

# EQUIVALENT PHOTONS METHOD and EIKONAL APPROXIMATION in the THEORY of TRANSITION RADIATION

N. Shul'ga<sup>1</sup>, V.Syshchenko<sup>2</sup>, S. Shul'ga<sup>1</sup>

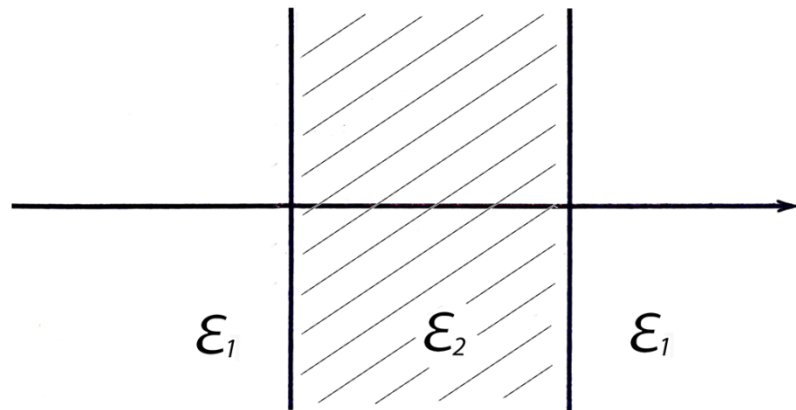
<sup>1</sup> Akhiezer Institute for Theoretical Physics of National Science Center  
“Kharkov Institute of Physics and Technology”  
Kharkov, Ukraine

<sup>2</sup> Belgorod State University, Belgorod, Russia

e-mail: [shulga@kipt.kharkov.ua](mailto:shulga@kipt.kharkov.ua)

## Transition Radiation for Sharp and Smooth Function $\epsilon(\mathbf{r})$

$\epsilon(x)$  – sharp function of  $x$



Joining of solutions for  $\epsilon_1$  and  $\epsilon_2$

Problem

$$\frac{\omega_p^2}{\omega} l \geq 1 \quad \text{for } \epsilon(\mathbf{r})$$

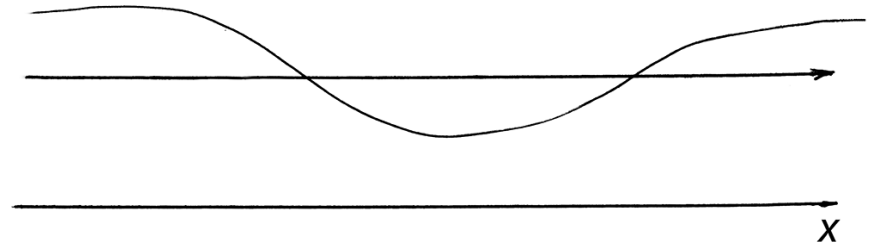
$\epsilon(\mathbf{r})$  – smooth function of  $\mathbf{r}$

1.  $|1 - \epsilon(\mathbf{r})| \ll 1$  approximation
2. Semiclassical approximation for  $\epsilon(x)$

*M.Ter – Mikaelian*(1980)

*L.Durand* (1975)

*J.Lepore, R.Riddell*(1976)



## Simplification of the Problem for $\varepsilon(\mathbf{r})$

1. Equivalent photons method (the field of relativistic electrons  $\rightarrow$  wave packet of free waves  $\rightarrow$  evolution of this wave packet with time)
2. Eikonal approximation  
(the changing of free waves packet due to  $\varepsilon(\mathbf{r}) \neq 1$  )

## Spectral-Angular Density of Transition Radiation

$$\left( \Delta + \omega^2 \varepsilon_\omega(\mathbf{r}) \right) \mathbf{E}_\omega = \text{grad div} \mathbf{E}_\omega - 4\pi e i \omega \frac{\mathbf{v}}{v} \delta(\dots) e^{i\omega z/v}$$

$$\frac{dE}{d\omega d\Omega} = \frac{\omega^2}{(8\pi^2)^2} |\mathbf{k} \times \mathbf{I}|^2$$

$$\mathbf{I} = \int d^3r e^{-i\mathbf{k}\mathbf{r}} (1 - \varepsilon(\mathbf{r})) \mathbf{E}_\omega(\mathbf{r})$$

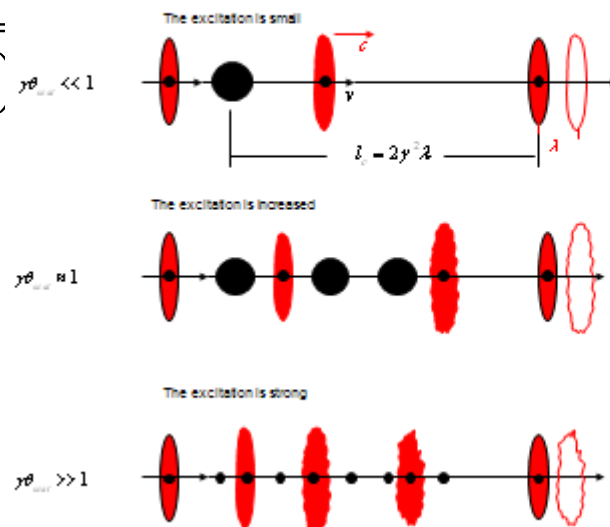
It is necessary to know  $E(\mathbf{r})$  in the region where  $\varepsilon_\omega(\mathbf{r}) \neq 1$

## Proper Field of Relativistic Electrons

$$\mathbf{E}_c(\mathbf{r}, t) = -e \left( \nabla_{\perp} + \frac{n}{\gamma^2} \frac{\partial}{\partial z} \right) \varphi_c(\mathbf{r}, t)$$

How Does Electron Radiate?

$$\varphi_c(\mathbf{r}, t) = \frac{1}{\sqrt{(z - vt)^2 + \gamma^2 \rho^2}} e^{-ik_{\perp} \rho} e^{-ik_z(z - vt)}$$



$$\frac{4\pi}{k_{\perp}^2 + k_z^2 / \gamma^2} e^{-ik_{\perp} \rho} e^{-ik_z(z - vt)}$$

$$\nu_{\text{rel}} = c - v$$

$$\lambda = (c - v) \Delta t \rightarrow \Delta t \approx 2\gamma^2 \lambda$$

$$l_c = 2\gamma^2 \lambda$$

$$\varepsilon = 5 \text{ GeV}$$

$$l_c \sim 10^{-3} \text{ cm for } \omega \sim 1 \text{ MeV}$$

$$l_c \sim 1 \text{ cm for } \omega \sim 1 \text{ KeV}$$

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$$c_k = \frac{4\pi\gamma}{k_{\perp}} e^{-i\mathbf{k}_{\perp} \cdot \mathbf{r}_{\perp}} \int_0^{\infty} dz \cos k_z z e^{-\gamma \rho} =$$

$$= 4\pi e^{-ik_{\perp} \rho} \int_0^{\infty} \rho d\rho J_0(k_{\perp} \rho) K_0(k_z \rho / \gamma)$$

$$z_{eff} \sim \lambda$$

$$\rho_{eff} \sim \gamma \lambda$$

## Equivalent Photons Method

$$\varphi(\mathbf{r}, t) = \frac{1}{\pi^2} \text{Re} \int d^3k e^{i(\mathbf{k}\mathbf{r} - kt)} \frac{\theta(k_z)}{k_{\perp}^2 + k_z^2/\gamma^2} \quad c=1$$

$$\text{for } t=0 \quad \varphi(\mathbf{r}, 0) = \varphi_c(\mathbf{r}, 0)$$

this packet spreads for  $t > 0$ , but for  $\frac{kt}{\gamma^2} \ll 1$  the spreading is small!!!

$$\begin{cases} \varphi(\mathbf{r}, t) = \frac{e}{\sqrt{(z-t)^2 + \rho^2/\gamma^2}} = \int \frac{d^3k}{(2\pi)^3} e^{i\mathbf{k}\mathbf{r}} c_k, \\ c_k = e^{ik_z t} \frac{4\pi}{k_{\perp}^2 + k_z^2/\gamma^2} \end{cases} \quad t \ll \frac{\gamma^2}{k}$$

## Eikonal Approximation

The changing of free wave packet  $\mathbf{E}_\omega(\mathbf{r})$  is determined by

$$\left(\Delta + \varepsilon_\omega(r)\omega^2\right)\mathbf{E}_\omega = \nabla(\text{div}\mathbf{E}_\omega)$$

$$\mathbf{E}_\omega = e^{i\omega z} \Phi(\mathbf{r})$$

The second derivatives of  $\Phi$  are neglected

$$i\frac{\partial}{\partial z}\Phi_\perp = \frac{\omega}{2}(1 - \varepsilon_\omega(\mathbf{r}))\Phi_\perp$$

$$E_\omega^{(eik)}(\mathbf{r}) \approx E_\omega^0(\mathbf{r}) \exp\left\{-i\frac{\omega}{2}\int_{-\infty}^z dz' (1 - \varepsilon_\omega(\mathbf{r}))\right\}$$

$$E_\omega^0(\mathbf{r}) \approx \frac{2e\omega}{v^2\gamma} e^{i\omega z} \frac{\rho}{\rho} K_1(\omega\rho/v\gamma)$$

## Spectral Angular Density of Transition Radiation

$$\frac{dE}{d\omega d\Omega} \approx \frac{\omega^4}{(8\pi)^2} \left| \mathbf{I}_{\perp}^{(eik)} \right|^2$$

$$\mathbf{I}_{\perp}^{(eik)} = i \frac{4e}{v^2 \gamma} \int d^2 \rho \, e^{-i \mathbf{k}_{\perp} \rho} \frac{\boldsymbol{\rho}}{\rho} K_1 \left( \frac{\omega \rho}{v \gamma} \right) \left\{ e^{-i \frac{\omega}{2} \int_{-\infty}^{\infty} dz (1 - \epsilon_{\omega}(\mathbf{r}))} - 1 \right\}$$

*Validity conditions*

$$\gamma^{-1} \ll \frac{\omega_{\rho}}{\omega} \ll 1, \quad \theta \ll \frac{\omega_{\rho}}{\omega}$$

$$\omega \theta l \ll 1$$



# Conclusions

1. Description method for forward transition radiation for smooth function  $\varepsilon_{\omega}(\mathbf{r})$ .
2. Spreading of wave photon packet (equivalent photons method for proper field of relativistic electron)  
the wave packet doesn't spread during the time  $\Delta t < \gamma^2 / \omega$
3. Eikonal approximation for forward transition radiation (validity conditions)
4. Applications
  - thin plate
  - thin dielectric fiber