

1
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QUANTUM MECHANICS
and
NEUTRAL KAON INTERFEROMETRY

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"If there is any place where we have a chance to test the main principles of quantum mechanics in the purest way - does the superposition of amplitudes work or doesn't it? - this is it."

R. Feynman
"Lectures on Physics"

THE NEUTRAL KAON SYSTEM

1. $K^0 - \bar{K}^0$ basis \leftrightarrow Strangeness measurements

$$\left. \begin{array}{l} K^0 \sim d\bar{s} \quad S=+1 \\ \bar{K}^0 \sim s\bar{d} \quad S=-1 \end{array} \right\} \langle K^0 | \bar{K}^0 \rangle = 0$$

H^{weak} non-diag. $\leftrightarrow K^0 \rightleftharpoons \bar{K}^0$ oscillations in t

2. $K_S - K_L$ basis \leftrightarrow lifetime measurements

$$H^{\text{weak}} \text{ diag.} \sim \begin{pmatrix} m_S - \frac{i}{2}\Gamma_S & 0 \\ 0 & m_L - \frac{i}{2}\Gamma_L \end{pmatrix} \leftrightarrow \begin{aligned} |K_{S,L}(t)\rangle &= \\ &= e^{-i\lambda_{S,L}t} |K_{S,L}(0)\rangle \end{aligned}$$

↓

$$U(t) \sim \begin{pmatrix} e^{-\frac{\Gamma_S}{2}t} e^{i\frac{\Delta m}{2}t} & 0 \\ 0 & e^{-\frac{\Gamma_L}{2}t} e^{-i\frac{\Delta m}{2}t} \end{pmatrix}$$

no $K_S - K_L$ oscillations

$$\underline{\Delta m} \equiv m_L - m_S \sim \frac{1}{2}\Gamma_S \gg \underline{\Gamma_L}$$

$$\langle K_S | K_L \rangle = \frac{\epsilon + \epsilon^*}{1 + |\epsilon|^2} \sim 3 \cdot 10^{-3}$$

0 neglect CP

3. Bohr's complementarity

$$|K^0\rangle \approx \frac{1}{\sqrt{2}} |K_S\rangle + \frac{1}{\sqrt{2}} |K_L\rangle$$

$$|+\rangle_x = \frac{1}{\sqrt{2}} |+\rangle_z + \frac{1}{\sqrt{2}} |-\rangle_z$$

$$|\bar{K}^0\rangle \approx \frac{1}{\sqrt{2}} |K_S\rangle - \frac{1}{\sqrt{2}} |K_L\rangle$$

$$|-\rangle_x = \frac{1}{\sqrt{2}} |+\rangle_z - \frac{1}{\sqrt{2}} |-\rangle_z$$

S-measurement
 Γ - } "maximally" incompatible

$K^0 - \bar{K}^0$ OSCILLATIONS

1.- At $t=0$ $\pi p \rightarrow \Lambda K^0$

2.- Transition P_2 at $t \geq 0$

$$\left\{ \begin{aligned} |\langle K^0 | K^0(t) \rangle|^2 &= \frac{1}{4} \left\{ e^{-\Gamma_S t} + e^{-\Gamma_L t} \right\} \left\{ 1 + \underline{V_0(t)} \cos \Delta\Gamma t \right\} \\ |\langle \bar{K}^0 | K^0(t) \rangle|^2 &= \quad \quad \quad \left\{ 1 - \underline{V_0(t)} \cos \Delta\Gamma t \right\} \end{aligned} \right.$$

with $\underline{V_0(t)} = 1/\cosh \frac{1}{2} \Delta\Gamma t$ $\Delta\Gamma \equiv \Gamma_L - \Gamma_S$
 \uparrow "oscillation visibility"

$$\left\{ \begin{aligned} |\langle K_L | K^0(t) \rangle|^2 &= \frac{1}{2} |\langle K_L | K_L(t) \rangle|^2 = \frac{1}{2} e^{-\Gamma_L t} \\ |\langle K_S | K^0(t) \rangle|^2 &= \frac{1}{2} |\langle K_S | K_S(t) \rangle|^2 = \frac{1}{2} e^{-\Gamma_S t} \end{aligned} \right.$$

DOUBLE-SLIT ANALOGY

Particle-like: Trajectories!

"Which Way" \leftrightarrow "Which Width" $K_S - K_L$
 \uparrow \uparrow
 WW

Wave-like: Interference!

"Fringe Visibility" \leftrightarrow "Oscillation Visibility" $K^0 - \bar{K}^0$
 \uparrow

Duality: Acquisition of WW info } mutually
 observed visibility } exclusive

INTERFEROMETRIC DUALITY (I)

1.- B.-G. Englert 1996

$$\left. \begin{array}{l}
 \text{"Fringe visibility" } V_0 \\
 \text{vs} \\
 \text{Path "Predictability" } \mathcal{P} \\
 \text{(a priori WW knowledge)} \\
 \text{* Greenberger + Yasin 1988}
 \end{array} \right\} \text{ in } \mathcal{P} \equiv |w_I - w_{II}| \\
 w_I + w_{II} = 1$$

$$\left. \begin{array}{l}
 \text{Quantitative} \\
 \text{duality/complementarity}
 \end{array} \right\} V_0^2 + \mathcal{P}^2 \leq 1 \\
 \text{(pure states)} \\
 \text{(our case)}$$

2.- Strangeness oscillations (int)
for (undecayed) neutral kaons
 $K^0(0) \rightarrow K^0(t)$

$$\mathcal{P} = |w_S - w_L| = \left| \frac{e^{-\Gamma_S t} - e^{-\Gamma_L t}}{e^{-\Gamma_S t} + e^{-\Gamma_L t}} \right| = \tanh \left[\frac{1}{2} \overbrace{\Delta\Gamma}^{\Gamma_L - \Gamma_S} t \right] = \mathcal{P}(t)$$

$$V_0 = \dots = \frac{2}{\sqrt{2 + e^{\Delta\Gamma t} + e^{-\Delta\Gamma t}}} = \frac{1}{\cosh \left[\frac{1}{2} \Delta\Gamma t \right]} = V_0(t)$$

Thus, $V_0(t)^2 + \mathcal{P}(t)^2 = 1 \quad \forall t$ t-dependence!

3.- CLEAR data reinterpreted

Phys. Lett. B503(01)49
" " B444(98)38

CPLEAR DATA

 (Single kaon)

Strangeness or K^0 - \bar{K}^0 oscillations in time t
 P's $| \langle K^0 | K^0(t) \rangle |^2, | \langle \bar{K}^0 | K^0(t) \rangle |^2, \dots$

1. CPLEAR (semileptonic decays of K^0 or \bar{K}^0)

Phys. Lett. B444 (1998) 38

$$A_{\Delta m}^{\text{weak}} \left(\frac{t}{\tau_S} \right) = \underbrace{\frac{1}{\cosh \frac{\Delta \Gamma}{2} t}}_{V_0(t)} \cdot \cos \Delta m t$$

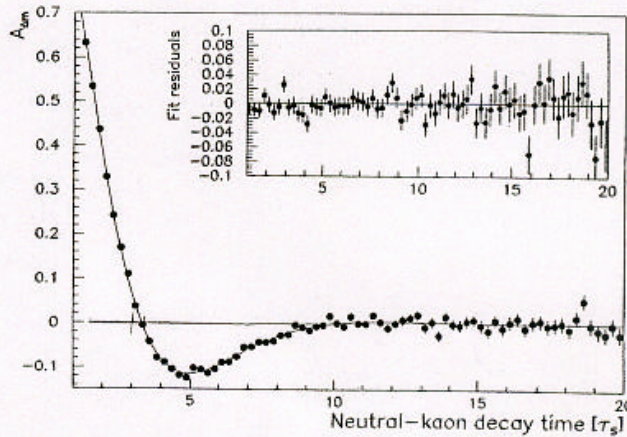


Fig. 2. The asymmetry $A_{\Delta m}$ versus the neutral-kaon decay time (in unit of τ_S). The solid line represents the result of the fit.

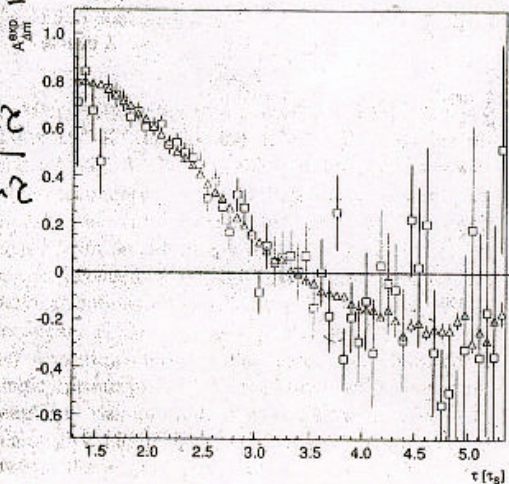
2. CPLEAR (strong K^0, \bar{K}^0 interactions on nuclei: $K^0 p \rightarrow K^+ n$ or $\bar{K}^0 n \rightarrow \Lambda \pi^0$)

Phys. Lett. B503 (2001) 49

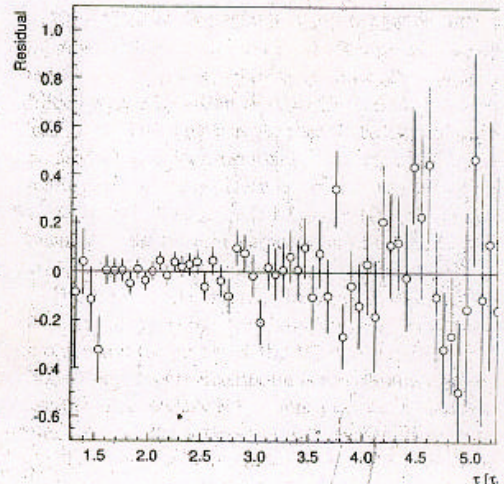
$\rightarrow p\pi^-$

$$A_{\Delta m}^{\text{strong}} \left(\frac{t}{\tau_S} \right) = \frac{2 V_0(t) \cos \Delta m t}{1 + V_0(t)^2 \cos^2 \Delta m t}$$

CPLEAR Collaboration / Physics Letters B 503 (2001) 49-57



(a)



(b)

Fig. 5. (a) Best fit to the data points $A_{\Delta m}^{\text{exp}}(\tau)$ (squares) with the simulated asymmetries (triangles) for the $(K^- + \Lambda)$ sample, see text. Distribution of the fit residuals.

A CONJECTURE

If $I(z) \propto \left\{ 1 + V_0(z) \cos[bz] \right\}$ then
 $V_0(z)^2 + P(z)^2 = 1$ is satisfied with

$$V_0(z) = \frac{1}{\cosh[az]} \quad P(z) = \tanh[az]$$

and

$$\# \text{ of fringes/oscillations} \sim 0.26 \frac{b}{a}$$

1.- Neutrons
(CPLEAR data)

$$\frac{b}{a} = \frac{\Delta m}{\frac{1}{2}\Delta E} \sim 1$$

2.- Wheeler + Bartell (80)
(double-slit + lens)

$\frac{b}{a}$ introduced

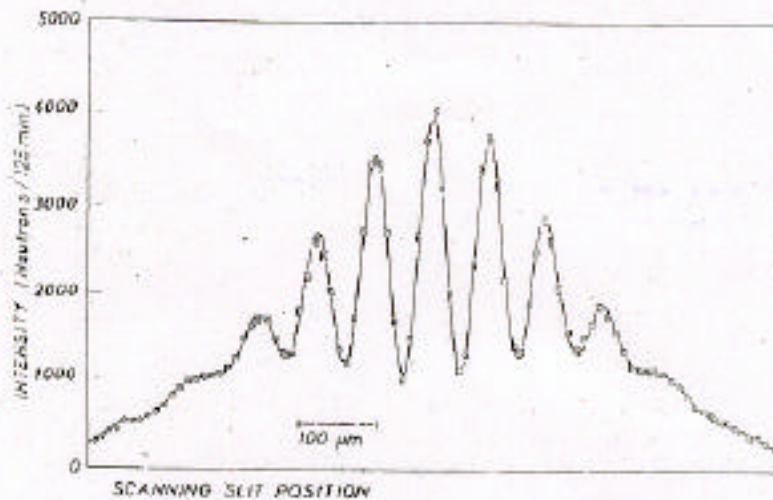
3.- Zeilinger et al (1991)
(double-slit for neutrons)

$$\frac{b}{a} \sim 10$$

4.- Mott: ${}_6\text{C}^{12} + {}_6\text{C}^{12} \rightarrow {}_6\text{C}^{12} + {}_6\text{C}^{12}$
(PRL 4 (1960) 365 data)
at $5 \text{ MeV} = E_{\text{cm}}$

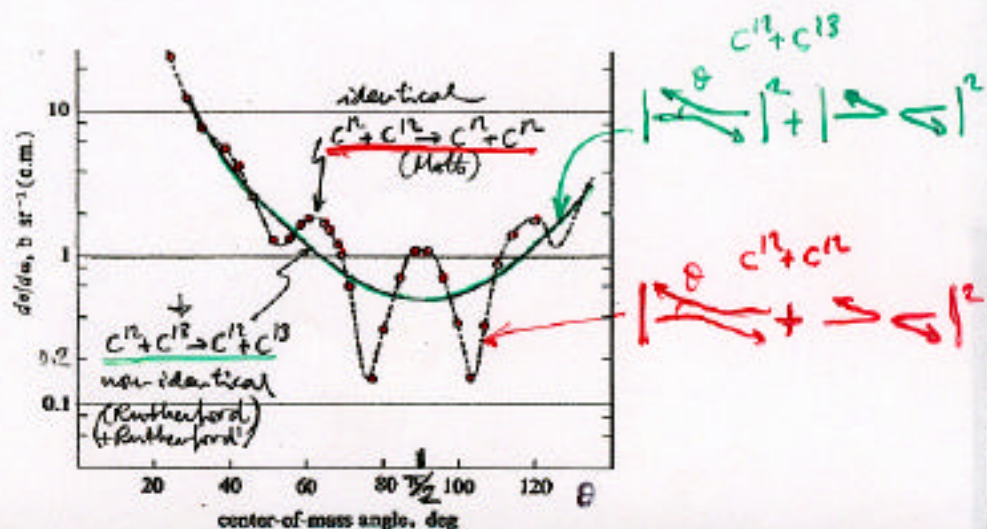
$$\frac{b}{a} = \frac{z_1 z_2 \alpha c}{v_{\text{cm}}} \sim 9$$

Neutron data in a double-slit.



$C^{12} + C^{12}$ elastic scattering (Mott). Data at 5 MeV = E_{CM}

$$\begin{aligned} \frac{d\sigma}{d\Omega} \Big|_{CM}^{Mott} &\propto \left(\frac{1}{\sin^2 \frac{\theta}{2}} + \frac{1}{\cos^2 \frac{\theta}{2}} + \frac{2}{\sin \frac{\theta}{2} \cos \frac{\theta}{2}} \cos \left[b \ln \tan^2 \frac{\theta}{2} \right] \right) \quad b = \frac{Z_1 Z_2 \alpha c}{v_{CM}} \\ &= \left(1 + \frac{1 - \cos \theta}{1 + \cos \theta} \cos \left[b \ln \tan^2 \frac{\theta}{2} \right] \right) \\ &= \left(1 + \frac{1}{\cosh x} \cos [bx] \right) \quad \text{where } e^{\pm x} = \frac{1 + \cos \theta}{1 - \cos \theta} \end{aligned}$$



INTERFEROMETRIC DUALITY II

1. - B-G. Englert 1996

"Fringe visibility" V [1 + V cos φ]
 vs
 "Path or WW knowledge" $K_W > P$
 (after a W-measurement, W either S or Γ)

Quantitative duality/complementarity $V(t)^2 + K_W(t)^2 \leq 1$
↑
pure state

2. - $\varphi \rightarrow K\bar{K}$ (left/right kasn = object/meter kasn)

$t=0 \quad |\phi(0)\rangle = \frac{1}{\sqrt{2}} K^0 \bar{K}^0 - \frac{1}{\sqrt{2}} \bar{K}^0 K^0 = \frac{1}{\sqrt{2}} K_L K_S - \frac{1}{\sqrt{2}} K_S K_L$

$\Delta t \equiv t_e - t_r$
 $|\phi(\Delta t)\rangle = \frac{1}{\sqrt{2} \sqrt{1 + e^{i\Delta\Gamma\Delta t}}} \left\{ \begin{array}{l} K^0 K_S - \bar{K}^0 K_L - \\ - e^{i\Delta m \Delta t} e^{i\frac{1}{2}\Delta\Gamma\Delta t} (K^0 K_L + \bar{K}^0 K_S) \end{array} \right\}$

S(left)
 S(right)

i) W-measurement:
 $\Gamma_{S,L}(\text{right}) \rightarrow K_r = 1, V_r = 0.$

$|\phi(\Delta t)\rangle = \frac{1}{\sqrt{2} \sqrt{1 + e^{i\Delta\Gamma\Delta t}}} \left\{ \begin{array}{l} (1 - e^{i\Delta m \Delta t} e^{i\frac{1}{2}\Delta\Gamma\Delta t}) (K^0 K^0 - \bar{K}^0 \bar{K}^0) + \\ + (1 + e^{i\Delta m \Delta t} e^{i\frac{1}{2}\Delta\Gamma\Delta t}) (K^0 \bar{K}^0 - \bar{K}^0 K^0) \end{array} \right\}$

ii) W-measurement:

$S(\text{right}) \rightarrow K_S(\Delta t) = \tanh | \frac{1}{2} \Delta\Gamma\Delta t |$
 $V_S(\Delta t) = [\cosh(\frac{1}{2} \Delta\Gamma\Delta t)]^{-1}$

S in both sides
QUANTUM ERASER
 (Scully + Drühl '82)

3. - Data

CPLEAR: $\Delta t = 0$
 $\Delta t = 1.2 \tau_S$
 PLB 422 (1999) 339

$\Delta\phi_{NE} : \dots$