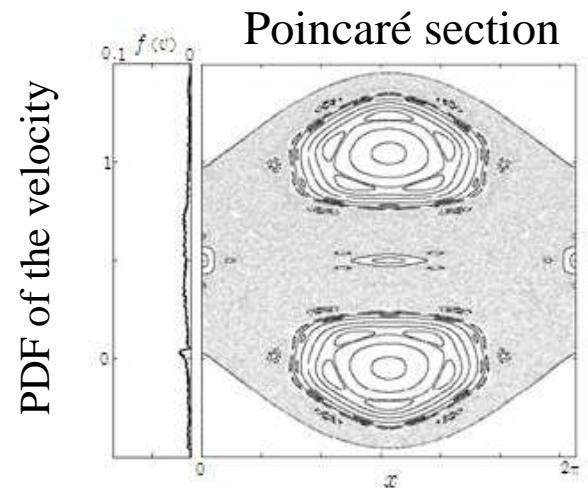


- 1) helix
- 2) electron gun
- 3) trochoïdal analyzer
- 4) antenna
- 5) vacuum chamber

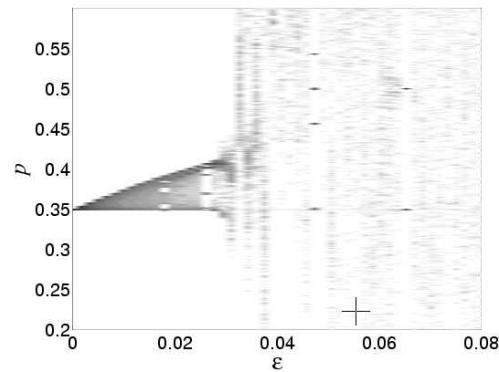


Experiment realized by F. Doveil, Y. Elskens, A. Macor  
at PIIM, Université de Provence

$$H_c = \frac{p^2}{2} + e_1 \cos(k_1 x - w_1 t + j_1) + e_2 \cos(k_2 x - w_2 t + j_2)$$



> PDF of the momenta of a trajectory

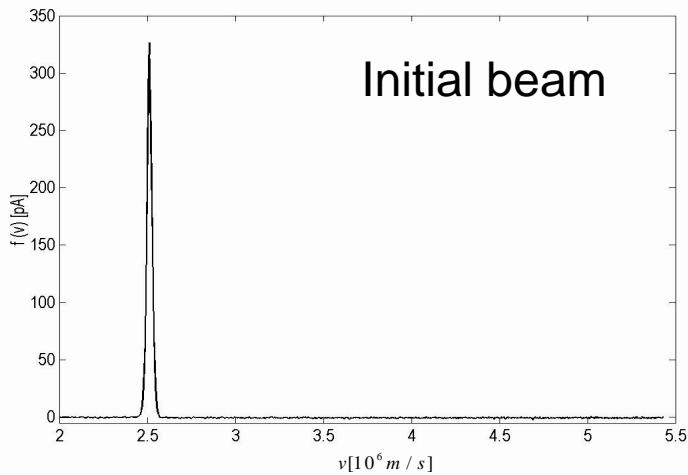


The new barriers prevent diffusion of electrons in position-momentum space

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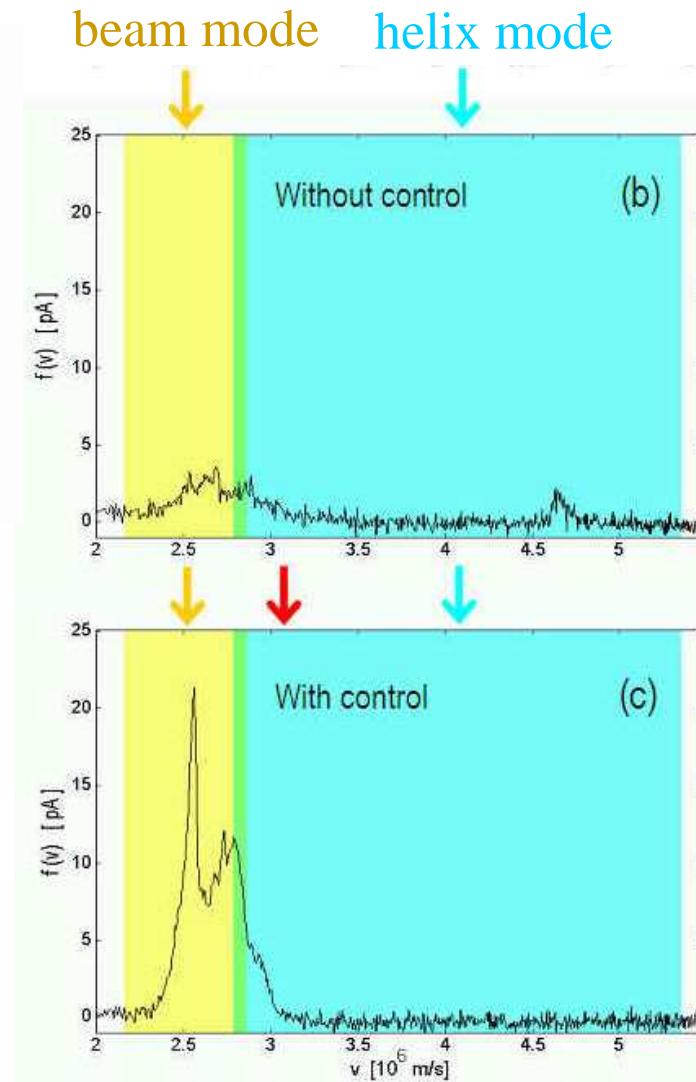
# *Experiment on a Travelling Wave Tube*

Beam velocity distribution



Initial beam

Energy : 
$$\frac{E_{\text{contr}}}{E_{\text{syst}}} \approx 0,1\%$$



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# *Experiment on a Travelling Wave Tube*

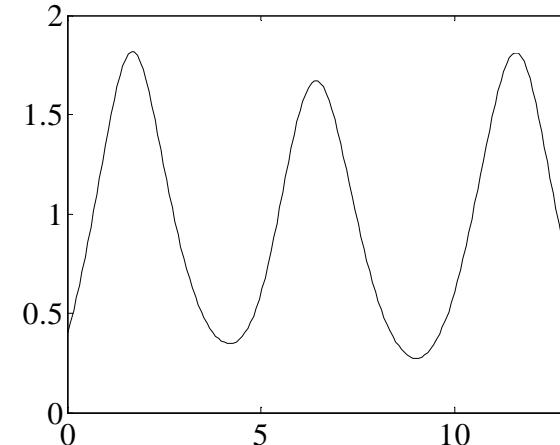
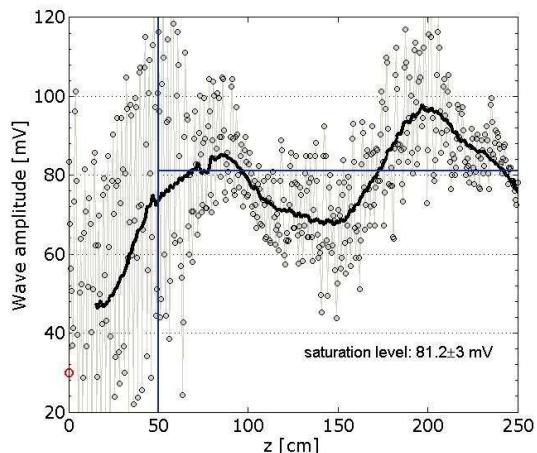
TWT with one wave at low intensity : test-particle regime

$$H = \ddot{\mathbf{a}}_j \frac{\alpha p_j^2}{2} + e_1 \cos(k_1 x_j - w_1 t + j_1) \frac{\ddot{\phi}}{\phi}$$

TWT with one wave at larger intensity : self-consistent regime

$$H = \ddot{\mathbf{a}}_j \frac{\alpha p_j^2}{2} + \sqrt{I} \cos(kx_j + j) \frac{\ddot{\phi}}{\phi}$$

→ The TWT is a nice test-bed for FEL experiments



- Methods of control of invariant structures of phase-space
  - cutting down aggregation of particles by destroying invariant structures
  - reducing diffusion of particles in phase-space by building barriers
- Perspectives/in progress
  - a method of control to limit sensitivity to initial conditions (jittering), based on a Lie algebra formalism
  - derivation of a Hamiltonian model for a beam of particles interacting with an electromagnetic field