Craig Dukes	Daphne 2004	University of Virginia
Search for	CP Violation in Hyperon Decays with	the HyperCP
	Spectrometer at Fermilab	
	Craig Dukes	
	University of Virginia	
	for the HyperCP collaboration	
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## Why Search for *CP* Violation in Hyperon Decays?

- After 40 years of intense experimental effort and many beautiful experiments we still know little about *CP* violation: the origin of *CP* violation remains unknown and there is little hard evidence that it is explained by the Standard Model.
- The asymmetry can be relatively large: up to  $1 \times 10^{-2}$ .
- The price is modest:
  - No new accelerators needed.
  - Apparatus is modest in scope and cost.
- Hyperons are sensitive to sources of CP violation that, for example, kaons are not.
- *CP* violation is too important, and experimental evidence is too meagre, not to examine every possible manifestation of the effect.

"We are willing to stake our reputation on the prediction that dedicated and comprehensive studies of CP violation will reveal the presence of New Physics."

Bigi and Sanda, CP Violation

## How to Search for *CP* Violation in $\Lambda$ Decays

Due to parity violation the proton likes to go in the direction of the  $\Lambda$  spin:

$$\Lambda \to p\pi^{-}: \qquad \frac{dN(p)}{d\cos\theta} = \frac{N_0}{2}(1 + \alpha_{\Lambda}P_{\Lambda}\cos\theta) \qquad \alpha = \frac{2Re(S^*P)}{|S|^2 + |P|^2}$$

Under *CP* the antiproton likes to go in the direction opposite to the  $\overline{\Lambda}$  spin:



## **Problem: Producing** $\Lambda$ 's of Known Polarization

 $\Lambda/\overline{\Lambda}$ 's of known polarization can be produced through the decay of **unpolarized**  $\Xi^{-}/\overline{\Xi}^{+}$ 's:

 $\Xi^- \to \Lambda \pi^ \overline{\Xi}^+ \to \overline{\Lambda} \pi^+$ 

If the  $\Xi$  is produced unpolarized — which can simply be done by targeting at 0 degrees — then the  $\Lambda$  is found in a helicity state:

 $\vec{P}_{\Lambda} = \alpha_{\Xi} \hat{p}_{\Lambda}$  $\vec{P}_{\overline{\Lambda}} = \overline{\alpha}_{\Xi} \hat{p}_{\overline{\Lambda}}$  $\frac{d\mathrm{N}(\bar{\mathrm{p}})}{d\cos\theta} = \frac{\mathrm{N}_0}{2} (1 + \overline{\alpha}_{\Lambda} \overline{\alpha}_{\Xi} \cos\theta)$  $\frac{dN(p)}{d\cos\theta} = \frac{N_0}{2}(1 + \alpha_\Lambda \alpha_\Xi \cos\theta)$ If CP is good, the slopes of the  $\Xi^- \to \Lambda \pi^- \to p \pi^- \pi^ \overline{\Xi}^+ \to \overline{\Lambda} \pi^+ \to \overline{p} \pi^+ \pi^+$ proton and antiproton  $\cos \theta$  disslope =  $\overline{\alpha}_{\Lambda} \overline{\alpha}_{\Xi}$ slope =  $\alpha_{\Lambda} \alpha_{\Xi}$ tributions are identical, and: dN dN dcosθ dcosθ  $\alpha \Xi \alpha_{\Lambda} = \overline{\alpha} \Xi \overline{\alpha}_{\Lambda}$ -1 0 +10 +1-1 cosθ cosθ

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### We are sensitive to both $\Xi$ and $\Lambda$ *CP* violation

$$A_{\Xi\Lambda} = \frac{\alpha_{\Xi}\alpha_{\Lambda} - \overline{\alpha}_{\Xi}\overline{\alpha}_{\Lambda}}{\alpha_{\Xi}\alpha_{\Lambda} + \overline{\alpha}_{\Xi}\overline{\alpha}_{\Lambda}} \approx A_{\Xi} + A_{\Lambda}$$

where:

What we experimentally measure is the slope of the proton (antiproton)  $\cos(\theta)$  distribution in the special  $\Lambda$  rest frame where the  $\Lambda$  momentum direction in the  $\Xi$  rest frame defines the polar axis: the Lambda Helicity Frame.

 $A_{\Xi} = \frac{\alpha_{\Xi} + \overline{\alpha}_{\Xi}}{\alpha_{\Xi} - \overline{\alpha}_{\Xi}}$ 



## Phenomenology of CP Violation in $\Xi$ and $\Lambda$ Decay

- CP violation in  $\Xi$  and  $\Lambda$  decays is manifestly direct with  $\Delta S = 1$ .
- Three ingredients are needed to get a non-zero asymmetry:
  - 1. At least two channels in the final state: the S-and P-wave amplitudes.
  - 2. The CP violating weak phases must be different in the two channels.
  - 3. There must be unequal final-state scattering phase shifts in the two channels.

$$A_{\Lambda} = (\alpha_{\Lambda} + \alpha_{\overline{\Lambda}})/(\alpha_{\Lambda} - \alpha_{\overline{\Lambda}}) \cong -\tan(\delta_P - \delta_S)\sin(\phi_P - \phi_S),$$
  
$$A_{\Xi} = (\alpha_{\Xi} + \alpha_{\overline{\Xi}})/(\alpha_{\Xi} - \alpha_{\overline{\Xi}}) \cong -\tan(\delta_P - \delta_S)\sin(\phi_P - \phi_S),$$
  
strong phases  $(\phi_P - \phi_S)$ ,  
weak phases

- Asymmetry greatly reduced by the small strong phase shifts.
  - The  $p\pi$  phase shifts have been measured to a precision of about one degree:

$$\Lambda \begin{cases} \delta_P = -1.1 \pm 1.0^{\circ} \\ \delta_S = 6.0 \pm 1.0^{\circ} \end{cases}$$

• The  $\Lambda \pi$  phase shifts can't be measured, and theoretical predictions disagree.

$$\Xi^{-} \left\{ \begin{array}{l} \delta_{P} = -2.7^{\circ} \\ \delta_{S} = -18.7^{\circ} \end{array} \right\} 1965 \qquad \begin{array}{l} = -1^{\circ} \\ = 0^{\circ} \end{array} \right\} \operatorname{recent} \chi PT$$

HyperCP has measured the  $\Lambda \pi$  phase shift:  $(4.6 \pm 1.4 \pm 1.8)^{\circ}$ 

### What do we Expect: Theoretical Predictions

• Standard Model predictions range from about  $10^{-4} - 10^{-5}$ . Beware

"Given our crude estimate of the hadronic matrix elements involved, all our numerical results should be viewed with caution."

He and Valencia, PRD52 (1995) 5257.

• Beyond-the-standard-model predictions larger, and not well contrained by kaon *CP* measurements: hyperon *CP* violation probes both parity conserving and parity violating amplitudes.

Recent preprint by Tandean (hep-ph/0311036) shows that the upper bound on  $A_{\Xi\Lambda}$  is  $\sim 100 \times 10^{-4}$ .

• For example, some supersymmetric models that do not generate  $\epsilon'/\epsilon$  can lead to  $A_{\Lambda}$  of  $O(10^{-3})$ .



## What is the experimental situation?

- There are no limits on  $A_{\Xi}$ .
- $A_{\Lambda}$  has been measured to  $2 \times 10^{-2}$ .

Mode	Method	Limit	Exp	Date
$A_{\Lambda}$	$p\overline{p} \to \Lambda X,  p\overline{p} \to \overline{\Lambda}X$	$-0.02 \pm 0.14$	R608	1985
$A_\Lambda$	$e^+e^- \to {\rm J}/\psi \to \Lambda\overline{\Lambda}$	$0.01\pm0.10$	DM2	1988
$A_{\Lambda}$	$p\overline{p} \to \Lambda\overline{\Lambda}$	$0.010 \pm 0.022$	PS185	1998

• There is a recent measurement of  $A_{\Xi\Lambda}$ , based on the HyperCP technique.

Mode	Method	Limit	Exp	Date
$A_{\Xi\Lambda}$	pN $\rightarrow \Xi \pi \rightarrow \Lambda \pi \rightarrow p \pi \pi$	$0.012 \pm 0.014$	E756	2000

- This measurement of  $A_{\Xi\Lambda}$  can be used with measurements of  $A_{\Lambda}$  to infer a limit on  $A_{\Xi}$ .
- None of these measurements is in the regime of testing theory.
- HyperCP is pushing two orders of magnitude beyond the best limit, to  $\sim 10^{-4}$ .



- 800 GeV/c incident proton beam.
- $\bullet$  10–15 MHz, 167 GeV/c secondary beam.
- Low-mass, high-rate, narrow-pitch wire chambers.
- Very high-rate DAQ:
  - $\Rightarrow$  50-80 KHz evts/spill-s to tape.
  - $\Rightarrow$  27 MB/s on 27 Exabyte 8705 tape drives.

• Simple, low-bias trigger using hodoscopes and calorimeter.

 $\mathrm{SS}(\geq\!\!1~\mathrm{hit}){\cdot}\mathrm{OS}(\geq\!\!1~\mathrm{hit}){\cdot}\mathrm{Cal}(\geq\!\!40~\mathrm{G~eV})$ 

• Muon system for rare and forbidden hyperon and kaon decays.

Craig Dukes			Daphne 2004		University of Virgini
		king we record	0	<b>ds</b> est data samples ev nd 119.5 TB data.	er by a particle
	E	vents		x 10 <sup>5</sup> 1400	
		ear		1200 -	$\Xi^-$ , 2.07 billion $\Xi^+$ , 0.46 billion $\sigma = 1.6 \text{ MeV/c}^2$
Trigger	1997	1999	Total		
Cascade	$39 \times 10^9$	$81 \times 10^{9}$	$120 \times 10^9$	Λπ <sup>±</sup>	
All	$58 \times 10^9$	$173 \times 10^{9}$	$231 \times 10^9$	t set t	
	Reconstru	ucted Events	3	++ 200 [I] 0	
	Channeled b	peam polarity			1.320 1.325 1.330 1.335 triant mass of $p\pi\pi$ (GeV/c <sup>2</sup> )
Type	+	_	Total	x 10 <sup>2</sup> 10000	
$\Xi \to \Lambda \pi$	$458 \times 10^{6}$	$2032 \times 10^{6}$	$2490 \times 10^{6}$	2000	$\Omega^{-}$ , 14.6 million $\overline{\Omega}^{+}$ , 5.0 million
$K \to \pi \pi \pi$	$391 \times 10^{6}$	$164 \times 10^{6}$	$555 \times 10^{6}$		$\sigma = 1.5 \text{ MeV/c}^2$
$\Omega \to \Lambda K$	$4.9 \times 10^{6}$	$14.1 \times 10^{6}$	$19.0 \times 10^{6}$	$\xrightarrow{\text{Events / 0.24 MeV/c}^2} MK^{\pm}$	P A
				2000	
					A A A A A A A A A A A A A A A A A A A
				1.660 1.665	5 1.670 1.675 1.680 iant mass of Kp $\pi$ (GeV/c <sup>2</sup> )
				Invar	Tant mass of Kph (GeV/C)

### Care Taken to Mimimize Differences in + and – Running

- Target length changed to equalize channeled beam rates.
  - + polarity: 2.0 cm Cu
  - - polarity: 6.0 cm Cu



- When flipping polarity, field magnitude kept within  $\sim 2 \times 10^{-4}$ .
- This corresponds to a  $\sim 0.3$  mm deflection difference at 10 m for the lowest momentum ( $\sim 10 \text{ GeV}/c \text{ pions}$ ).





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### Two Different CP Analyses Being Done

### West Coast (LBL): Hybrid Monte Carlo Method

- Compare corrected  $\cos \theta$  distributions.
- Take a real Ξ → Λπ, Λ → pπ event, discard proton and pion, generate 10 new unpolarized Λ decays.
- Advantage: Absolute measurement of  $\alpha_{\Lambda}\alpha_{\Xi}$ .
- Disadvantage: Monte Carlo must be very, very good, and fast: ~20 billion events needed.

### East Coast (UVa): Weighting Method

- Compare uncorrected  $\cos(\theta)$  distributions.
- Force the Ξ<sup>-</sup> and Ξ<sup>+</sup> events to have similar momentum and spatial distributions by appropriate weighting.
- Advantage: No Monte Carlo needed to measure apparatus acceptance, smaller statistical error.
- Disadvantage: inflexible, event-size dependent analysis.



## Weighting Technique

x 10<sup>3</sup>

4000

2000

1500

1000

- Problem: Geometical acceptance identical for  $\Xi^-$  and  $\overline{\Xi}^+$  decay products only if parent  $\Xi^-$ → 4000 3500 3000 and  $\overline{\Xi}^+$  have same momentum and inhabit the **v** 2500 same phase space exiting the collimator.
- Solution: Weight the  $\Xi^-$  and  $\overline{\Xi}^+$  events to force the two distributions to be identical.



# $\Xi^-$ and $\overline{\Xi}^+$ x Slopes and Positions not Weighted

- Distributions almost identical  $\Rightarrow$  cut out regions where they are not.
- $\Xi^-$ : Solid lines
- $\overline{\Xi}^+$ : Dashed lines



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# Extracting the CP Asymmetry

• Determine weighted proton and weighted antiproton  $\cos\theta$  distributions.

$$\frac{dN_{-}}{d\cos\theta_{-}} = A_{-}\frac{N_{-}}{2}(1+\alpha_{\Xi}\alpha_{\Lambda}\cos\theta_{-})$$

• Take the ratio of  $\cos \theta$  distributions and fit to:

 $R(\theta, \delta) = C \frac{1 + \alpha_{\Xi} \alpha_{\Lambda} \cos \theta}{1 + (\alpha_{\Xi} \alpha_{\Lambda} - \delta) \cos \theta}$ 

to extract asymmetry  $\delta$ :

$$\delta = \alpha_{\Xi} \alpha_{\Lambda} - \overline{\alpha}_{\Xi} \overline{\alpha}_{\Lambda}$$
$$A_{\Xi\Lambda} = \frac{\delta}{\alpha_{\Xi} \alpha_{\Lambda} + \overline{\alpha}_{\Xi} \overline{\alpha}_{\Lambda}} \approx \frac{\delta}{2\alpha_{\Xi} \alpha_{\Lambda}}$$

• Note: acceptances cancel out!



 $\frac{dN_+}{d\cos\theta_+} = A_+ \frac{N_+}{2} (1 + \overline{\alpha}_{\Xi} \overline{\alpha}_{\Lambda} \cos\theta_+)$ 



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### The CP Asymmetry $A_{\Xi\Lambda}$ from Weighting Method

- Data broken up into 18 sets of positive and negative data.
- No acceptance corrections.
- No efficiency corrections.

from Analysis Set 1

• No background subtraction.







Weighted average of all 18 data sets:

$$\delta = (-1.3 \pm 3.0) \times 10^{-4}$$
  
 $A_{\Xi\Lambda} = (2.2 \pm 5.1) \times 10^{-4}$   
 $\chi^2 = 24$ 

### **Background Subtraction Has Little Effect**

- Triple Gaussian fit with fourth-order polynomial for background.
- Background fraction:
  - $\Xi^{-}: 0.43\%$  (lines)
  - $\overline{\Xi}^+$ : 0.41% (circles)



Low mass:  $\delta = (-2.2 \pm 0.5) \times 10^{-2}$ High mass:  $\delta = (-3.8 \pm 0.7) \times 10^{-2}$ 

• Weighted background asymmetry:

$$A_{\Xi\Lambda}(bs) = (0.0 \pm 5.1) \times 10^{-4}$$



### Helicity Frame Analysis Naturally Minimizes Biases

 The helicity frame axes changes from event to event since we always define the polar axis to be the direction of the Λ momentum in the Ξ rest frame.



• Acceptance differences localized in a particular part of the apparatus do **not** map into a particular part of the proton (antiproton)  $\cos \theta$  distribution.

Important! Overall acceptance differences do not cause any biases.

# Weighted Analysis Bias Error Summary

Systematic	Method	$\delta A_{\Xi\Lambda}(10^{-4})$
Analyzing Magnets field uncertainties	Data	2.8
Calorimeter inefficiency uncertainty	Data	2.1
Validation of analysis code	CHMC	1.9
Collimator exit $x$ slope cut	Data	1.4
Collimator exit $x$ position cut	Data	1.2
PWC inefficiency uncertainty	CHMC	1.0
Hodoscope inefficiency uncertainty	Data	0.3
Particle/antiparitle interaction differences	MC	0.9
Momentum weights bin size	Data	0.4
Background subtraction uncertainty	Data	0.3
Error on $\alpha \alpha_{PDG}$	Data	0.03
Polarization	MC	negligible
Earth's magnetic field	CHMC	negligible
Total systematic error		4.2

### **Results from** CP Violation Search

### Weighting Technique:

- $\sim 10\%$  total data sample
- $\bullet$  selected from end of 1999 run
- 118.6 million  $\Xi^-$
- 41.9 million  $\overline{\Xi}^+$
- no acceptance or efficiency corrections

 $A_{\Xi\Lambda} = [0.0 \pm 5.1 (\text{stat}) \pm 4.2 (\text{syst})] \times 10^{-4}$ 

### Check with HMC Technique:

- ~ 5% of the total data sample
- $\bullet$  prescaled selection of 1997 and 1999
- 15 million  $\Xi^-$
- 30 million  $\overline{\Xi}^+$

 $A_{\Xi\Lambda} = [-7 \pm 12(\text{stat}) \pm 6.2(\text{syst})] \times 10^{-4}$ 

 $\Rightarrow 20 \times$  improvement on previous result.



## Measuring the $\Lambda$ - $\pi$ Strong Phase Shift

 $\vec{p}_{\text{beam}}$ 

- This is done by analyzing the  $\Lambda$  decay distribution from **polarized**  $\Xi^{-}$ 's.
- Polarized  $\Xi^{-}$ 's are produced by targeting at nonzero angles.
- A polarized  $\Xi^-$  decay produces a  $\Lambda$ with three components of polarization:

$$\vec{P}_{\Lambda} = \frac{(\alpha_{\Xi} + \vec{P}_{\Xi} \cdot \hat{p}_{\Lambda})\hat{p}_{\Lambda} + \beta_{\Xi}(\vec{P}_{\Xi} \times \hat{p}_{\Lambda}) + \gamma_{\Xi}(\hat{p}_{\Lambda} \times (\vec{P}_{\Xi} \times \hat{p}_{\Lambda}))}{(1 + \alpha_{\Xi}\vec{P}_{\Xi} \cdot \hat{p}_{\Lambda})}$$

where:

$$\alpha = \frac{2\text{Re}(S^*P)}{|S|^2 + |P|^2} \ \beta = \frac{2\text{Im}(S^*P)}{|S|^2 + |P|^2} \ \gamma = \frac{|S|^2 - |P|^2}{|S|^2 + |P|^2}$$

• Measuring  $\beta_{\Xi}$  allows the phase shift to be extracted:

$$\frac{\beta_{\Xi}}{\alpha_{\Xi}} = \tan(\delta_P - \delta_S) \approx \delta_P - \delta_S$$



### Measurement of the $\Lambda$ - $\pi$ Phase Shift



 $\beta_{\Xi} = -0.037 \pm 0.011 (\text{stat}) \pm 0.010 (\text{syst})$ 

- $\gamma_{\Xi} = 0.888 \pm 0.0004 (\text{stat}) \pm 0.006 (\text{syst})$
- Using the known value of  $\alpha_{\Xi}$ (-0.456±0.008), the strong phase shift is:

$$\delta_P - \delta_S = \tan^{-1} \left( \frac{\beta_{\Xi}}{\alpha_{\Xi}} \right) = (4.6 \pm 1.4 \pm 1.2)^\circ$$

- First non-zero measurement of phaseshift.
- This is about the same magnitude as the p- $\pi$  phase shift:
  - $\Rightarrow$  CP equally likely to be seen in  $\Xi \to \Lambda \pi$  decays.
  - $\Rightarrow$  CP predictions underestimated,
  - $\Rightarrow \chi \text{PT}$  calculations off.



### Search for Parity Violation in $\Omega^- \to \Lambda K^-$ Decays

 $\Omega^- \to \Lambda K^- \qquad \Lambda \to p\pi^-$ 

• Although spin-3/2,  $\Omega^- \to \Lambda K^-$  decay goes much like the other hyperon two-body decays:

$$\frac{dP}{d\cos\theta} = \frac{1}{2}(1 + \alpha_{\Omega}P_{\Omega}\cos\theta)$$

• Here:

$$\alpha_{\Omega} = \frac{2\text{Re}(P^*D)}{|P|^2 + |D|^2}$$

- A non-zero  $\alpha_{\Omega}$  indicates parity violation.
- All other hyperons have non-zero α parameters; only the Ω<sup>-</sup> has resisted efforts to find an asymmetrical decay distribution.
- HyperCP is measuring  $\alpha_{\Omega}$  using unpolarized  $\Omega^{-}$ 's through the polarization given to the daughter  $\Lambda$ , which is  $\alpha_{\Omega}$ :

$$\frac{dP}{d\cos\theta} = \frac{1}{2}(1 + \alpha_{\Omega}\alpha_{\Lambda}\cos\theta)$$

• Large data sample, little background.



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# **Conclusions and Outlook**

- Analyzing by far the largest sample of hyperons ever recorded, we have assulted *CP* violation from a different direction than the kaon and B experiments.
- We find no evidence of CP violation in  $\Xi^{\pm}$  and  $\Lambda$  decays:

$$\delta A_{\Xi\Lambda} = (0.0 \pm 5.1 \pm 4.2) \times 10^{-4}$$

- This result is constraining some SUSY models of *CP* violation.
- Shortly we will push our statistical limit to HyperCP  $\delta A_{\Xi\Lambda} \approx 2 \times 10^{-4}$  Limit

two orders of magnitude better than the present limit.

- We have the first evidence of parity violation in  $\Omega^- \to \Lambda K^-$  decays.
- A preliminary analysis shows no evidence of CP violation in  $\Omega^- \to \Lambda K^- / \overline{\Omega}^+ \to \overline{\Lambda} K^+$  decays.

