Summary talk

Paul Hoyer University of Helsinki

First Workshop on Quark-Hadron Duality and the Transition to pQCD

Laboratori Nazionali di Frascati June 6-8 2005



Amor Sacro e Amor Profano

Quark and Hadron Dynamics in QCD



Click on picture for QCD explanation

(Keynote version only)

Art or Science?

Henk Blok:

Duality: Belief or hold? $p vs. n, L vs. T, \sigma_{3/2} vs. \sigma_{1/2}$, semi-incl., excl.? global, local?resonance vs. background

Rolf Ent:

If one integrates over all resonant and non-resonant states, quark-hadron duality should be shown by any model. This is simply unitarity. However, quark-hadron duality works also, for $Q^2 > 0.5$ (1.0) GeV², to better than 10 (5) % for the F₂ structure function in both the N- Δ and N-S₁₁ region! (Obviously, duality does not hold on top of a peak! -- One needs an appropriate energy range)

Why does local quark-hadron duality work so well, at such low energies? Paul Hoyer Frascati 8.6 2005

Confinement is local

Many acknowledgements to:

Quark-Hadron Duality in Electron Scattering

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Abstract

The duality between partonic and hadronic descriptions of physical phenomena is one of the most remarkable features of strong interaction physics. A classic example of this is in electron-nucleon scattering, in which low-energy cross sections, when averaged over appropriate energy intervals, are found to exhibit the scaling behavior expected from perturbative QCD. We present a comprehensive review of data on structure functions in the resonance region, from which the global and local aspects of duality are quantified, including its flavor, spin and nuclear medium dependence. To interpret the experimental findings, we discuss various theoretical approaches which have been developed to understand the microscopic origins of quark-hadron duality in QCD. Examples from other reactions are used to place duality in a broader context, and future experimental and theoretical challenges are identified.

200 pages400 references

24 Jan 2005 arXiv:hep-ph/0501217 v1

Duality in $e^+e^- \rightarrow hadrons$



Resonances build the parton subprocess cross section because of a separation of scales between hard and soft processes.

QCD Sum Rules,... SVZ



The softer is confinement physics, the smaller Q_0 can be chosen, and the more local is duality.

The "higher twist" contributions to an inclusive cross section are naturally of order $(Q_0/Q)^2$, in the absence of another scale.

 $\Rightarrow \text{Local duality (low } Q_0)$ comes with small higher twist.

These features are built into quantum "mechanical resonance models. Wally Melnitchouk

Duality arises as consequence of the separation of time scales between hard and soft physics.



More detailed features such as spin dependent distributions appear to require higher Q^2 and broader averaging.



Comparison of $\Gamma_1^{res}/\Gamma_1^{DIS}$ in same x range

Use PDF



Bianchi, Fantoni & Liuti, PR D69 (2004) 014505

Duality in F₂

Local duality indicates that o_{DIS} and the form factors $F(p \rightarrow N^*)$ are governed by the same hard subprocess and by the same target Fock states.



However, it is frequently held that

- DIS occurs via incoherent scattering on a single quark in a general Fock state
- Form factors are governed by coherent scattering on compact Fock states

Brodsky - Lepage

Form factor dynamics



pQCD calculations show that most of the proton form factor arises from gluon exchanges within the hard exclusive subprocess which are much softer than the photon: $k^2 \ll Q^2$ P. Kroll et al

The transverse size of the Fock states is then $r_{\perp} >> 1/Q$ Hence the virtual photon scatters incoherently on the constituents. Wally Melnitchouk:



Moreover:

There are fast transitions between compact Fock states and non-compact ones, where one quark nearly carries all the momentum. Both states have short lifetimes, similar to the virtual photon.



If the virtual photon scatters from the $x \sim 1$ quark the scattering is incoherent and qualitatively similar to DIS at large x.

Summing up:

The QCD dynamics of hadron form factors is still controversial

Isgur - Llewellyn-Smith Radyushkin

I. Niculescu et al, Phys. Rev. Lett. 85, 1182 (2000) • $Q^2 = 0.20 (GeV/c)$ • $Q^2 = 0.45 (GeV/c)$ 0.4 = 0.85 (GeV)= 1.4 (GeV/c0.3 = 2.4 (GeV/c) $Q^2 = 3.3 (GeV/c)^2$ NMC 10 0.2 (a) 0.1 0.8 = 0.20 (GeV/c)= 0.45 (GeV) 0.85 (GeV 0.6 1.4 (GeV/c 2.4 (GeV/c 3.3 (GeV 0.4 0.2 (b) 0.9 0.2 0.3 0.5 0.6 0.8 04 07

Duality gives important clues, and indicates that "end-point" regions are important. These involving transversally large Fock states where one quark carries most of the momentum.

The photon scatters incoherently off such states, hence interference terms $\propto e_q e_q$, do not contribute. This makes local duality possible for both proton and neutron targets.

Cynthia Keppel Nuclear target dependence

Duality in F_2 ...let the nucleus do the averaging

- Data in resonance region, spanning Q² range 0.7 - 5 GeV²
- GRV curve
- The nucleus does the averaging
- For larger A, resonance region indistinguishable from DIS



Cynthia Keppel Duality and the EMC Effect

Red = resonance region data

Blue, purple, green = deep inelastic data from SLAC, EMC

Medium modifications to the structure functions *are the same* in the resonance region as in the DIS

Cross-over can be studied with new data



J. Arrington, et al., submitted

Coherent vs Incoherent Rescattering in Nuclei

The inclusive DIS cross section (including resonance region) is influenced only by rescattering which happens within the coherence length of the virtual photon

$$L_I = \frac{1}{Q} \cdot \frac{\nu}{Q} = \frac{\nu}{Q^2} = \frac{1}{2mx_B} \approx 0.1 \text{ fm at } x_B = 1$$

Thus the nucleon G_E and G_M extracted by Rosenbluth separation should not be affected by rescattering in the nuclear environment.

Hadrons observed after traversal through the nucleus, as in SIDIS $e A \rightarrow e+h+X$, are affected by incoherent rescattering.

Thus G_E and G_M measured through spin transfer to the final nucleon can be sensitive to spin flips in nuclear rescattering.

Xin-Nian Wang



Duality works in both $\sigma_{\rm L}$ and $\sigma_{\rm T}$

Rolf Ent

E94-110 : Precise Measurement of Separated Structure Functions in Nucleon Resonance Region



- The resonance region is, on average, well described by NNLO QCD fits.
- This implies that Higher-Twist (FSI) contributions cancel, and are on average small.
- The result is a smooth transition from Quark Model Excitations to a Parton Model description, or a smooth quark-hadron transition.
- This explains the success of the parton model at relatively low W^2 (=4 GeV²) and Q^2 (=1 GeV²).

"The successful application of duality to extract known quantities suggests that it should also be possible to use it to extract quantities that are otherwise kinematically inaccessible."

(CERN Courier, December 2004)

σ_L is not suppressed at large x

In the Bjorken limit: $Q^2 \rightarrow \infty$ at fixed x σ_L is of order α_s and arises when γ_L couples to a virtual quark.

For $x \rightarrow 1$ the initial quark goes off-shell. When $Q^2(1-x)$ is held fixed the quark virtuality is commensurate with the photon resolution, and $\sigma_L \sim \sigma_T$ is not suppressed.



In local duality $Q^2(1-x) = M_{N*}^2 - M_p^2$ is fixed: This is not the Bj limit! \Rightarrow There is no twist expansion, hence one cannot properly talk about "higher twist corrections"



Experimental evidence for dominance of σ_{I} at high x

In $\pi N \rightarrow \mu^+ \mu^- X$ the angular distribution of the $\mu^+ \mu^- \propto 1 + \lambda \cos^2 \theta$ reveals directly the polarization of the (timelike) virtual photon:

 $\lambda = +1$: Transverse polarization $\lambda = -1$: Longitudinal polarization

For $x_F(\mu^+\mu^-) \rightarrow 1$ the quark in the pion carries $x \rightarrow 1$ and goes off-shell. In effect the (scalar) pion scatters coherently via a longitudinal photon.





Local duality for the pion

Duality for a pion "target" can be studied by comparing $f_{q/\pi}(x)$ measured in Drell-Yan (which contains σ_L) with the pion e.m. form factor (which only has contributions from σ_L).

In pQCD, $f_{q/\pi}(x) \propto (1-x)^2$ for σ_T in the Bj limit. The E615, $f_{q/\pi}(x) \sim (1-x)^1$ indicating substantial σ_L (as seen above).

Insofar as σ_L dominates in DY at large x, this is another succesful test of local duality, now for the pion pole.



Local duality at $Q^2 = 0$

Alexander Donnachie

The spin-averaged photoproduction total cross section obeys the duality of soft hadron scattering: Resonances are dual to Regge exchange "Background" is dual to "Pomeron" exchange

Regge duality was a very succesful phenomenology in the 60-70's, and led to the Veneziano model, whence to string theory.

The QCD origins of Regge duality is more obscure than that of Bloom-Gilman duality, since it involves only soft physics.



Resonances dual to valence quark distribution

At high x, the contribution from gluons and sea quarks is small.

At low x and Q^2 , the relation between resonances and DIS is close to Regge duality, where two-component duality has been shown to hold.

The success of Bloom-Gilman duality in the valence quark distribution at low x may thus be a consequence of Regge duality.



Rolf Ent

Duality in Meson Electroproduction



Requires non-trivial cancellations of decay angular distributions If duality is not observed, factorization is questionable

Duality and factorization possible for $Q^2, W^2 \le 3 \text{ GeV}^2$ (Close and Isgur, Phys. Lett. B509, 81 (2001))

Unpolarized SIDIS: new test

Simple picture in valence region:

 $\sigma_{p}(\pi^{+}) = 4u(x)D^{+}(z) + d(x)D^{-}(z)$ $\sigma_{p}(\pi^{-}) = 4u(x)D^{-}(z) + d(x)D^{+}(z)$ $\sigma_{d}(\pi^{+}) = [u(x)+d(x)] [4D^{+}(z) + D^{-}(z)]$ $\sigma_{d}(\pi^{-}) = [u(x)+d(x)] [4D^{-}(z) + D^{+}(z)]$

$$\mathsf{R}_{pd} = [\sigma_p(\pi^+) + \sigma_p(\pi^-)] / [\sigma_d(\pi^+) + \sigma_d(\pi^-)]$$

= [4u(x) + d(x)] / 5[u(x)+d(x)]

= \(\sigma_p(x)/\(\sigma_d(x))\)

R_{pd} should be independent of z and p_t

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x-dependence of
R<sub>pd</sub> should be same
as
inclusve σ<sub>p</sub>(x)/σ<sub>d</sub>(x)
in valence region
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P. Bosted, Duality Workshop, June 2005 Peter Bosted

Unpolarized SIDIS: new test Peter Bosted



Extending duality: Semi-exclusive processes

C. Carlson et al S. Brodsky et al

Replace short-distance probe of DIS ($\gamma^* q \rightarrow q$) by some other compact subprocess, e.g., $\gamma^* u \rightarrow \pi^+ d$, with hard scale given by t and/or Q^2 .



The cross sections for $\gamma^{(*)} p \rightarrow \gamma Y$, $\gamma^{(*)} p \rightarrow \pi^+ Y$, ... can then be predicted in the limit where $\Lambda^2_{QCD} \ll M^2_Y \ll W^2$.

Note: Data on semi-exclusive processes such as $\gamma^* p \rightarrow \gamma Y$ is very welcome!

Using duality as in DIS ($W^2 \rightarrow M_N^2$) one may try to extrapolate $M_Y^2 \rightarrow M_N^2$ The calculated result is below the data by \geq factor 10.

D. M. Scott P. Eden et al



Failure may be due to duality extrapolation or to an incorrect dynamics for the semi-exclusive process, or...

Either way, we are learning something!



Summary Conclusions

- Duality is a vibrant and promising field, which is driven by experiment.
- Duality is hinting that QCD is scale (Q^2) invariant
 - $\Rightarrow Conformal Field Theory? Dimensional scaling rules. Guy de Teramond$ $Yet in confinement region, the <math>\Lambda_{OCD}$ scale is relevant Arkady Vainshtein
- QCD dynamics of hadron form factors should resemble DIS dynamics
 ⇒ Incoherent scattering on single quark?
- Two limits: 1. $Q^2 \rightarrow \infty$ with x fixed, then $x \rightarrow 1$ (Bj, twist expansion) 2. $Q^2 \rightarrow \infty$ with $Q^2(1-x)$ fixed (No twist expansion)

Local duality for a given resonance requires limit 2.

• Duality suggests principal features of confinement region, but our field theory tools are limited in this region.

Lattice measurements $\alpha_{MS}(M_Z)$ Howard Trottier

Schwinger-Dyson approximations

Limited applications

- ⇒ Need to develop description of relativistic bound states moving in color fields – using hints from duality.
- Duality developments likely to be driven by data also in the future

Spin dependence may yet be a challenge also to PQCD

• Progress requires close collaboration between experiment and theory