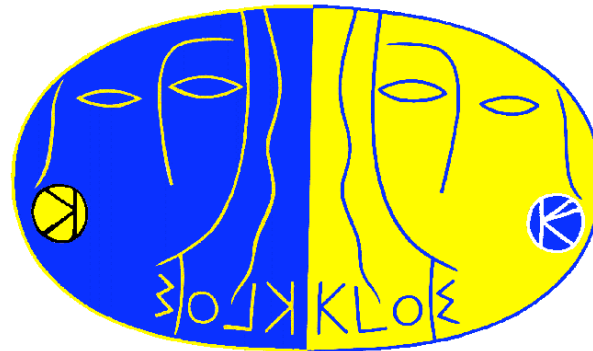


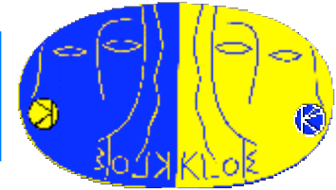
*Measurement of $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$
at DAΦNE with the radiative return*

Stefan E. Müller
Institut für Exp. Kernphysik,
Universität Karlsruhe
(for the KLOE collaboration)

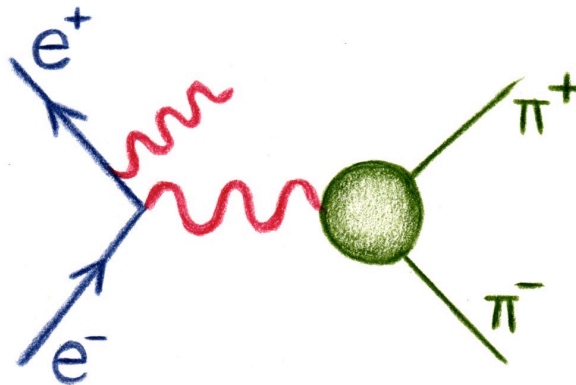


Workshop on
 e^+e^- in the 1-2 GeV range
Alghero, 10.-13. September 2003

$\sigma(e^+e^- \rightarrow \pi^+\pi^-)$ with ISR:



Particle factories have the opportunity to measure the cross section $\sigma(e^+e^- \rightarrow \text{hadrons})$ as a function of the hadronic c.m.s energy M_{hadrons}^2 by using the radiative return.



$$\sigma_{\text{had}}^2 \frac{d\sigma(e^+e^- \rightarrow \text{hadrons} + \gamma)}{dM_{\text{hadrons}}^2} = \sigma(e^+e^- \rightarrow \text{hadrons}) H(M_{\text{had}}^2, \cos\theta_{\text{min}})$$

This method is a complementary approach to the standard energy scan.

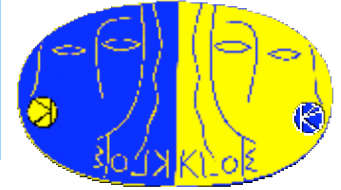
It requires precise calculations of the radiator H .

→ EVA + Phokhara MC Generator

(S. Binner, J.H. Kühn, K. Melnikov, Phys. Lett. B 459, 1999)

(H. Czyz, A. Grzelinska, J.H. Kühn, G. Rodrigo, hep-ph/0308312)

Signal selection:



Pion tracks are measured at angles
 $50^\circ < \theta < 130^\circ$

Photons are required to be within
 $\theta < 15^\circ$ or $\theta > 165^\circ$

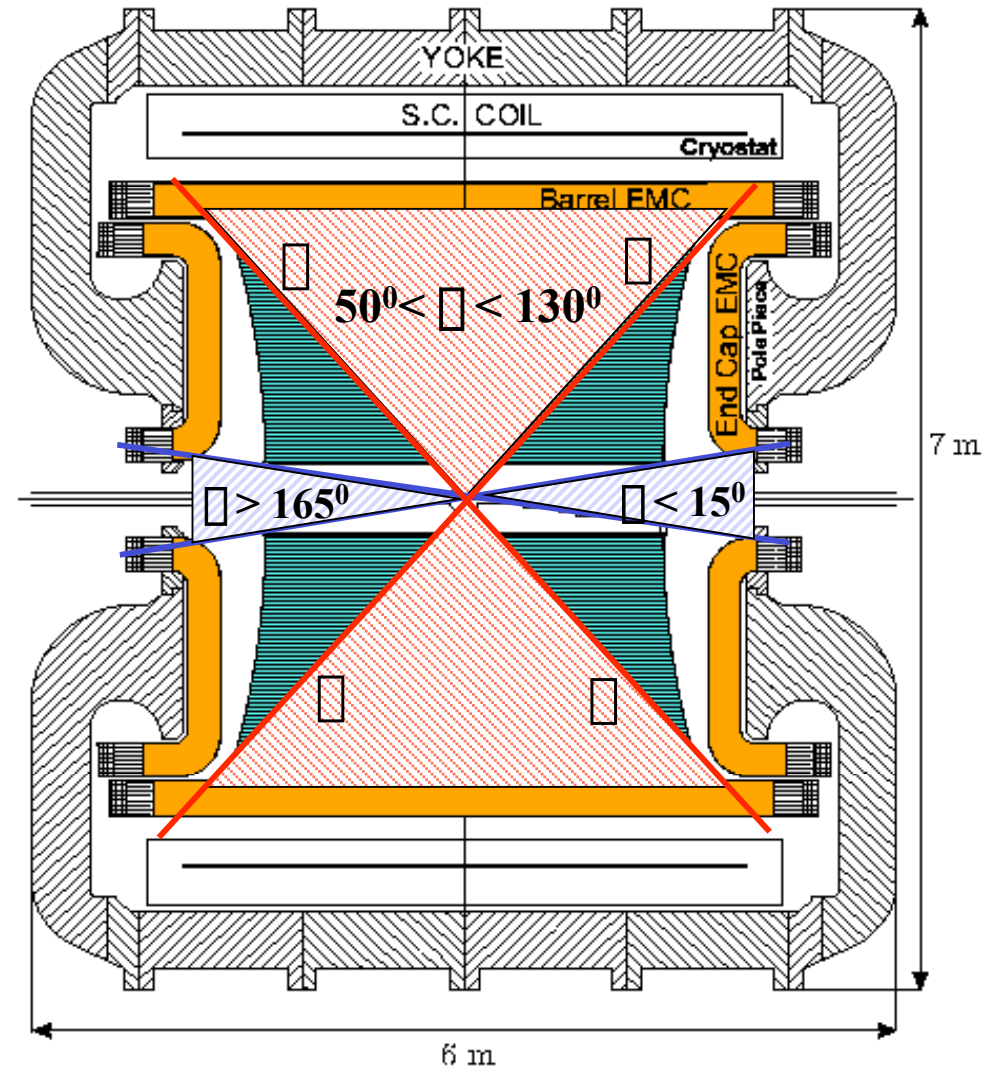
Untagged measurement in which we cut on the direction of the missing momentum

$$\vec{p}_\square = \square \vec{p}_{\text{miss}} = \square (\vec{p}_+ + \vec{p}_\square)$$

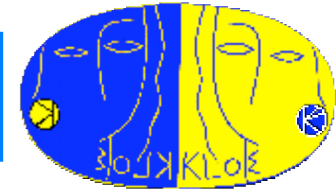
The choice of this kinematical region was motivated by:

- **small relative contribution of FSR**
- **reduced background contamination:**

- $e^+e^- \rightarrow e^+e^- \square$
- $e^+e^- \rightarrow \square + \square \square$
- $e^+e^- \rightarrow \square \square \square + \square \square \square$



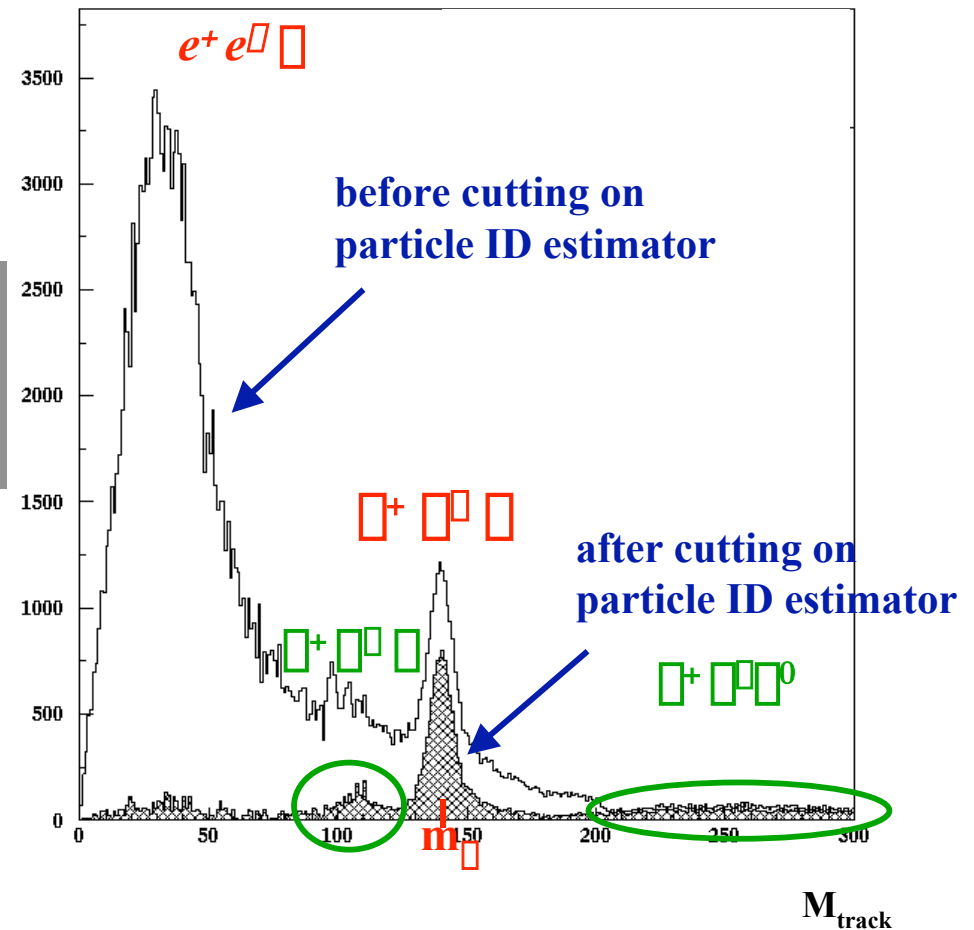
Background rejection:



To reduce Bhabha contamination, a π -e-separation is performed using a particle ID estimator based on:

- TOF of charged clusters in EMC
- Shape and energy deposition of the cluster

The event is selected if one of the charged tracks is identified to be a pion.



Background rejection:

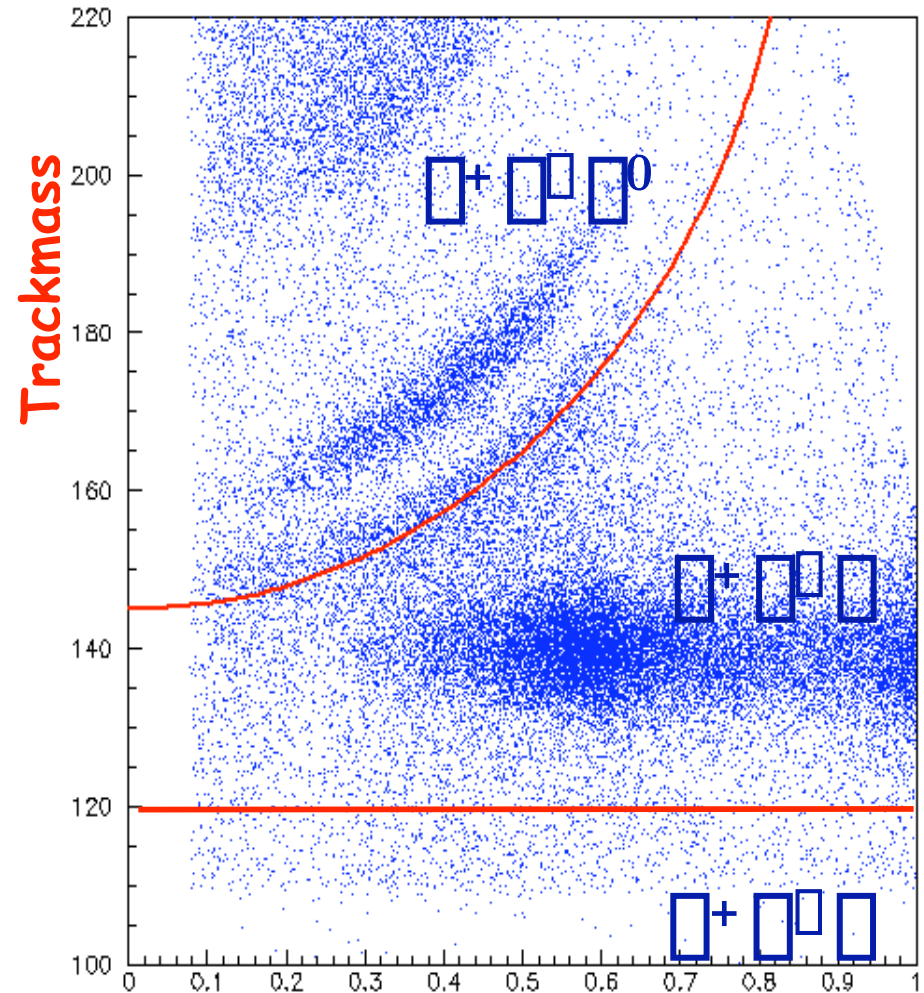


The signal is further selected by performing a cut in the kinetical variable **trackmass** in order to reduce $\pi^+ \pi^0 \pi^0$ **background**

$\pi^+ \pi^0 \pi^0$ background ($M_{\text{track}} \approx 105 \text{ MeV}$) is rejected by a cut on $M_{\text{track}} = 120 \text{ MeV}$

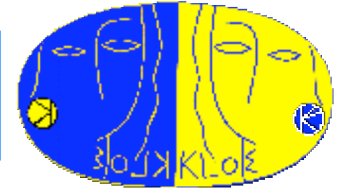
The trackmass is the particle mass for the two tracks obtained by using the 4-momentum-conservation and the assumption that both particles have the same mass M_{trk} :

$$q_{\square}^2 = \left[M_{\square} \sqrt{p_1^2 + M_{\text{trk}}^2} \right]^2 + \left[M_{\square} \sqrt{p_2^2 + M_{\text{trk}}^2} \right]^2 - (\vec{p}_1 + \vec{p}_2)^2 = 0$$



$M_{\pi^0}^2$

$M_{\mu\mu}^2$ - Spectrum:



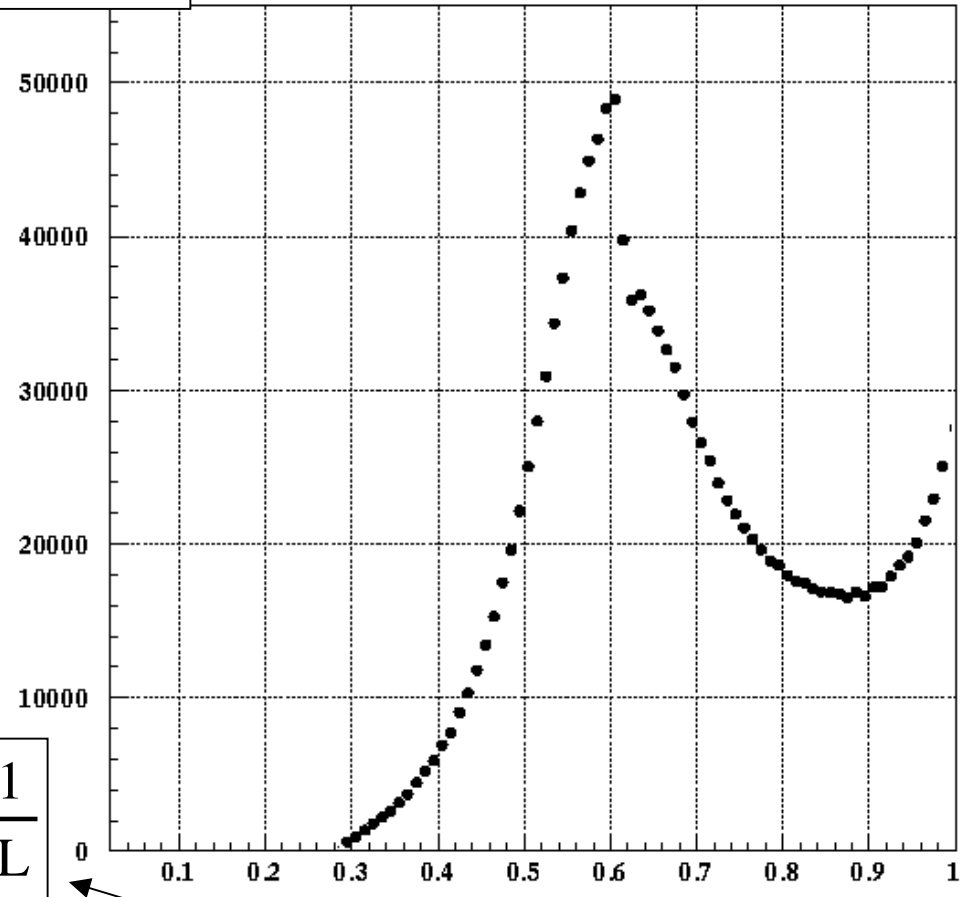
140 pb⁻¹ of 2001 data were analyzed according to the items discussed.

After selection: **1 500 000 events**
(11000 evts/pb⁻¹)

statistical error/bin < 1%
for $M_{\mu\mu}^2 > 0.45 \text{ GeV}^2$

$N_i/0.01\text{GeV}^2$

1548036 events



$$\frac{dN_{\mu\mu}}{dM_{\mu\mu}^2} = \frac{N^{\text{obs}} - N^{\text{bkg}}}{\epsilon M_{\mu\mu}^2} \times \frac{1}{\epsilon_{\text{Select.}}} \times \frac{1}{L}$$

Signal

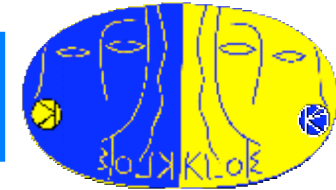
Background

Selection efficiency

Luminosity

Acceptance: $|\eta_{\mu}| < 15^\circ$ ($|\eta_{\mu}| > 165^\circ$), $50^\circ < \phi_{\mu\mu} < 130^\circ$, $E_{\mu} > 10 \text{ MeV}$

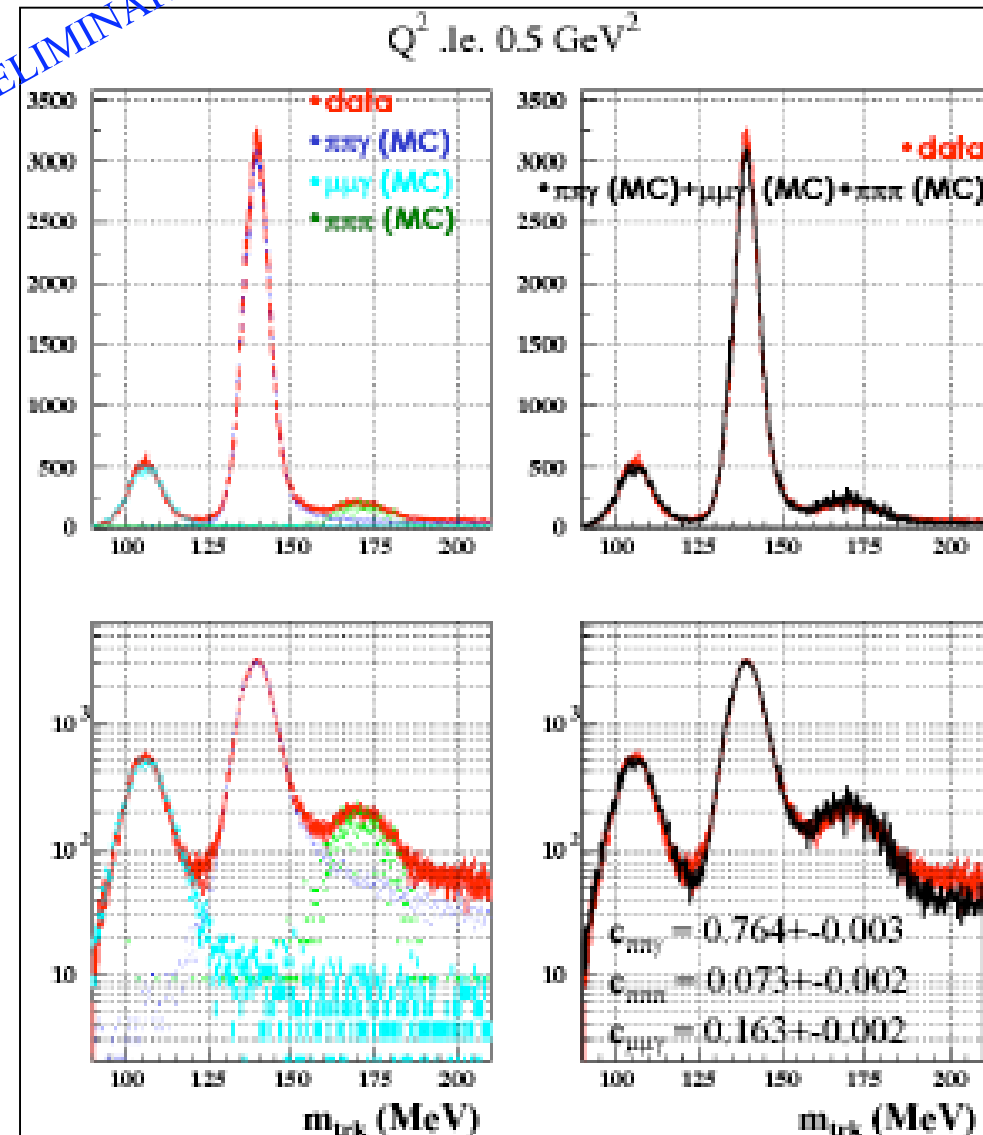
Background subtraction:



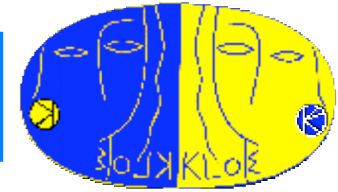
Remaining contaminations from $\pi^+\pi^-\pi^0$ and $\pi\pi\pi$ are measured by fitting the shape of signal and background in the trackmass distribution for different bins of $M_{\pi\pi}^2$.

The estimated number of background events is then subtracted from the spectrum.

PRELIMINARY



Luminosity:

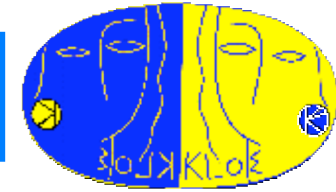


- Luminosity measured with Large Angle Bhabhas: $55^\circ < \theta_e < 135^\circ$
- 2 independent generators used for radiative corrections:
 - BABAYAGA (Pavia group): $\sigma_{\text{eff}} = (428.8 \pm 0.3_{\text{stat}}) \text{ nb}$
 - BHAGENF (Berends modified): $\sigma_{\text{eff}} = (428.5 \pm 0.3_{\text{stat}}) \text{ nb}$

- ◆ Systematics from generator claimed to be 0.5%
- ◆ Experimental systematic error determined by comparing data and MC angular and momentum distributions

Systematics on Luminosity	
Theory	0.5 %
Acceptance	0.3 %
Background (□□□+□□□)	0.1 %
Trigger+Track+Clustering	0.2 %
Knowledge of s run-by-run	0.1 %
TOTAL	0.5 % theory ⊕ 0.4% exp = 0.6 %

Efficiencies:



Trigger

including Cosmic Veto Eff.

Reconstr. Filter

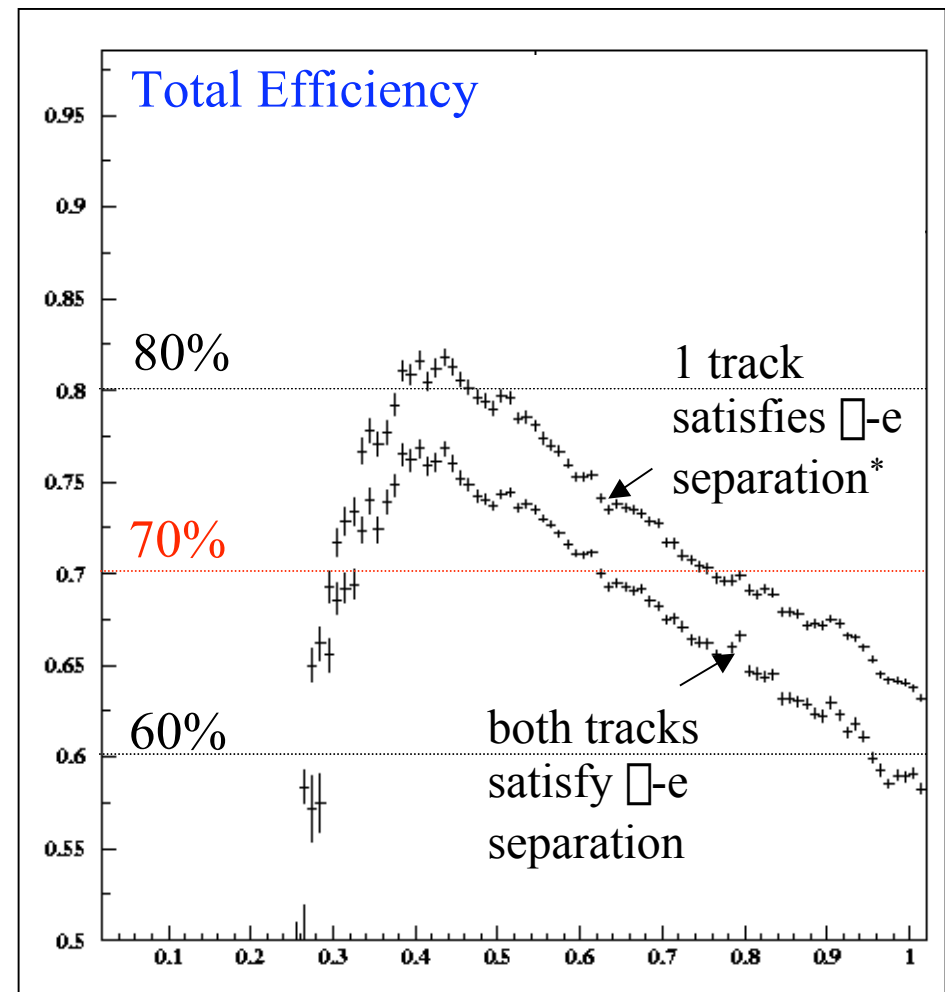
Tracking / Vertex

μ -e separation

Trackmass

blue = estimated from data and/or indep. control samples $\mu^+\mu^-\mu^0$, $\mu^+\mu^-$
Kinematics simulated by MC

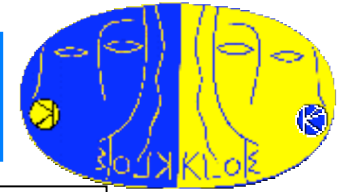
red = estimated from MC and compared with data



$M_{\mu\mu}^2(\text{GeV})$

* used in this analysis

$\sigma(e^+e^- \rightarrow \pi^+\pi^0)$:



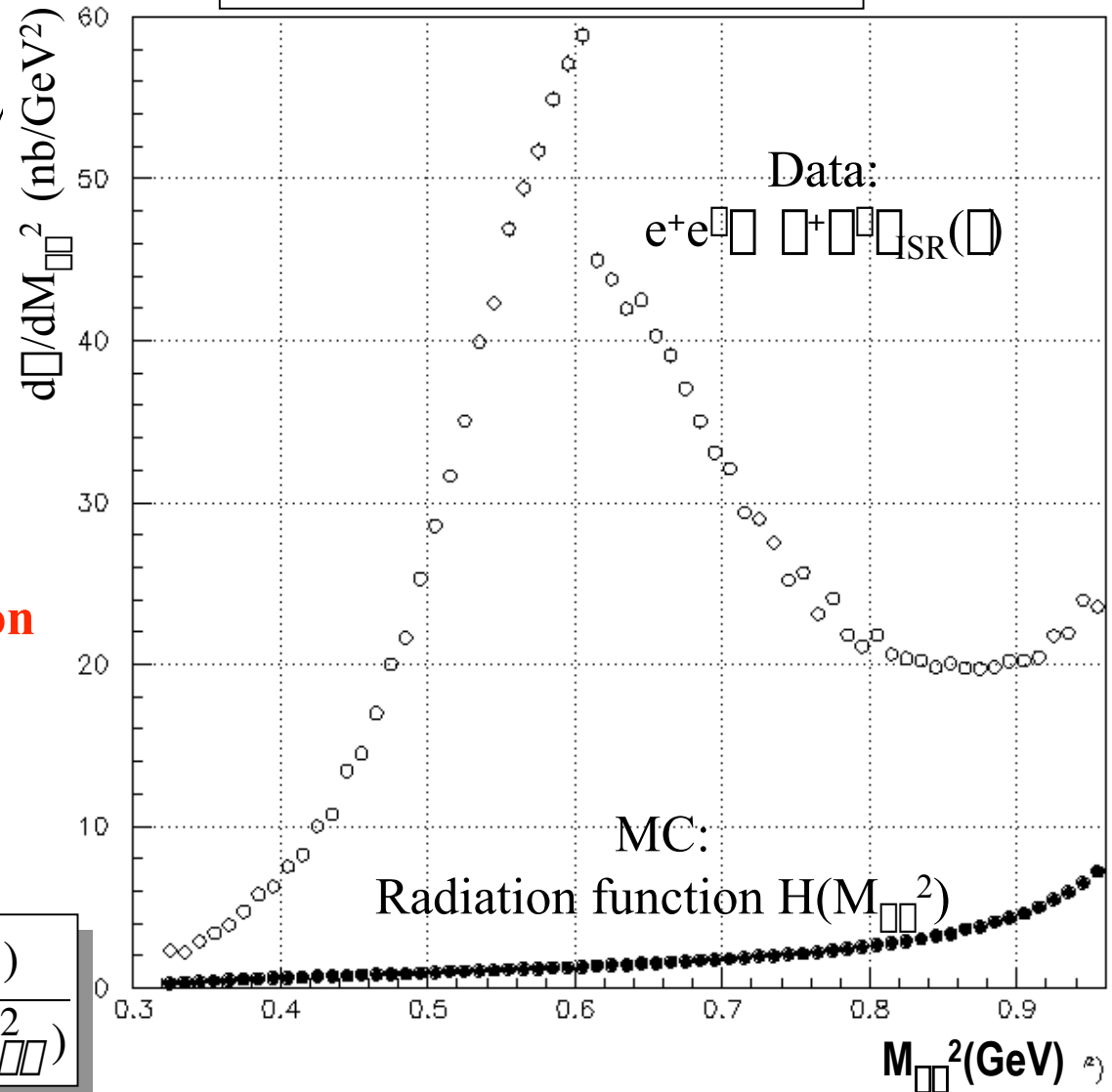
After subtracting the residual background, correcting for efficiencies, dividing for luminosity and unfolding the detector resolution, we arrive to

$\sigma(e^+e^- \rightarrow \pi^+\pi^0)$ in bins of $M_{\pi\pi}^2$

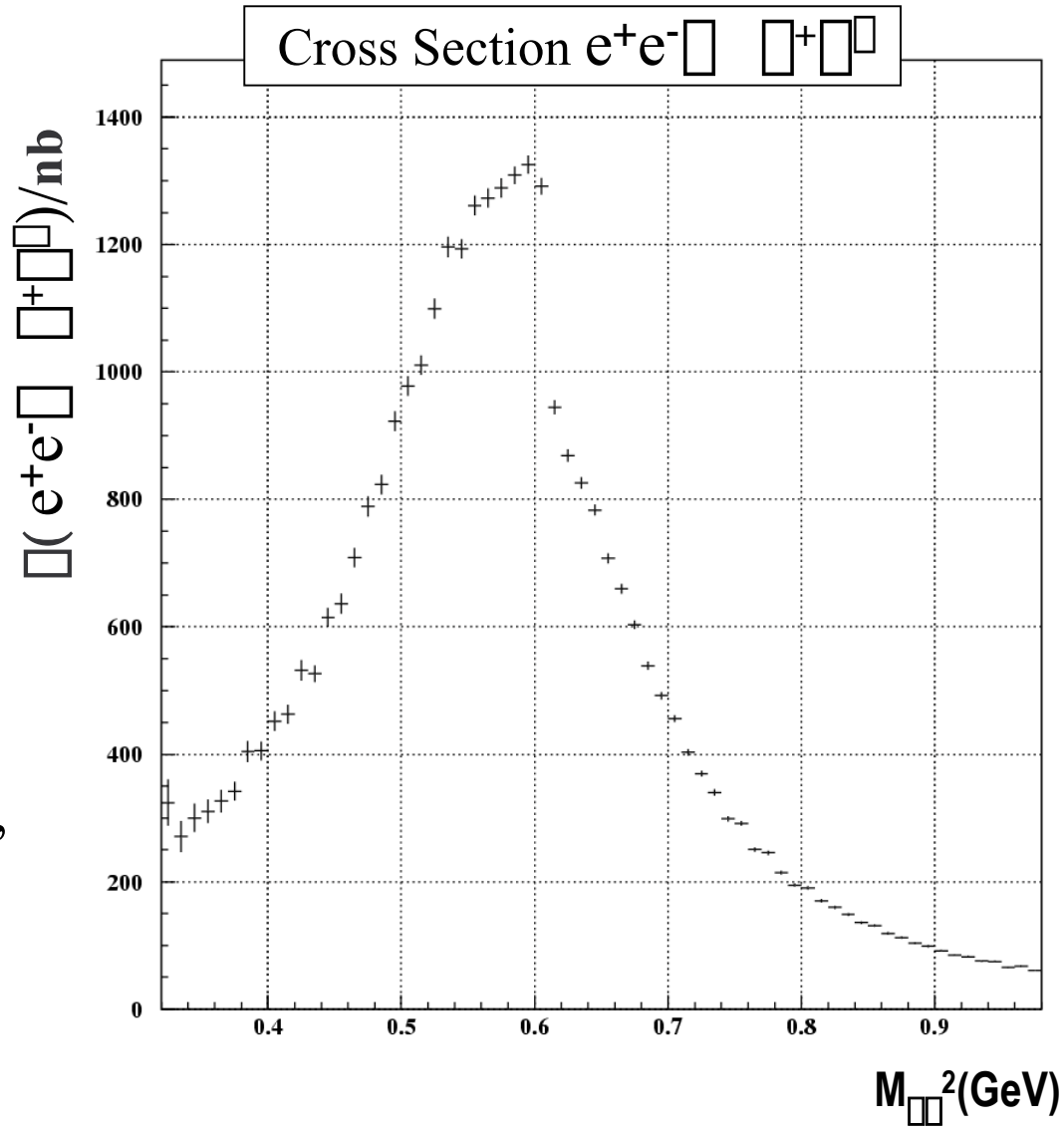
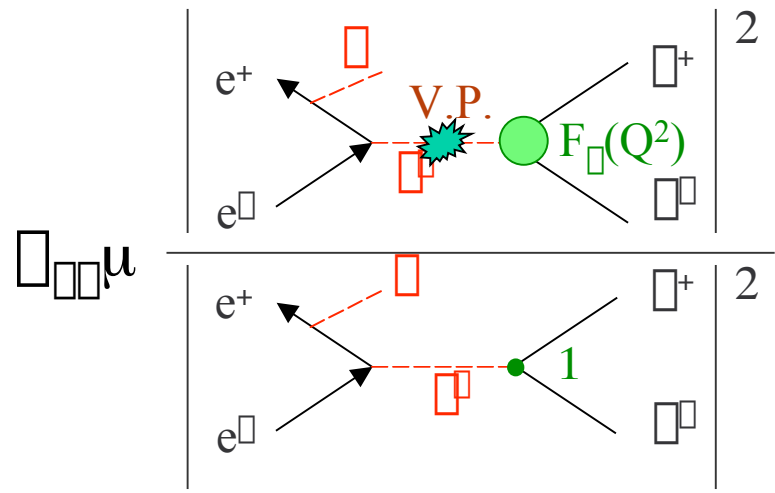
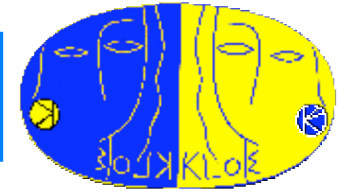
To get the cross section for $e^+e^- \rightarrow \pi^+\pi^-$ we divide the $\pi^+\pi^0$ cross section by the cross section $\pi^+\pi^0$ for “pointlike” pions which is obtained technically from the MC generator by setting $F_\pi = 1$:

$$\sigma_{\pi\pi} |F_\pi(M_{\pi\pi}^2)|^2 = \frac{d\sigma_{\pi\pi\pi^0}(M_{\pi\pi}^2)}{d\sigma_{\pi\pi\pi^0, F_\pi=1}(M_{\pi\pi}^2)}$$

Cross Section $e^+e^- \rightarrow \pi^+\pi^0$

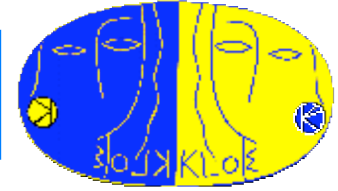


$\sigma(e^+e^- \rightarrow \mu^+\mu^-)$:



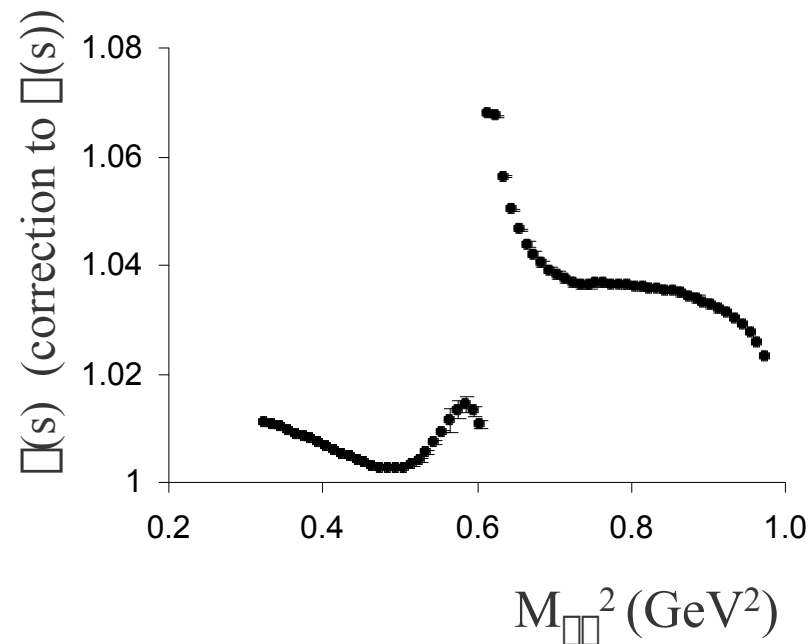
After dividing by the radiation function $H(M_{\mu^2}^2)$, one gets the cross section $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$.

Rad. corrections: Vac. pol.

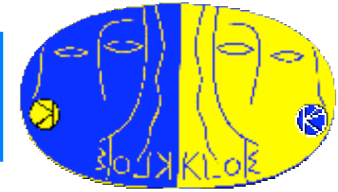


The “bare cross section” has to be used for the evaluation of the hadronic contribution to a_μ in the dispersion integral, i.e. the cross section has to be divided for the running of the fine structure constant α

$$\alpha(s) \quad \alpha^2(s) = \frac{\alpha_0^2}{1 - \Pi_{lep}(s) - \Pi_{had}(s)} \equiv \alpha(s) \cdot \alpha_0^2 \quad \alpha_{bare}(s)$$

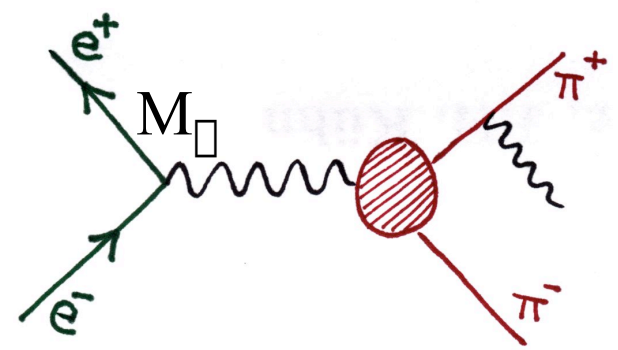


Rad. corrections: FSR (LO)

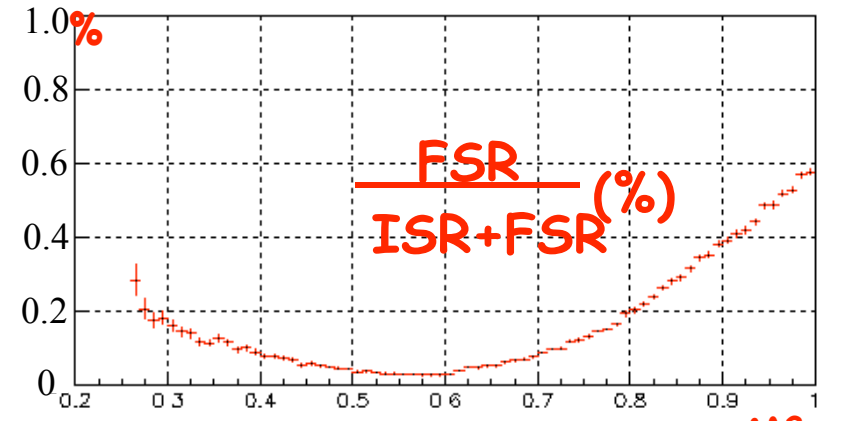
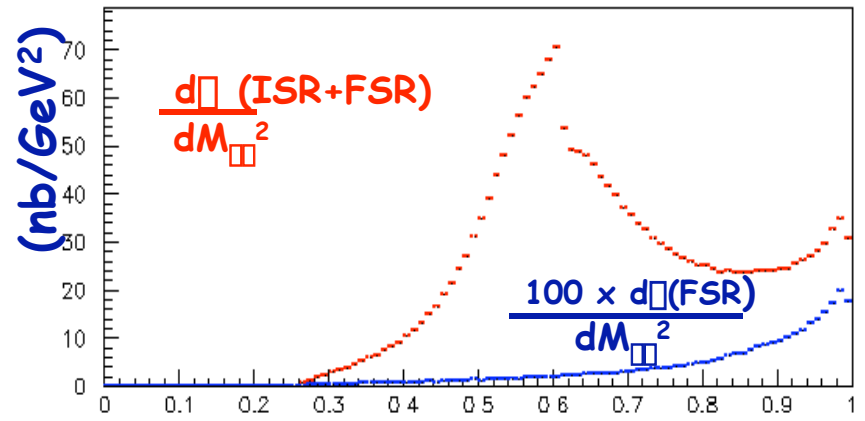


The cross section has to be corrected with respect to **Final State Radiation** (FSR).

At LO final state radiation, there is no initial state radiation and the e^+ and the e^- collide at the energy M_{\square} :



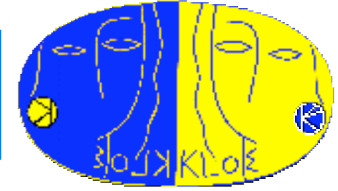
This process has been studied with the **EVA** MC program:



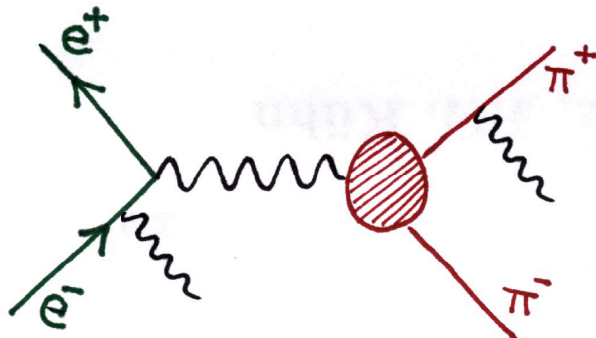
Contribution of **FSR(lo)** <1% for our selection cuts

M_{\square}^2

Rad. corr.: FSR (LO+NLO)

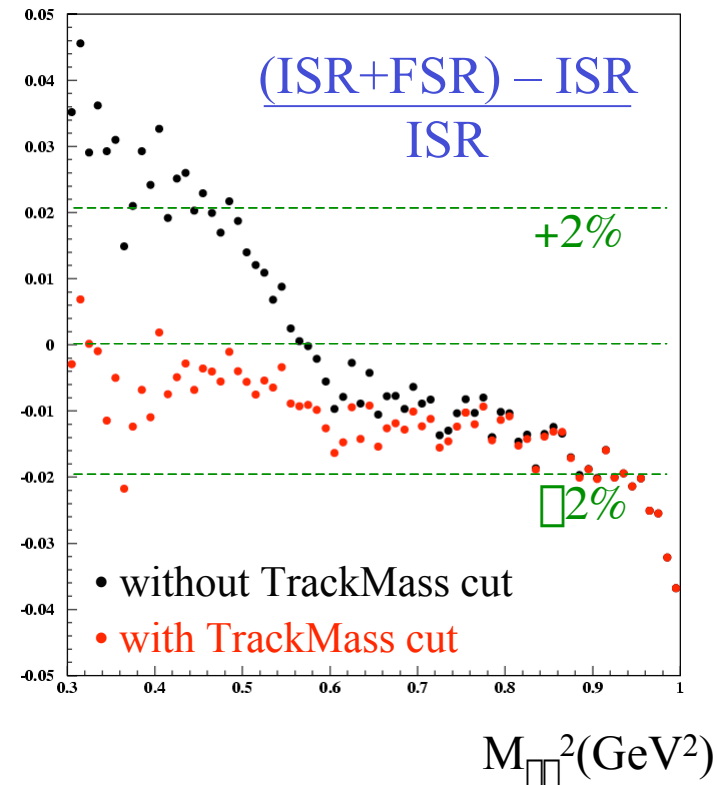


Just recently we got a **new version of Phokhara** which also simulates events with the **presence of 1 ISR- and 1 FSR-photon**:

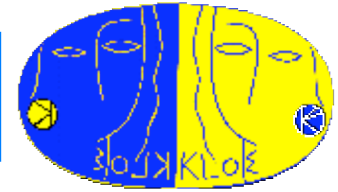


A preliminary check shows that **the FSR contribution is at most 1-2%**.

As of now, we *do not* apply any correction for FSR and add a **contribution of 2% to the systematic error**



Systematic error:



The **systematic error** is under **final evaluation!**
Contributions from:

Theory	
– Radiator Function H	0.5%
– Vacuum Polarization	0.1%
– Luminosity	0.6%
TOTAL	0.8%

Experiment	
– Acceptance	0.3%
– Trigger	0.2%
– Tracking	0.3%
– Vertex	0.8%
– Rec. Filter	0.6%
– Likelihood	0.1%
– Track Mass	0.2%
– BKG subtr.	0.5%
– Unfolding	0.6%
TOTAL	1.4%

FSR	2.0%
□	< 1%

□ 1%

Systematic error can be reduced to • in a short time scale

Conclusions:



- Using **initial state radiation** as a means to measure the cross section for $e^+e^- \rightarrow \mu^+\mu^-$ from threshold to the full collider energy has been proven very effective
- **Analysis almost finished:**
Final checks, especially on **Final State Radiation Correction**
- **Next Steps:**
 - publish final results
 - study events at large photon angles to access lower $M_{\mu\mu}^2$ regions
 - use $\mu\mu\mu$ events for cross checking vacuum polarisation, FSR and additional ISR effects